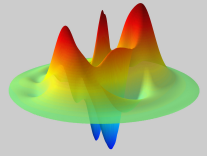


Quantum measurement with "trapped" microwave photons in a SRF cavity

Michel Brune



Paris, France



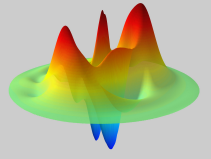
« Ridiculous » quantum phenomena

Schrödinger 1952 :

« one never experiments with just **one** electron, **one** atom or **one** molecule. In thought experiments we sometimes assume that we do, this invariably entails **ridiculous consequences**... »

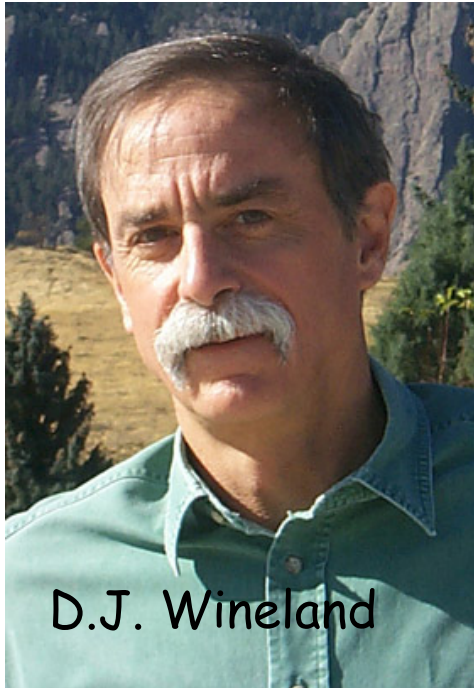
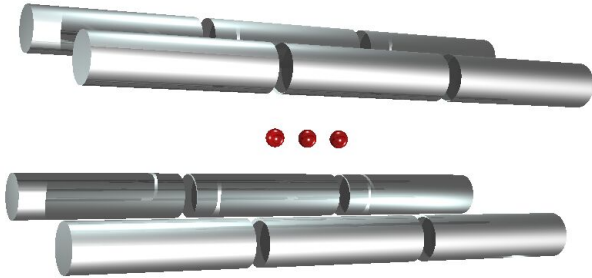
(British Journal of the Philosophy of Sciences, vol 3, 1952)





Two “spin and spring” experiments

- Ion traps

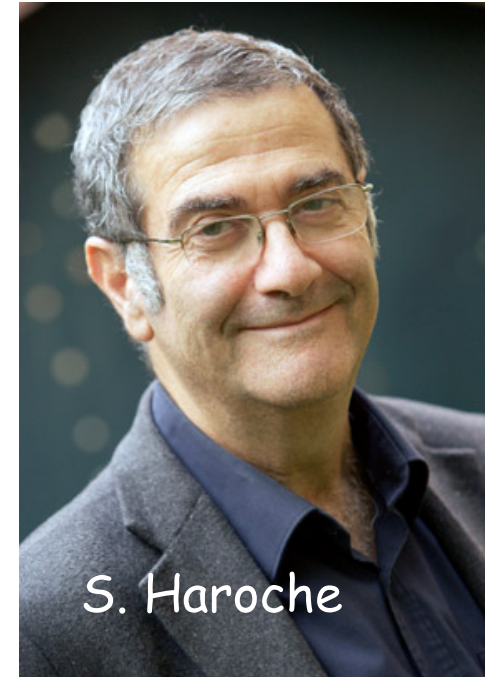
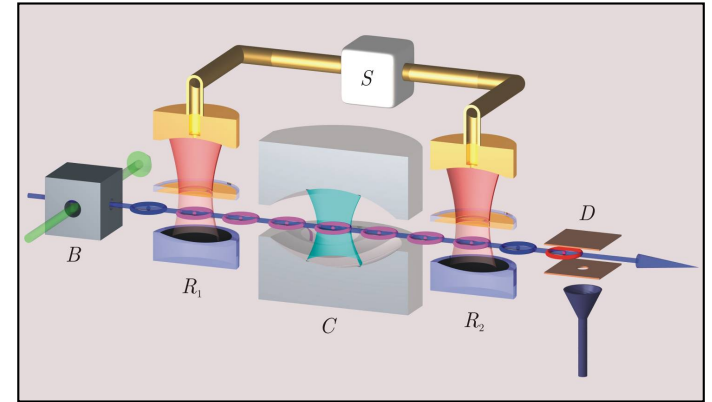


D.J. Wineland

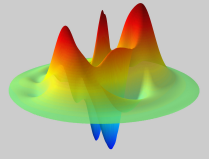
Different systems
with strong
similarities



- Atoms and cavities



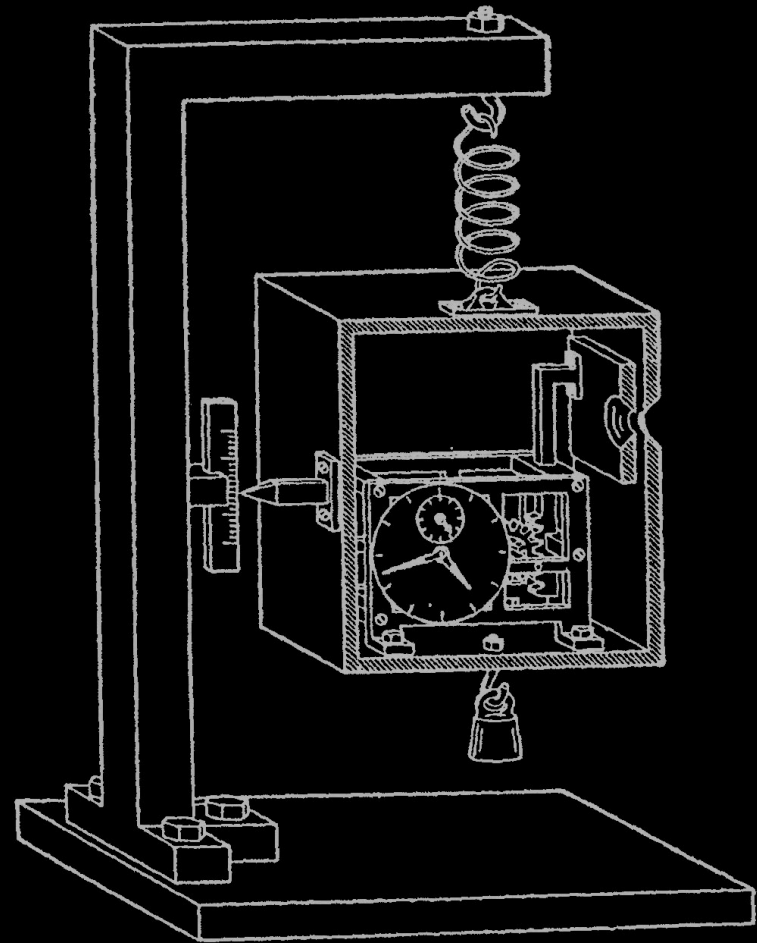
S. Haroche

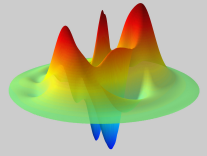


The « photon box » thought experiment



N. Bohr A. Einstein





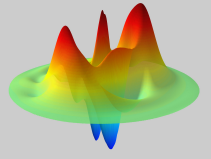
The « photon box » thought experiment



N. Bohr A. Einstein

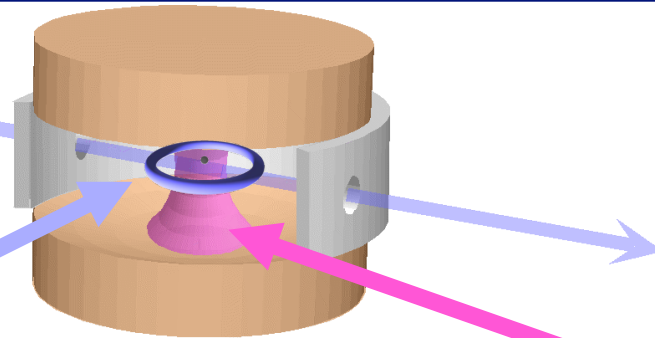


Experiments on QND counting
of trapped photons



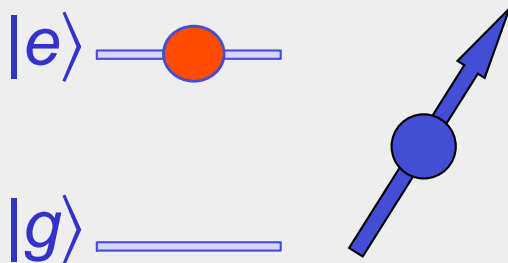
The tools: spin and spring

A cavity QED experiment



The SPIN:

One atom, two levels

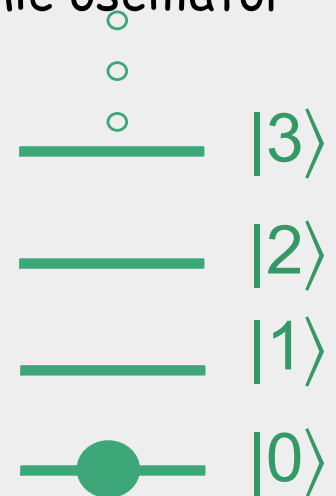
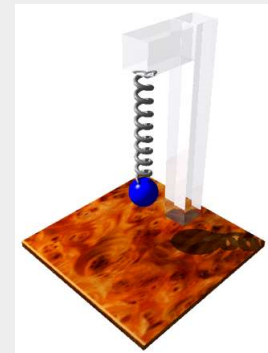


Electric dipole coupling

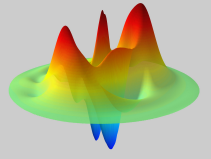
$$\Omega_0$$

The SPRING:

One high Q cavity mode as an harmonic oscillator

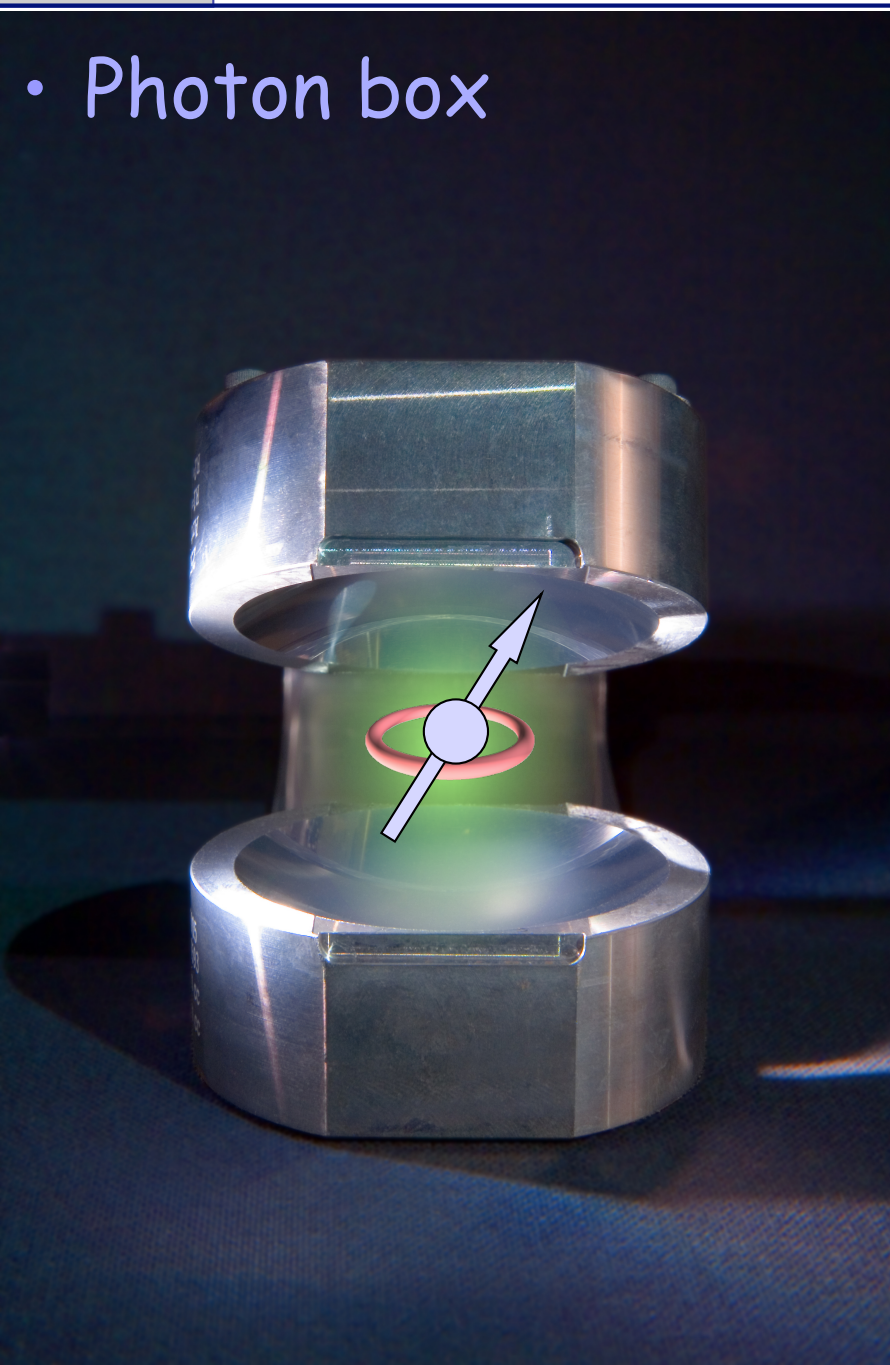


➔ Nearly ideal realization of a simple generic system

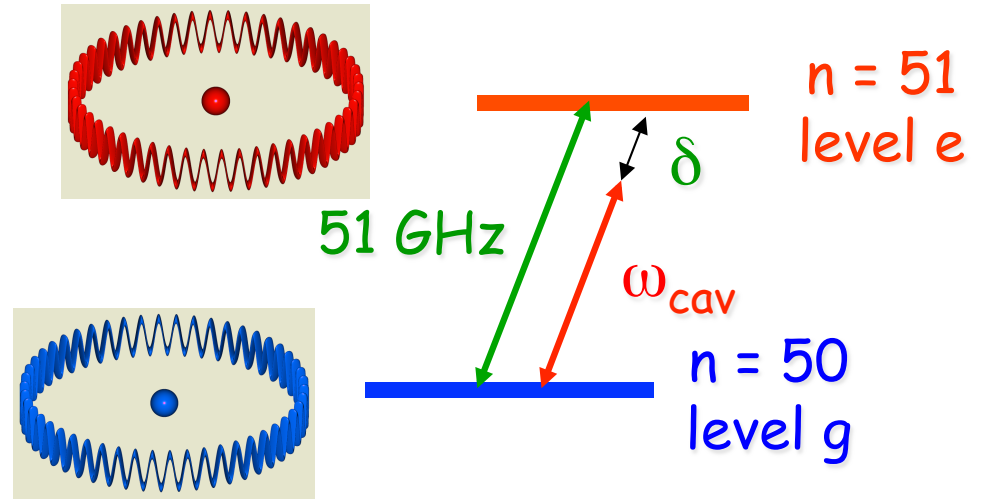


The "Spin"

