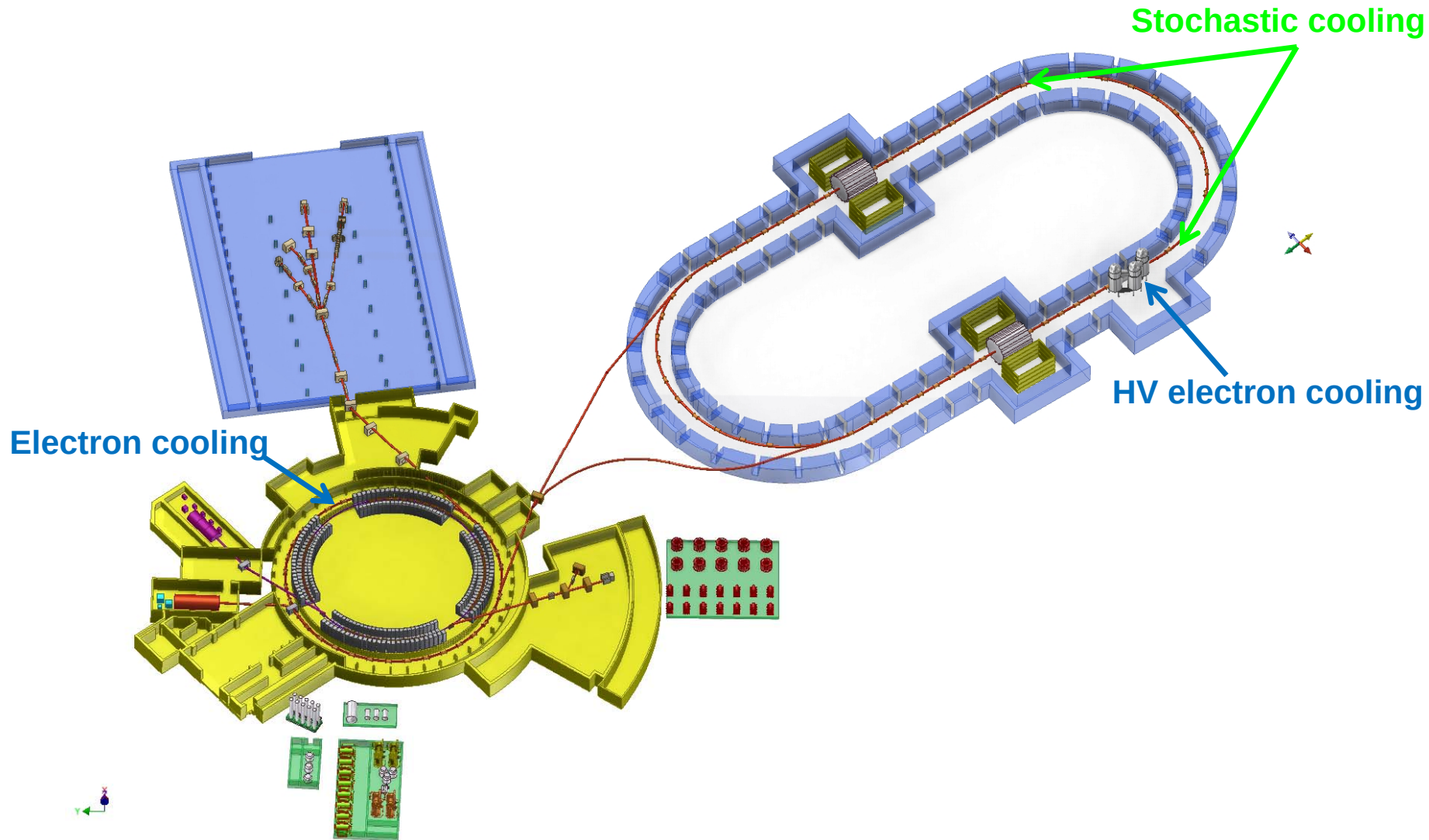


Application of cooling methods at NICA project

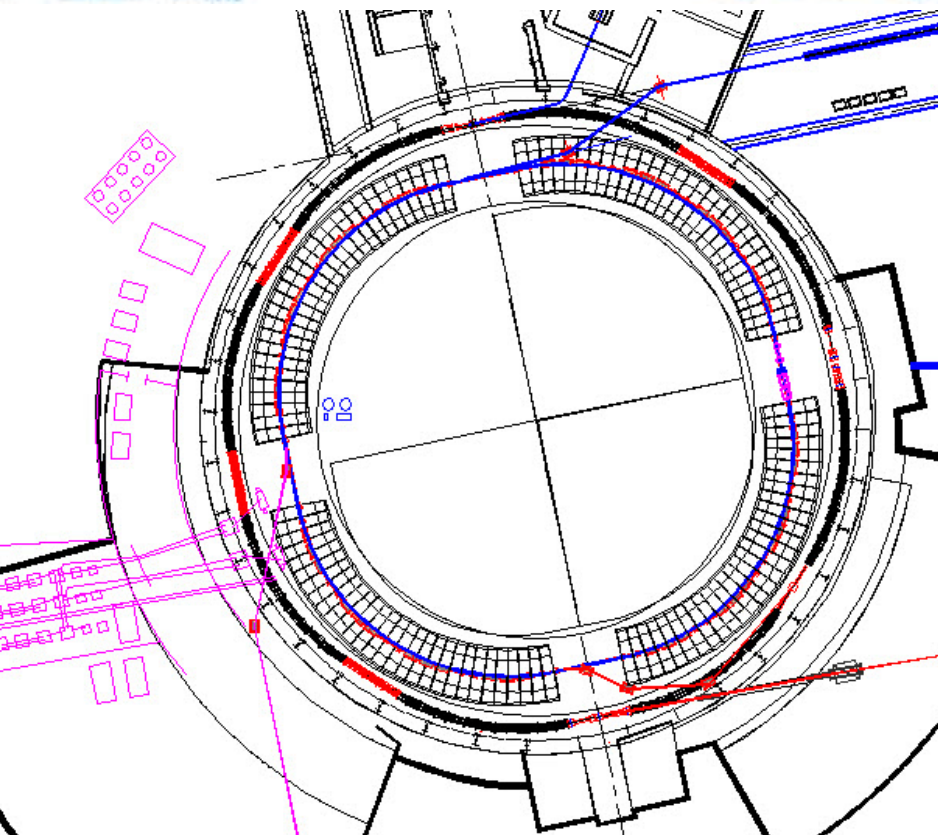
G.Trubnikov on behalf of team

*RUPAC-12 conference
St.Petersburg, 25 September 2012*

NICA complex



Booster synchrotron

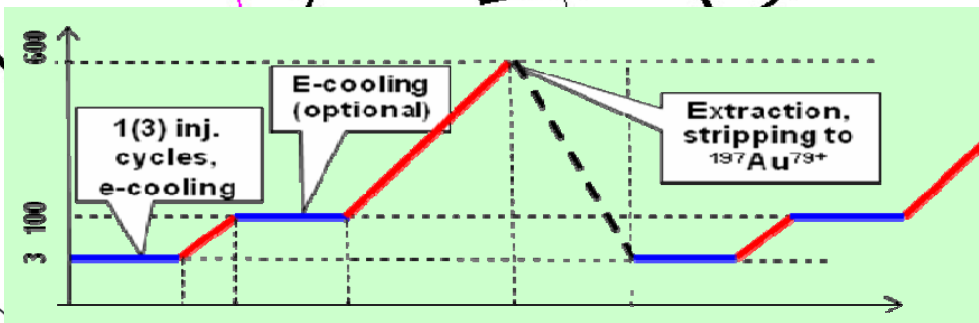
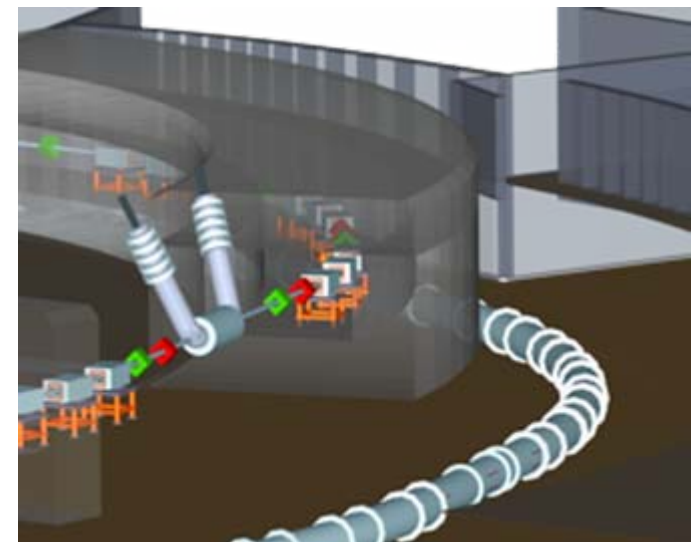
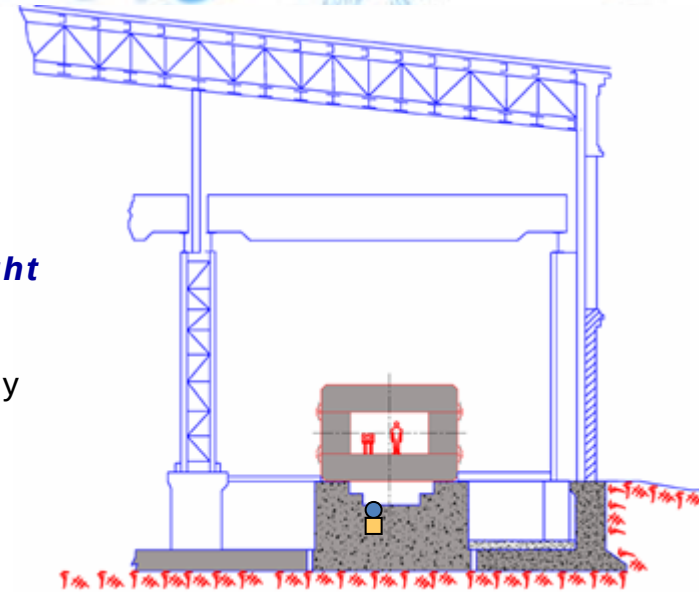


