



# Theory of Multilayer Coating for proof-of-concept experiments

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**KEK**  High Energy Accelerator Research Organization, Tsukuba, Japan

# KUBO, Takayuki

久保

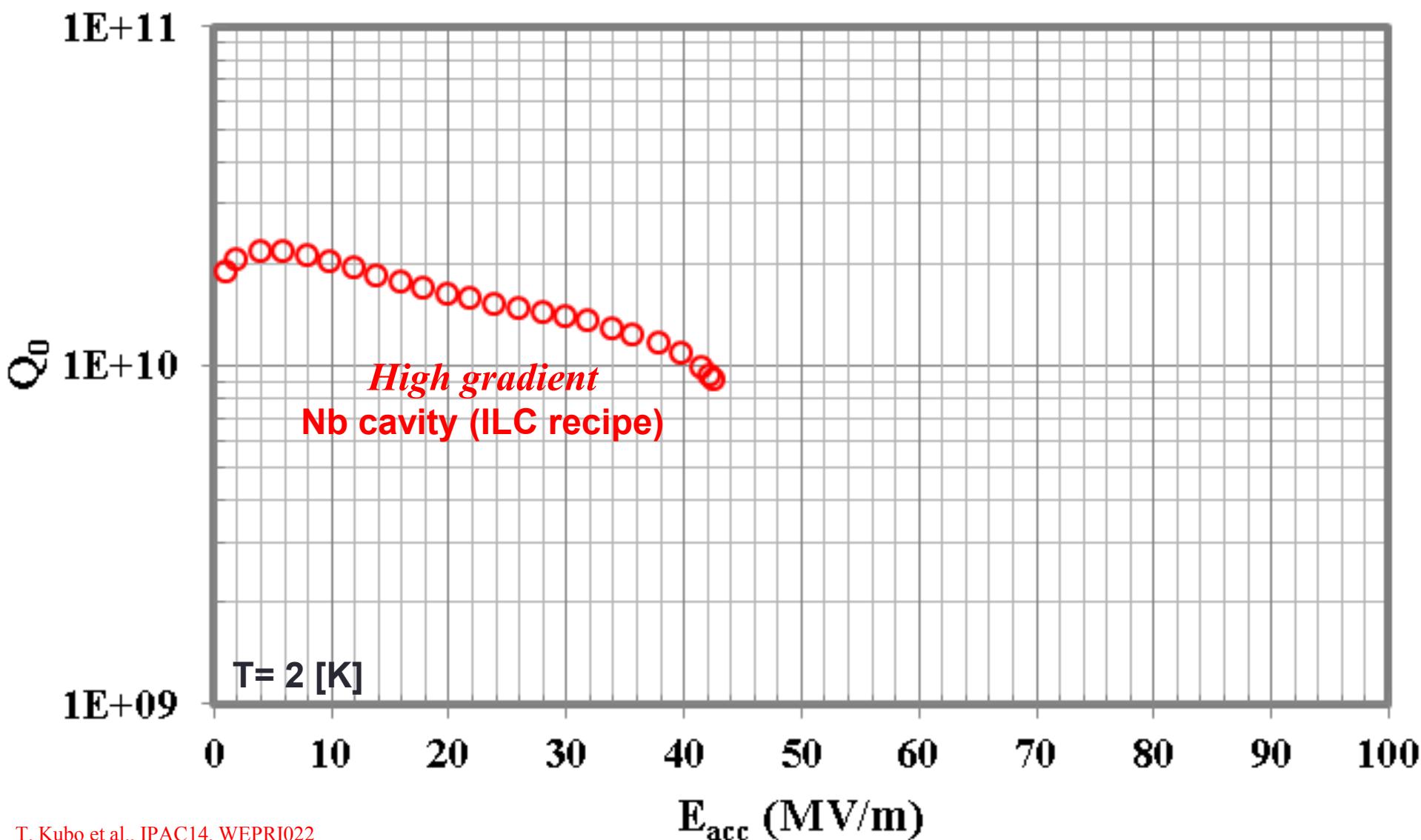
毅幸

<http://researchmap.jp/kubotaka/>

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Grant-in-Aid for Young Scientists (B), Number 26800157,  
Grant-in-Aid for Challenging Exploratory Research, Number 26600142,  
and Photon and Quantum Basic Research Coordinated Development Program from the Ministry of Education, Culture, Sports, Science and Technology, Japan.

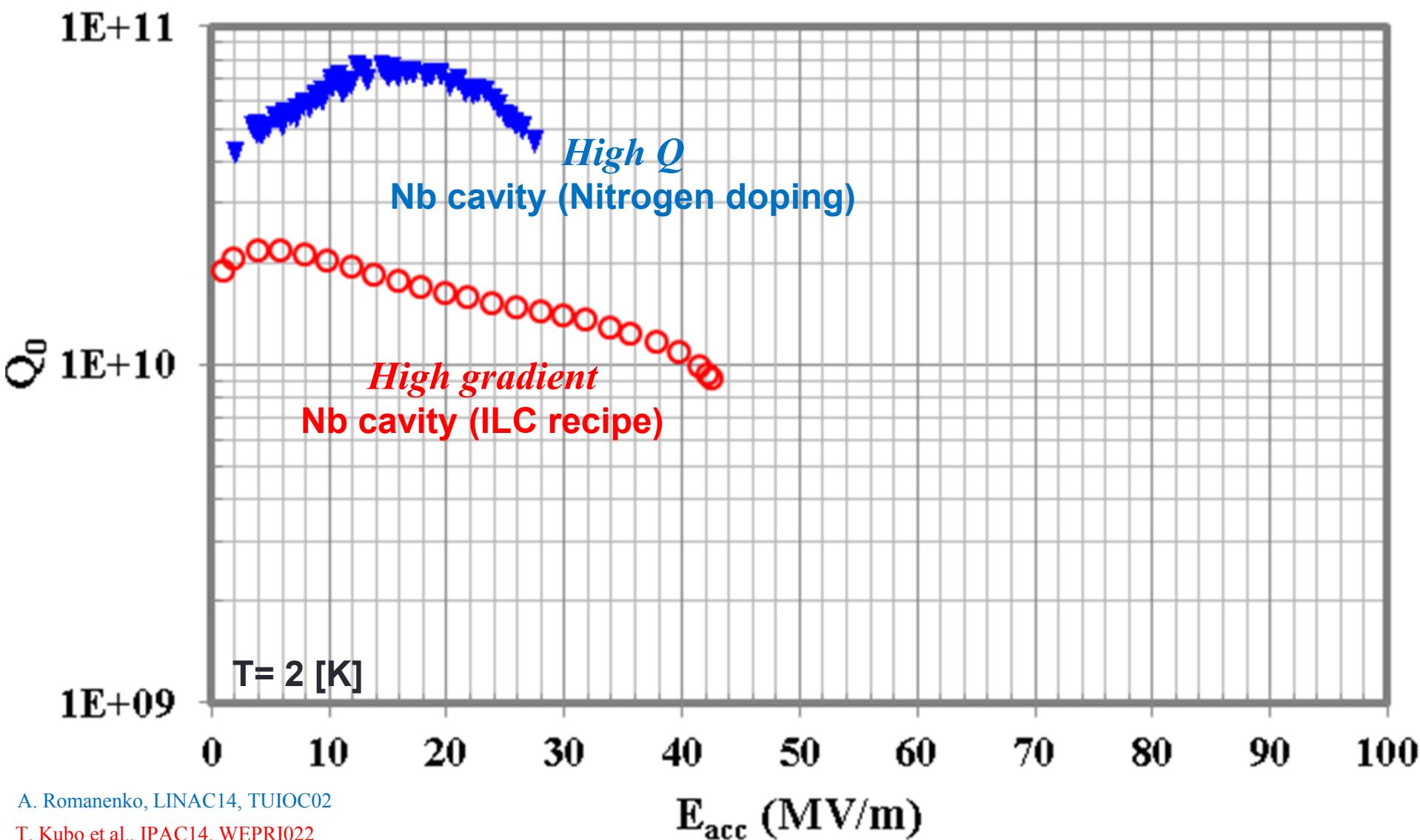


# Nb cavity processed by the ILC recipe

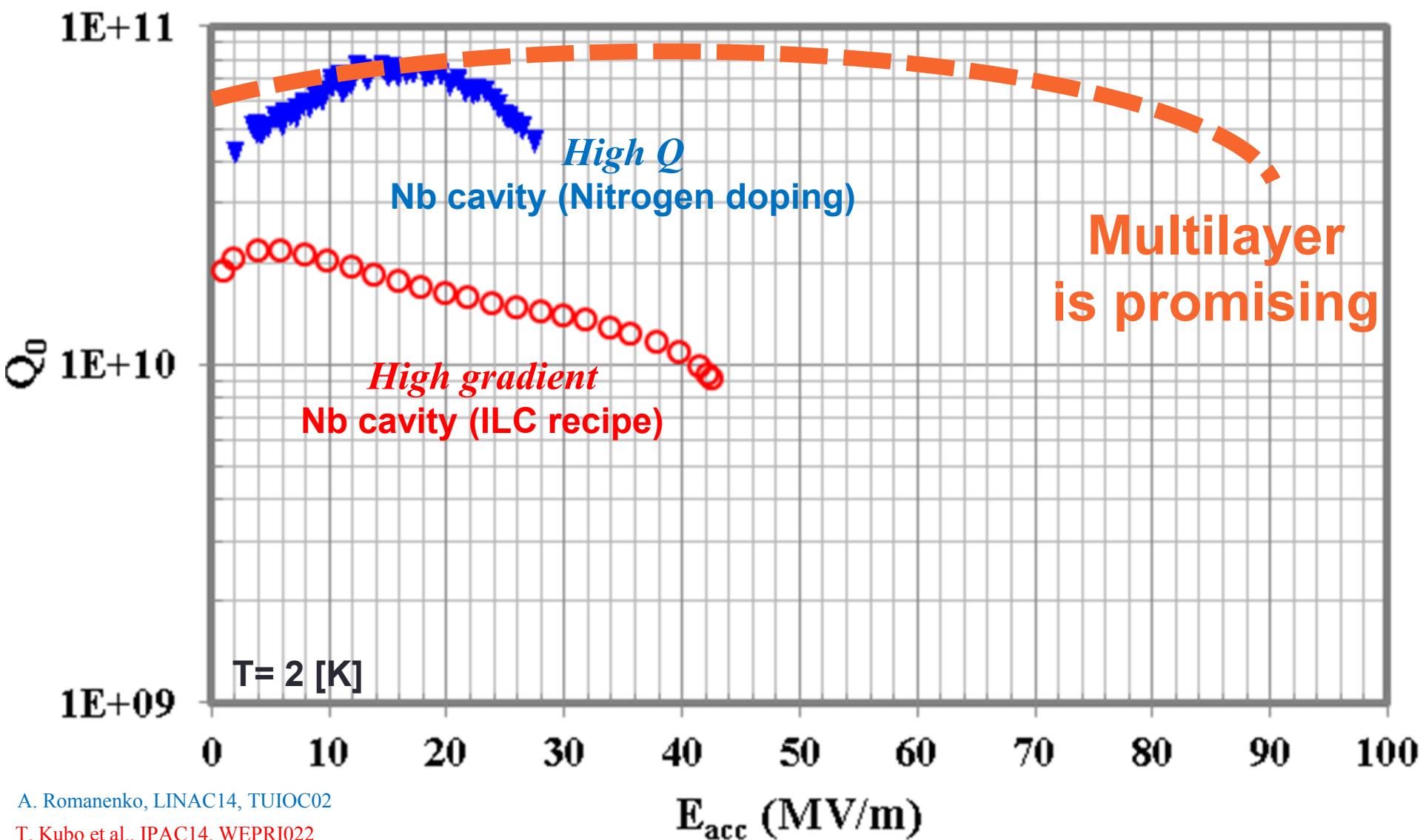


# Breakthrough by the nitrogen doping

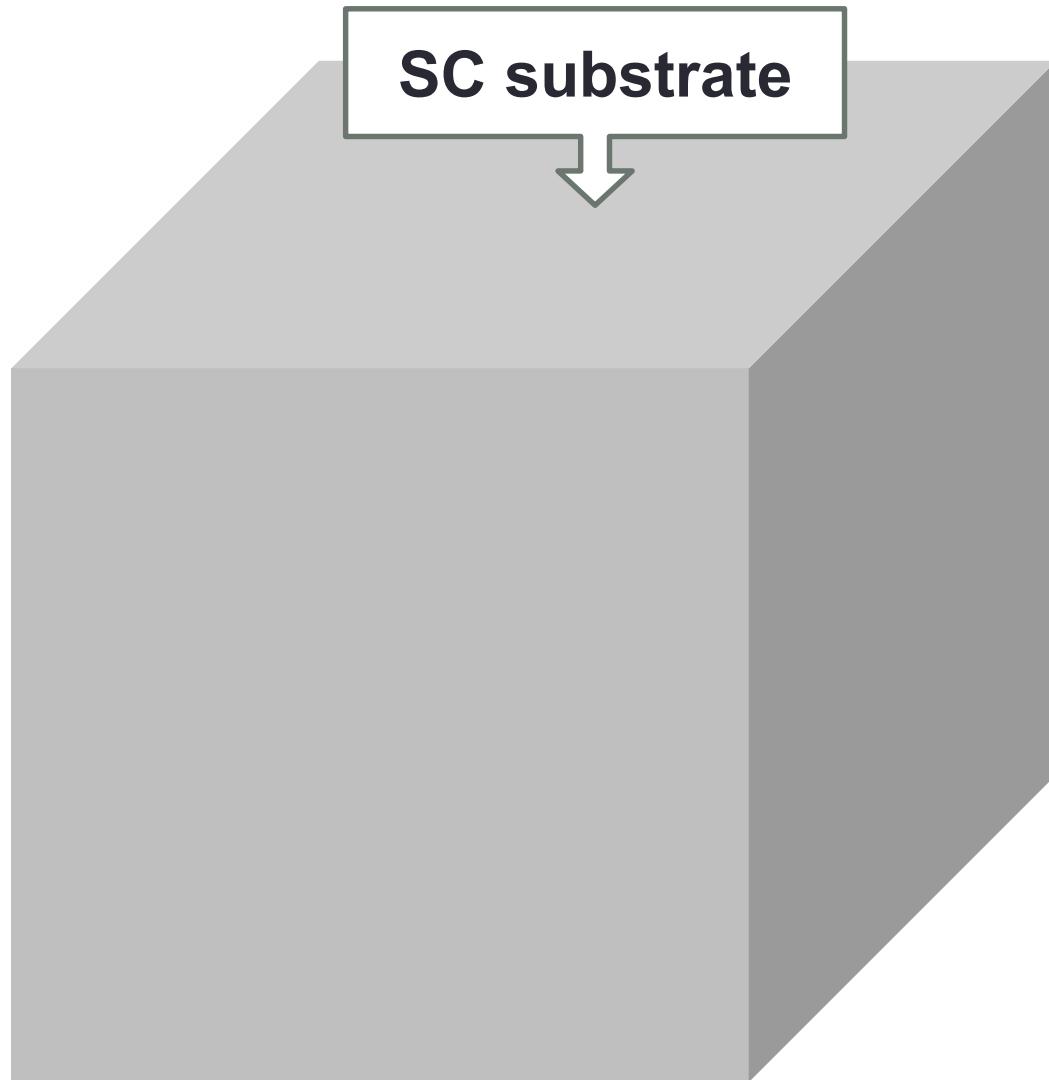
A. Grassellino et al, Supercond. Sci. Technol. **26**, 102001 (2013)