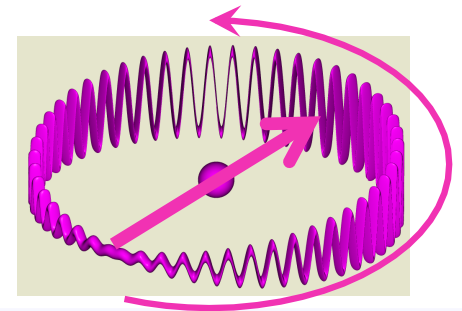
- Photon box

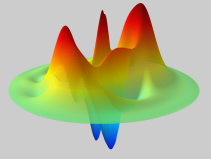


- Photon probes
Circular Rydberg atoms



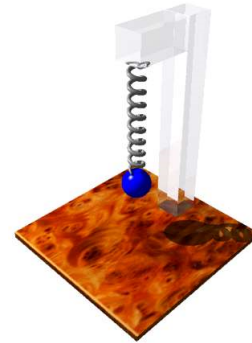
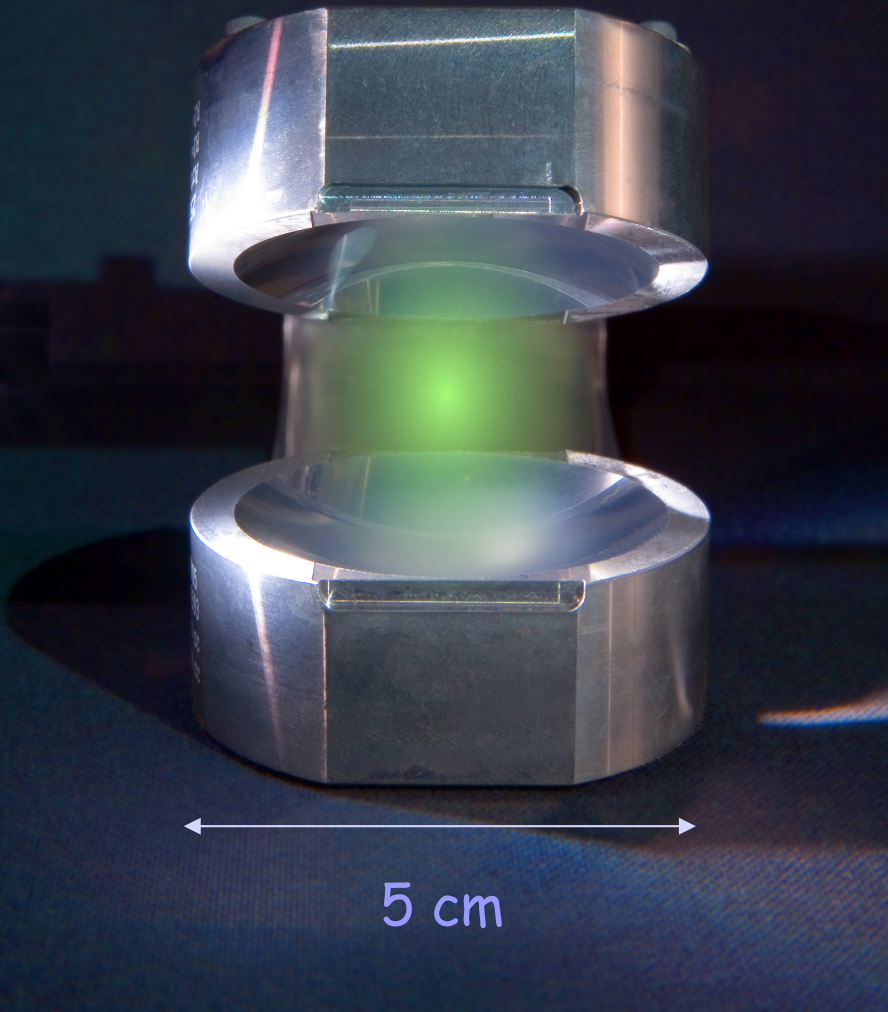
Superposition
State
 $|e\rangle + |g\rangle$





The “spring”

- Photon box
Superconducting mirrors



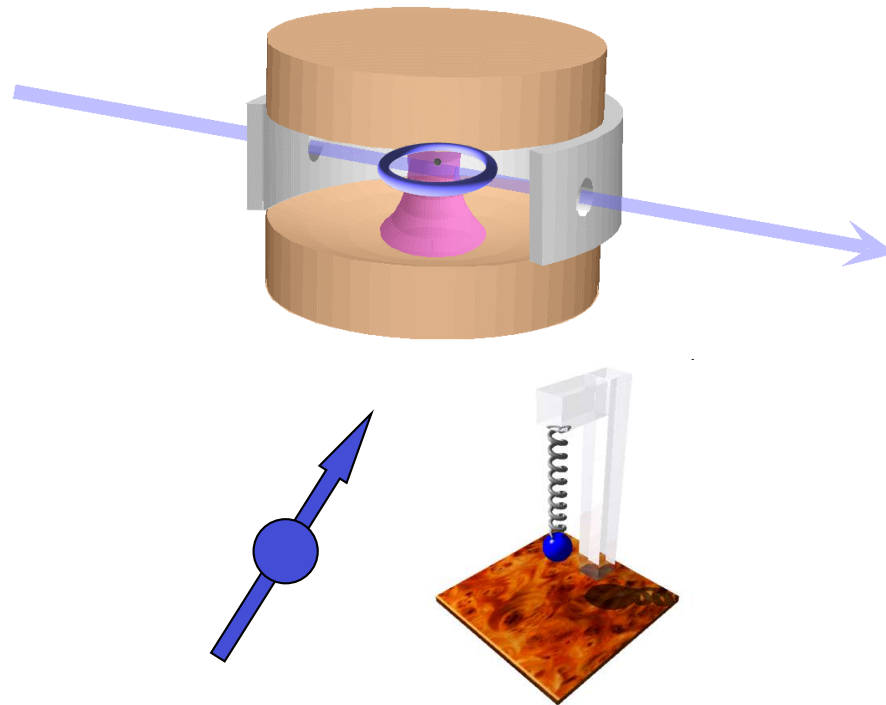
- Storing microwave photons
@ $\nu_{\text{cav}} = 51 \text{ GHz}$

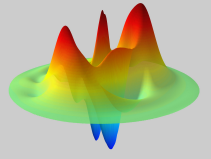
- Photon lifetime:

$T_{\text{cav}} = \text{from } 2 \mu\text{s} \dots$

$\dots \text{ to } 0.130\text{s}$

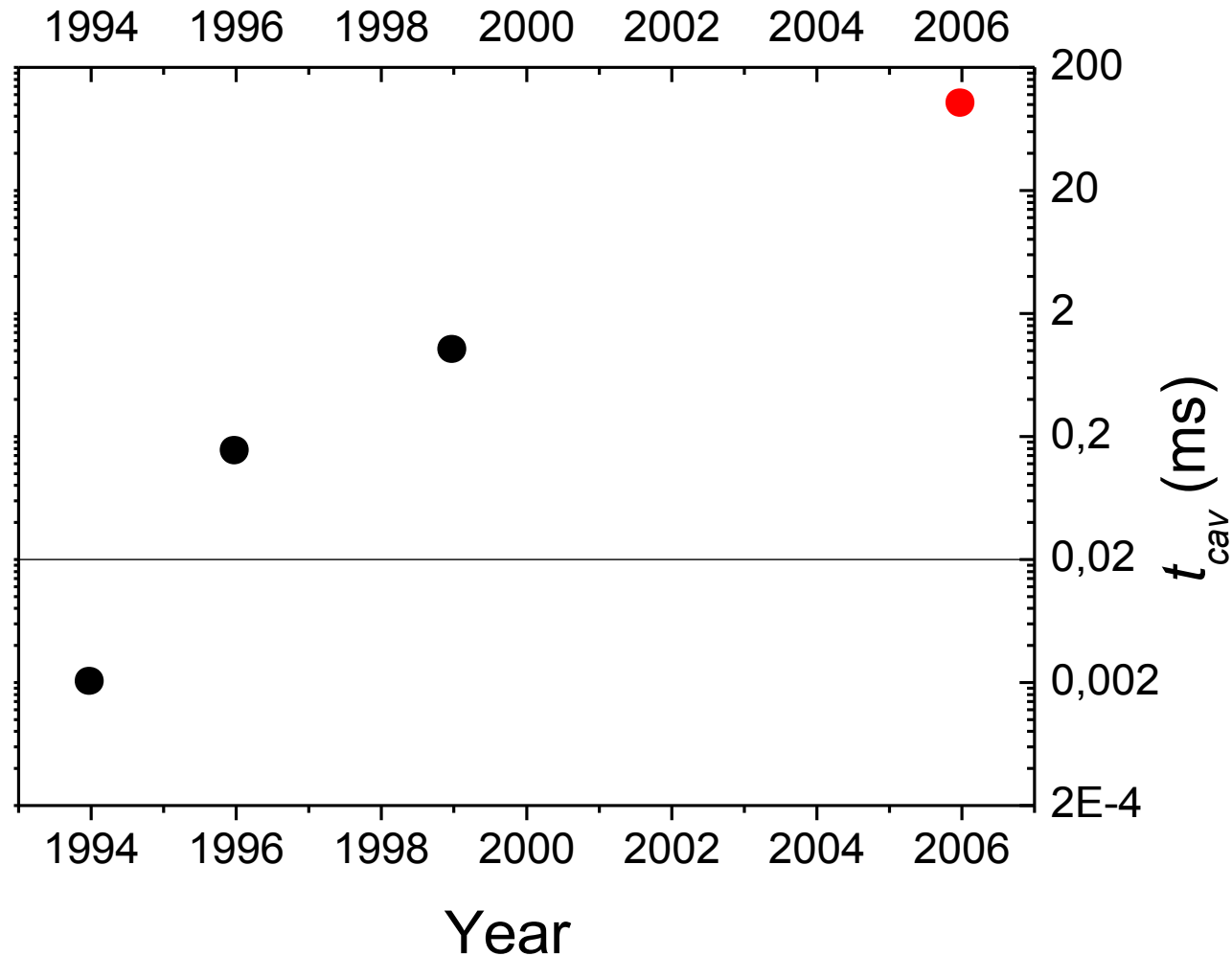
I. Cavity QED spin and spring: Brief history

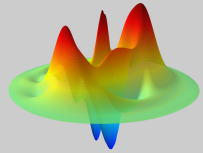




Cavity technology

- Our version of Moore's law:





The beginning of the story ...

VOLUME 65, NUMBER 8

PHYSICAL REVIEW LETTERS

20 AUGUST 1990

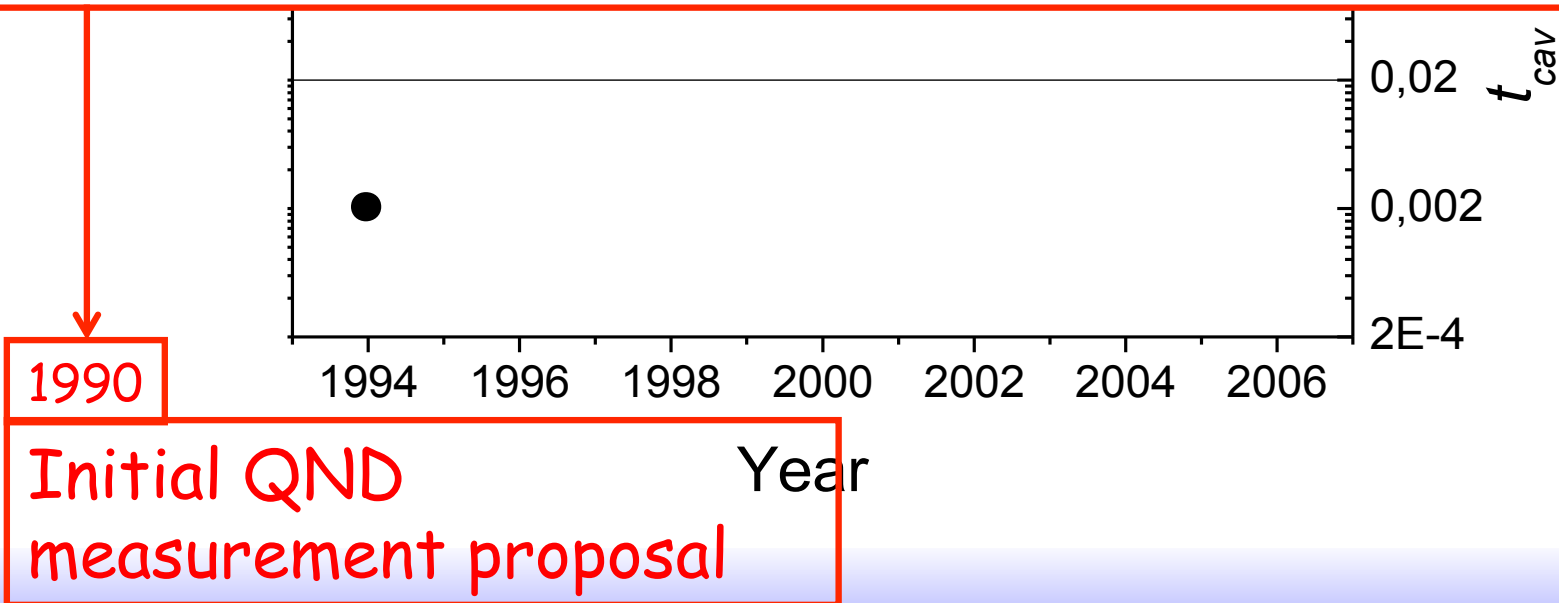
Quantum Nondemolition Measurement of Small Photon Numbers by Rydberg-Atom Phase-Sensitive Detection

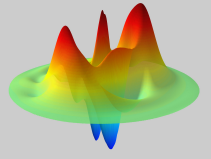
M. Brune, S. Haroche, V. Lefevre, J. M. Raimond, and N. Zagury^(a)

*Département de Physique de l'Ecole Normale Supérieure, Laboratoire de Spectroscopie Hertzienne,
24 rue Lhomond, F-75231 Paris CEDEX 05, France*

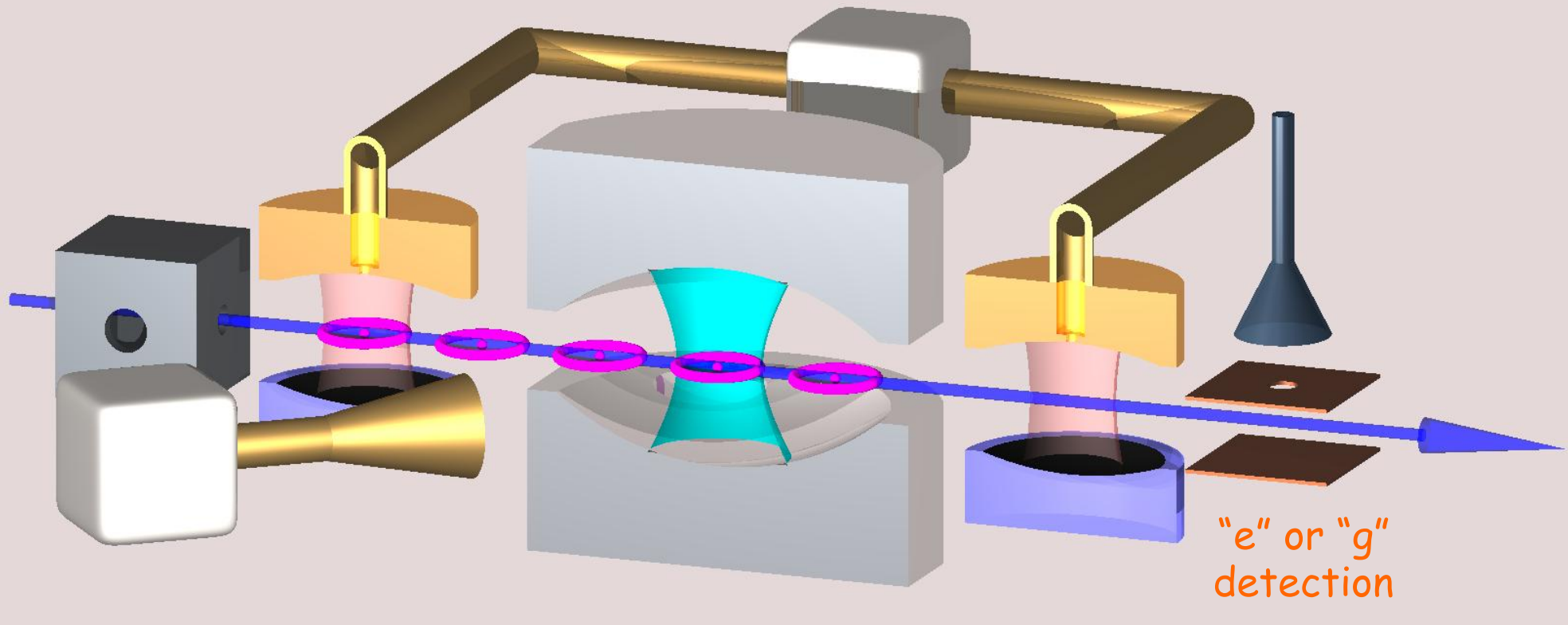
(Received 18 April 1990)

We describe a new quantum nondemolition method to monitor the number N of photons in a microwave cavity. We propose coupling the field to a quasiresonant beam of Rydberg atoms and measuring the resulting phase shift of the atom wave function by the Ramsey separated-oscillatory-fields technique. The detection of a sequence of atoms reduces the field into a Fock state. With realistic Rydberg atom-cavity systems, small-photon-number states down to $N=0$ could be prepared and continuously monitored.