Injector

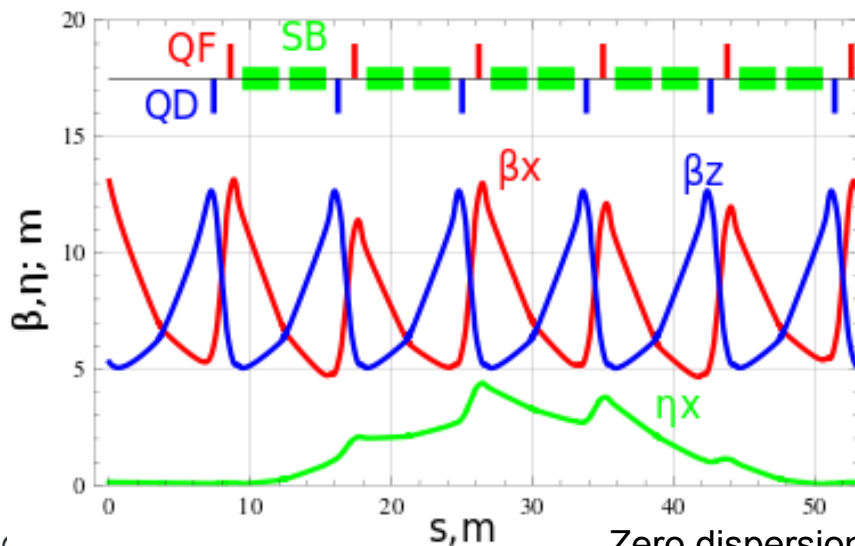
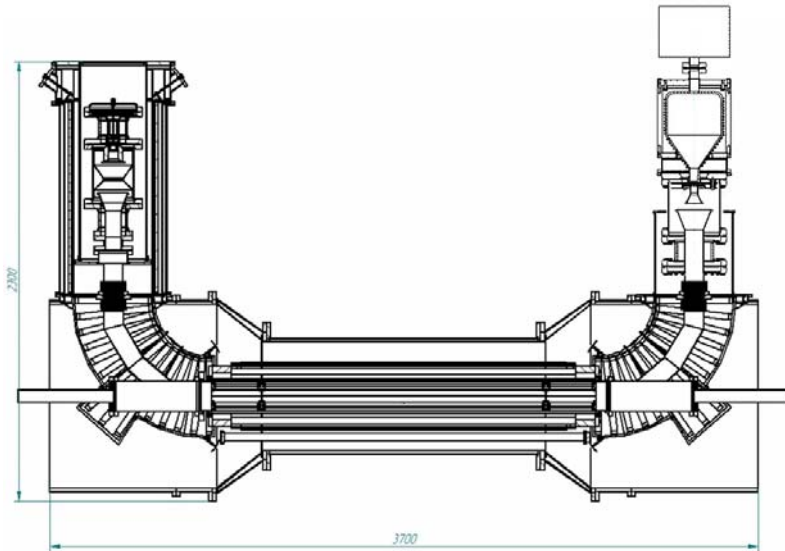
3MeV/nucleon
Au³¹⁺ RFQ

4 large straight sections

1. Injection
2. Two RF-cavity
3. One turn extraction
4. E-cooler



Booster electron cooling system



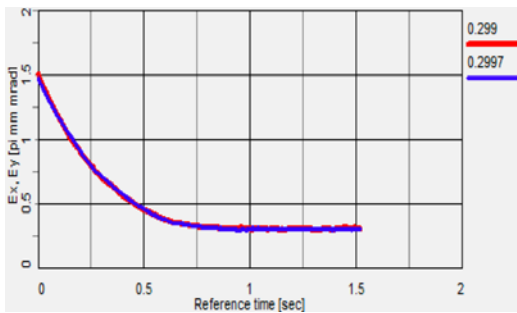
Ions	$^{197}\text{Au}^{31+ (65+)}$
Booster circumference, m	211.2
Injection/extraction energy, MeV/u	3/600
Max. dipole field, T	1.8
Ion number	2×10^9
Beta functions in cooling section, m	10-6 / 6-10
Dispersion in cooling section, m	0.6
Maximum electron energy, keV	50.0
Electron beam current, A	$0 \div 1.0$
Cooler overall length, m	5.7
Eff. length of the cooling section, m	2.5
Magnetic field in the e-cooler, kG	1.5
Magnetic field inhomogeneity in the cooling section, $\Delta B/B$	$1 \cdot 10^{-4}$
Electron beam radius, cm	2.5
Trans. electron temperature, meV	200
(long) electron temperature, meV	0.5
Cooling time, s ($I_e=0.05\text{A}$)	0.4
Residual gas pressure, Torr	10^{-11}

Simulation of cooling process with BETACOOl

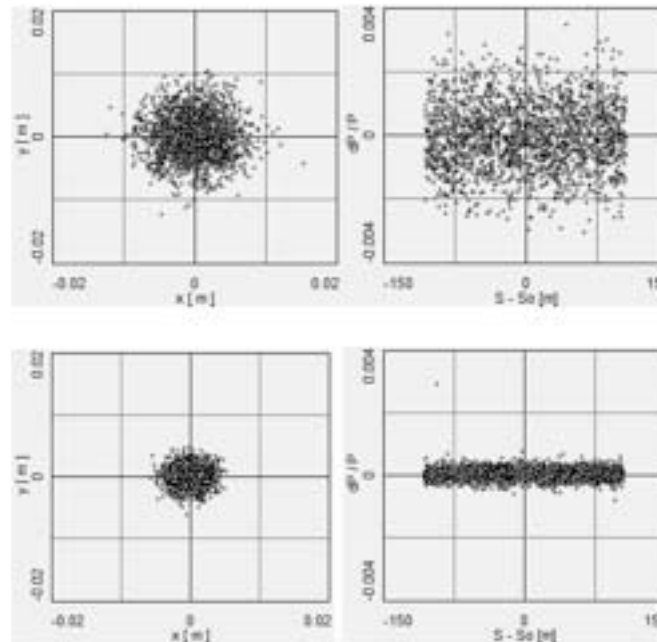
Evolution of the bunched ion beam parameters during the cooling process (Ion energy = 3 MeV/u).

Initial parameters of the cooling

Horizontal and vertical emittances

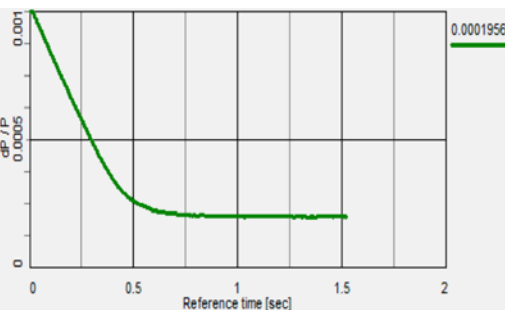


Phase space diagrams



Ion energy, MeV	3
Ion kind	$^{197}\text{Au}^{31+}$
Particle number	2×10^9
Initial Tr_emittance, π mm mrad (rms)	1.5
Initial momentum spread	1×10^{-3}
RF voltage, kV	10
Electron beam current, A	0.05
Electron beam temp. long/trans, meV	200 / 0.5

ion momentum spread



Transverse emittance evolution in the collider injection chain

From HILac $6\sigma \leq 10 \pi \cdot \text{mm} \cdot \text{mrad}$ (upper limit for vertical emittance)

Horizontal Booster acceptance $\sim 120 \pi \cdot \text{mm} \cdot \text{mrad}$

(upper limit for the horizontal 6σ emittance in the case of multiturn injection)

In Booster: Kinematic decrease + Electron cooling (if necessary)

At 65 MeV/u rms $\varepsilon_h \leq 4.21 \pi \cdot \text{mm} \cdot \text{mrad}$ $\varepsilon_v \leq 0.35 \pi \cdot \text{mm} \cdot \text{mrad}$