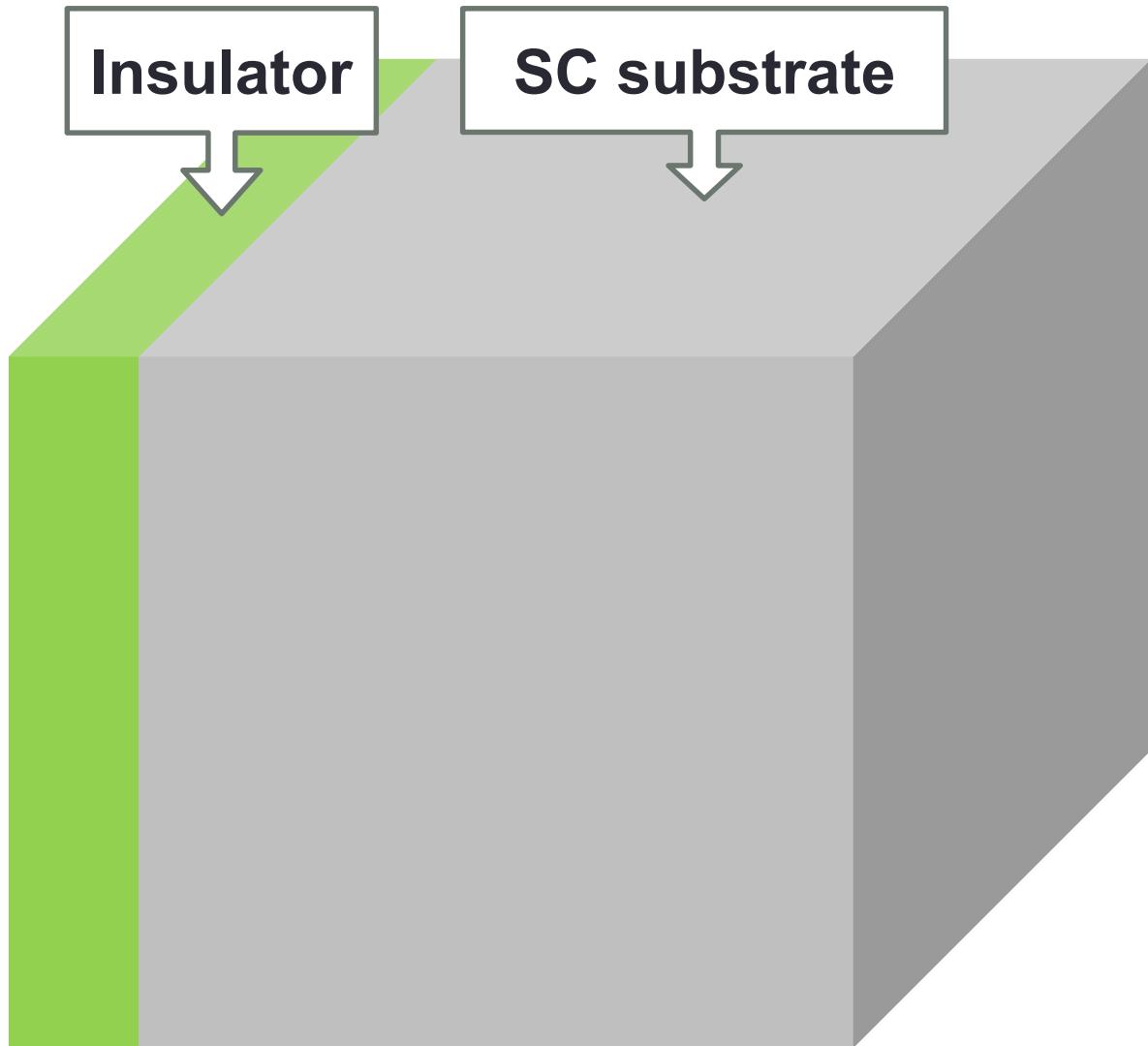
# We want to go beyond Nb!



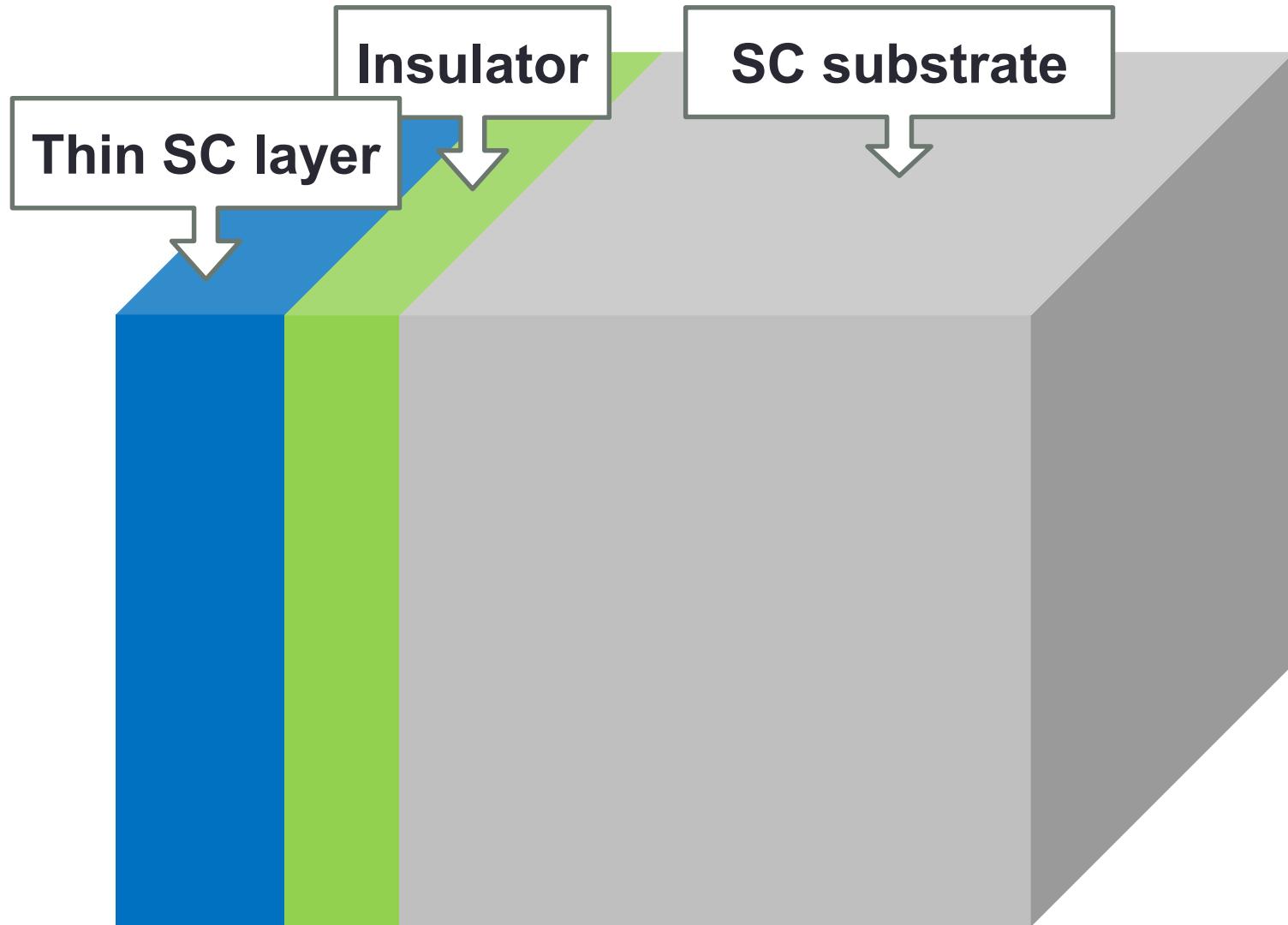
The multilayer coating was proposed by A. Gurevich, [Appl. Phys. Lett. 88, 012511 \(2006\)](#)



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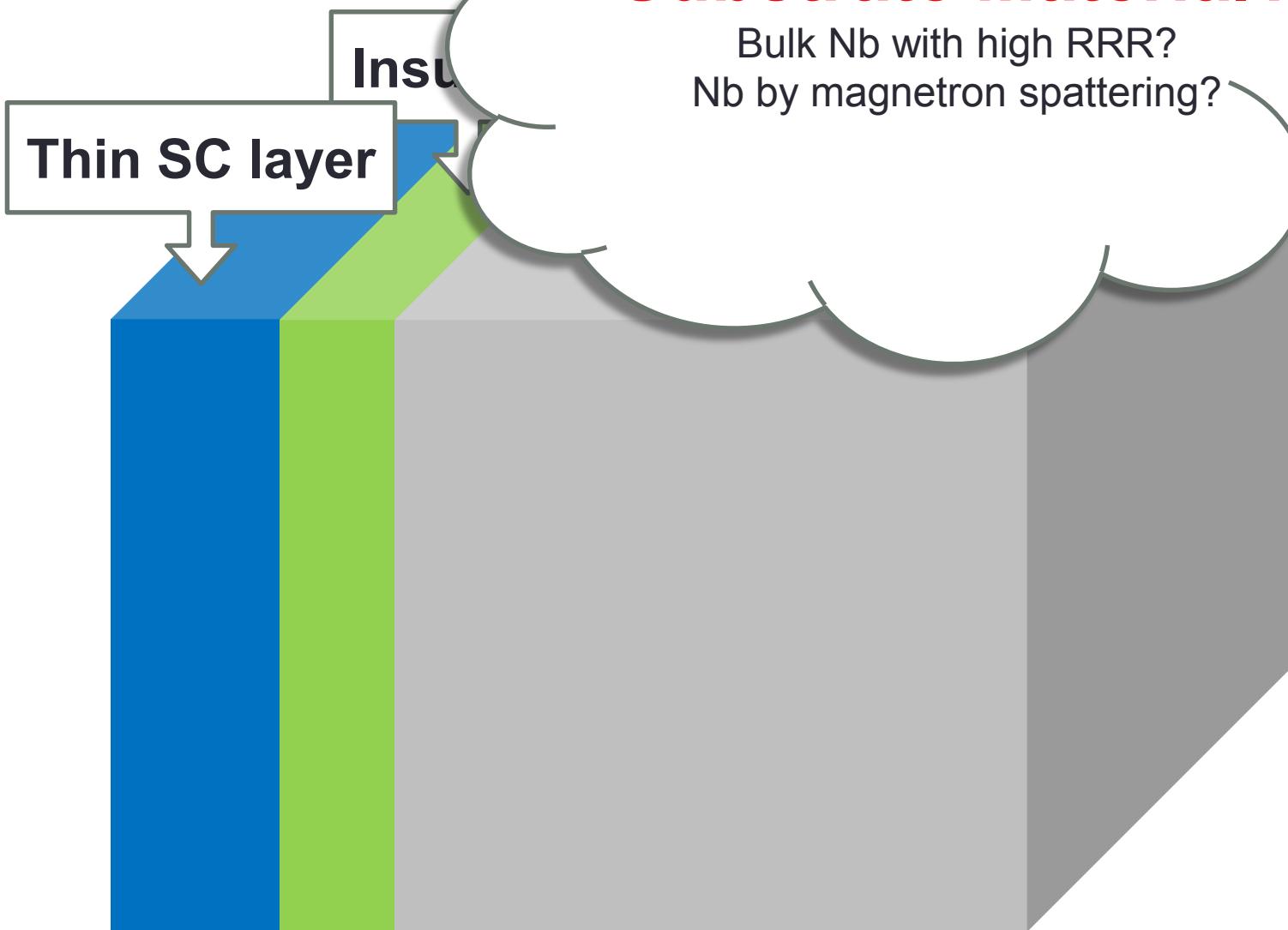


The multilayer coating is proposed by A

## Substrate Material?

Bulk Nb with high RRR?

Nb by magnetron sputtering?



The n

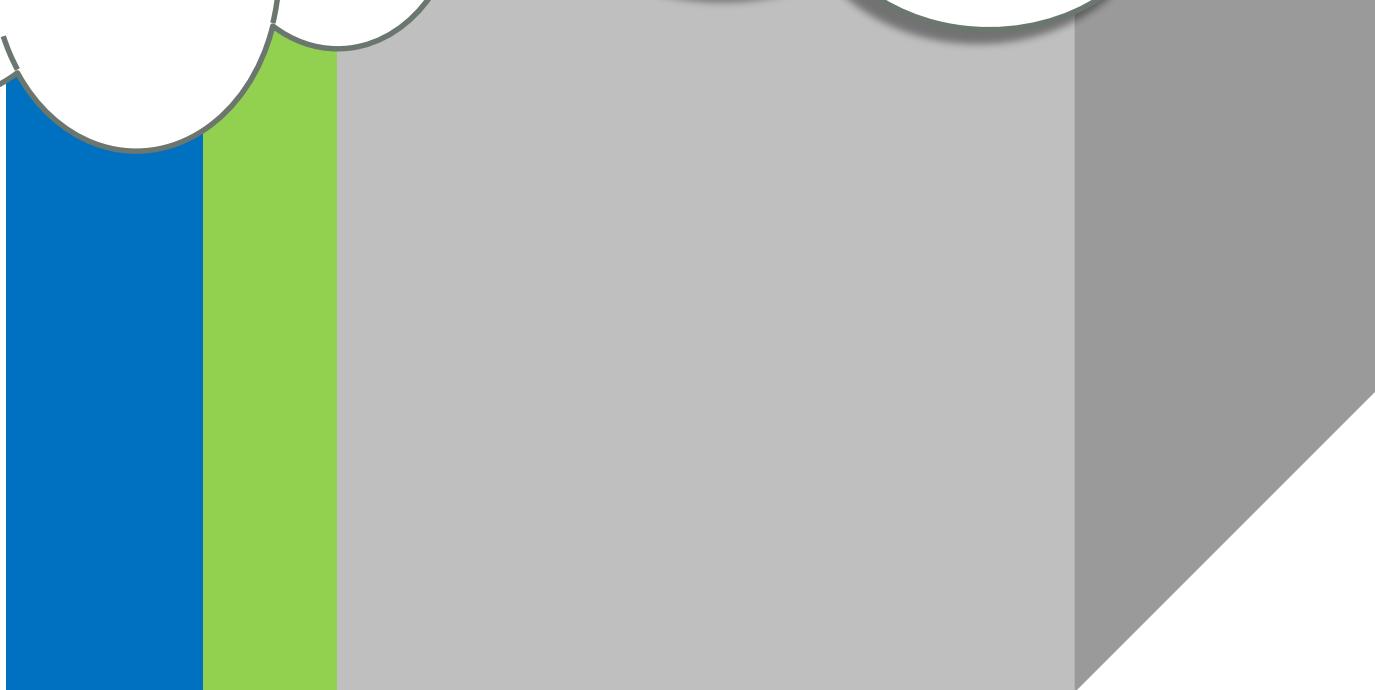
**How to fix  
insulator  
thickness  $d_I$  ?**

10nm?