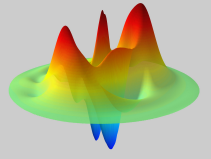




... with a simple experimental scheme

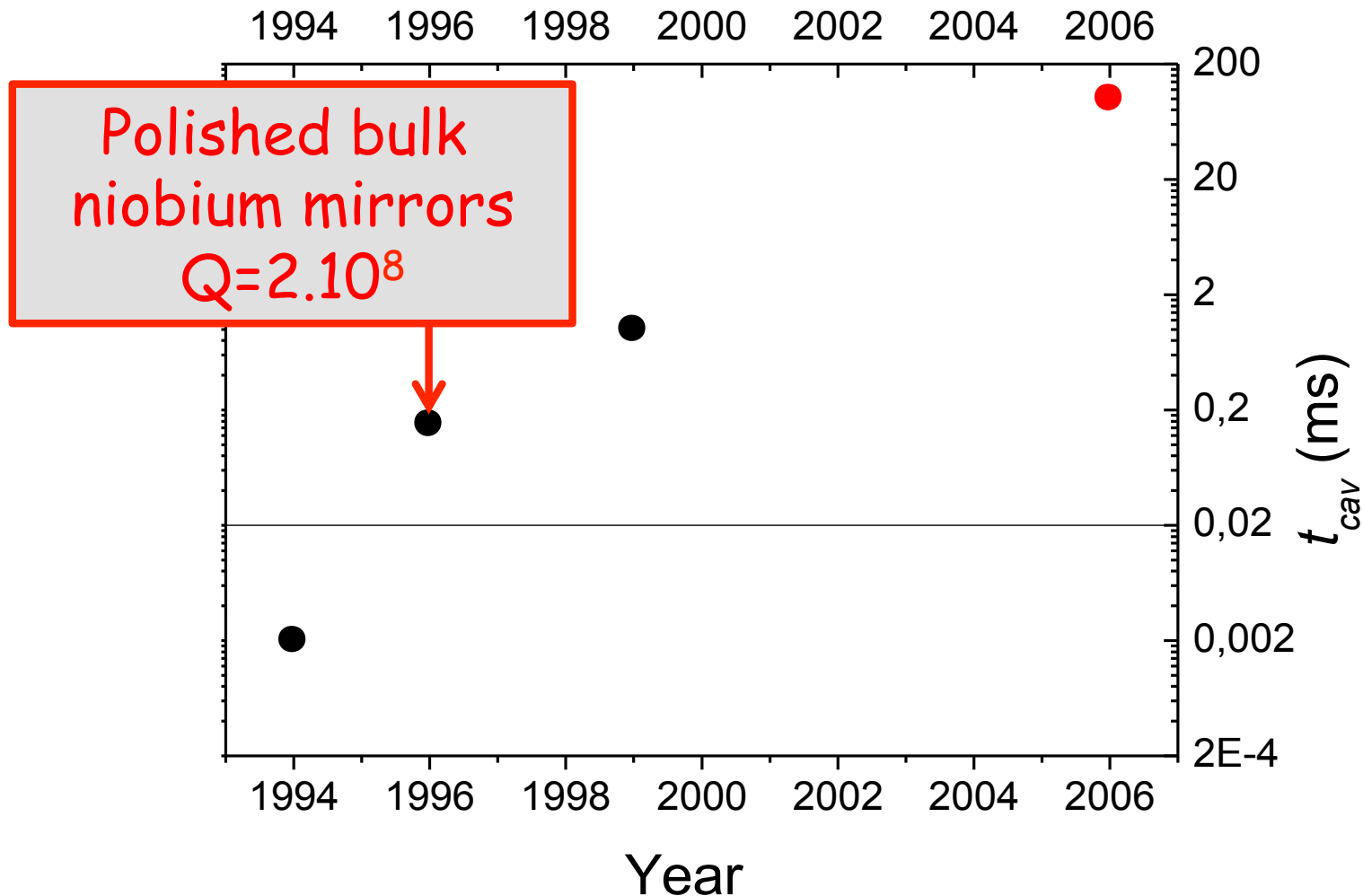


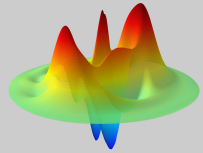
- An atomic clock (Ramsey setup) made of Rydberg for probing light-shifts induced by "trapped" photons
- State selective detection of atoms by field ionization: Atoms detected on "e" or "g" one by one



Cavity technology

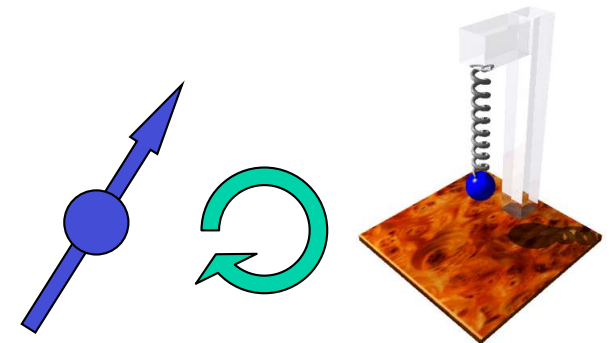
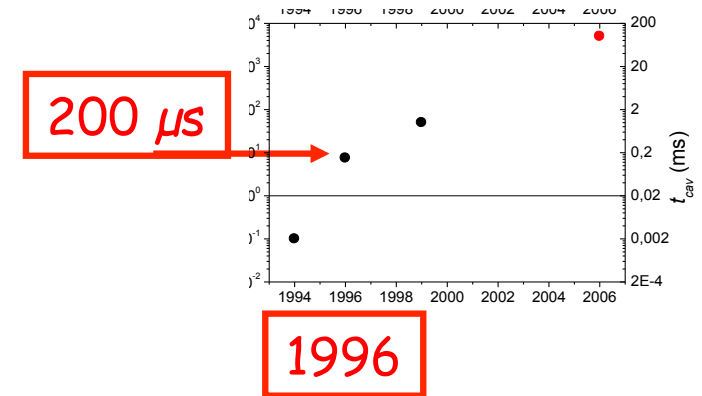
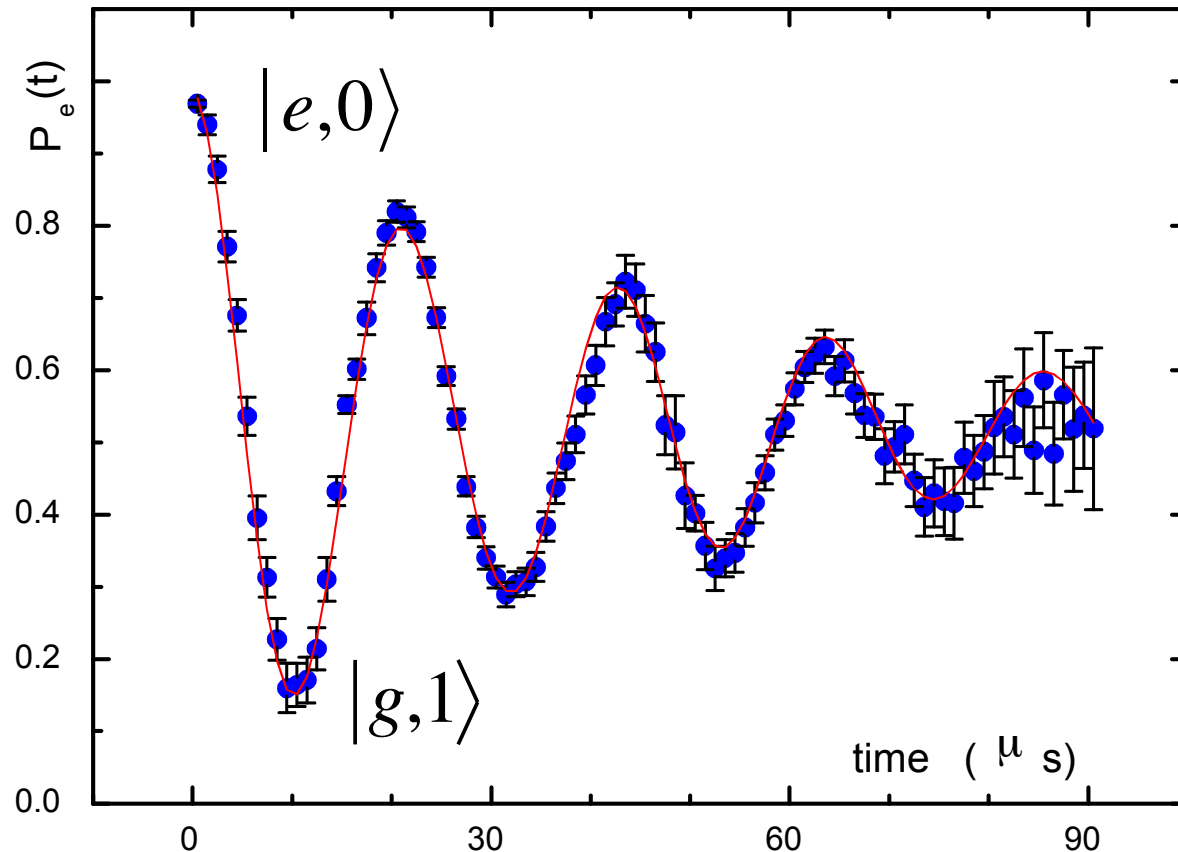
- Our version of Moore's law:



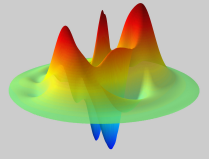


Strong coupling: quantum Rabi oscillations

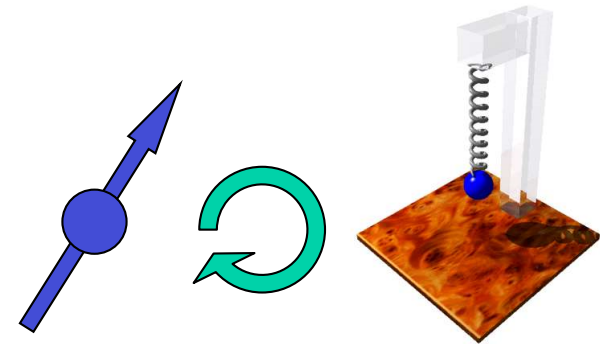
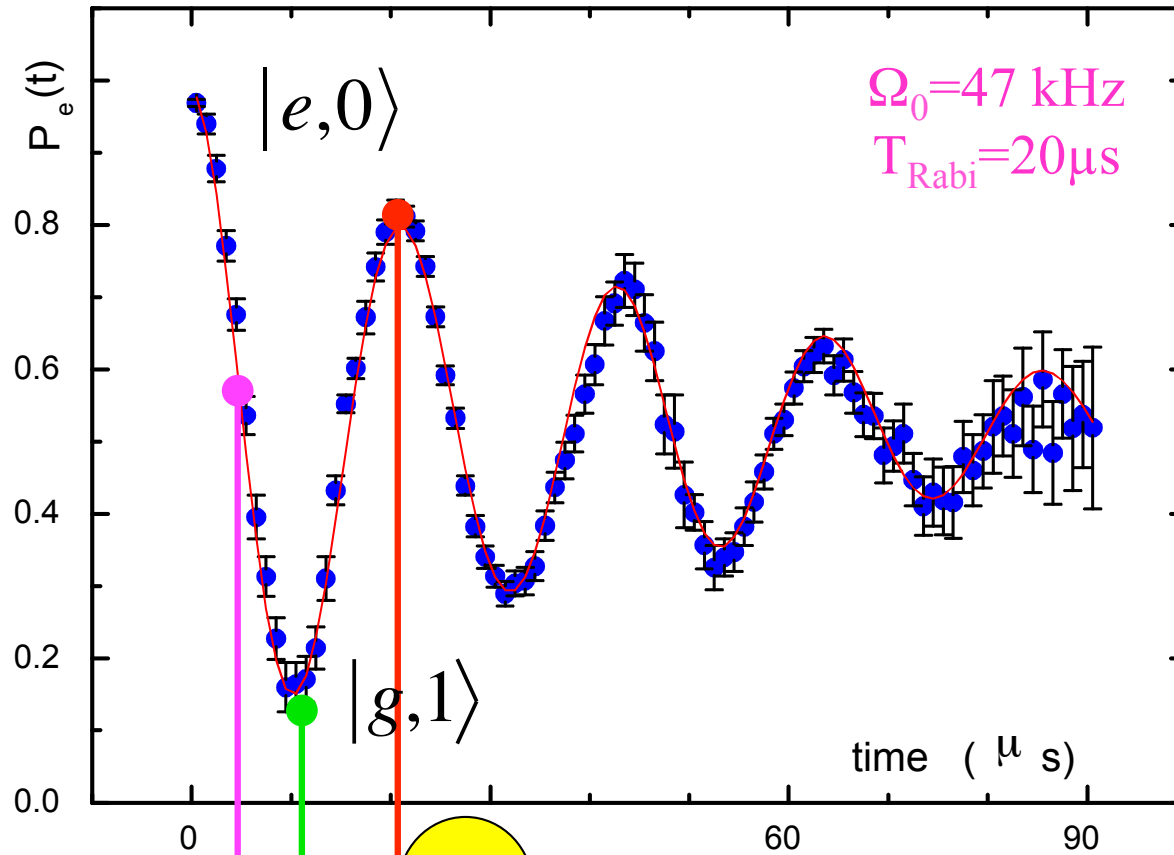
- Resonant atomic emission in an empty cavity



Resonant energy exchange like with coupled oscillators



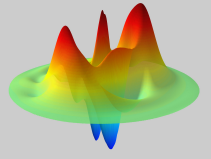
Quantum logic operations



2π • Phase gate, QND detection of a single photon (but not twice)