At 600 MeV/u (without cooling) rms $\varepsilon_h \leq 1.23 \pi \cdot \text{mm} \cdot \text{mrad}$ $\varepsilon_v \leq 0.102 \pi \cdot \text{mm} \cdot \text{mrad}$

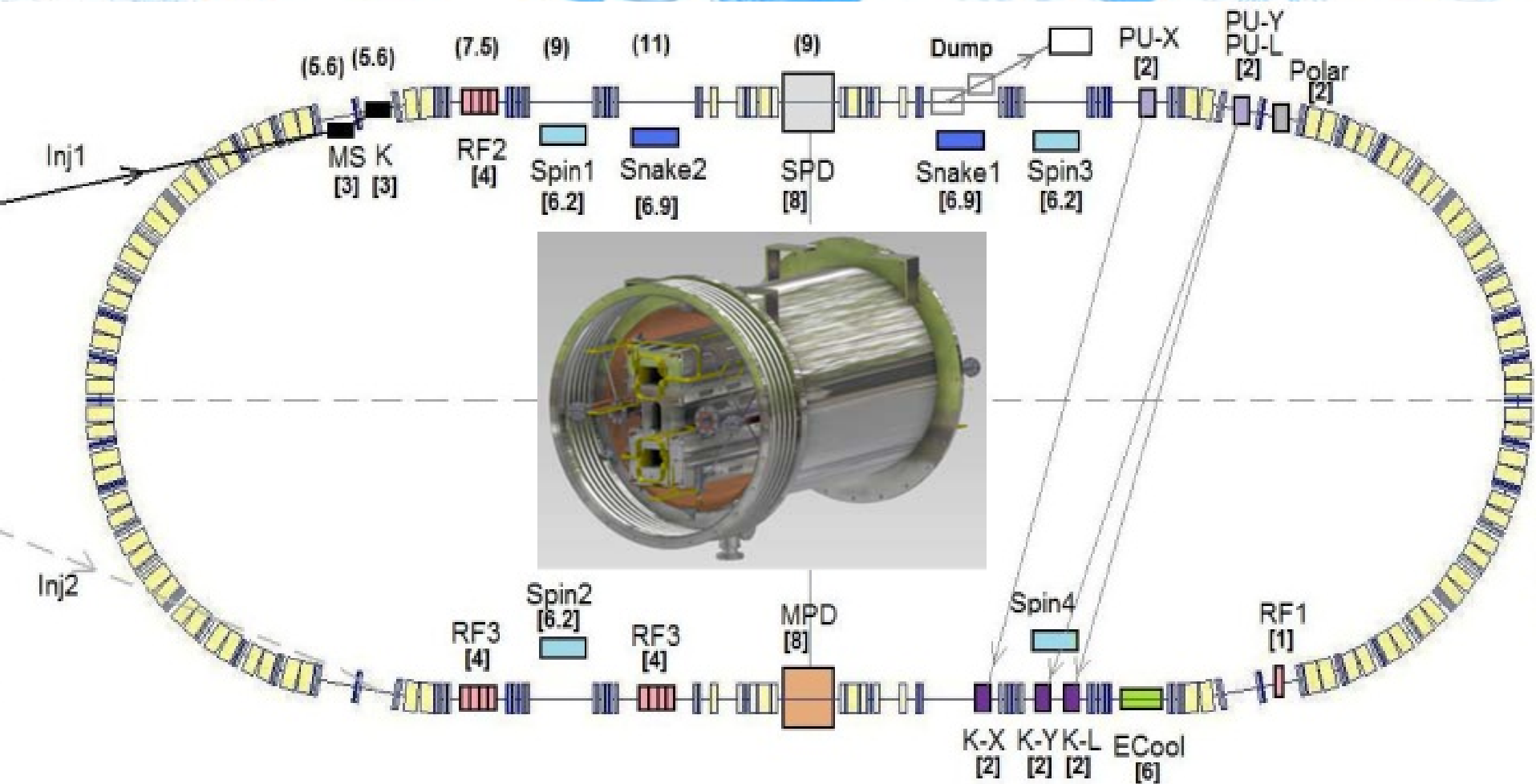
Increase in transfer line from Booster to Nuclotron:

- Stripper
- Coupling
- Mismatch

If $\varepsilon_h \gg \varepsilon_v$, $\varepsilon_h \sim \text{constant}$, ε_v increases by about 3 times

At injection into Nuclotron rms $\varepsilon_h \leq 1.23 \pi \cdot \text{mm} \cdot \text{mrad}$ $\varepsilon_v \leq 0.3 \pi \cdot \text{mm} \cdot \text{mrad}$

NICA collider



Collider parameters

Ring circumference, m	503,04		
Number of bunches	23		
Rms bunch length, m	0.6		
Beta-function in the IP, m	0.35		
Ring acceptance (FF lenses)	$40 \pi \text{ mm mrad}$		
Long. acceptance, dp/p	± 0.010		
Gamma-transition, γ_{tr}	7.091		
Ion energy, GeV/u	1.0	3.0	4.5
Ion number per bunch	$2.75 \cdot 10^8$	$2.4 \cdot 10^9$	$2.2 \cdot 10^9$
Rms momentum spread, 10^{-3}	0.62	1.25	1.65
Rms beam emittance, h/v, (unnormalized), $\pi \cdot \text{mm} \cdot \text{mrad}$	1.1/ 1.01	1.1/ 0.89	1.1/ 0.76
Luminosity, $\text{cm}^{-2}\text{s}^{-1}$	1.1e25	1e27	1e27
IBS growth time, sec	186	702	2540

Peak luminosity can be estimated as:

$$L = \frac{N_b^2}{4\pi\epsilon\beta^*} F_{coll} f_{HG} \left(\frac{\sigma_s}{\beta^*} \right)$$

The collision repetition rate:

$$F_{coll} = \frac{\beta c}{l_{bb}}, \quad l_{bb} = \frac{C_{Ring}}{n_{bunch}}$$

Hour-glass effect ~ 1 (because in our case $\sigma_s \ll \beta^*$):

$$f_{HG} \left(\frac{\sigma_s}{\beta^*} \right) = \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} \frac{\exp(-u^2) du}{\left[1 + \left(\frac{u\sigma_s}{\beta^*} \right)^2 \right]}$$

Maximum luminosity is reached when the bunch phase volume corresponds to the ring acceptance

Collider parameters

To reach maximum peak luminosity one needs to meet the following evident requirements:

- minimum beta function in the IP;
- maximum collision repetition rate (that corresponds to bunch number in rings as maximum as possible);
- maximum bunch intensity;
- minimum beam emittance;
- minimum bunch length.

$$L = \frac{N_b^2}{4\pi\epsilon\beta^*} F_{coll} f_{HG} \left(\frac{\sigma_s}{\beta^*} \right)$$

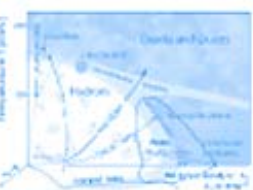
$$f_{HG} = \frac{C_{Ring}}{n_{bunch}}$$

Hour-glass effect ~ 1 because in our case $\sigma_s \ll \beta^*$:

$$f_{HG} \left(\frac{\sigma_s}{\beta^*} \right) = \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} \frac{\exp(-u^2) du}{1 + \left(\frac{u\sigma_s}{\beta^*} \right)^2}$$

FF lenses aperture (radius) : 40mm

Proposed chromaticity correction scheme provides the transverse dynamic aperture of about 120 $\mu\text{m}\cdot\text{mrad}$ and dynamic aperture on the relative momentum deviation of about $\pm 1\%$ when the bunch phase volume corresponds to the ring acceptance



The problems we have met and the solutions

The problems:

- ✓ Beam space charge effects
- ✓ Beam-beam effects
- ✓ IntraBeam scattering (IBS) ⇒
⇒ luminosity degradation

- ✓ Recombination in e-cooler

- ✓ Bunch halo formation ⇒
⇒ parasitic collisions,
background growth

The solutions:

- Electron cooling application
- Stochastic cooling application
- Beam parameters choice
- Operation scenario optimisation

➤ Scrappers and collimators



Role of beam cooling

1. Beam stacking using BB system:

Injection repetition is 10 sec, the cooling times has to be short enough

2. Longitudinal cooling during beam bunching

3. During experiment:

- in IBS dominated mode - Suppression of IBS;
- in SC dominated mode – providing optimum phase volume



Beam stacking with BB system and cooling

$E > 3 \text{ GeV/u}$:

stacking with Stochastic cooling

The cooling time is proportional to the bunching factor
for “almost” coasting beam in BB the cooling times $\sim 10 \text{ sec}$
(T.Katayama previous MAC)