100nm?

**Substrate Material?**

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Nb by magnetron sputtering?



The n

**How to fix  
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10nm?

100nm?

**Substrate Material?**

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Nb by magnetron sputtering?

**How to fix  $S$  layer  
thickness  $d_S$  ?**

10nm?

100nm?



## § 1

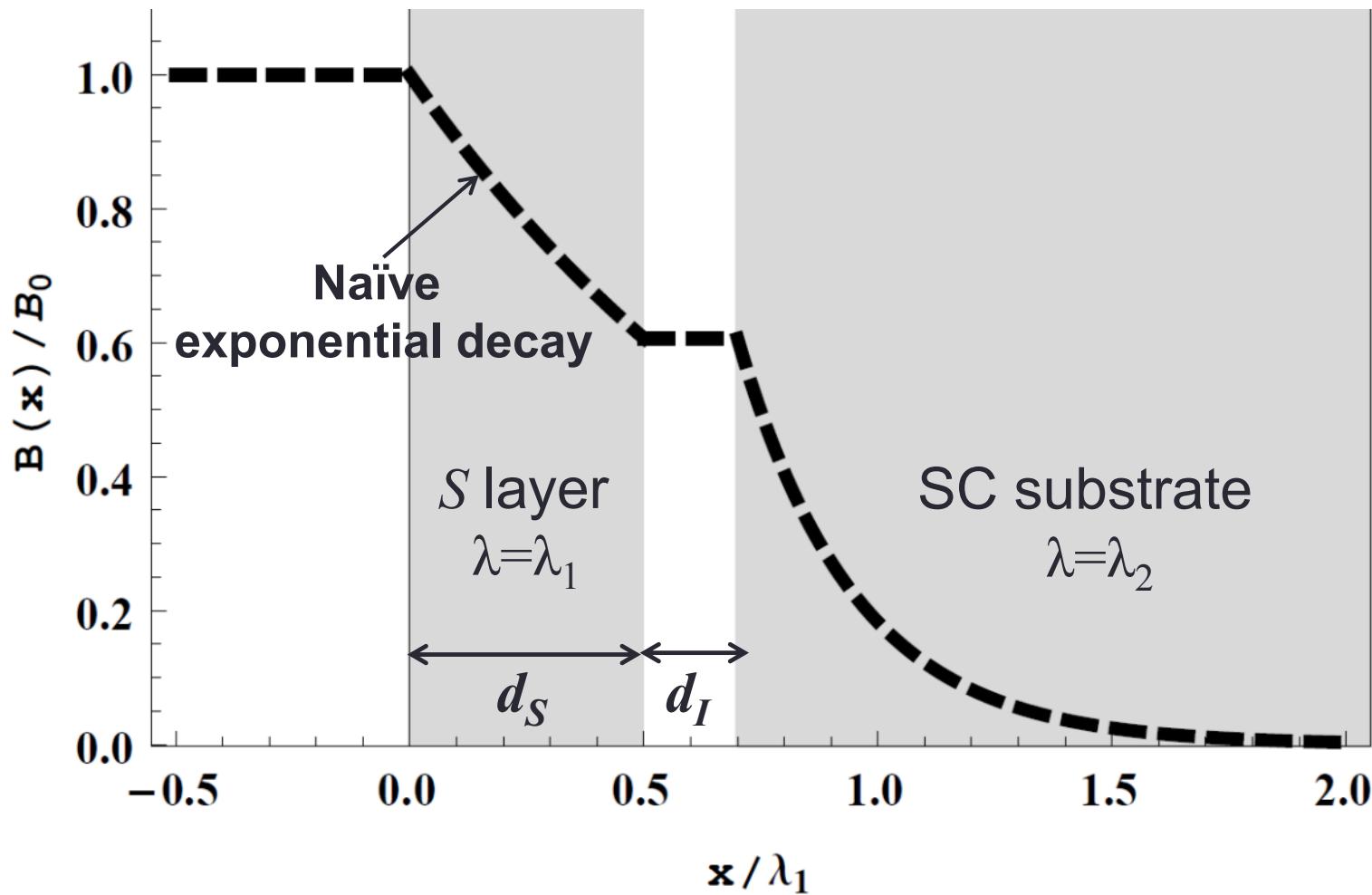
*The optimum parameters*

1. The magnetic field distribution (and thus the screening current distribution  $\mathbf{J} \propto d\mathbf{B}/dx$ ) in the  $S$  layer is different from the naïve exponential decay.
2. When  $d_S$  and  $d_I$  are thin enough and  $\lambda_1 > \lambda_2$ , the screening current in the  $S$  layer is suppressed, and the surface field can exceed the superheating field of the  $S$  layer.
3. However, an extremely thin  $d_S$  can not protect the SC substrate. Thus the  $S$  layer must have some thickness to decay the magnetic field and protect the SC substrate. The optimum thickness of  $d_S$  exists.

1. **T. Kubo** et al., Appl. Phys. Lett. **104**, 032603 (2014) [submitted to arXiv on April 2013; published on January 2014]
2. **S. Posen** et al., in proceedings of SRF2013, p. 788, WEIOC04 [Sep.2013].
3. **A. Gurevich**, AIP Advances **5**, 017112 (2015) [submitted on Sep.2014; published on Jan.2015]
4. **S. Posen** et al., arXiv:1506.08428v1[June 2015].

## Important!

1. The magnetic field distribution (and thus the screening current distribution  $J \propto dB/dx$ ) in the S layer is different from the naïve exponential decay.



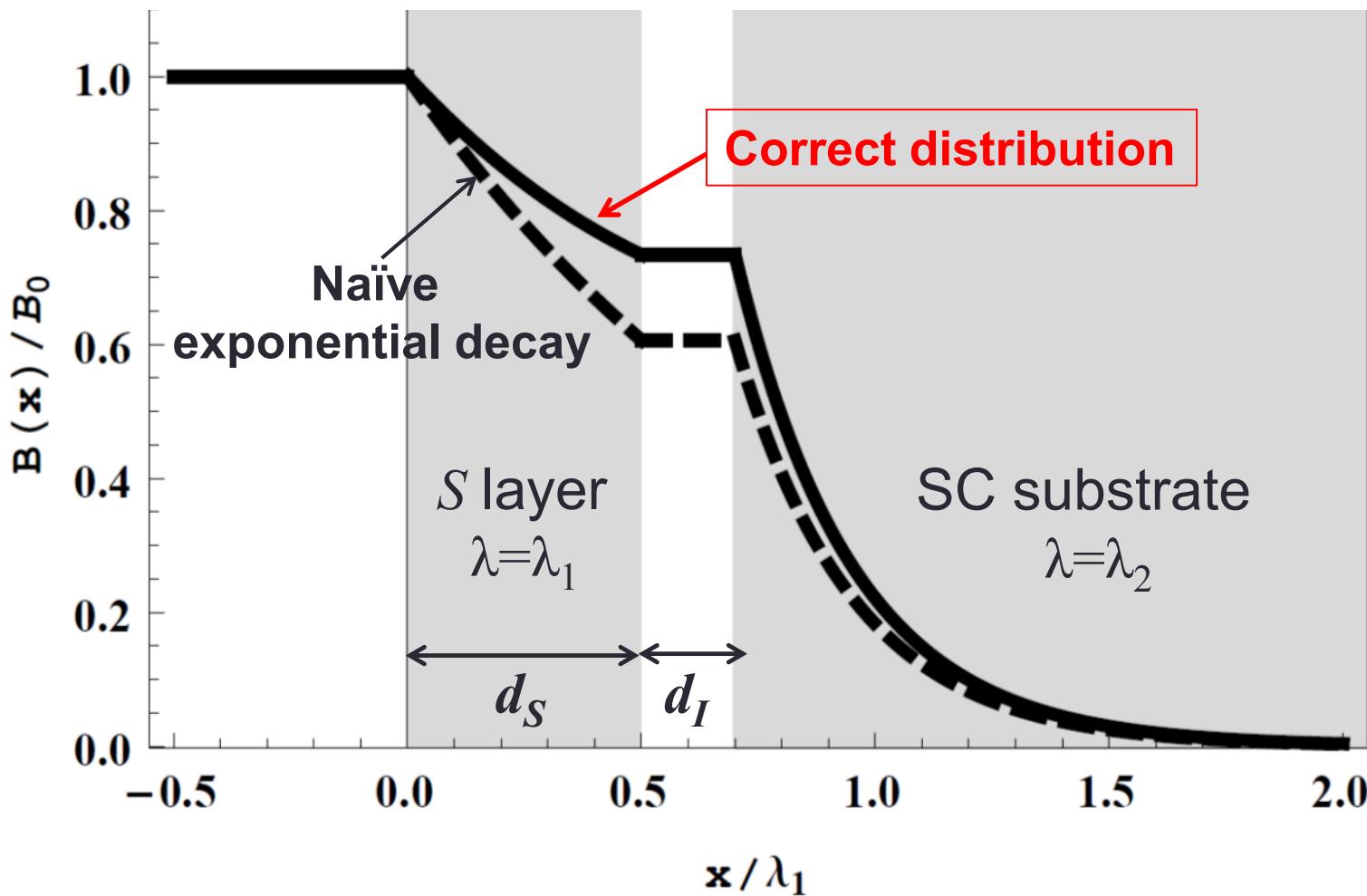
T. Kubo, Y. Iwashita, and T. Saeki, *Appl. Phys. Lett.* **104**, 032603 (2014)

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The derivation processes are explained in detail in proceedings of IPAC13, p. 2343, WEPWO014 [May 2013]

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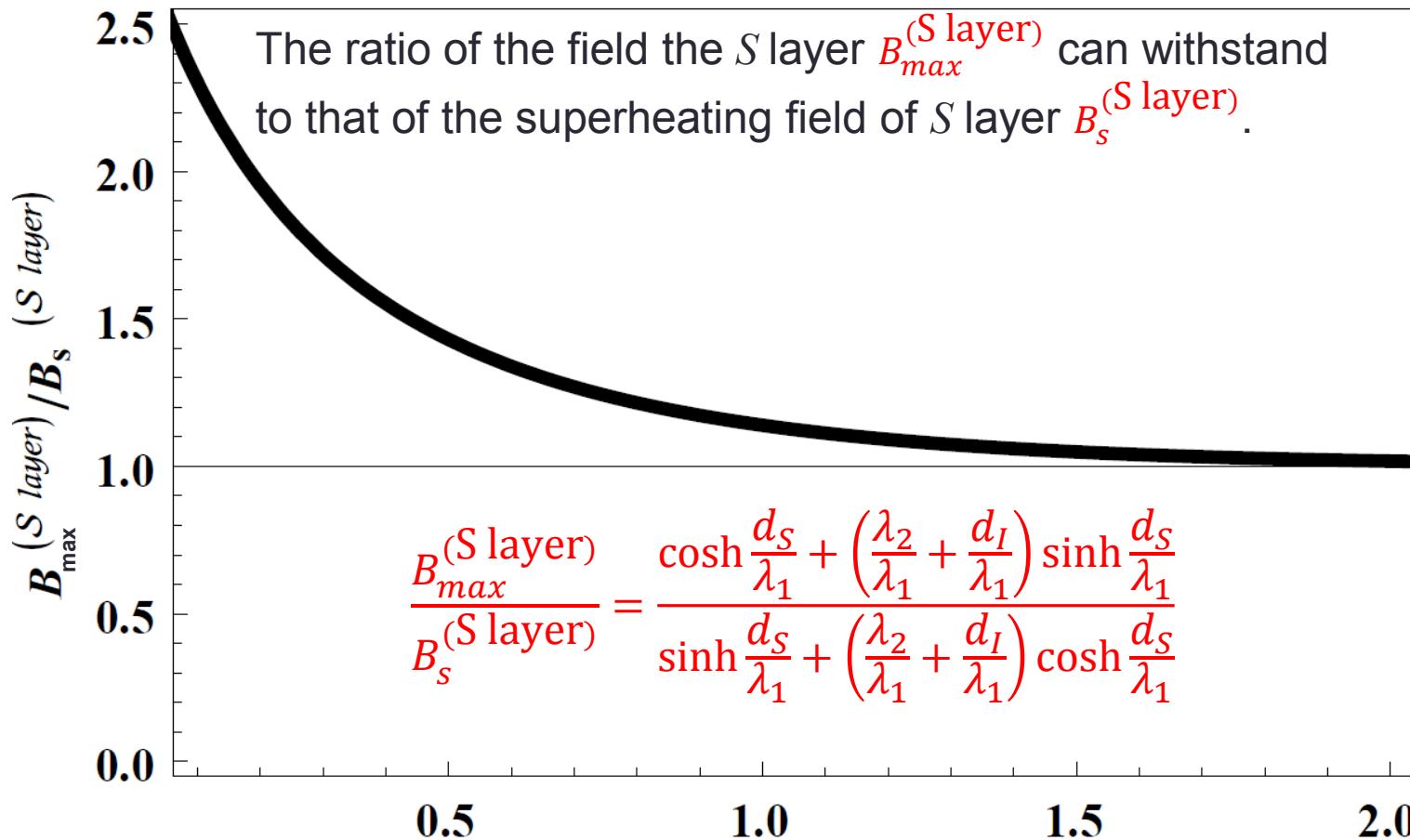


T. Kubo, Y. Iwashita, and T. Saeki, *Appl. Phys. Lett.* **104**, 032603 (2014)

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2. When  $d_S$  and  $d_I$  are thin enough and  $\lambda_1 > \lambda_2$ , the screening current in the  $S$  layer is suppressed, and the surface field can exceed superheating field of the  $S$  layer.