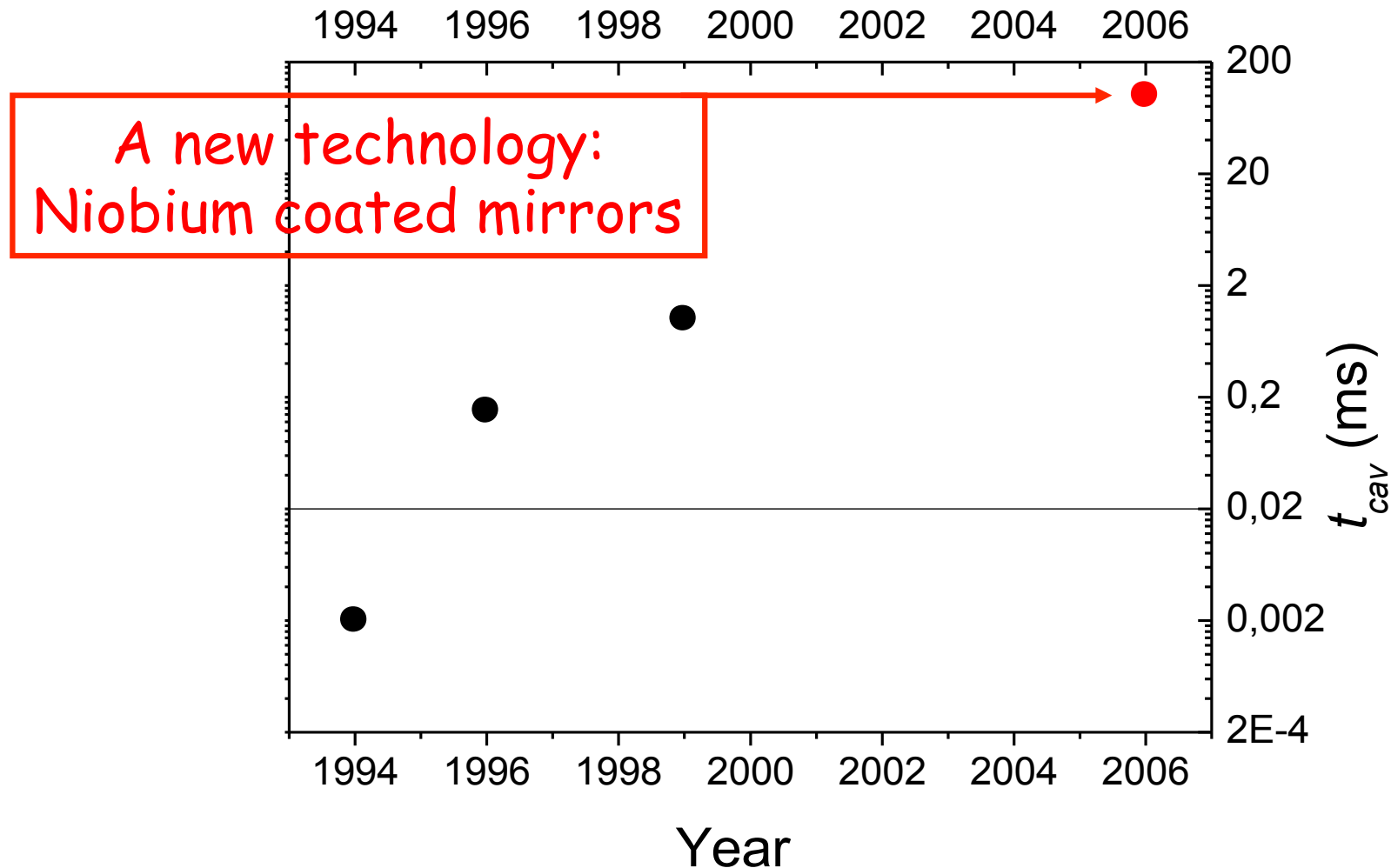
π • Atom-field state exchange

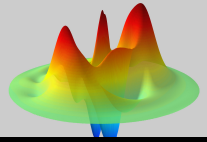
$\pi/2$ • EPR pair preparation



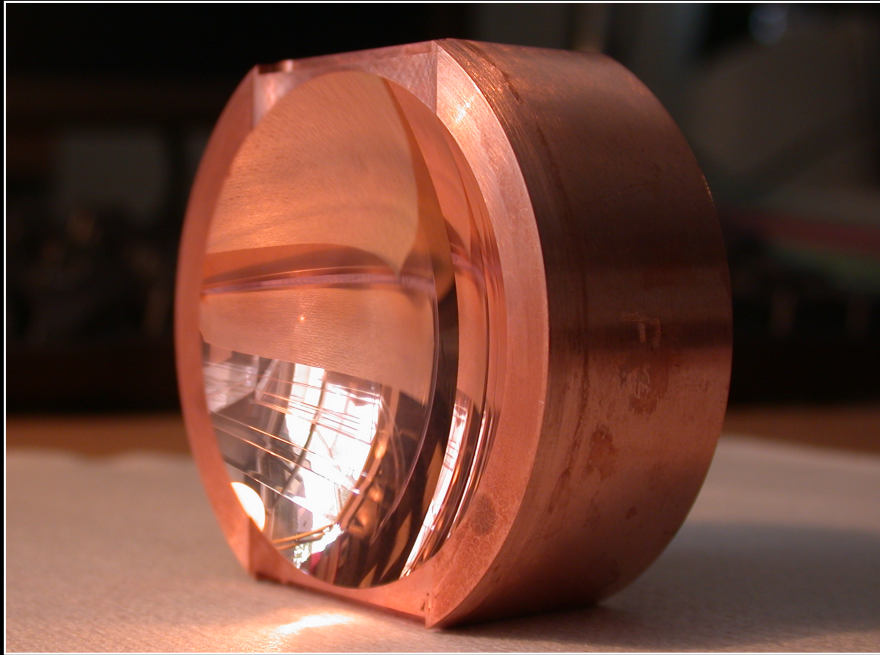
Cavity technology

- Our version of Moore's law:



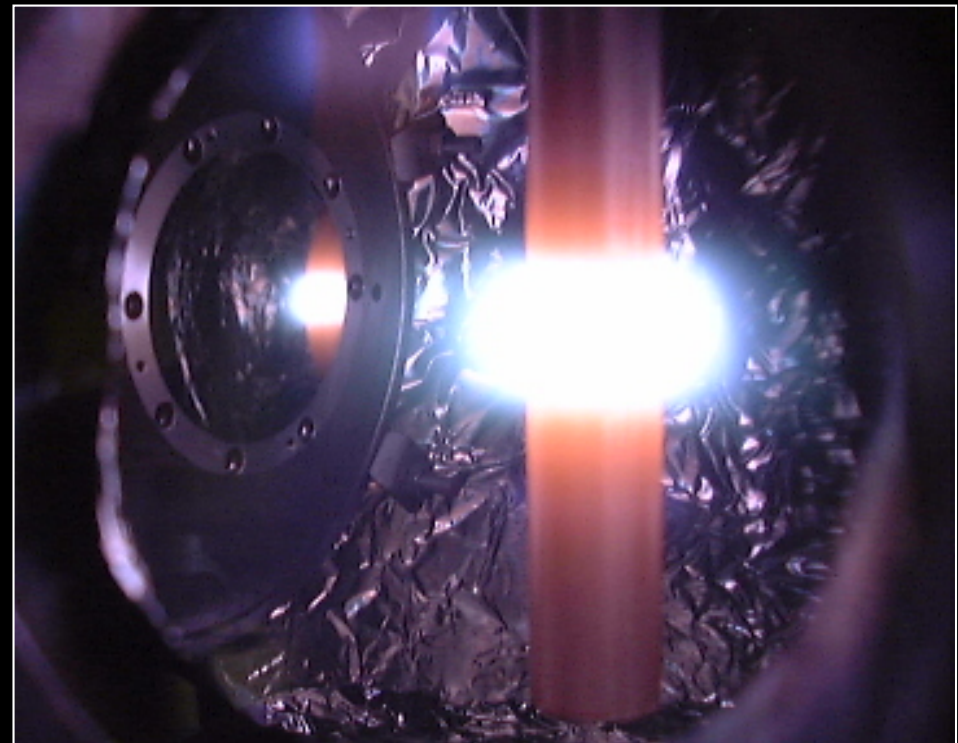


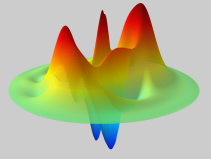
Niobium coated copper mirrors



- Copper mirrors
 - Diamond machined
 - $\sim 1 \mu\text{m}$ pTV form accuracy
 - $\sim 10 \text{ nm}$ roughness
- Toroidal $\rightarrow$ single mode

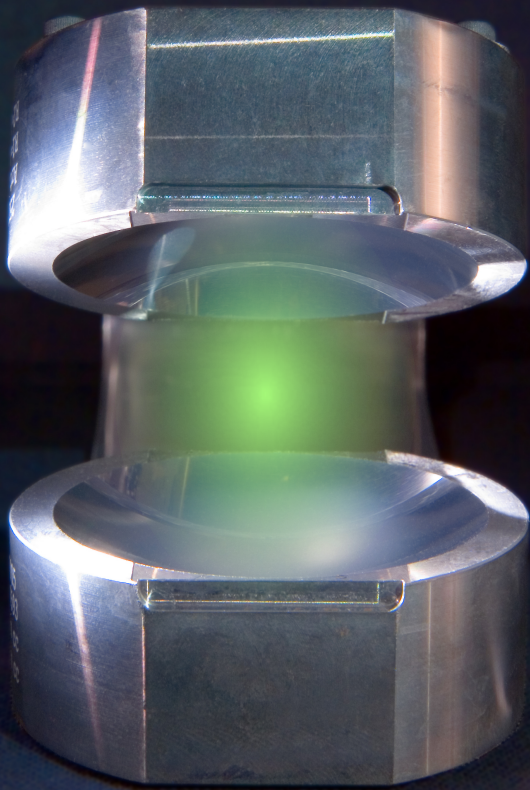
- Sputter $12 \mu\text{m}$ of Nb
 - Particles accelerator technique
 - Process done at CEA, Saclay
 - [E. Jacques, B. Visentin, P. Bosland]*





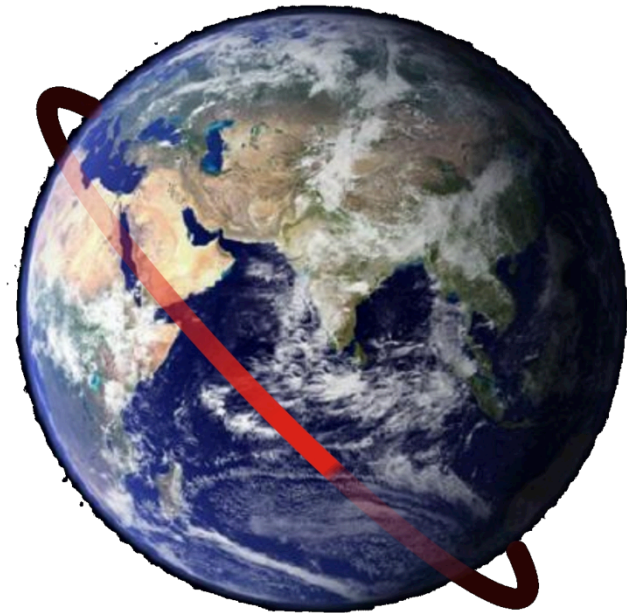
The photon box

- Photon box
Superconducting cavity



$$T_{\text{cav}} = 130 \text{ ms}$$

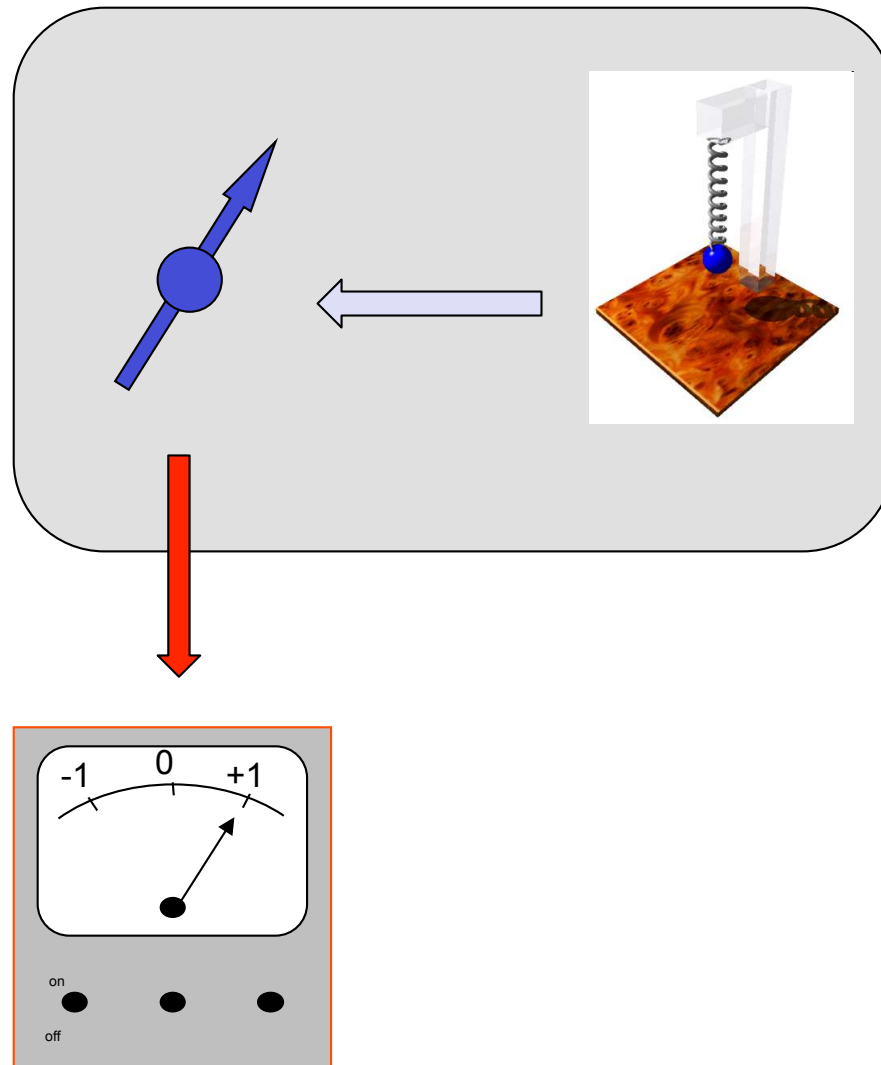
- Resonance @ $\nu_{\text{cav}} = 51 \text{ GHz}$
- Q factor = $4.2 \cdot 10^{10}$
- finesse = $4 \cdot 10^9$

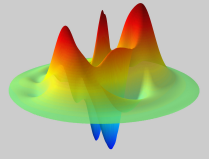


Photons running for 39 000 km
in the box before dying!