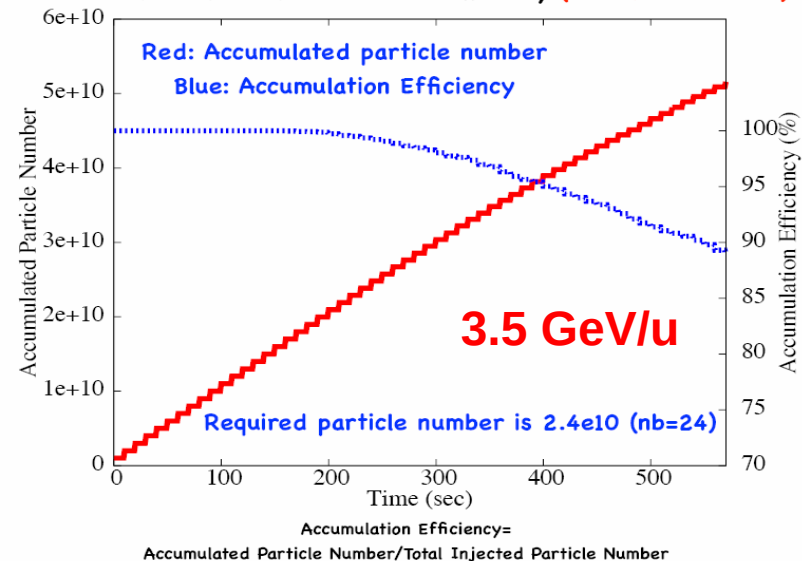
$E < 3 \text{ GeV/u}$:

Stacking with Electron cooling

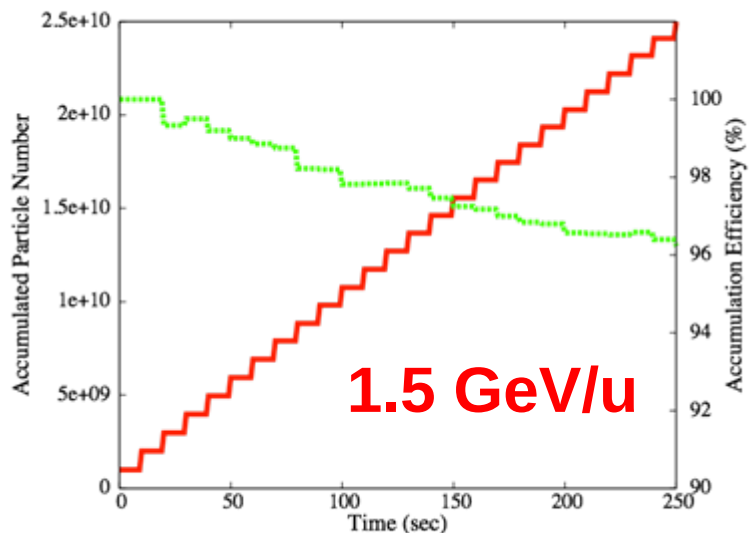
Problem – cooling time strongly depends on energy
and does not depend on bunching factor

The cooling power sufficient for experiment can be
insufficient for effective stacking

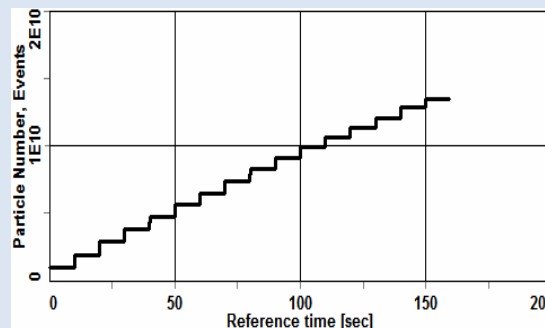
Accumulated Particle Number & Efficiency (Initial Gain=120dB)



Beam stacking with electron cooling

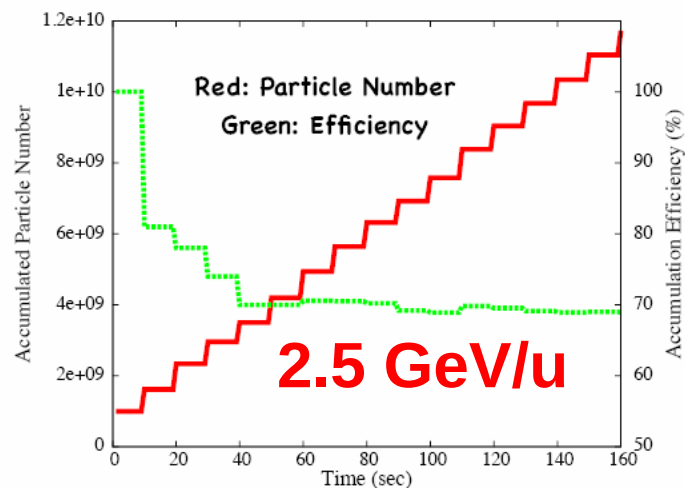
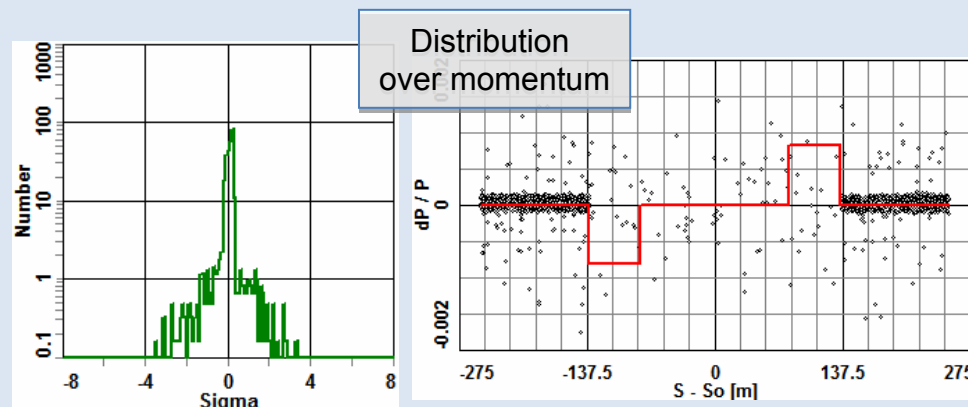


Stacking efficiency (A.Smirnov)



$E = 1.5 \text{ GeV/u}$
Stacking efficiency 86%

$E_{\text{Init}} = 1\pi \cdot \text{mm} \cdot \text{mrad}$



Stacking with electron cooling and beam bunching

Effective stacking can be realized below $\sim 2 \text{ GeV/u}$

Possible solution:

- Stacking at 1 GeV/u with efficiency closed to 100%
- Slow acceleration to the experiment energy

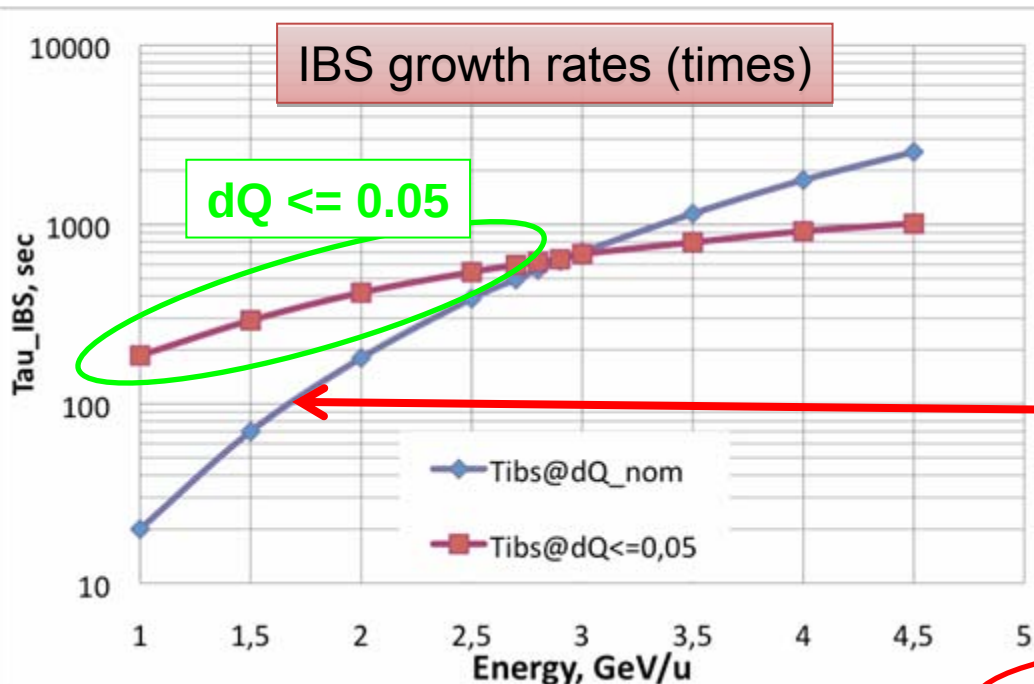
At constant longitudinal emittance ($\sigma_p \times \sigma_s = \text{const}$) minimum threshold current of microwave instability corresponds to coasting beam

$I_{\text{th}} \sim \sigma_p^2$ the momentum spread is inversely proportional to bunching factor