T. Kubo, Y. Iwashita, and T. Saeki, *Appl. Phys. Lett.* **104**, 032603 (2014)

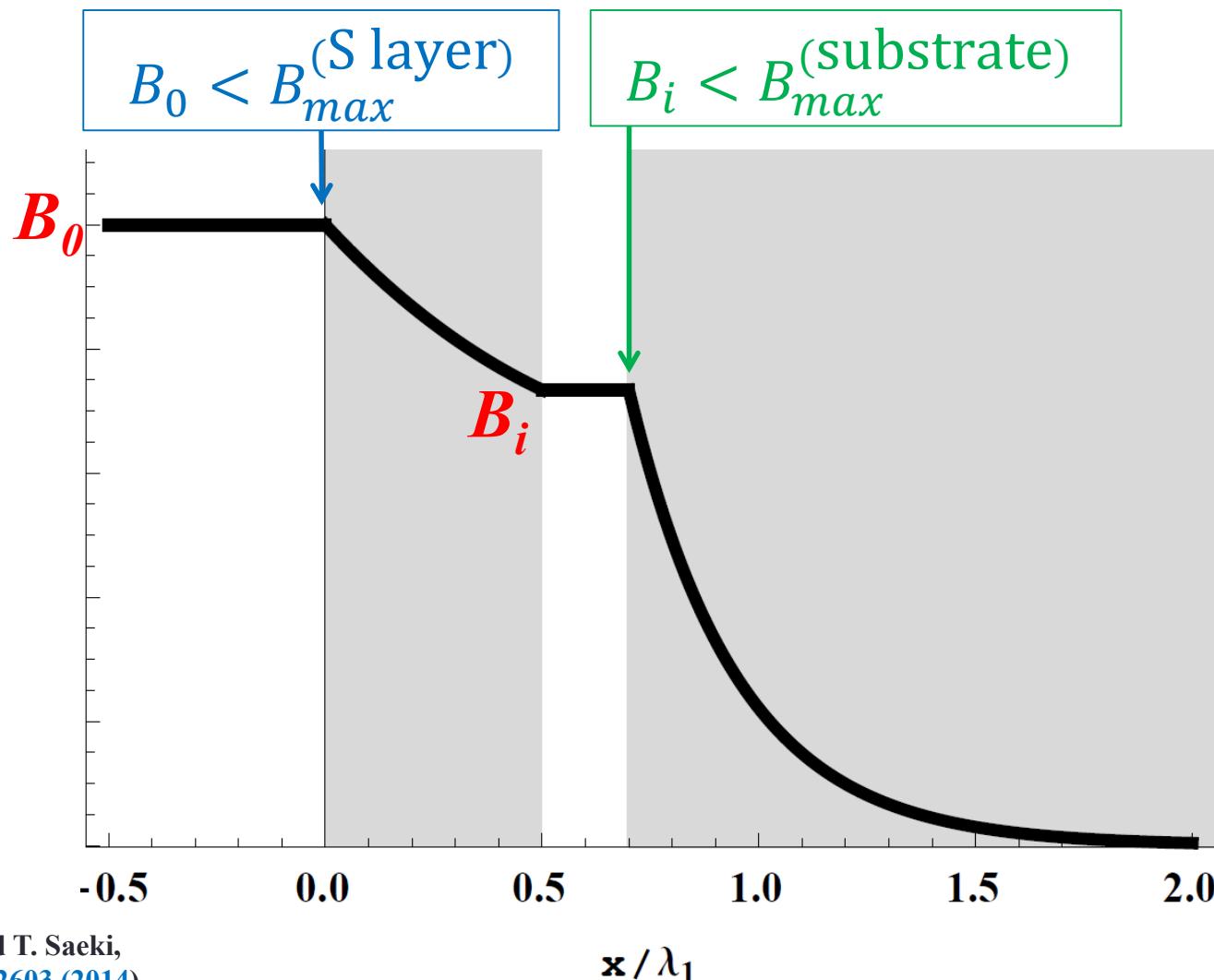
[submitted to arXiv on April 2013; published on January 2014]

The derivation processes are explained in detail in proceedings of SRF2013, p. 430, TUP007 [Sep. 2013]

$$\frac{d_I}{\lambda_1} = 0.1 \quad \frac{\lambda_2}{\lambda_1} = 0.25$$

are assumed here.

3. However, an extremely thin  $d_S$  can not protect the SC substrate. Thus the *S* layer must have some thickness to decay the magnetic field and protect the SC substrate.



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## Formula for the maximum screening field of the multilayer

$$B_{\max}^{(\text{multilayer})} = \begin{cases} B_{\max}^{(\text{S layer})} & (\text{if } \gamma B_{\max}^{(\text{S layer})} < B_{\max}^{(\text{substrate})}) \\ \gamma^{-1} \times B_{\max}^{(\text{substrate})} & (\text{if } \gamma B_{\max}^{(\text{S layer})} \geq B_{\max}^{(\text{substrate})}) \end{cases}$$

where

$$\left. \begin{array}{l} B_{\max}^{(\text{S layer})} = B_s^{(\text{S layer})} \frac{\cosh \frac{d_S}{\lambda_1} + \left( \frac{\lambda_2}{\lambda_1} + \frac{d_I}{\lambda_1} \right) \sinh \frac{d_S}{\lambda_1}}{\sinh \frac{d_S}{\lambda_1} + \left( \frac{\lambda_2}{\lambda_1} + \frac{d_I}{\lambda_1} \right) \cosh \frac{d_S}{\lambda_1}} \quad B_s^{(\text{S layer})} = 0.84 B_c^{(\text{S layer})} \\ \gamma = \frac{1}{\cosh \frac{d_S}{\lambda_1} + \left( \frac{\lambda_2}{\lambda_1} + \frac{d_I}{\lambda_1} \right) \sinh \frac{d_S}{\lambda_1}} \\ B_{\max}^{(\text{substrate})} = 170\text{mT} - 240\text{mT} \quad (\text{if the substrate is Nb}) \end{array} \right\}$$

3. However, an extremely thin  $d_S$  can not protect the SC substrate. Thus **the S layer must have some thickness to decay the magnetic field** and protect the SC substrate.

## Formula for the optimum thickness of the S layer

$$d_S = \lambda_1 \ln \left[ \frac{\lambda_1}{\lambda_1 + \lambda_2 + d_I} \frac{B_s^{(\text{S layer})}}{B_{max}^{(\text{substrate})}} + \sqrt{\left( \frac{\lambda_1}{\lambda_1 + \lambda_2 + d_I} \frac{B_s^{(\text{S layer})}}{B_{max}^{(\text{substrate})}} \right)^2 + \frac{\lambda_1 - \lambda_2 - d_I}{\lambda_1 + \lambda_2 + d_I}} \right]_{d_I \lesssim 0(10) \text{ nm}}$$

**T. Kubo (2015),**

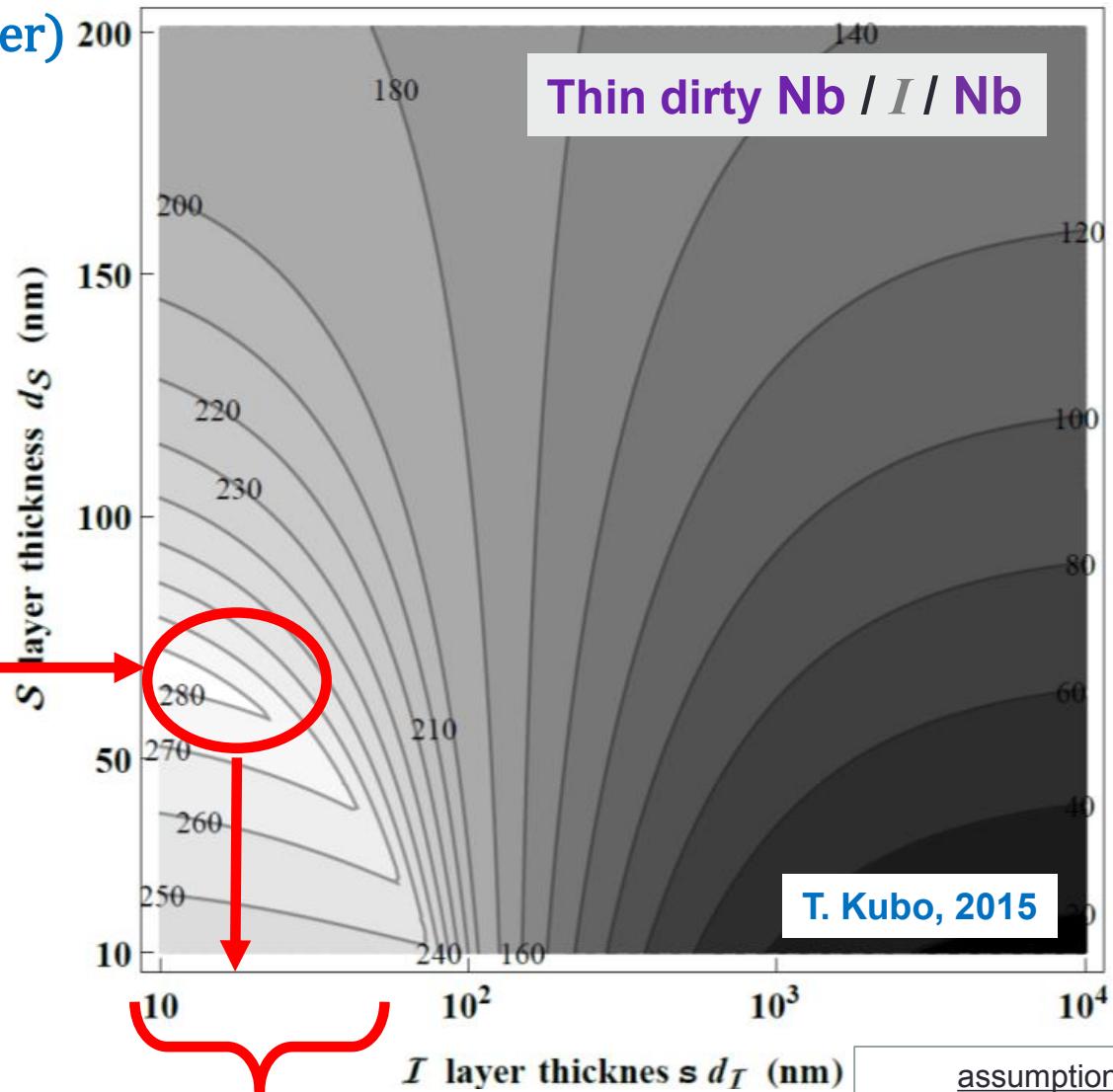
Obtained by using [T. Kubo et al., Appl. Phys. Lett. 104, 032603 \(2014\)](#) and [A. Gurevich, AIP Advances 5, 017112 \(2015\)](#)

- The formulae can be derived by using the discussion of A. Gurevich, AIP Advances 5, 017112 (2015) and are described by using the superheating field of the quasi-classical theory and thus valid even at  $T \ll T_c$ .
- The formulae are generalized version of the Gurevich's formulae. The formulae includes effects of insulator layer with a finite thickness. **When  $d_I \ll \lambda_1$ , the formulae are reduced to the Gurevich's formulae** [A. Gurevich, AIP Advances 5, 017112 (2015)].

## Contour plot of $B_{\max}^{(\text{multilayer})}$

**Optimum  $d_s \sim 70\text{nm}$**

see also A Gurevich, the 6<sup>th</sup> international workshop on thin films and new ideas for RF superconductivity, October 2014, Italy and A. Gurevich, AIP Advances 5, 017112 (2015)



**$d_I$  should be thin.  
Up to several tens of nm acceptable.**

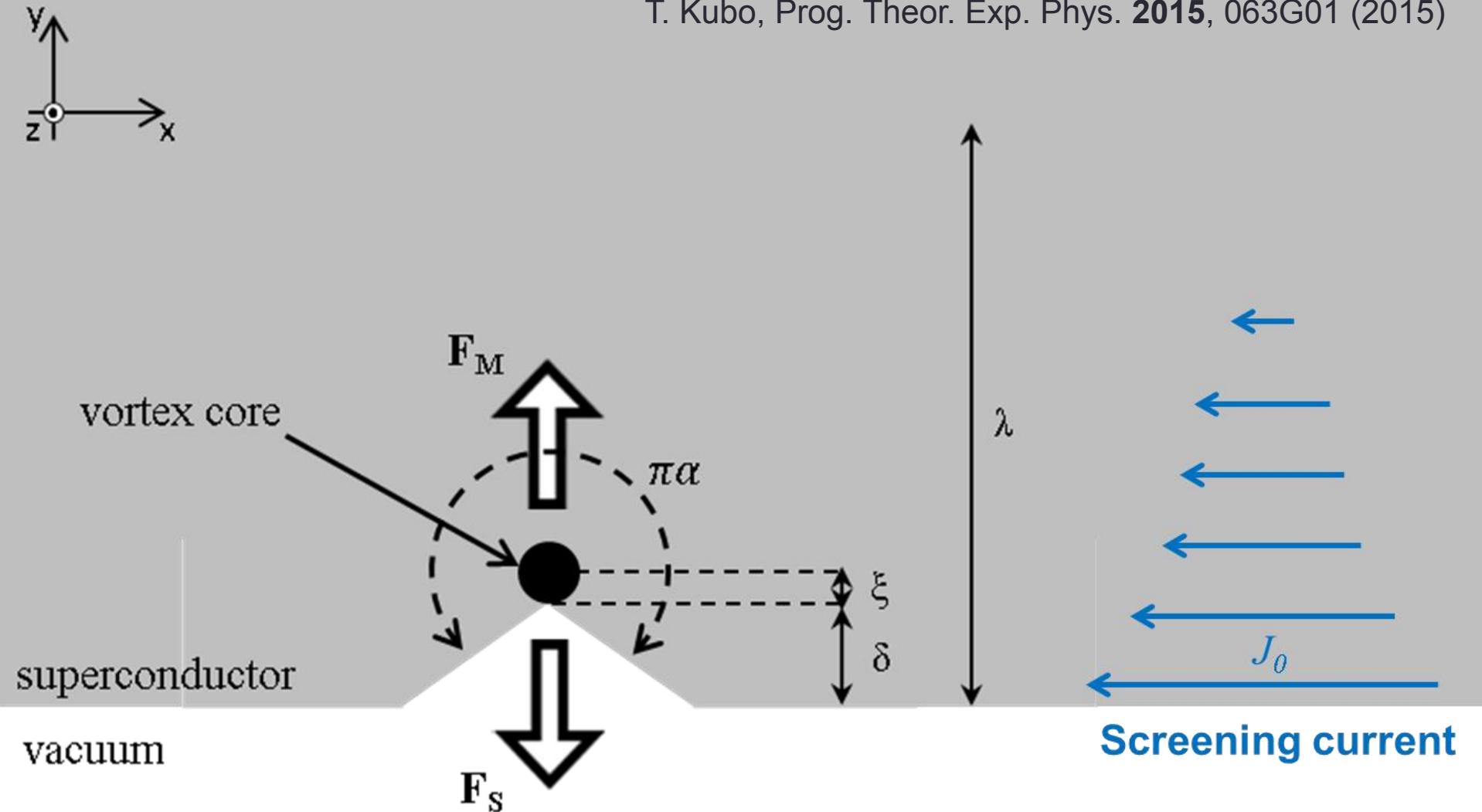
assumption  
 S layer: Dirty Nb  
 $B_c^{(\text{Nb})}=200\text{mT}$   
 $\lambda_1=\lambda(\text{dirtyNb})=180\text{nm}$   
 SC substrate: clean Nb  
 $B_{\max}^{(\text{Nb})}=240\text{mT}$   
 $\lambda_2=\lambda^{(\text{Nb})}=40\text{nm}$