II - QND photon counting in the box

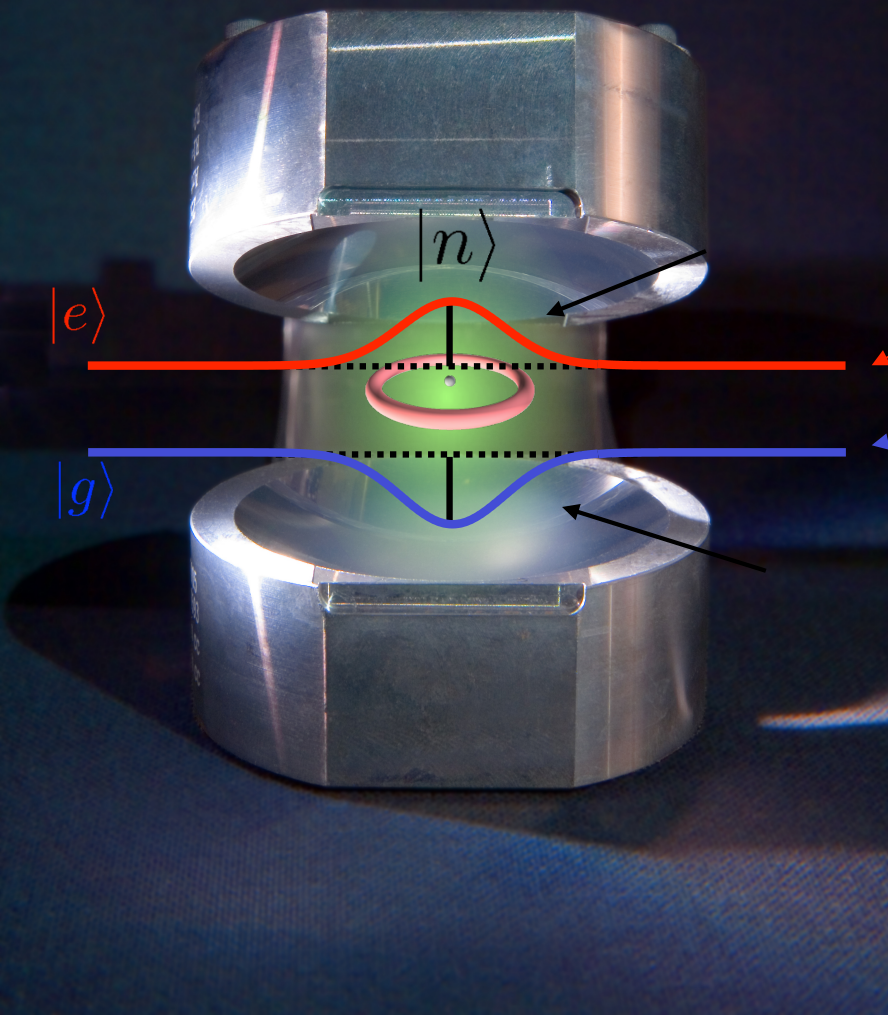
Measuring the "spring" with the "spin"





QND detection of photons: the tools

- Photon box
Superconducting cavity



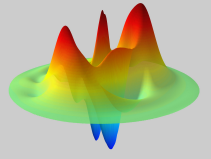
- Photon probes
Circular Rydberg atoms
- Non-resonant interaction
 $\Rightarrow$ light shifts

$$\Delta E_e = \hbar \frac{\Omega_0^2}{4\delta} (n + 1)$$

$$\Delta E_g = -\hbar \frac{\Omega_0^2}{4\delta} n$$

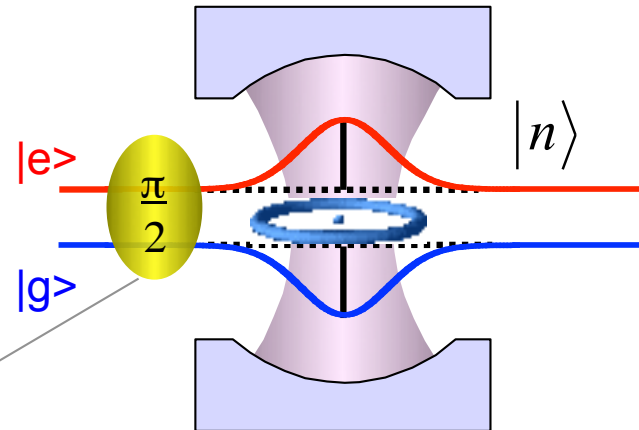
Atoms used as clock
for counting n by
measuring light shifts



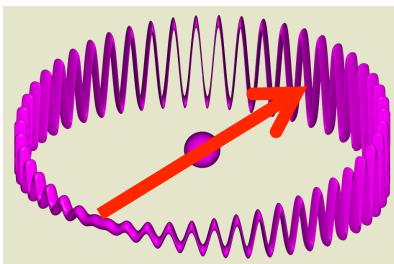
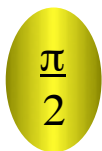
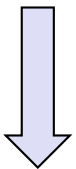
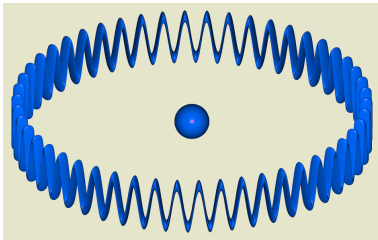


Detecting 0 or 1 photon

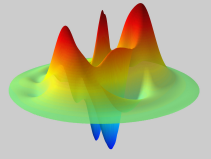
1. Trigger of the clock.



g

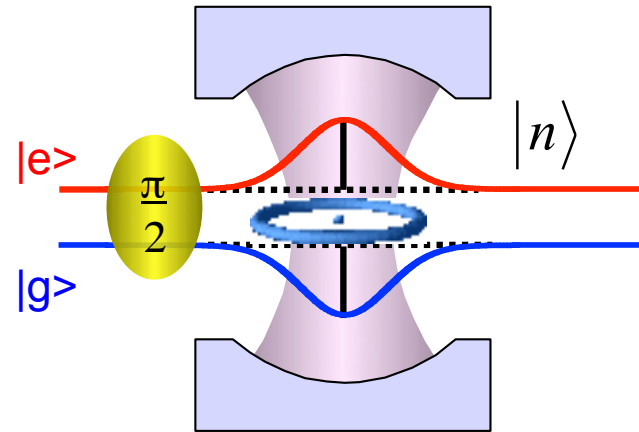


The $\pi/2$ pulse prepares a superposition states of "e" and "g". This corresponds to an induced atomic dipole oscillating at the atomic frequency.

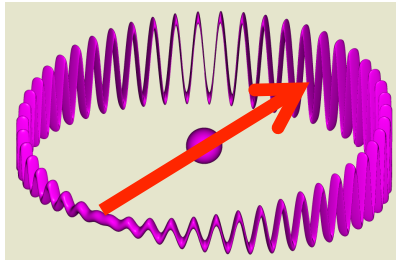
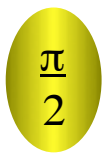
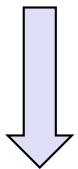
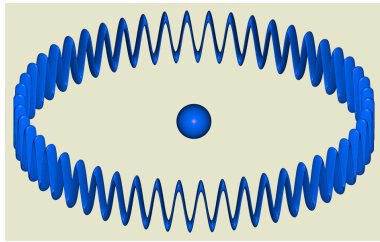


Detecting 0 or 1 photon

1. Trigger of the clock.
2. Let the dipole precess

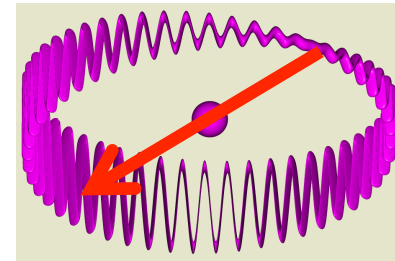


g

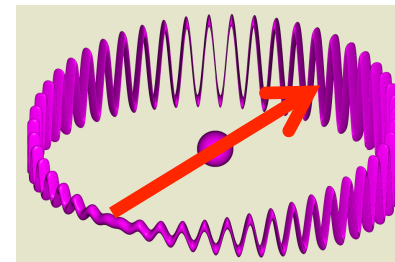


Single photon
phase shift

$$\Phi_0 = \pi$$

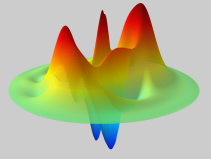


1 photon



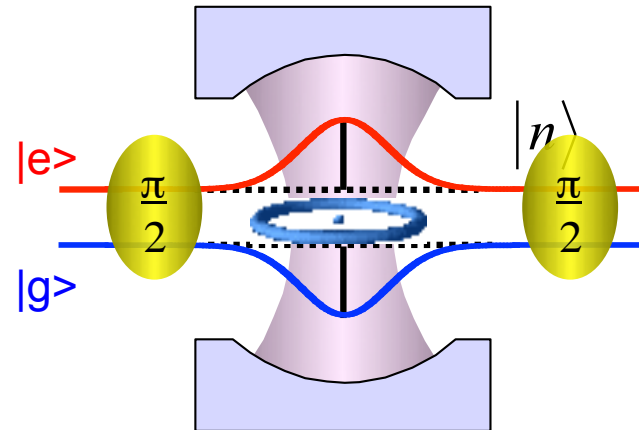
0 photon

Due to "light-shifts", dispersive interaction with the cavity delays the dipole oscillation.

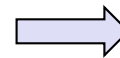
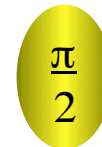
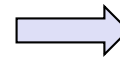
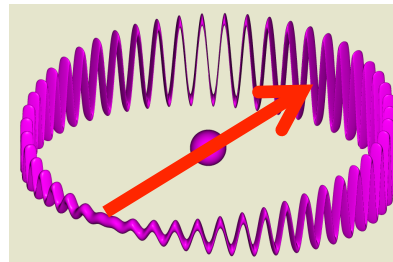
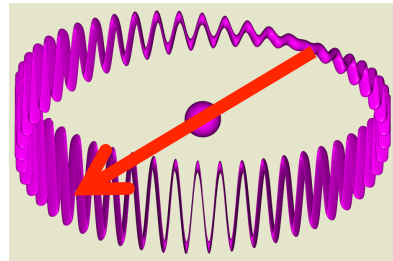
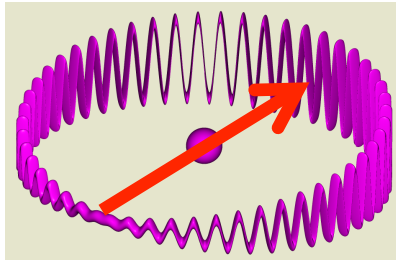
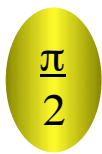
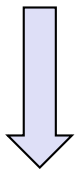
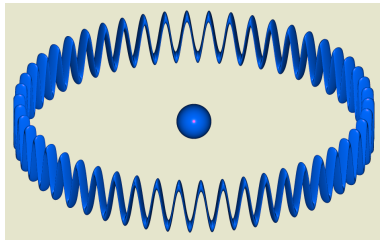


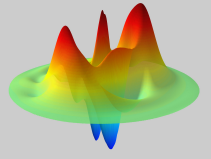
Detecting 0 or 1 photon

1. Trigger of the clock.
2. Let the dipole precess
3. Detect the atom



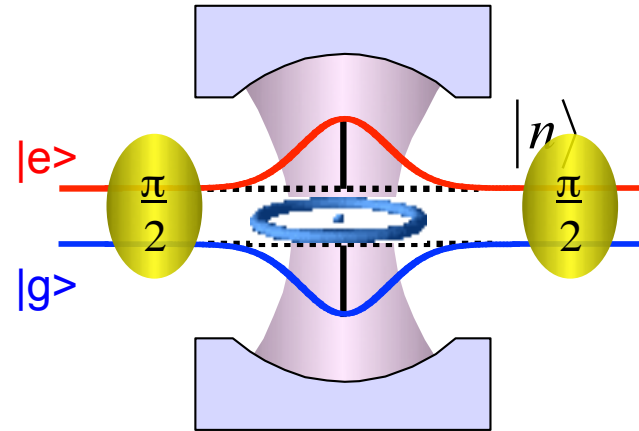
g



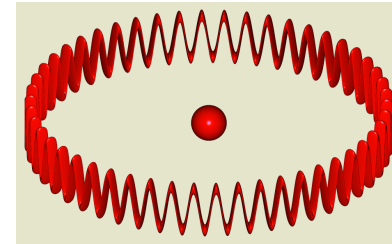
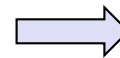
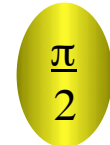
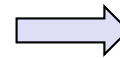
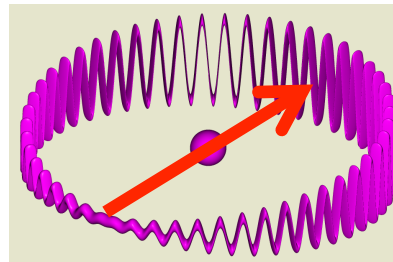
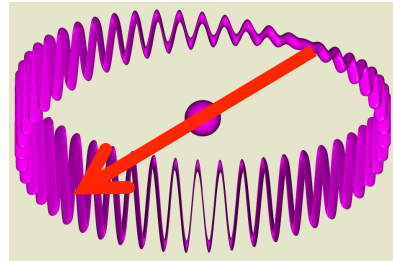
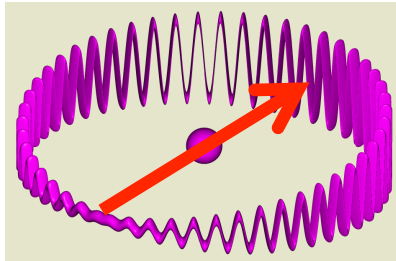
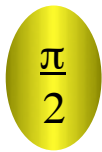
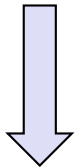
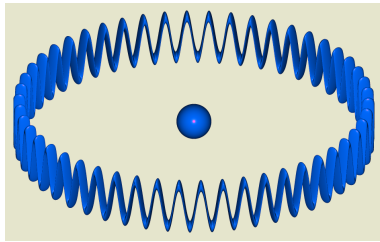


Detecting 0 or 1 photon

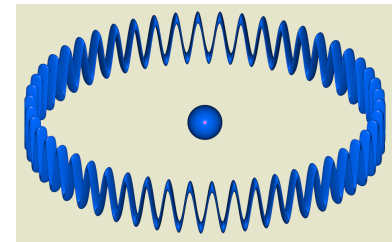
1. Trigger of the clock.
2. Let the dipole precess
3. Detect the atom



g

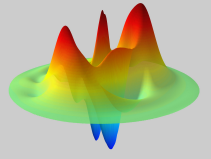


e
1 photon



g
0 photon

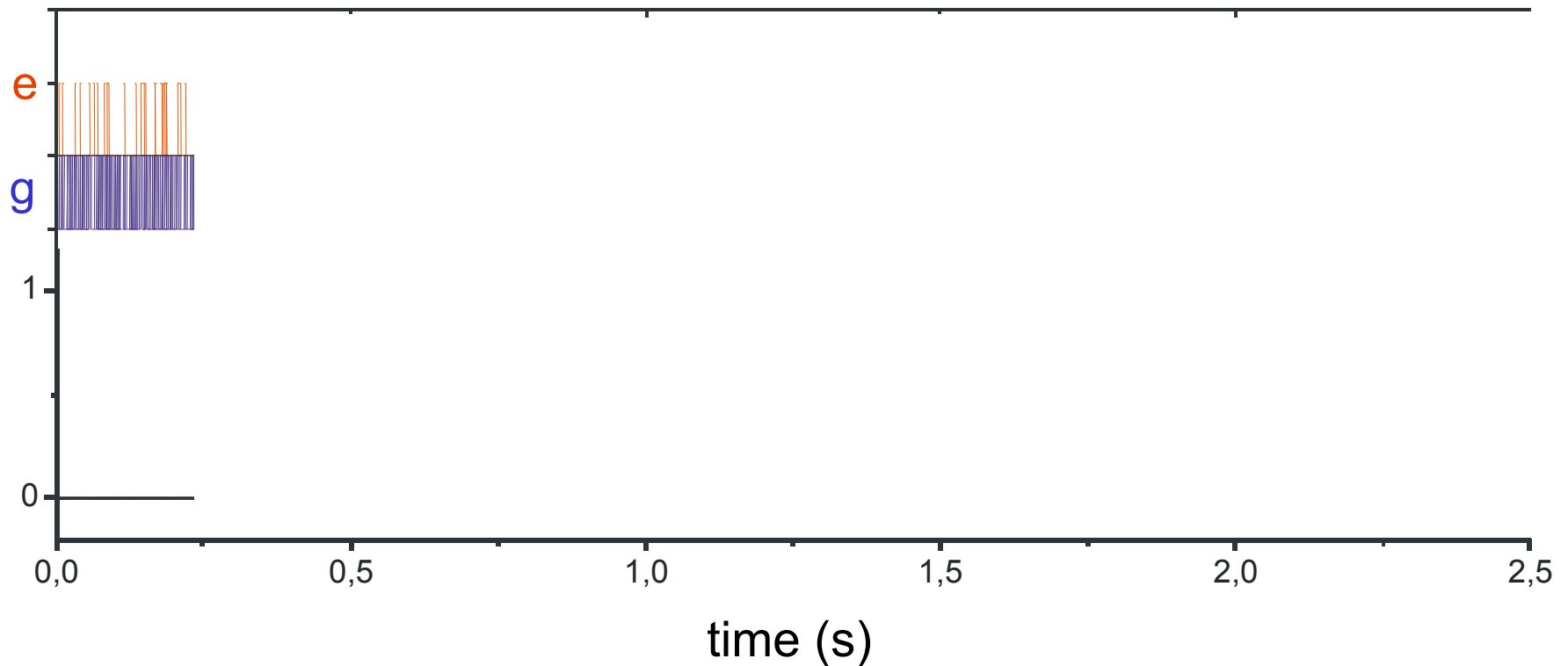
Second $\pi/2$ pulse converts phase information into atomic energy state