Formation of required long emittance by a few steps:

- At storage of coasting beam formation σ_p required for the stability (long emit. is larger than required at collisions by about bunching factor ~ 13)
- Bunching at 24 harmonics + cooling of long emittance by about 6 times, Formation of the bunch of about 1.2 m, required to recapture in $h=72$
- Bunching + cooling at $h = 72$, formation of required bunch length and momentum spread

IBS calculations

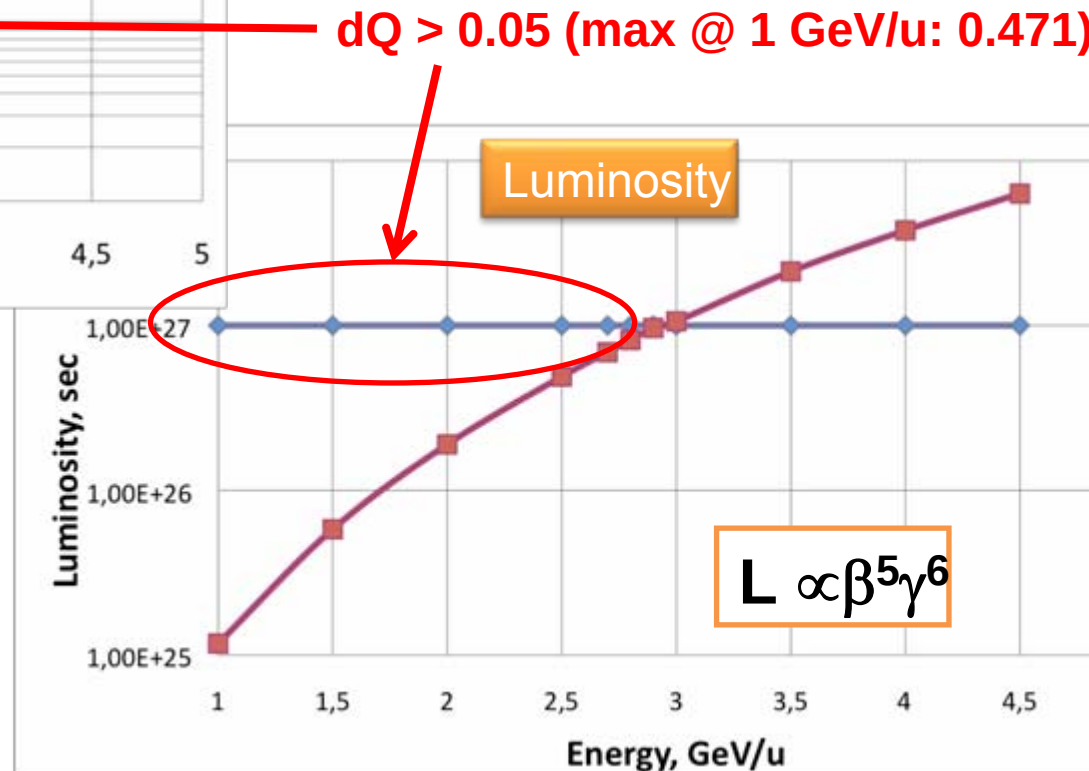


Strategy:

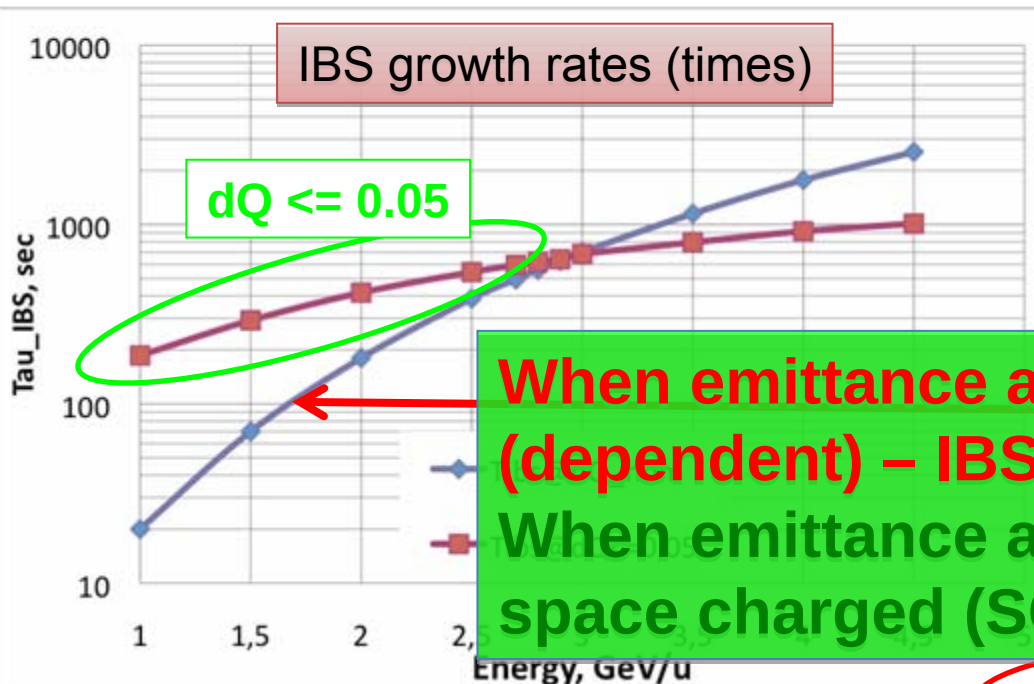
1. $\varepsilon_x = 1.1 \pi \text{ mm mrad}$ (due to $6\sigma_x = 40$)
2. Equal heating rates of all degrees
3. $dP/P \sim (1-1.5)e-3$ is acceptable
(from bunch coherent stability condition)
4. $L \leq 1e27$

“IBS dominated regime”:
bunch parameters are determined
by equilibrium between IBS
and beam cooling.

$$L = 8\pi^2 \beta^5 \gamma^6 \Delta Q^2 \frac{A^2}{Z^4} \cdot \frac{\varepsilon C}{r_p^2 \beta^* l_{bb}} \cdot \left(\frac{\sigma_s}{C_{Ring}} \right)^2 \cdot f_{HG}$$