## § 2

*A further step forward*

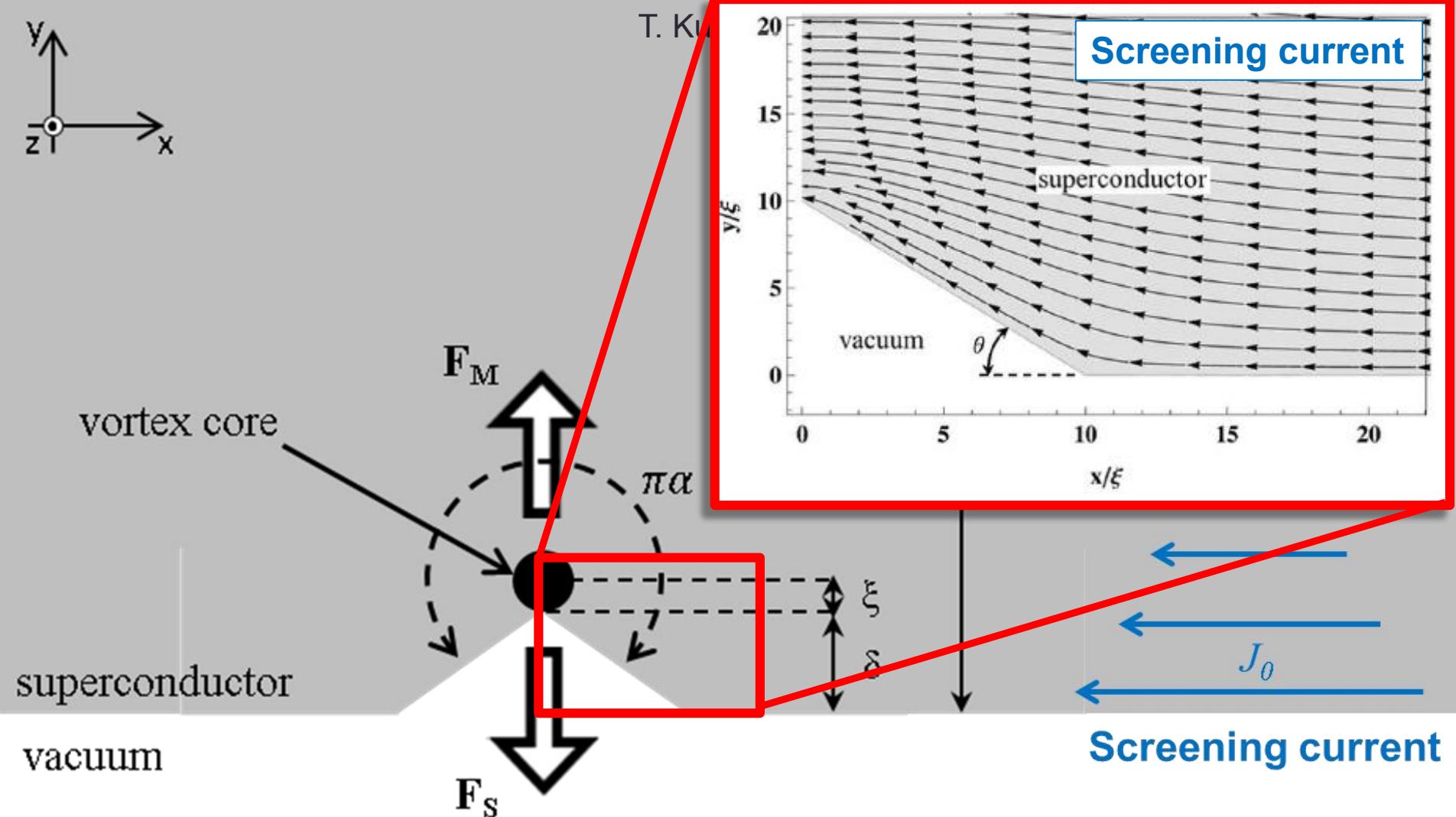
*incorporate non-ideal surfaces*

T. Kubo, Prog. Theor. Exp. Phys. 2015, 063G01 (2015)



We assume  $\xi < \delta < \lambda$ .

**Such small defects almost continuously distribute on the surface.**



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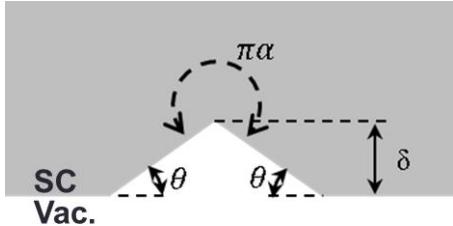
Such small defects almost continuously distribute on the surface.

# The superheating field is suppressed due to the enhanced screening current.

Suppression factor

$$\tilde{B}_s = \eta B_s$$

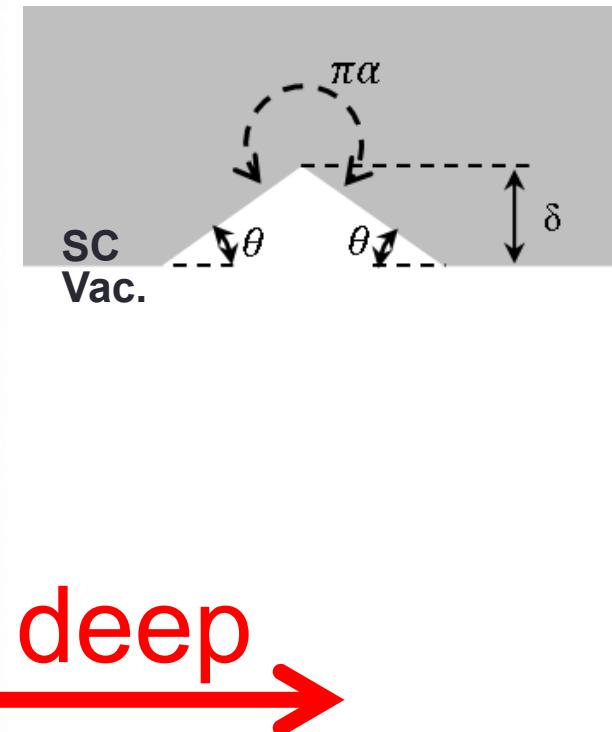
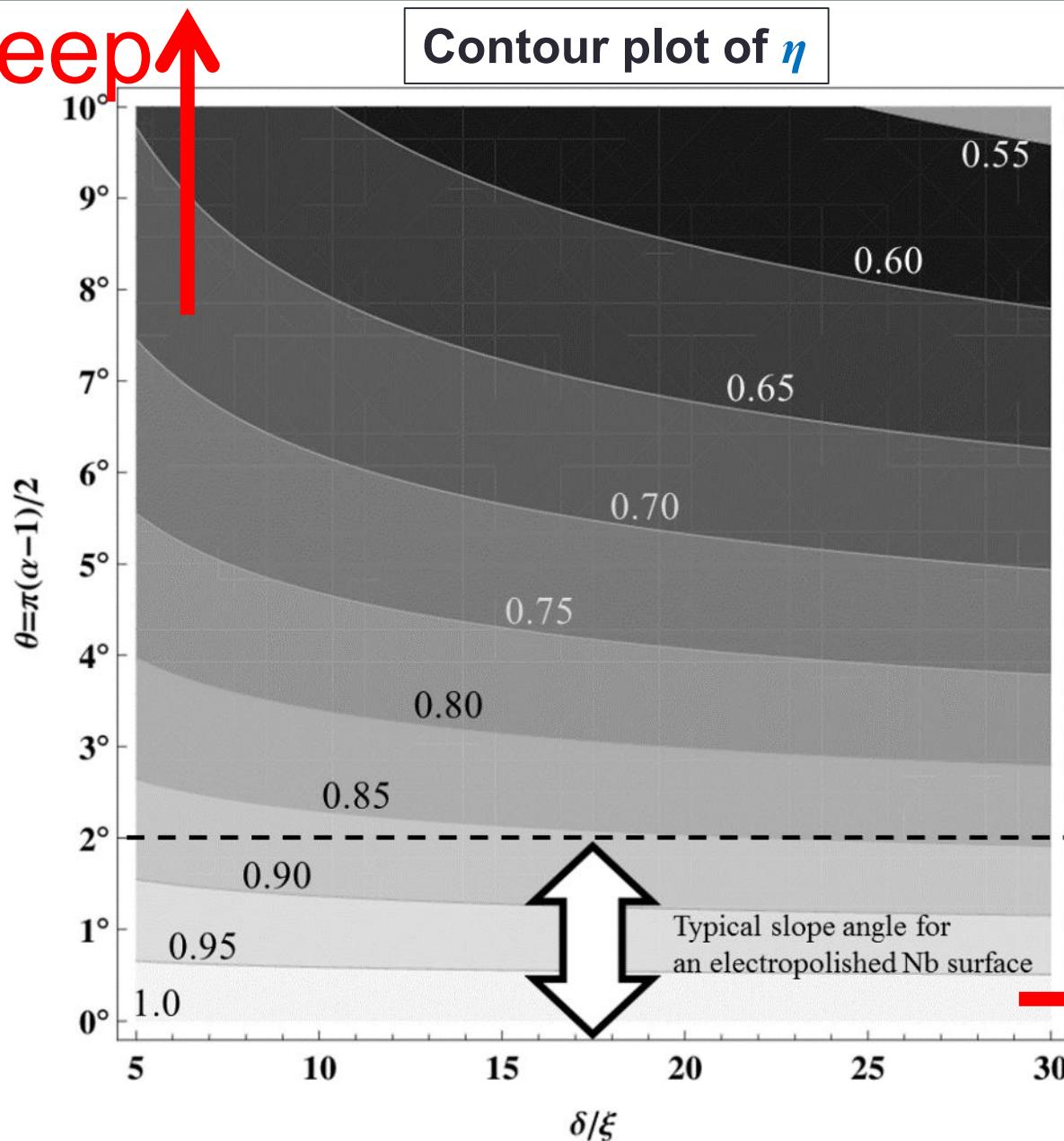
$$\eta = \frac{1}{\alpha} \left( \frac{\Gamma\left(\frac{\alpha}{2}\right) \Gamma\left(\frac{3-\alpha}{2}\right) \alpha \sin \frac{\pi(\alpha-1)}{2} \xi}{\sqrt{\pi}} \frac{\delta}{\xi} \right)^{\frac{\alpha-1}{\alpha}}$$



T. Kubo, Prog. Theor. Exp. Phys. 2015, 063G01 (2015)

We can evaluate a “**suppression factor  $\eta$** ” for materials, if we have **data of surface topographic studies**  
 (see for example “C. Xu et al. Phys. Rev. ST Accel. Beams 14, 123501 (2011)”).

steep↑



# The formula for the maximum screening field of the multilayer

$$B_{\max}^{(\text{multilayer})} = \begin{cases} B_{\max}^{(\text{S layer})} & (\text{if } \gamma B_{\max}^{(\text{S layer})} < B_{\max}^{(\text{substrate})}) \\ \gamma^{-1} \times B_{\max}^{(\text{substrate})} & (\text{if } \gamma B_{\max}^{(\text{S layer})} \geq B_{\max}^{(\text{substrate})}) \end{cases}$$

where

$$\left. \begin{aligned} B_{\max}^{(\text{S layer})} &= B_s^{(\text{S layer})} \frac{\cosh \frac{d_S}{\lambda_1} + \left( \frac{\lambda_2}{\lambda_1} + \frac{d_I}{\lambda_1} \right) \sinh \frac{d_S}{\lambda_1}}{\sinh \frac{d_S}{\lambda_1} + \left( \frac{\lambda_2}{\lambda_1} + \frac{d_I}{\lambda_1} \right) \cosh \frac{d_S}{\lambda_1}} \times \eta \\ \gamma &= \frac{1}{\cosh \frac{d_S}{\lambda_1} + \left( \frac{\lambda_2}{\lambda_1} + \frac{d_I}{\lambda_1} \right) \sinh \frac{d_S}{\lambda_1}} \\ B_{\max}^{(\text{substrate})} &= 170\text{mT} - 240\text{mT} \end{aligned} \right\}$$

$\eta$  can be estimated by

$$\eta = \frac{1}{\alpha} \left( \frac{\Gamma\left(\frac{\alpha}{2}\right) \Gamma\left(\frac{3-\alpha}{2}\right) \alpha \sin \frac{\pi(\alpha-1)}{2}}{\sqrt{\pi}} \xi \right)^{\frac{\alpha-1}{\alpha}}$$

T. Kubo (2015),

Obtained by using T. Kubo et al., Appl. Phys. Lett. **104**, 032603 (2014), A. Gurevich, AIP Advances **5**, 017112 (2015), and T. Kubo, Prog. Theor. Exp. Phys. **2015**, 063G01 (2015)

# The formula for the optimum thickness of the S layer

$$d_S = \lambda_1 \ln \left[ \frac{\lambda_1}{\lambda_1 + \lambda_2 + d_I} \frac{\eta B_s^{(\text{S layer})}}{B_{max}^{(\text{substrate})}} + \sqrt{\left( \frac{\lambda_1}{\lambda_1 + \lambda_2 + d_I} \frac{\eta B_s^{(\text{S layer})}}{B_{max}^{(\text{substrate})}} \right)^2 + \frac{\lambda_1 - \lambda_2 - d_I}{\lambda_1 + \lambda_2 + d_I}} \right]_{d_I \lesssim \mathcal{O}(10) \text{ nm}}$$

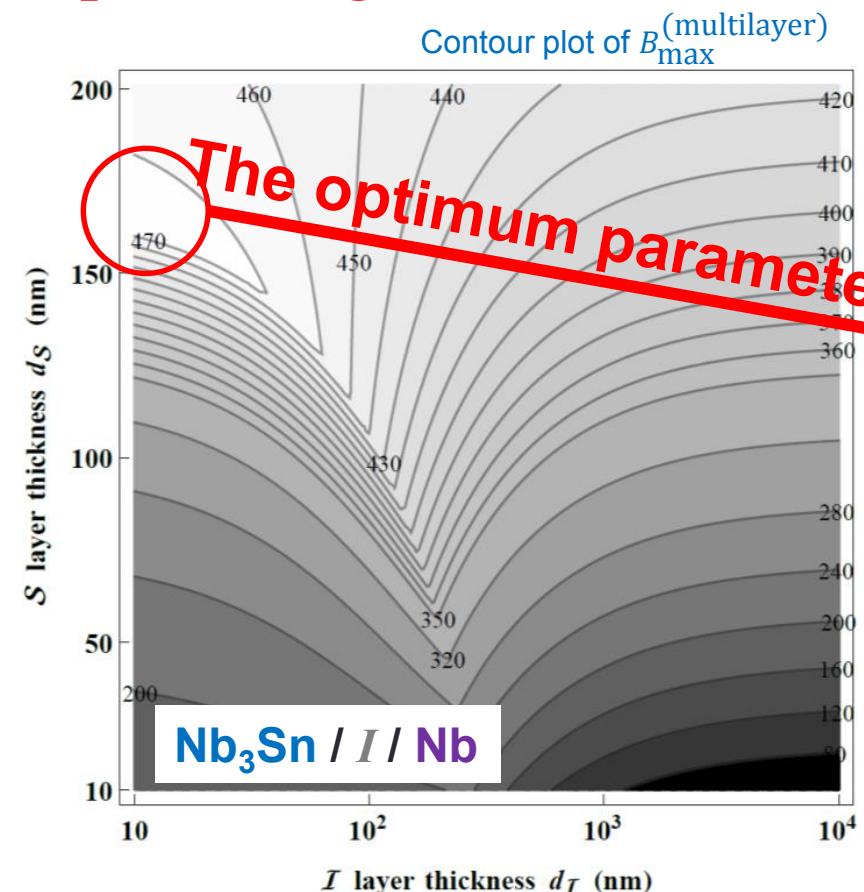
where

$$\left. \begin{array}{l} B_s^{(\text{S layer})} = 0.84 B_c^{(\text{S layer})} \\ B_{max}^{(\text{substrate})} = 170 \text{ mT} - 240 \text{ mT} \quad (\text{if the substrate is Nb}) \\ \eta \text{ can be estimated by} \quad \eta = \frac{1}{\alpha} \left( \frac{\Gamma\left(\frac{\alpha}{2}\right) \Gamma\left(\frac{3-\alpha}{2}\right) \alpha \sin \frac{\pi(\alpha-1)}{2}}{\sqrt{\pi}} \xi \right)^{\frac{\alpha-1}{\alpha}} \end{array} \right\}$$