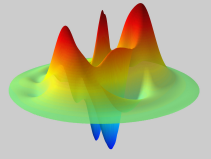
Detecting blackbody photons

$g \rightarrow$ field projected on $|0\rangle$

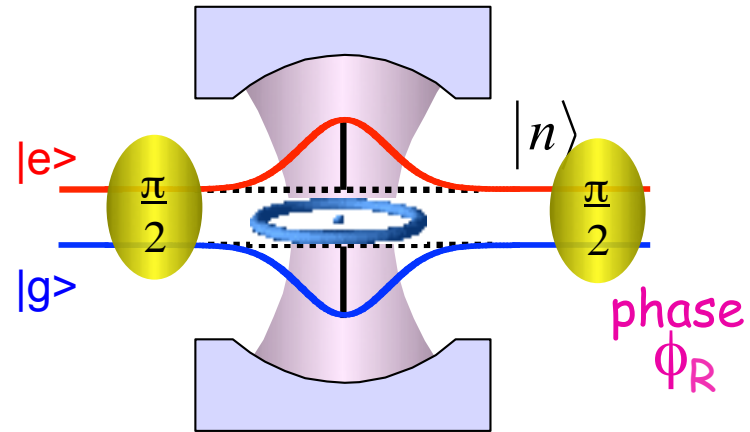
$e \rightarrow$ field projected on $|1\rangle$



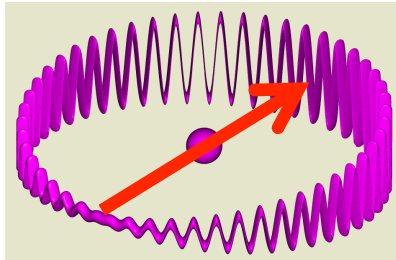
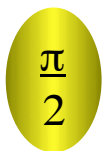
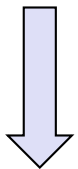
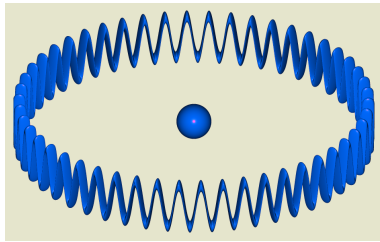
$$T = 0.8 \text{ K} \rightarrow n_{th} = 0.05 \quad (\text{proba. of } n=2 \text{ is negligible})$$



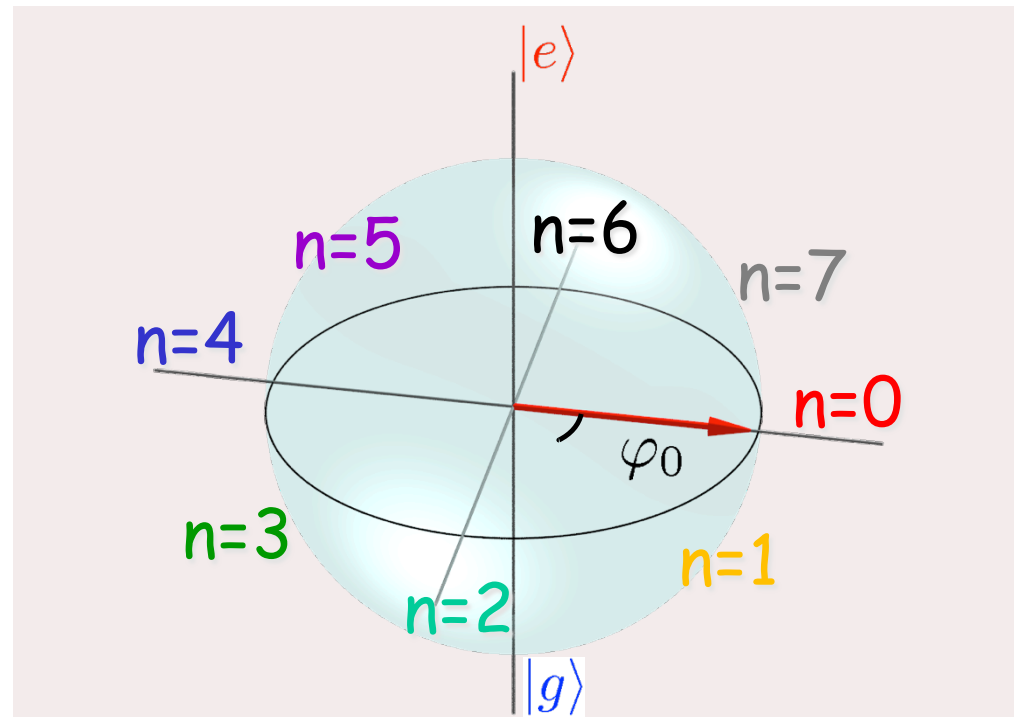
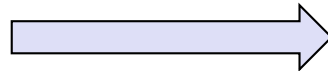
"Seeing" more photons



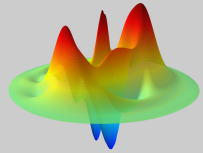
g



$$\Phi_0 = \pi/4$$

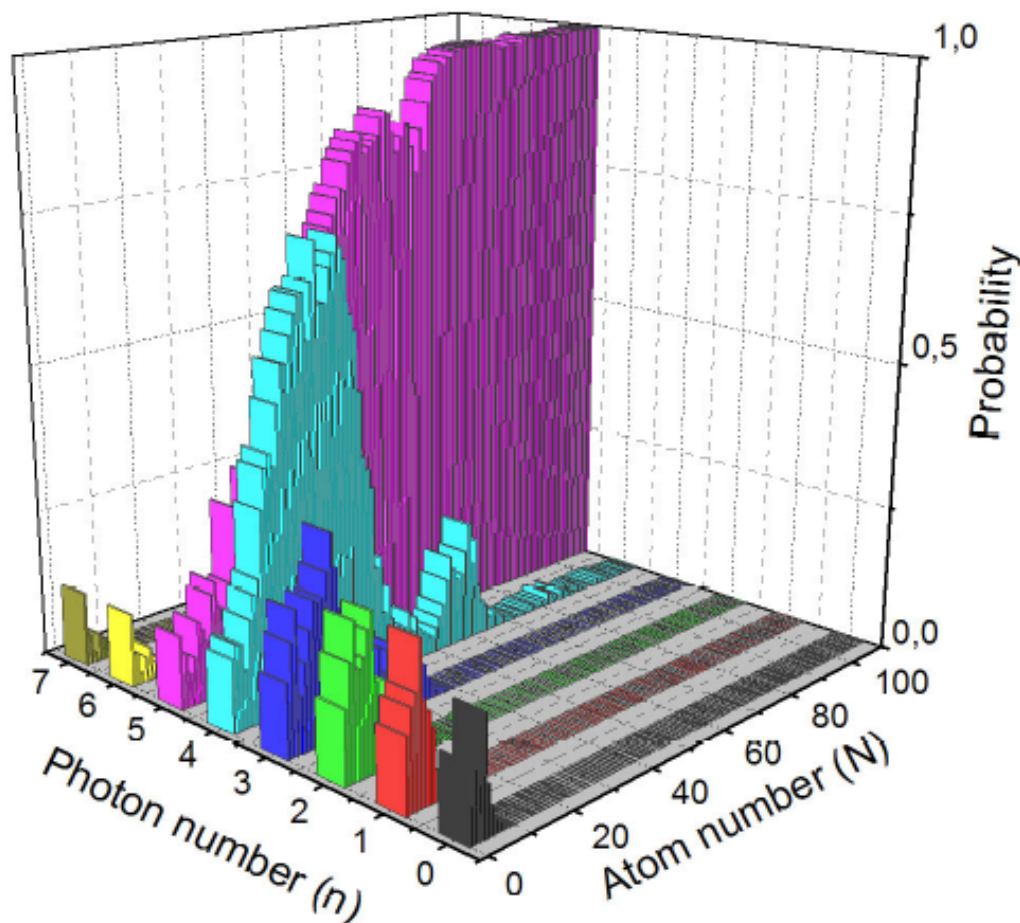


Depending on the photon number, the atomic dipole is the hand of a clock which can point in 7 different direction.



Atom by atom analysis of the measurement process

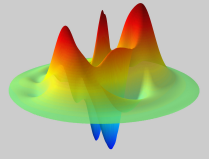
For each detected atom, one projects the field state according to the measurement result e or g



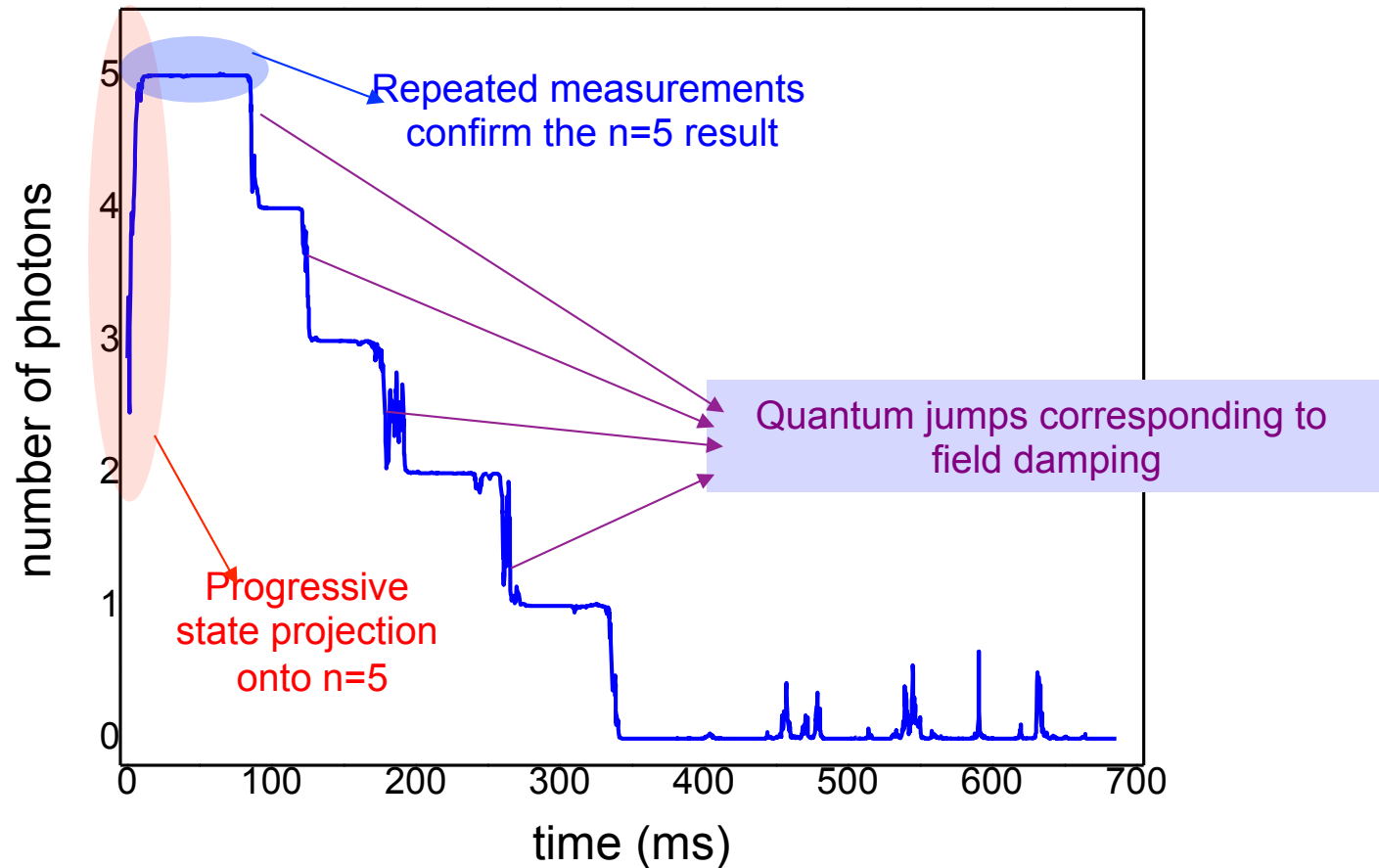
Measurement of a
coherent field
 $\langle n \rangle = 3.7 (\pm 0.008)$

Progressive collapse of
the field state on $n=5$

Initial knowledge of the
photon number
distribution is not
needed



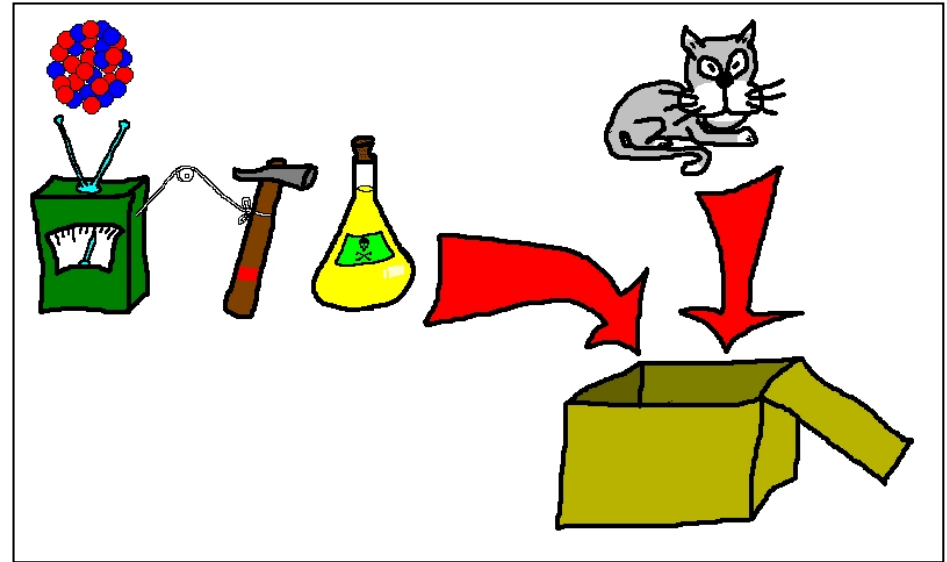
Repeated measurements: evolution of a continuously monitored field