IBS calculations

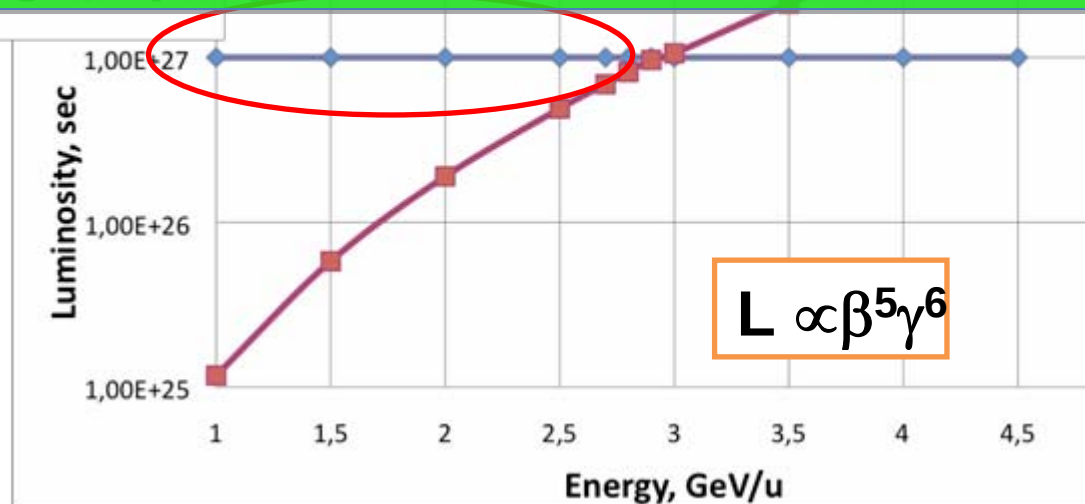


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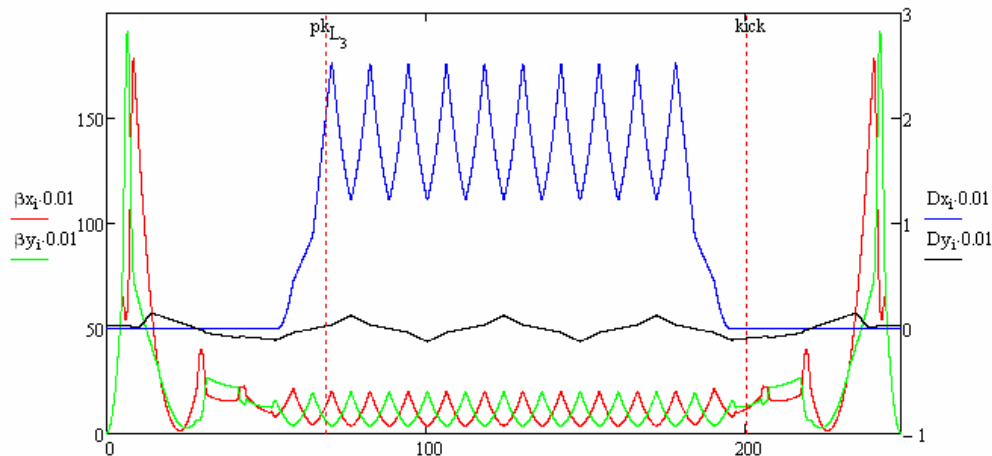
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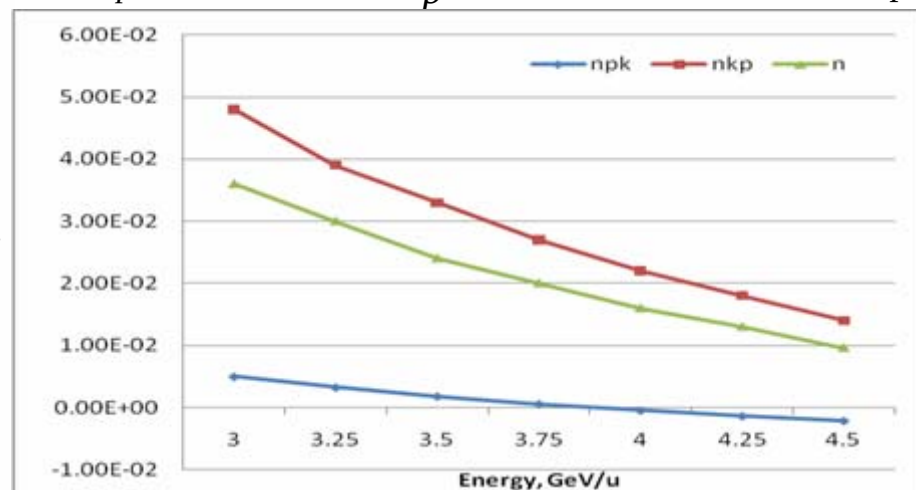


Stochastic cooling

$$\frac{1}{\tau} = \frac{W}{N} \frac{(1 - 1/M_{pk}^2)^2}{M} \quad N_{eq} = N \frac{C}{\sqrt{2\pi}\sigma_c} \quad M_{pk} = \frac{1}{2(f_{max} + f_{min})\eta_{pk}T_{pk}} \frac{\Delta p}{p} \quad f_{max} \leq \frac{1}{2\eta_{pk}T_{pk}} \frac{\Delta p}{p} \quad M_{kp} = \frac{1}{2(f_{max} - f_{min})\eta_{kp}T_{kp}} \frac{\Delta p}{p}$$

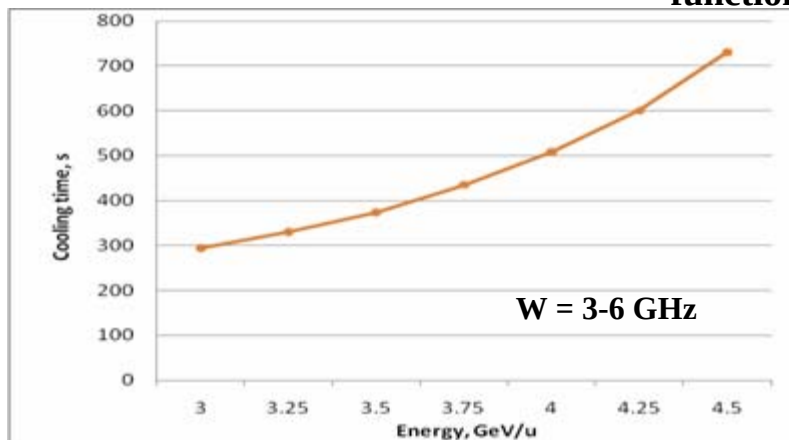


Kicker - 48 meters upstream the IP-point
PU - 132 meters upstream the Kicker



Total and partial slip-factors of the ring as the function of ion energy.

At such position of the kicker the condition gives for the acceptable upper frequency of the band the value of about 20 GHz (at the momentum spread equal to the ring dynamic aperture ± 0.01). The luminosity of $1 \cdot 10^{27} \text{ cm}^{-2}\text{s}^{-1}$ corresponds to about $2.3 \cdot 10^9$ ions per bunch, the effective ion number is about $8 \cdot 10^{11}$. To provide required cooling time the cooling bandwidth can be chosen from 3 to 6 GHz



“Slice” overlapping
(by D.Moehl)

3..6GHz: Tsc ~ 0,5Tibs
 2..4 GHz: Tsc ~ Tibs

Electron cooling

Beam emittances @ equilibrium state. Rates ($\tau_{\varepsilon_x} = \tau_{\varepsilon_y} = \tau_{\sigma_p}$) - from IBS calculations for lattice.

Luminosity is fitted to $1e27$, ε_x is fitted to 1.1π mm mrad

$$\rho_F = -V \frac{\rho 4Z^2 e^4 n_e L_P}{m} \frac{1}{(V^2 + \Delta_{e,eff}^2)^{3/2}} \quad L_P = \ln \left(\frac{\rho_{\max} + \rho_{\min} + \rho_{\perp}}{\rho_{\min} + \rho_{\perp}} \right) \quad \rho_{\min} = \frac{Ze^2}{m} \frac{1}{V^2 + \Delta_{e,eff}^2}$$

$$\rho_{\max} = \frac{v_i}{1/\tau_{flight} + \omega_p} \quad \Delta_{e,eff} = 0,0046 \text{ eV}$$

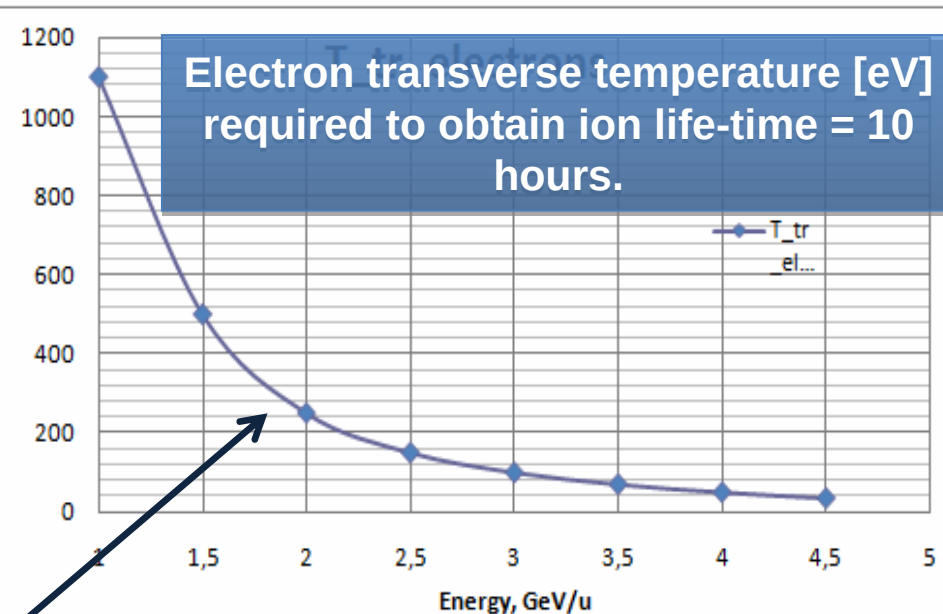
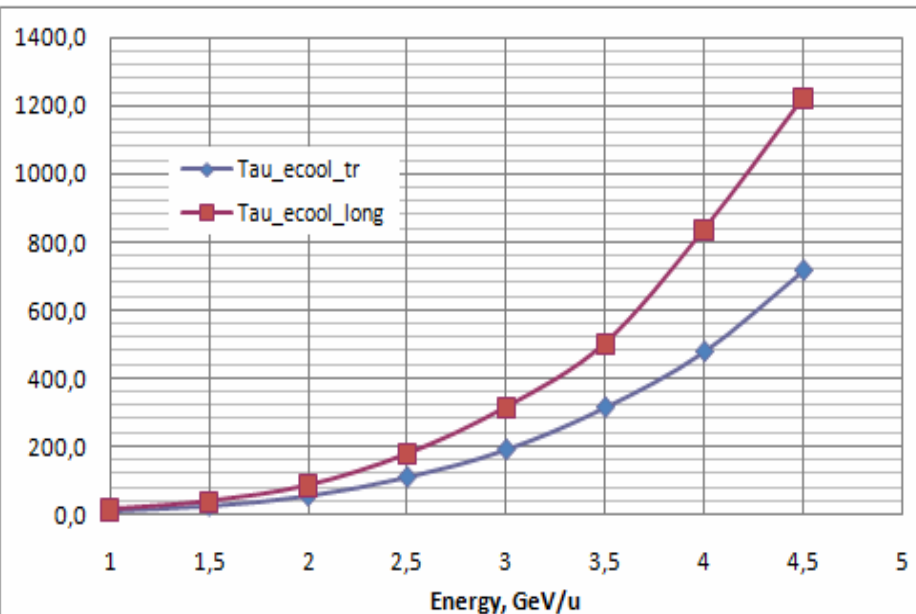
Angular spread [rad] = $2e-5$

Parkhomchuk model. $\beta_x = \beta_y \approx 20\text{m}$ @ cooling section, $L = 6\text{m}$, $B=1\text{T}$ (required mainly to provide adiabatic transport of the electron beam from HV source to the cooling section), $I_{\text{electron}} = 0,5\text{A}$. T_{tr_e} - chosen at all energies to the value in order to have τ_{life} (due_to_recombination) ≥ 10 hours (36000 seconds: recombination rate limit = $2,7E-5$).