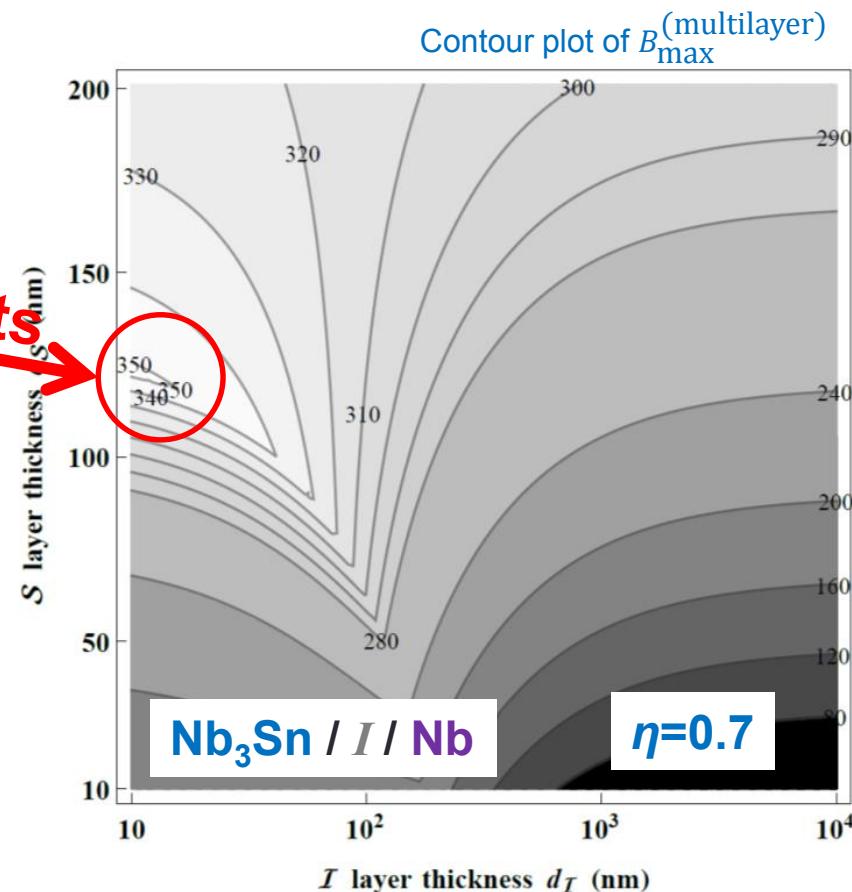
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The optimum thickness becomes thin in order to compensate **the suppressed superheating field.**



Optimum  $d_s \sim 170\text{nm}$   
Maximum field  $\sim 470\text{mT}$



Optimum  $d_s \sim 130\text{nm}$   
Maximum field  $\sim 350\text{mT}$

assumption  
S layer:  $\text{Nb}_3\text{Sn}$  (moderately dirty)  
 $B_c^{(\text{Nb}_3\text{Sn})}=540\text{mT}$   
 $\lambda_1=\lambda^{(\text{Nb}_3\text{Sn})}=120\text{nm}$   
SC substrate: clean Nb  
 $B_{\max}^{(\text{Nb})}=B_{c1}^{(\text{Nb})}=\underline{170\text{mT}}$   
 $\lambda_2=\lambda^{(\text{Nb})}=40\text{nm}$

# Summary

The optimum  $S$  layer thickness is a function of

- $\lambda_1$ , penetration depth of the  $S$  layer
- $\lambda_2$  ( $\lambda_1 > \lambda_2$ ), penetration depth of the substrate
- $d_I$  (< several tens of nm), thickness of the insulator
- $B_c^{(S\text{ layer})}$ , thermodynamic critical field of the  $S$  layer
- $B_{max}^{(\text{substrate})}$  (170-240mT for Nb).
- $\eta$  ( $0 < \eta < 1$ ), superheating field suppression factor

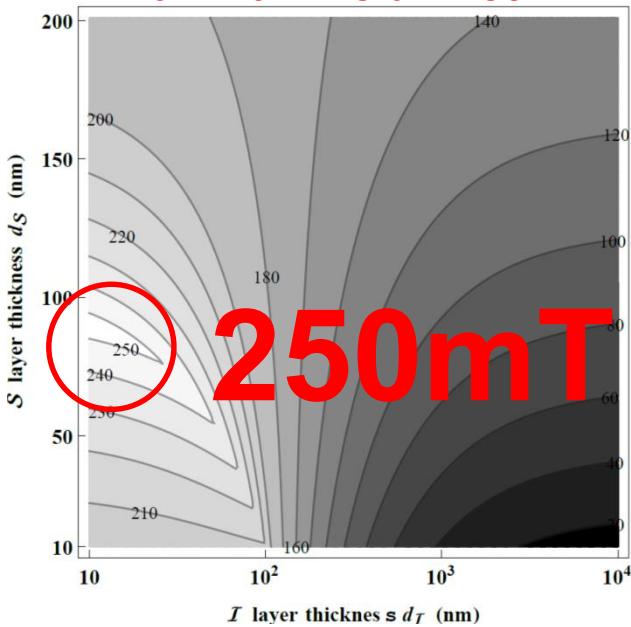
Material parameters and surface topographic studies are necessary in order to obtain the optimum  $S$  layer thickness.

The defect model and the formula for  $\eta$  may be useful to extract  $\eta$  from surface topographic studies.

# Let us go beyond Nb by using the optimum parameters!

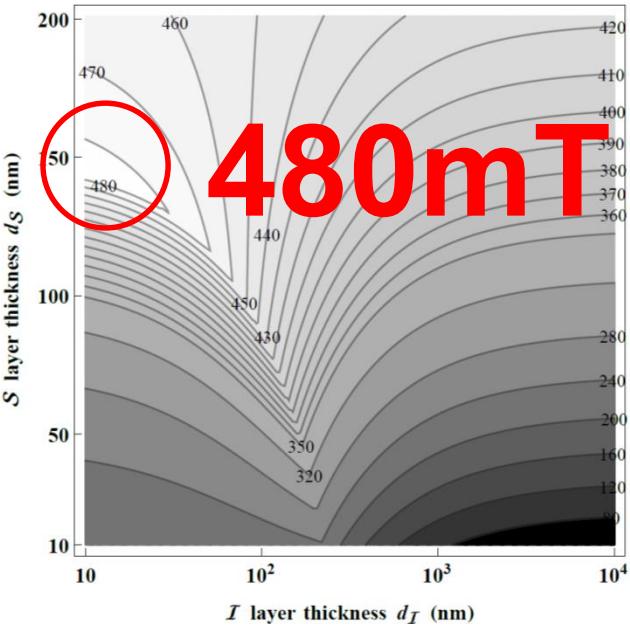
## Dirty Nb / I / Nb

Optimum  $d_s \sim 90\text{nm}$   
Maximum field  $\sim 250\text{mT}$



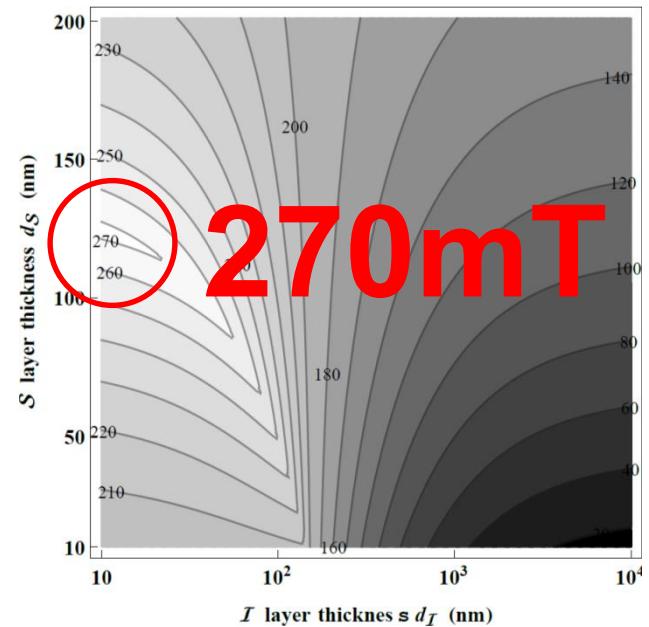
## Nb<sub>3</sub>Sn / I / Nb

Optimum  $d_s \sim 150\text{nm}$   
Maximum field  $\sim 480\text{mT}$



## NbN / I / Nb

Optimum  $d_s \sim 130\text{nm}$   
Maximum field  $\sim 270\text{mT}$



# Appendix for multilayer researchers

The optimum parameters for

**Dirty Nb /  $I$  / Nb**

**Nb<sub>3</sub>Sn /  $I$  / Nb**

**NbN /  $I$  / Nb**

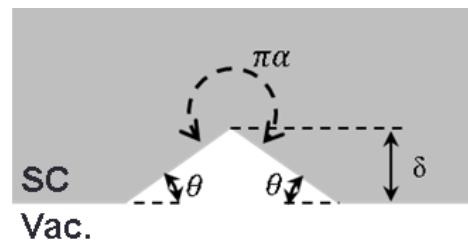
are given below!

# The formula for the maximum screening field

$$B_{\max}^{(\text{multilayer})} = \begin{cases} B_{\max}^{(\text{S layer})} & (\text{if } \gamma B_{\max}^{(\text{S layer})} < B_{\max}^{(\text{substrate})}) \\ \gamma^{-1} \times B_{\max}^{(\text{substrate})} & (\text{if } \gamma B_{\max}^{(\text{S layer})} \geq B_{\max}^{(\text{substrate})}) \end{cases}$$

where

$$\left\{ \begin{array}{l} B_{\max}^{(\text{S layer})} = B_s^{(\text{S layer})} \frac{\cosh \frac{d_s}{\lambda_1} + \left( \frac{\lambda_2}{\lambda_1} + \frac{d_I}{\lambda_1} \right) \sinh \frac{d_s}{\lambda_1}}{\sinh \frac{d_s}{\lambda_1} + \left( \frac{\lambda_2}{\lambda_1} + \frac{d_I}{\lambda_1} \right) \cosh \frac{d_s}{\lambda_1}} \times \eta \\ \gamma = \frac{1}{\cosh \frac{d_s}{\lambda_1} + \left( \frac{\lambda_2}{\lambda_1} + \frac{d_I}{\lambda_1} \right) \sinh \frac{d_s}{\lambda_1}} \\ B_{\max}^{(\text{substrate})} = 170\text{mT} - 240\text{mT} \end{array} \right. \quad \left. \begin{array}{l} B_s^{(\text{S layer})} = 0.84 B_c^{(\text{S layer})} \\ \eta = \frac{1}{\alpha} \left( \frac{\Gamma \left( \frac{\alpha}{2} \right) \Gamma \left( \frac{3-\alpha}{2} \right) \alpha \sin \frac{\pi(\alpha-1)}{2}}{\sqrt{\pi}} \xi \right)^{\frac{\alpha-1}{\alpha}} \end{array} \right.$$



# The formula for the optimum thickness of the *S* layer

$$d_S = \lambda_1 \ln \left[ \frac{\lambda_1}{\lambda_1 + \lambda_2 + d_I} \frac{\eta B_s^{(\text{S layer})}}{B_{max}^{(\text{substrate})}} + \sqrt{\left( \frac{\lambda_1}{\lambda_1 + \lambda_2 + d_I} \frac{\eta B_s^{(\text{S layer})}}{B_{max}^{(\text{substrate})}} \right)^2 + \frac{\lambda_1 - \lambda_2 - d_I}{\lambda_1 + \lambda_2 + d_I}} \right]_{d_I \lesssim \mathcal{O}(10) \text{ nm}}$$

## Necessary parameters

*S* layer:  $B_s^{(\text{S layer})} = 0.84 \times B_c^{(\text{S layer material})}$

$\lambda_1 = \lambda^{(\text{S layer material})}$

$\eta$  (suppression factor)

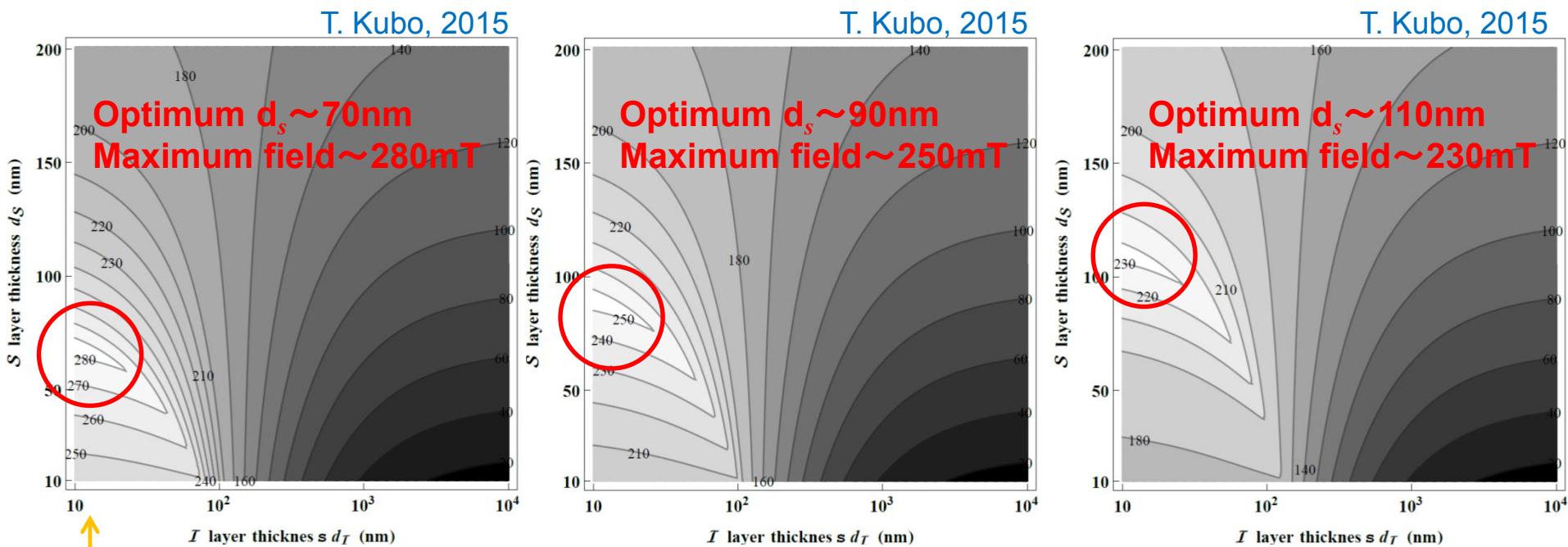
*I* layer:  $d_I \sim 10-100 \text{ nm}$

Substrate:  $B_{max}^{(\text{substrate})} = 170 \text{ mT} - 240 \text{ mT}$  (if the substrate is Nb)

$\lambda_2 = \lambda^{(\text{substrate})} = 40 \text{ nm}$  (if the substrate is Nb)

$$\eta=1$$

The optimum  $d_s$  and the maximum field for Dirty Nb / I / Nb system, where Nb substrate is assumed to withstand up to 240, 200, and 170mT.