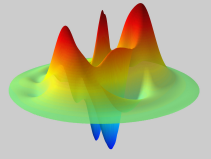


Field evolution due to cavity damping: not to QND measurement

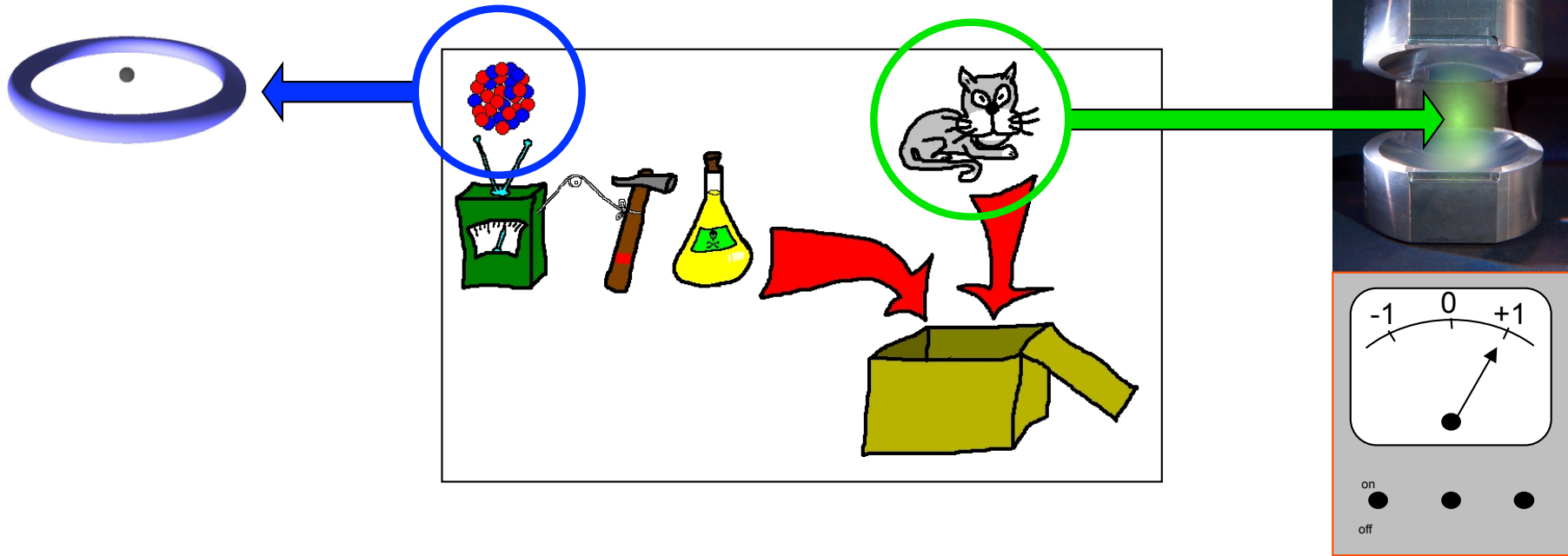
- Exhibits all features of quantum theory of measurement:
 - ❑ State collapse
 - ❑ Random result
 - ❑ repeatability

III. Preparing a Schrödinger cat state





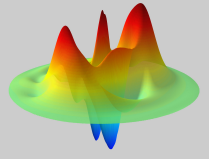
Measuring one atom **with the field**



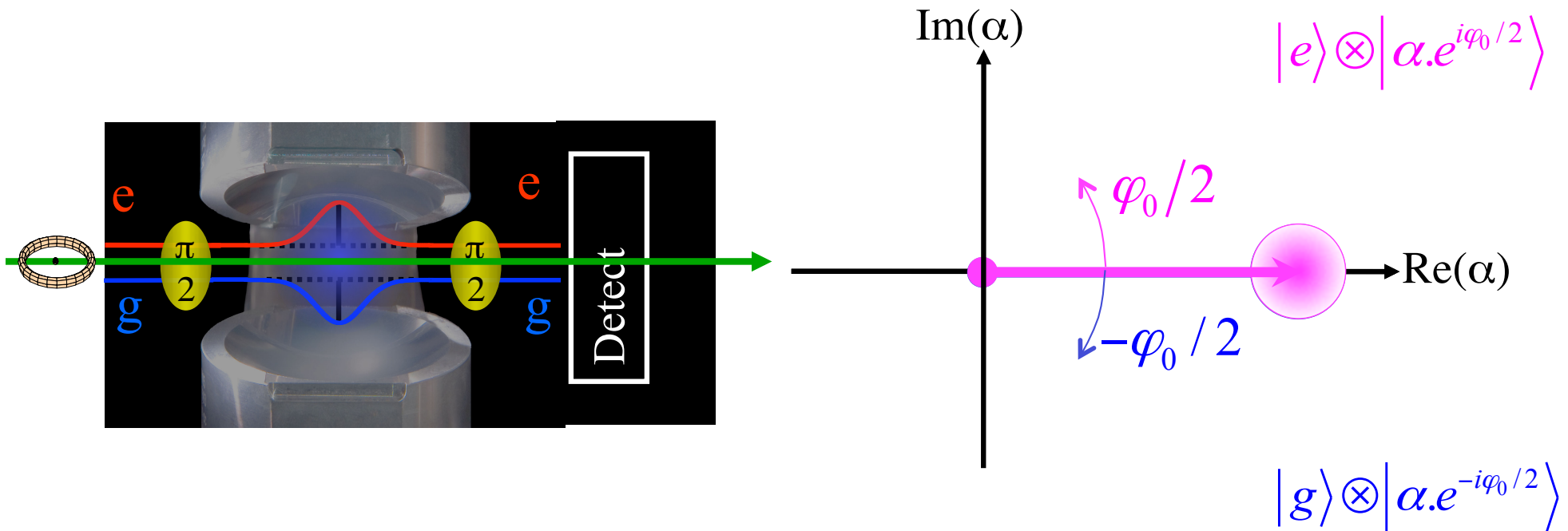
Distribution of roles:

- The radioactive atom: one Rydberg atom
- The cat: a mesoscopic field in the cavity.

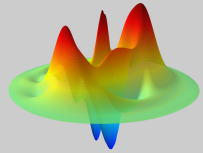
The "size" (i.e. the photon number) of the cat can be varied for exploring the quantum to classical transition



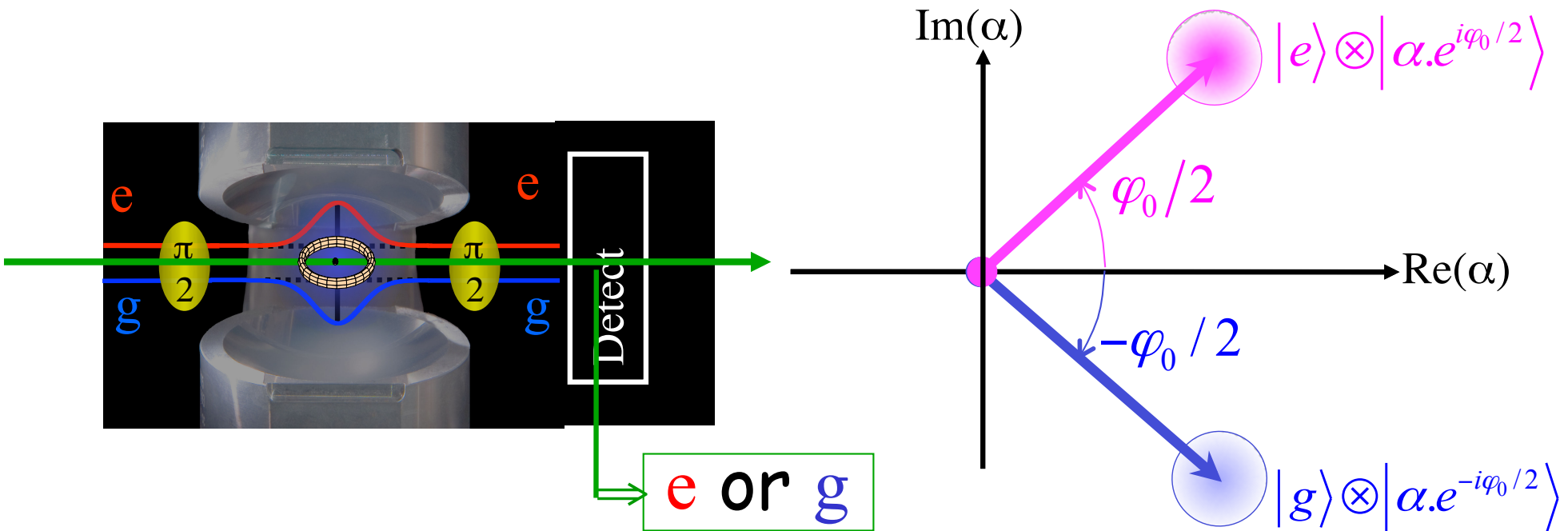
Preparing a phase Schrödinger cat state



$\Rightarrow$ the field phase is controlled by the atomic energy state



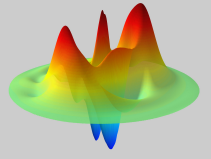
Preparing a phase Schrödinger cat state



Detection of the atom projects the field on:

$$\frac{1}{\sqrt{2}} \left(|\alpha.e^{i\varphi_0/2}\rangle \pm |\alpha.e^{-i\varphi_0/2}\rangle \right)$$