Radius_electron_beam chosen to have $T_{\text{ecool}} = \min$ (same at all energies)

The cooling rate is determined mainly by longitudinal electron temperature (that is dominated by HV generator stability) and logarithmically depends on the transverse one

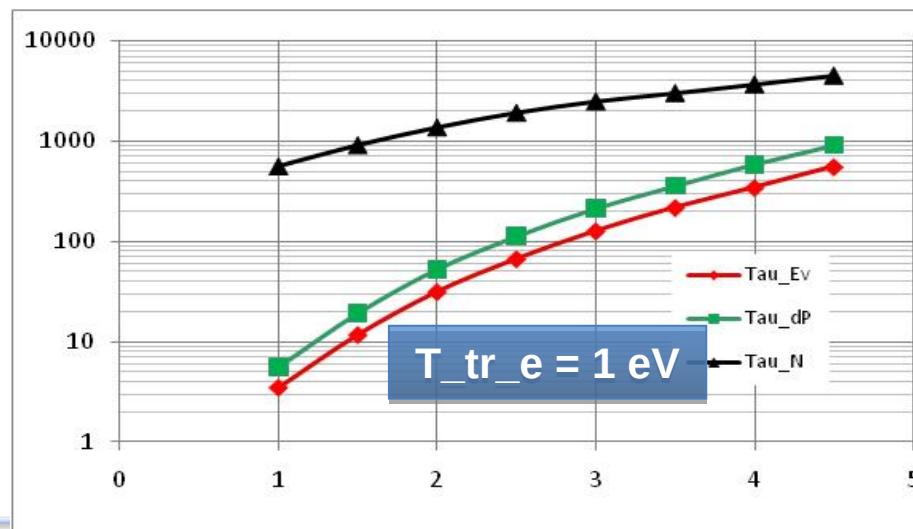
Electron cooling



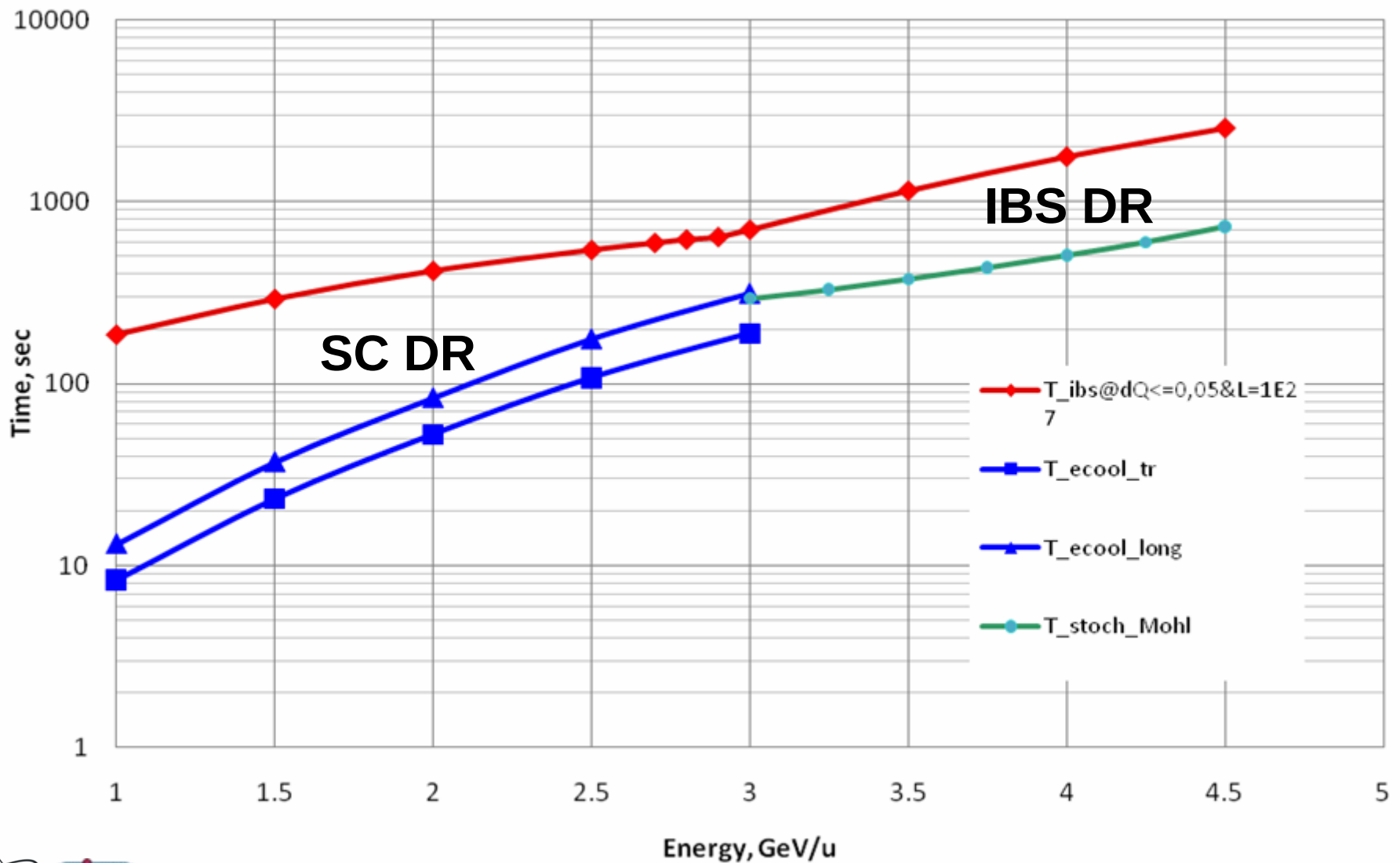
Dependence of the cooling times for transverse and longitudinal degrees of freedom

Recombination suppression:
a) Increasing T_{tr_e}
b) "Shift" of electron energy

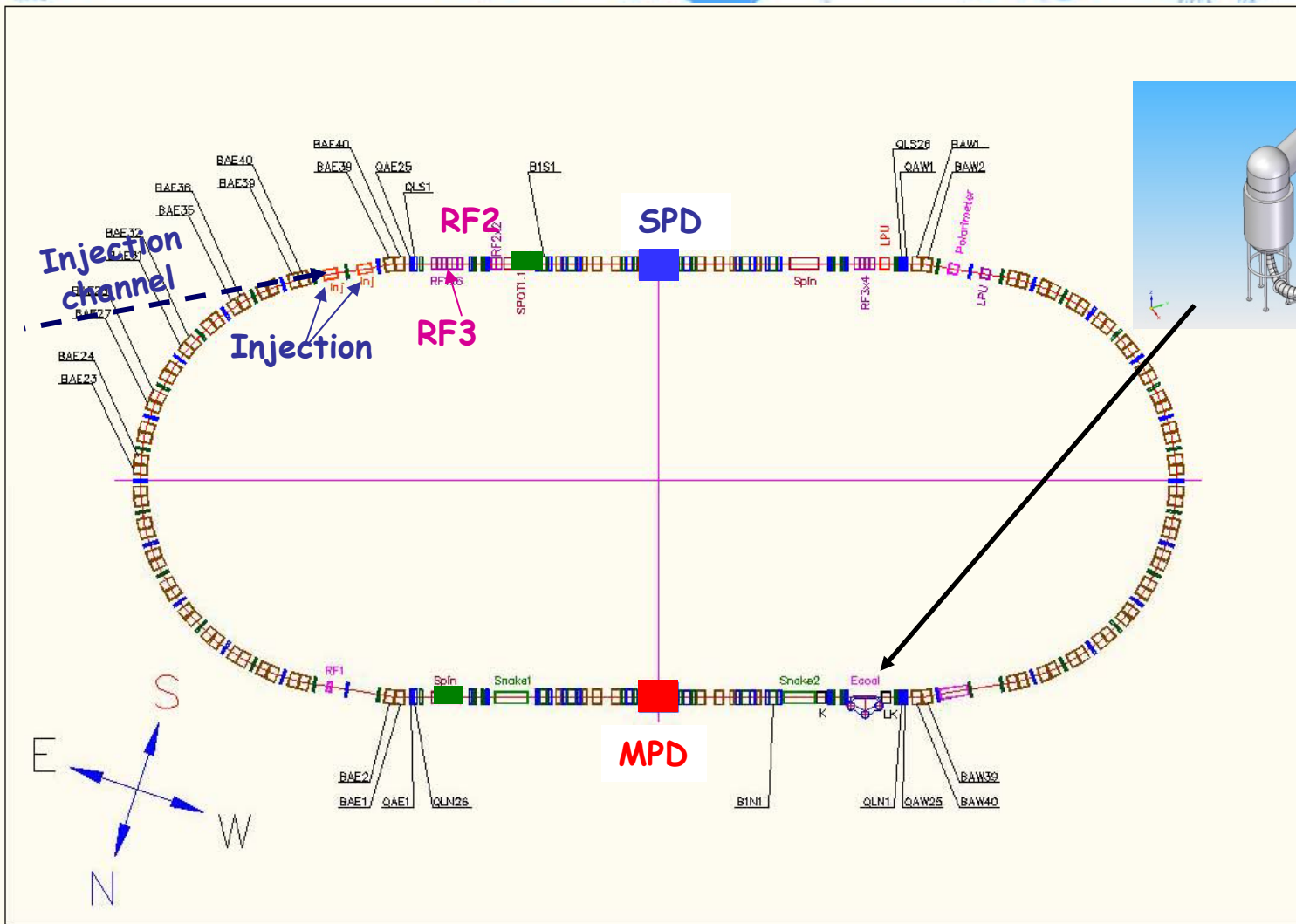
Conclusions: $T_{ecool} \sim 0,05$ Tibs at 1 GeV/u