Consistent with the Gurevich's recent result

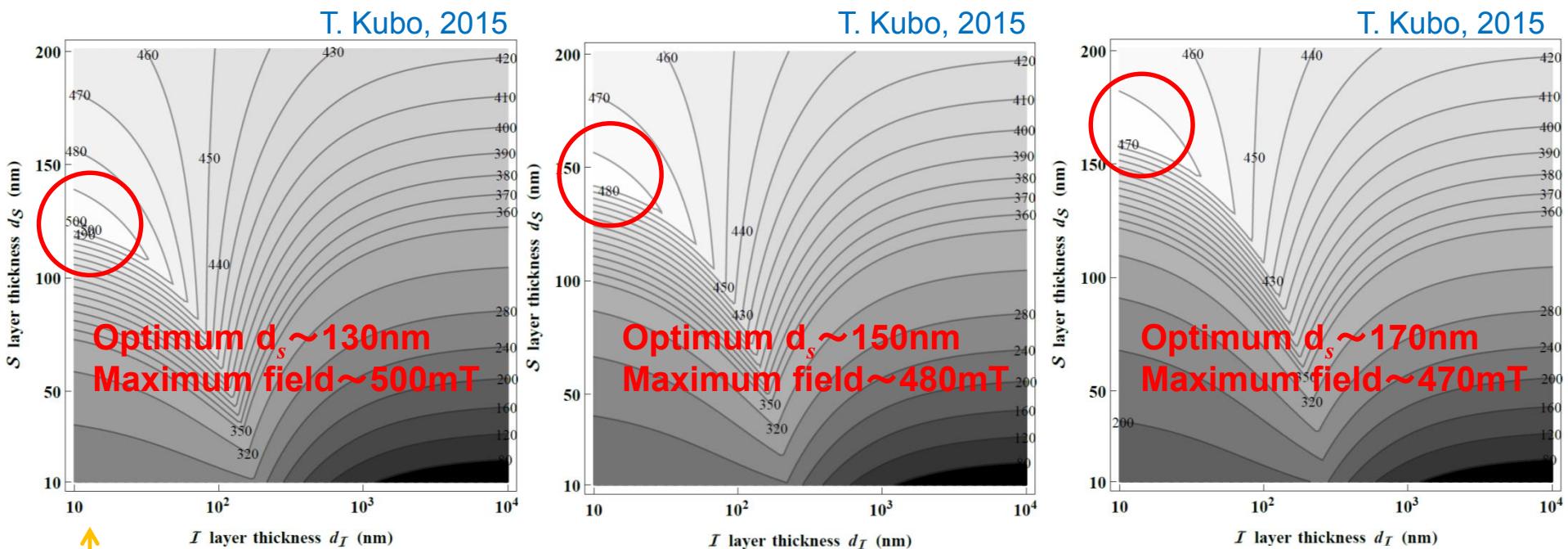
A. Gurevich, AIP Advances 5, 017112 (2015) and 6<sup>th</sup> thin film workshop at Italy

assumption  
**S layer: Dirty Nb**  
 $B_c^{(Nb)}=200$ mT  
 $\lambda_1=\lambda^{(\text{dirtyNb})}=180$ nm  
**SC substrate: clean Nb**  
 $B_{\max}^{(Nb)}=\underline{\underline{240}}\text{mT}$   
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 $B_{\max}^{(Nb)}=B_{c1}^{(Nb)}=\underline{\underline{170}}\text{mT}$   
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S layer:  **$\text{Nb}_3\text{Sn}$**  (moderately dirty)  
 $B_c^{(\text{Nb}_3\text{Sn})}=540\text{mT}$   
 $\lambda_1=\lambda^{(\text{Nb}_3\text{Sn})}=120\text{nm}$

SC substrate: clean **Nb**  
 $B_{\max}^{(\text{Nb})}=\underline{\underline{240\text{mT}}}$   
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 $B_c^{(\text{Nb}_3\text{Sn})}=540\text{mT}$   
 $\lambda_1=\lambda^{(\text{Nb}_3\text{Sn})}=120\text{nm}$

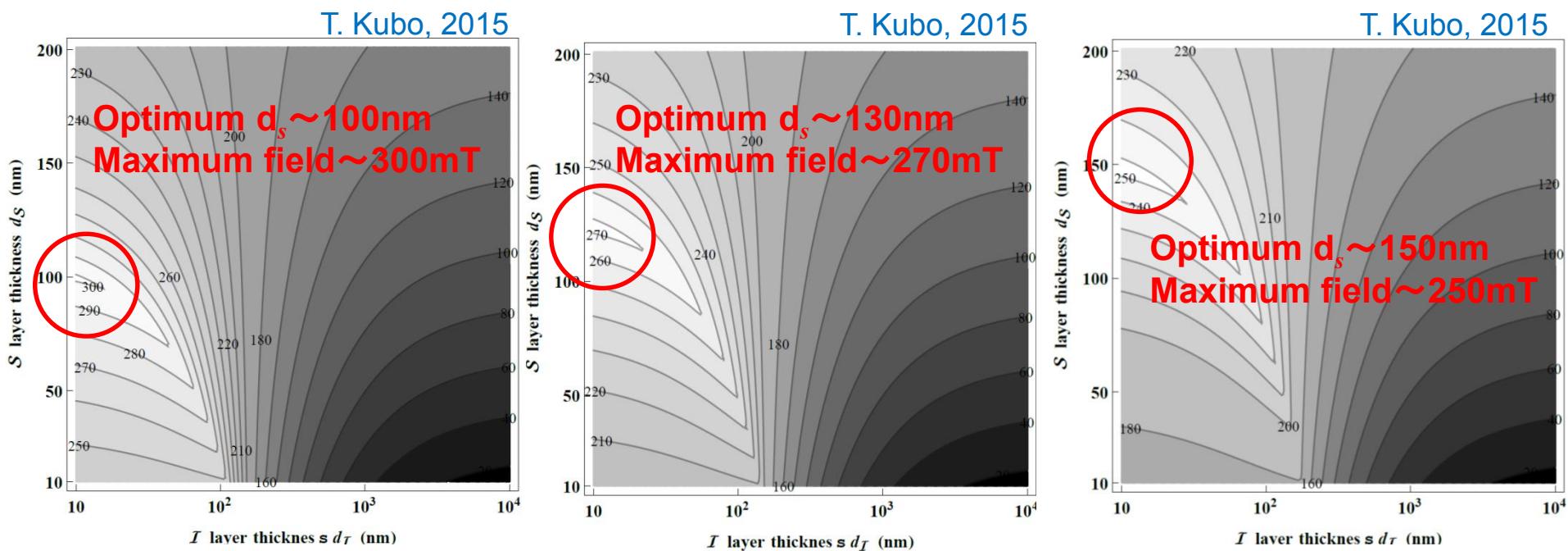
SC substrate: clean **Nb**  
 $B_{\max}^{(\text{Nb})}=\underline{\underline{200\text{mT}}}$   
 $\lambda_2=\lambda^{(\text{Nb})}=40\text{nm}$

assumption

S layer:  **$\text{Nb}_3\text{Sn}$**  (moderately dirty)  
 $B_c^{(\text{Nb}_3\text{Sn})}=540\text{mT}$   
 $\lambda_1=\lambda^{(\text{Nb}_3\text{Sn})}=120\text{nm}$

SC substrate: clean **Nb**  
 $B_{\max}^{(\text{Nb})}=B_{c1}^{(\text{Nb})}=\underline{\underline{170\text{mT}}}$   
 $\lambda_2=\lambda^{(\text{Nb})}=40\text{nm}$

The optimum  $d_s$  and the maximum field for NbN / I / Nb system, where Nb substrate is assumed to withstand up to 240, 200, and 170mT.



assumption

S layer: NbN  
 $B_c^{(NbN)}=230$ mT  
 $\lambda_1=\lambda^{(NbN)}=200$ nm

SC substrate: clean Nb  
 $B_{\max}^{(Nb)}=\underline{240}$ mT  
 $\lambda_2=\lambda^{(Nb)}=40$ nm

assumption

S layer: NbN  
 $B_c^{(NbN)}=230$ mT  
 $\lambda_1=\lambda^{(NbN)}=200$ nm

SC substrate: clean Nb  
 $B_{\max}^{(Nb)}=\underline{200}$ mT  
 $\lambda_2=\lambda^{(Nb)}=40$ nm

assumption

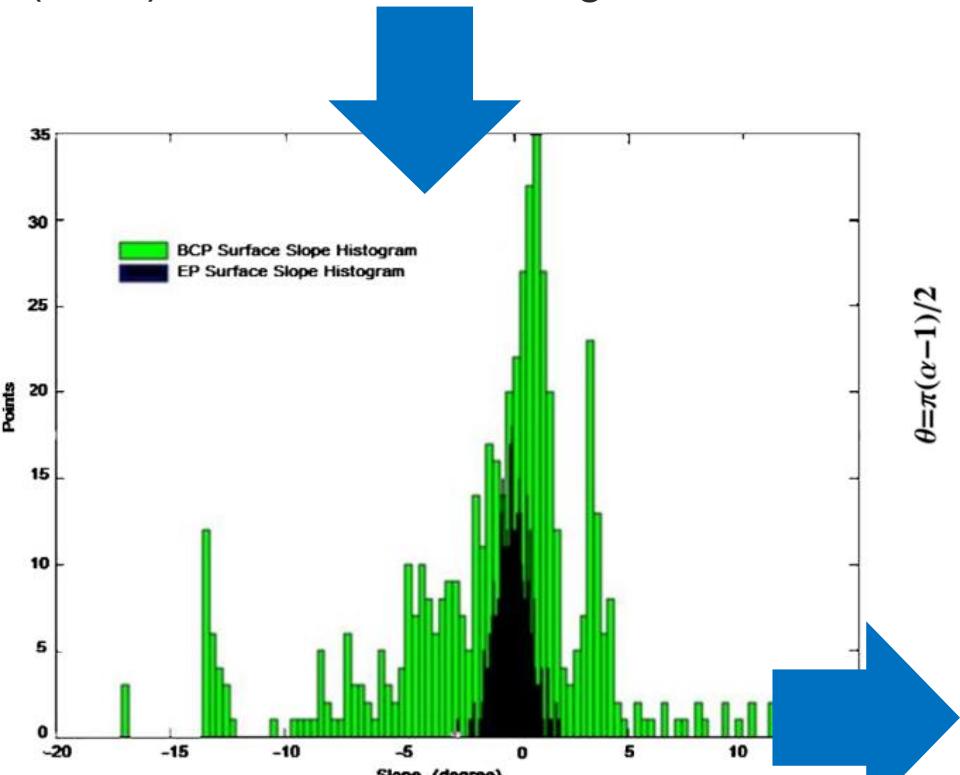
S layer: NbN  
 $B_c^{(NbN)}=230$ mT  
 $\lambda_1=\lambda^{(NbN)}=200$ nm

SC substrate: clean Nb  
 $B_{\max}^{(Nb)}=B_{c1}^{(Nb)}=\underline{170}$ mT  
 $\lambda_2=\lambda^{(Nb)}=40$ nm