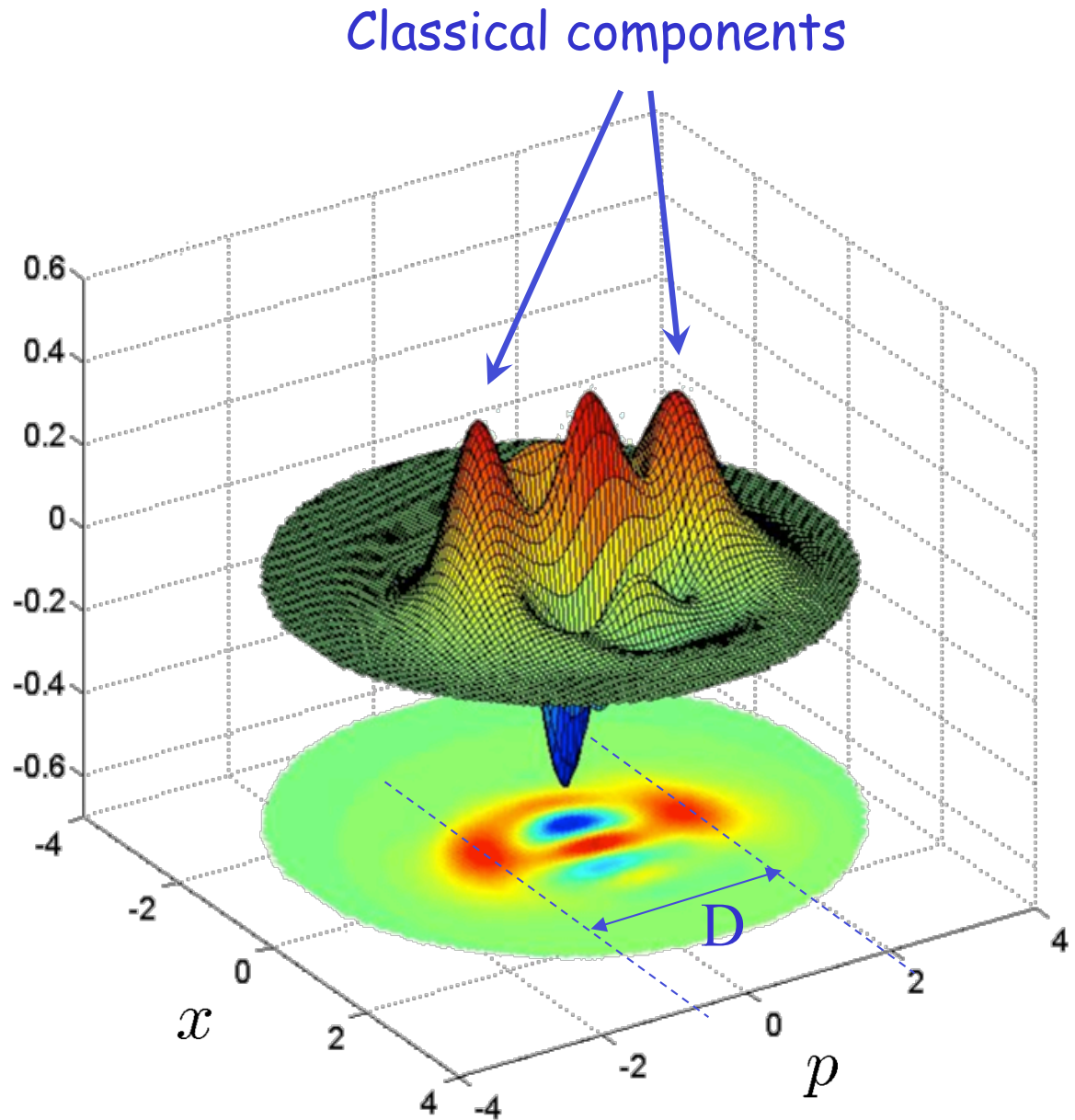
$\pm$: depends on detected state e or g

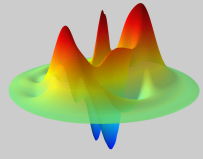


Reconstructed Wigner function

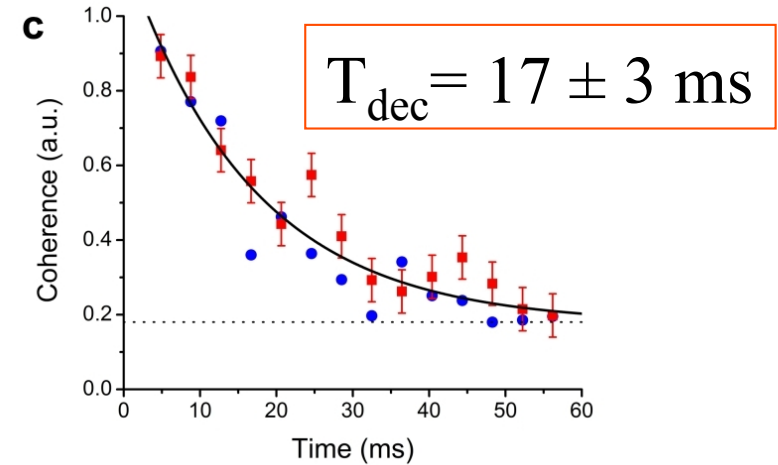
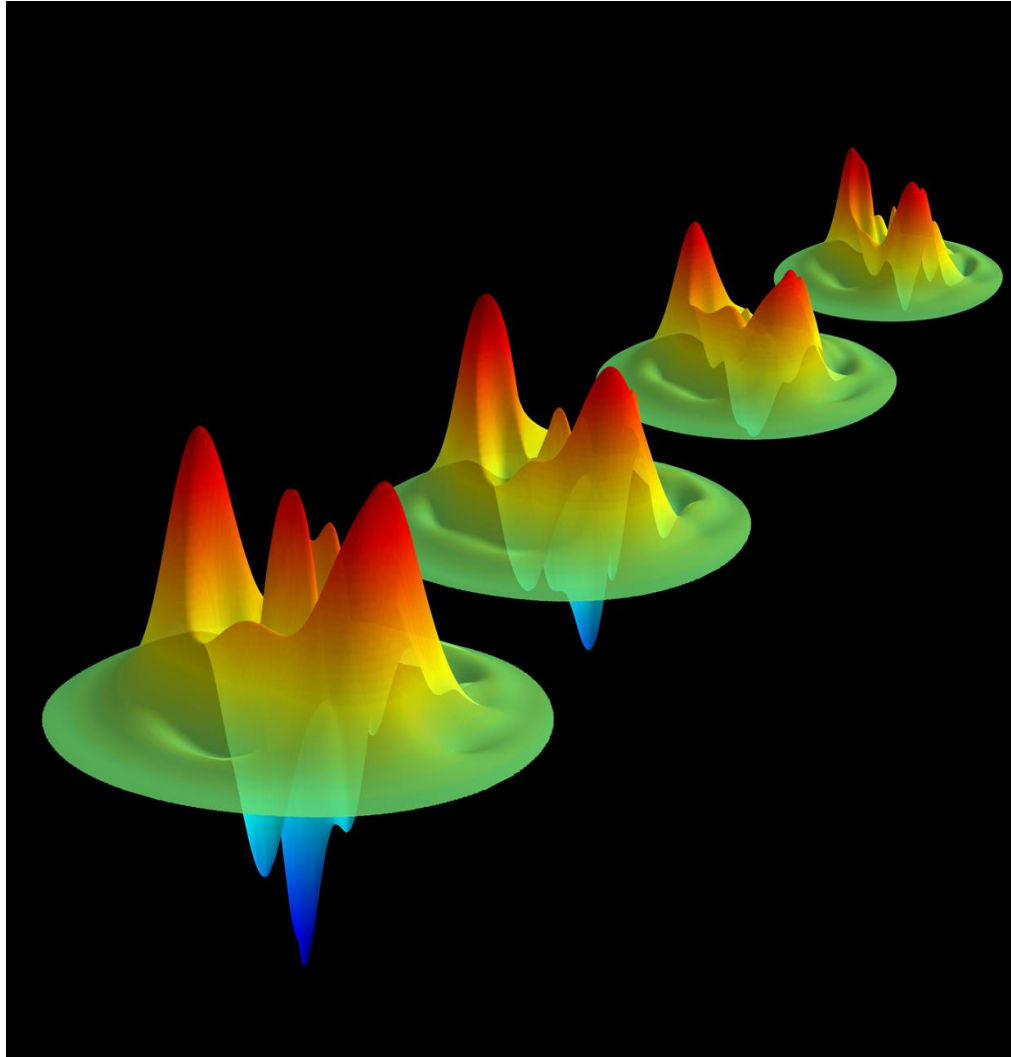
cat size:
 $D^2 \approx 7.5$ photons

coherent
components are well
separated:
 $D^2 \gg 1$





Decoherence of a $D^2=11.8$ photon cat state



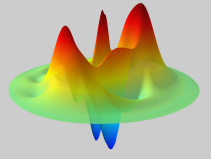
$$T_{\text{meas}} \approx 4 \text{ ms} < T_{\text{dec}, \bar{n}_{\text{th}}}$$

Theory:

$$T_{\text{dec}} = 2T_{\text{cav}}/D^2 = 22 \text{ ms}$$

+ small blackbody
contribution @ 0.8 K

The future?

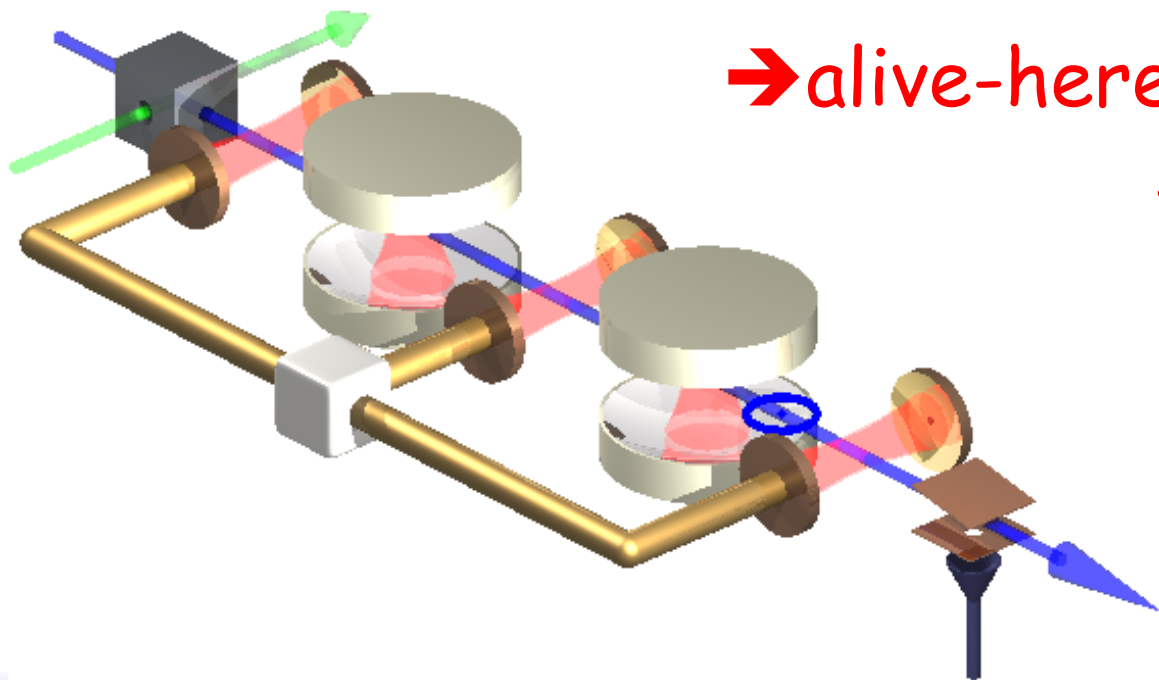


Perspectives

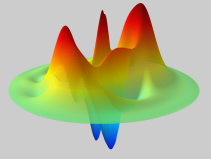
Experiments with 2 cavities

$$\frac{1}{\sqrt{2}} (|\alpha\rangle|0\rangle + |0\rangle|-\alpha\rangle)$$

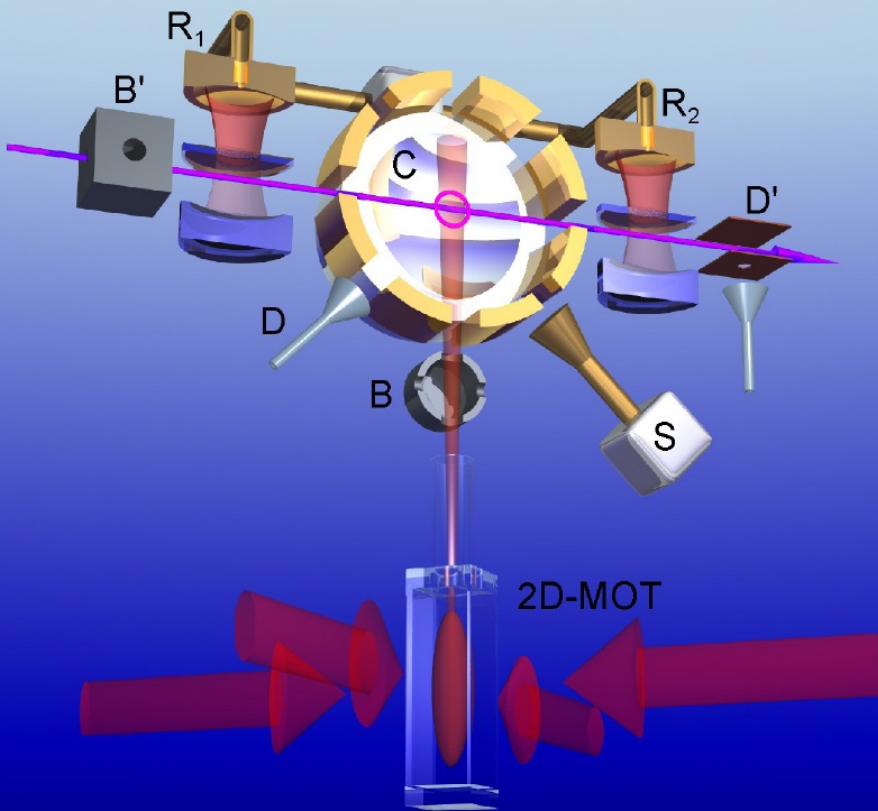
$$\frac{1}{\sqrt{2}} (| \text{alive-here} \rangle + | \text{dead-there} \rangle)$$



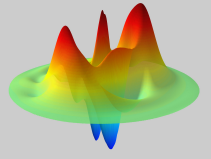
→ alive-here-and-dead-there state



Stationnary atoms in a cavity



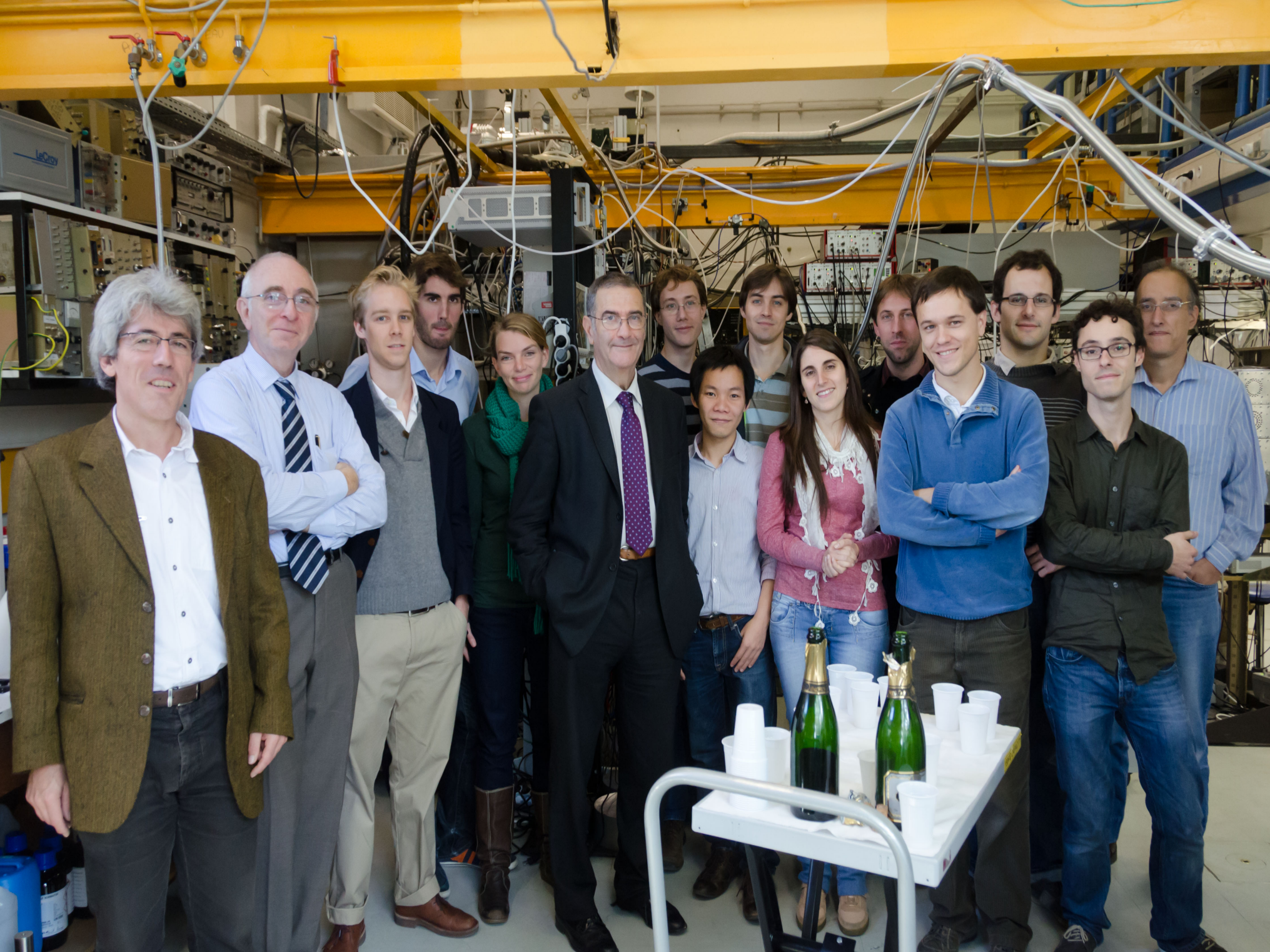
- Very long interaction times (few ms)
 - ⇒ Large cats
 - ⇒ Decoherence metrology
 - ⇒ Quantum random walks
 - ⇒ Quantum Zeno dynamics
 - ⇒ Reservoir engineering



And things would not have been like that...



... without Serge and Jean-Michel



The team

- Starring,
by order of apparition

- ❑ Serge Haroche
- ❑ Michel Gross
- ❑ Claude Fabre
- ❑ Philippe Goy
- ❑ Pierre Pillet
- ❑ Jean-Michel Raimond
- ❑ Guy Vitrant
- ❑ Yves Kaluzny
- ❑ Jun Liang
- ❑ Michel Brune
- ❑ Valérie Lefèvre-Seguín
- ❑ Jean Hare
- ❑ Jacques Lepape
- ❑ Aephraim Steinberg
- ❑ Andre Nussenzveig
- ❑ Frédéric Bernardot
- ❑ Paul Nussenzveig
- ❑ Laurent Collot
- ❑ Matthias Weidemüller
- ❑ François Treussart
- ❑ Abdelamid Maali
- ❑ David Weiss
- ❑ Vahid Sandoghdar
- ❑ Jonathan Knight
- ❑ Nicolas Dubreuil
- ❑ Peter Domokos
- ❑ Ferdinand Schmidt-Kaler
- ❑ Jochen Dreyer

- ❑ Peter Domokos
- ❑ Ferdinand Schmidt-Kaler
- ❑ Ed Hagley
- ❑ Xavier Maître
- ❑ Christoph Wunderlich
- ❑ Gilles Nogues
- ❑ Vladimir Ilchenko
- ❑ Jean-François Roch
- ❑ Stefano Osnaghi
- ❑ Arno Rauschenbeutel
- ❑ Wolf von Klitzing
- ❑ Erwan Jahier
- ❑ Patrice Bertet
- ❑ Alexia Auffèves
- ❑ Romain Long
- ❑ Sébastien Steiner
- ❑ Paolo Maioli
- ❑ Philippe Hyafil
- ❑ Tristan Meunier
- ❑ Perola Milman
- ❑ Jack Mozley
- ❑ Stefan Kuhr
- ❑ Sébastien Gleyzes
- ❑ Christine Guerlin
- ❑ Thomas Nirrengarten
- ❑ Cédric Roux
- ❑ Julien Bernu

- ❑ Ulrich Busk-Hoff
- ❑ Andreas Emmert
- ❑ Adrian Lupascu
- ❑ Jonas Mlynek
- ❑ Igor Dotsenko
- ❑ Samuel Deléglise
- ❑ Clément Sayrin
- ❑ Xingxing Zhou
- ❑ Bruno Peaudecerf
- ❑ Raul Teixeira
- ❑ Sha Liu (post-doc)
- ❑ Theo Rybarczyk
- ❑ Carla Hermann
- ❑ Adrien Signolles
- ❑ Adrien Facon
- ❑ Eva Dietsche
- ❑ Stefan Gerlich (post-doc)
- ❑ Than Long Nguyen