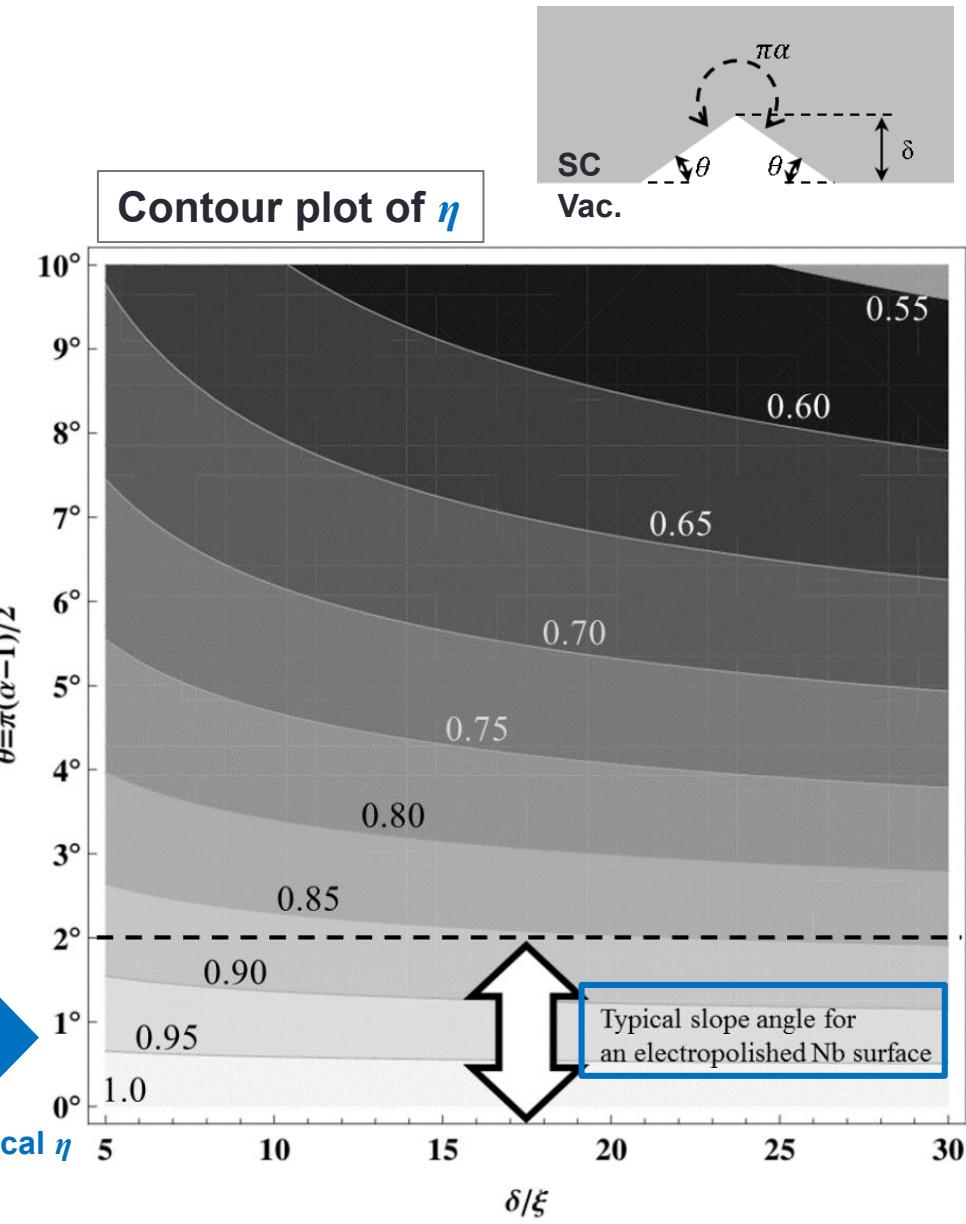
$$\eta < 1$$

## Example: electropolished Nb

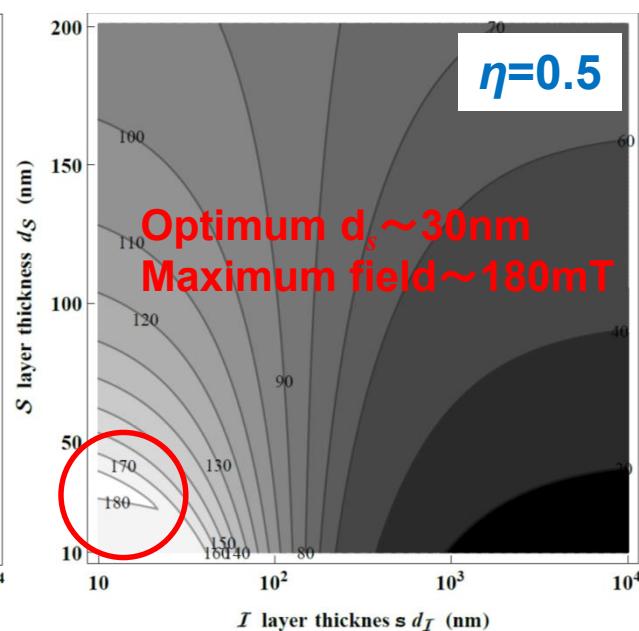
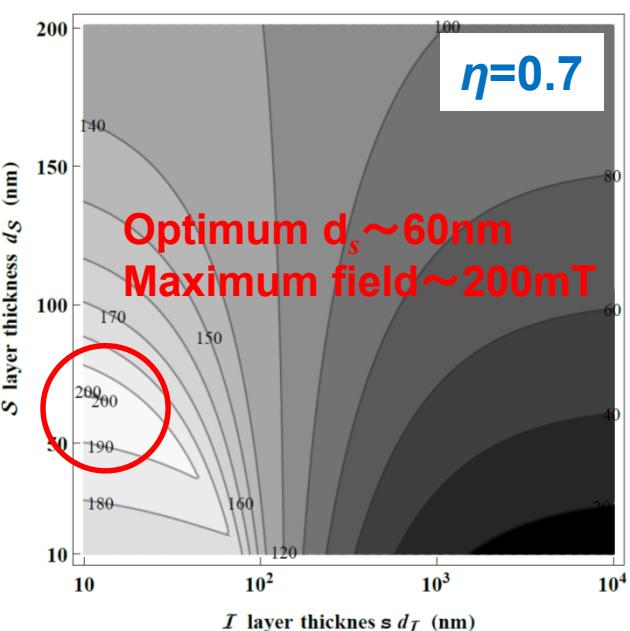
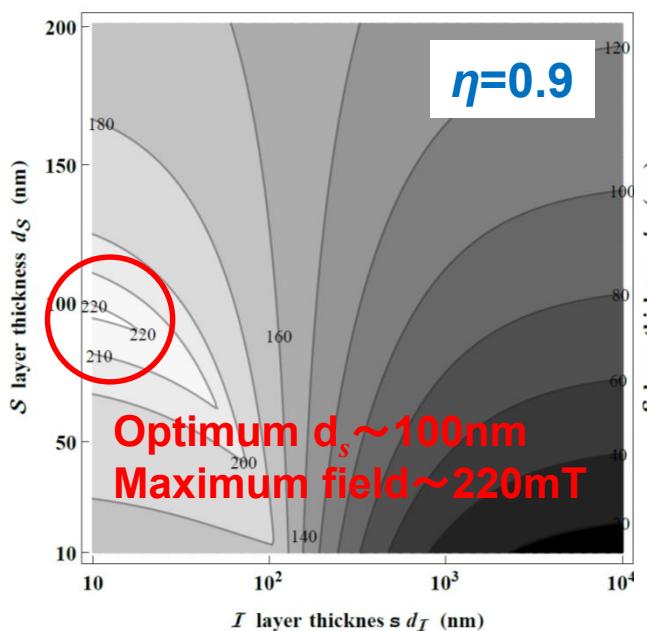
A surface after EP is studied by C.Xu, H.Tian, C.Reece, and M.Kelley, Phys. Rev. ST Accel. Beams **14**, 123501 (2011), which shows the figure below.



We can find typical  $\eta$  of EPed surface



The optimum  $d_s$  and the maximum field for Dirty Nb / I / Nb system, when the suppression factor due to nano-defects are given by  $\eta=0.9$ , 0.7, and 0.5. Nb substrate is assumed to withstand up to 170mT.

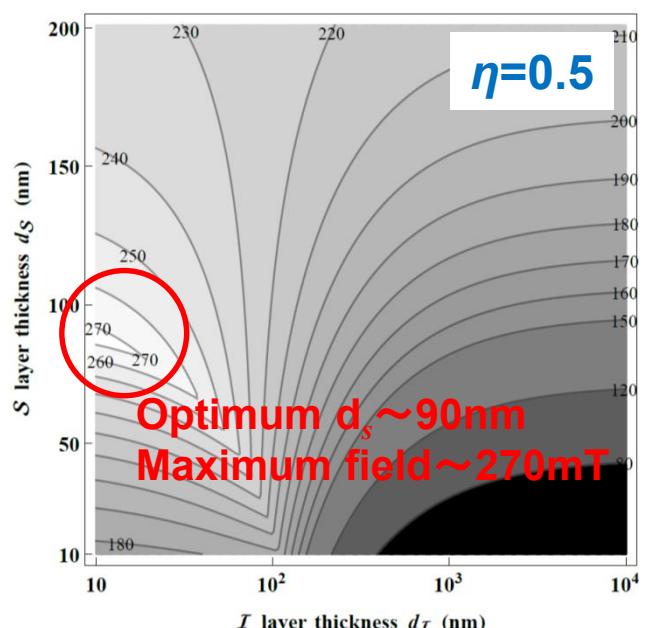
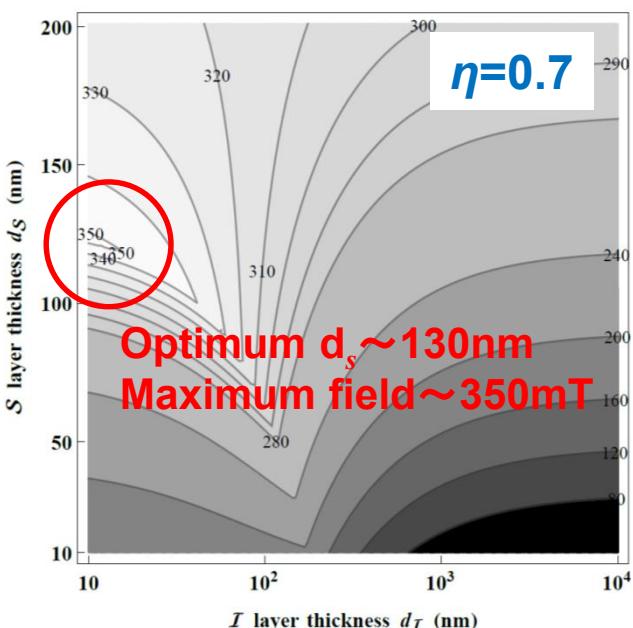
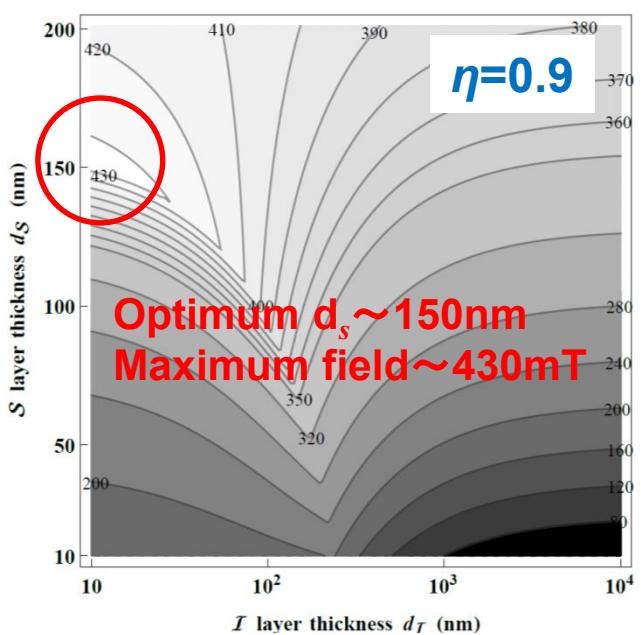


assumption

S layer: Dirty Nb  
 $B_c^{(Nb)}=200$ mT  
 $\lambda_1=\lambda^{(\text{dirtyNb})}=180$ nm

SC substrate: clean Nb  
 $B_{\max}^{(Nb)}=B_{c1}^{(Nb)}=\underline{170}$ mT  
 $\lambda_2=\lambda^{(Nb)}=40$ nm

The optimum  $d_s$  and the maximum field for  $\text{Nb}_3\text{Sn}$  / I / Nb system, when the suppression factor due to nano-defects are given by  $\eta=0.9$ , 0.7, and 0.5. Nb substrate is assumed to withstand up to 170mT.

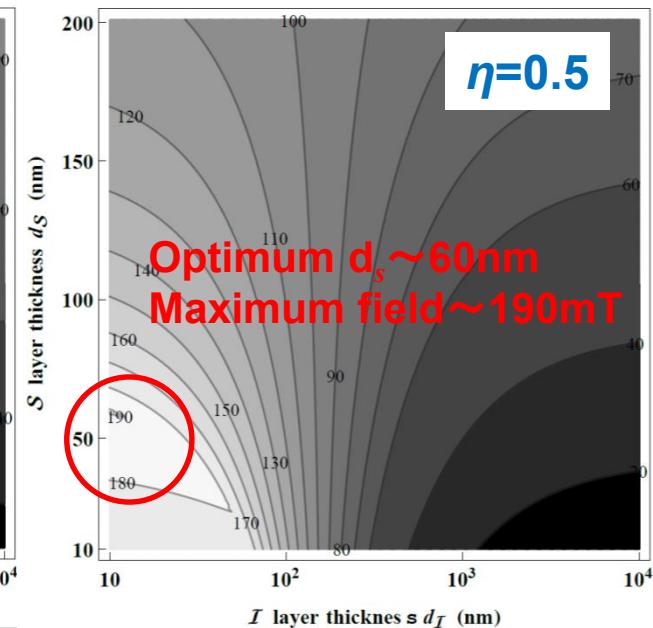
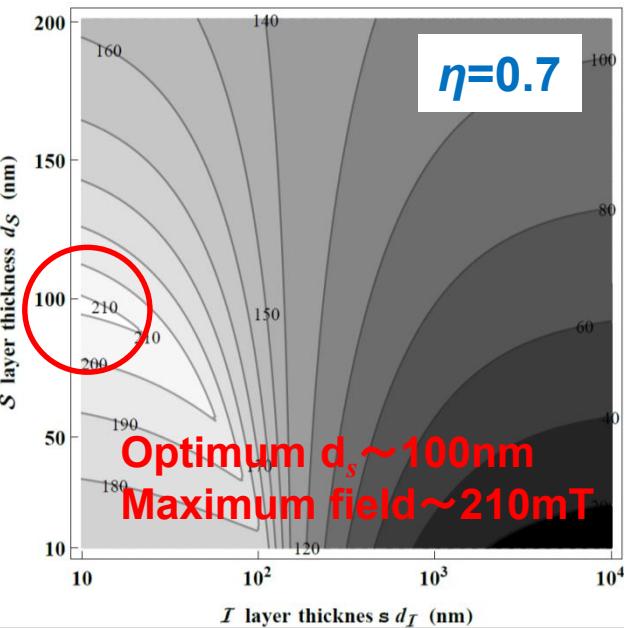
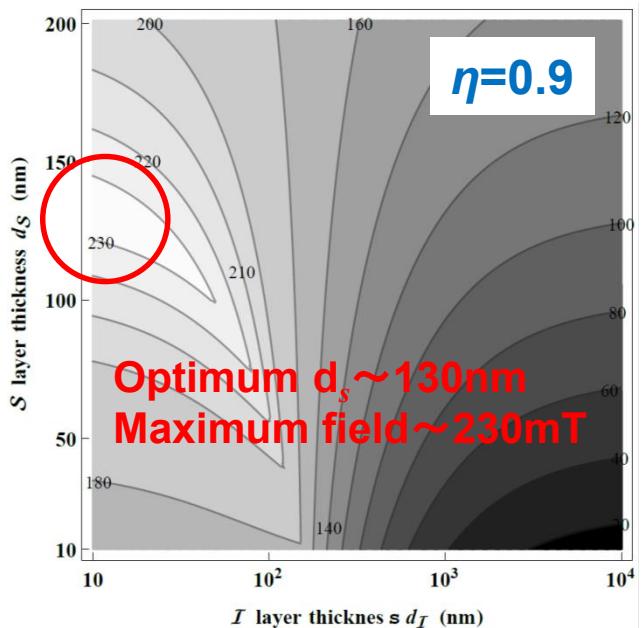


assumption

S layer:  $\text{Nb}_3\text{Sn}$  (moderately dirty)  
 $B_c^{(\text{Nb}_3\text{Sn})}=540$  mT  
 $\lambda_1=\lambda^{(\text{Nb}_3\text{Sn})}=120$  nm

SC substrate: clean Nb  
 $B_{\max}^{(\text{Nb})}=B_{c1}^{(\text{Nb})}=170$  mT  
 $\lambda_2=\lambda^{(\text{Nb})}=40$  nm

The optimum  $d_s$  and the maximum field for NbN / I / Nb system, when the suppression factor due to nano-defects are given by  $\eta=0.9$ , 0.7, and 0.5. Nb substrate is assumed to withstand up to 170mT.



### assumption

S layer: NbN

$$B_c^{(\text{NbN})} = 230 \text{ mT}$$

$$\lambda_1 = \lambda^{(\text{NbN})} = 200 \text{ nm}$$

SC substrate: clean Nb

$$B_{\max}^{(\text{Nb})} = B_{c1}^{(\text{Nb})} = 170 \text{ mT}$$

$$\lambda_2 = \lambda^{(\text{Nb})} = 40 \text{ nm}$$