Summary final

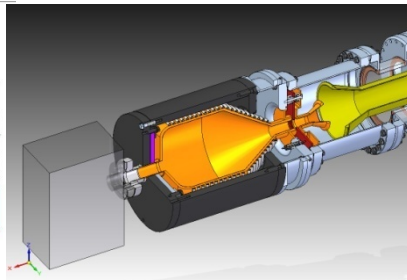
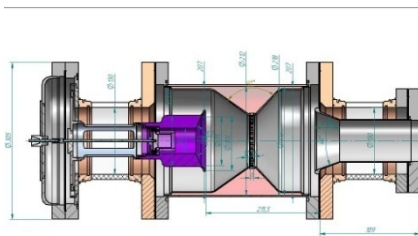
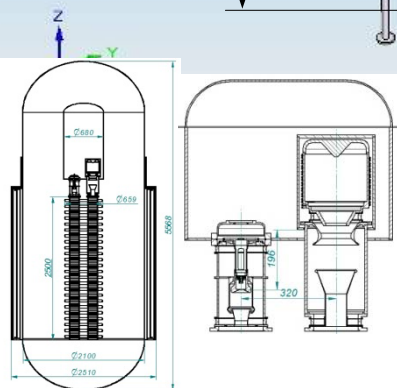


HV electron cooling system





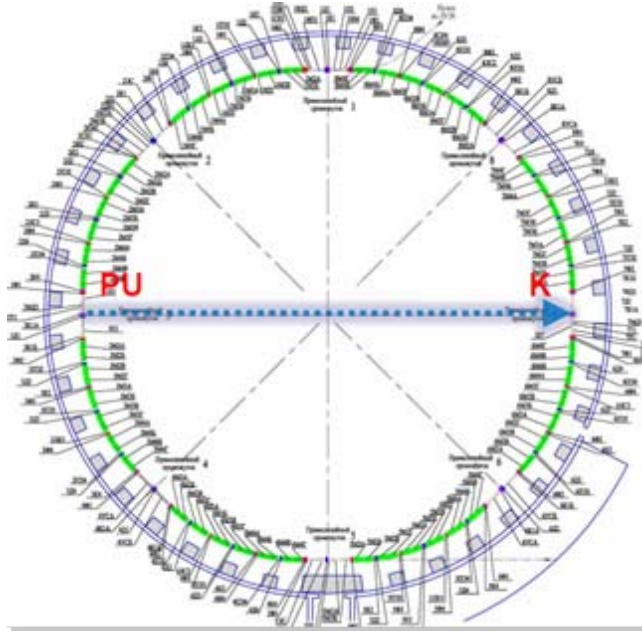
HV electron cooling system



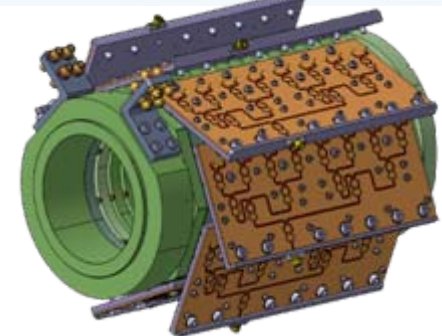
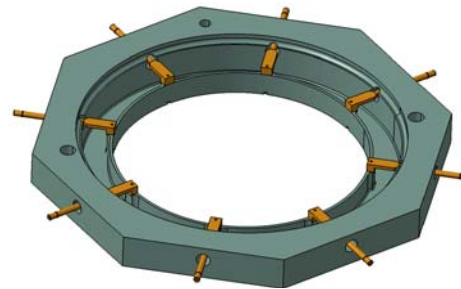
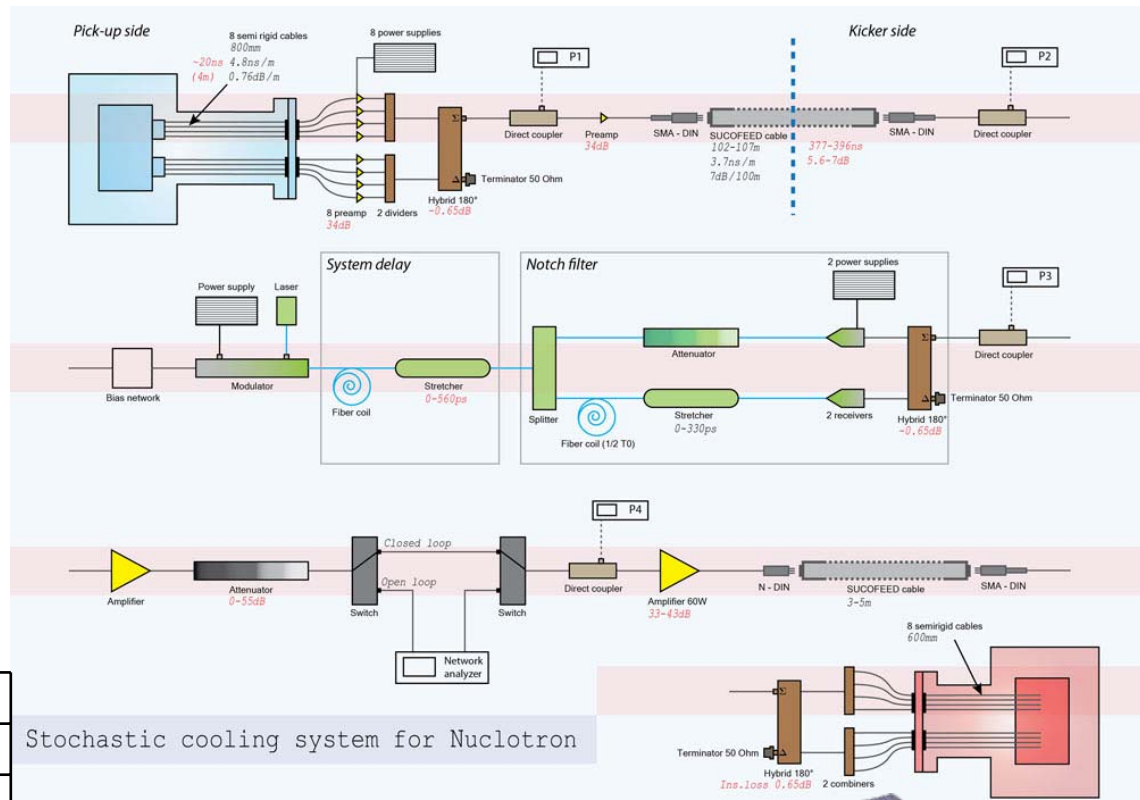
HV Generator prototype U=250 kV, I=1mA



SC experiment at Nuclotron

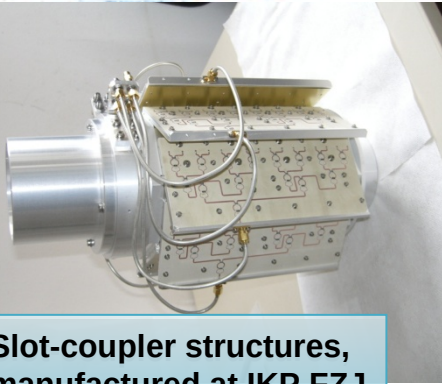


Circumference, m	251.5
Ions	up to A=56
Energy, GeV	3.5
Rev.frequency, MHz	1.2
Vacuum, Torr	10^{-10}
Intensity	$10^{11}(\text{p})$ - $10^9(\text{C}^{12})$
Ring slippage factor	0,0322
dp/p	10^{-3}



Nuclotron-NICA

Stochastic cooling system prototype at Nuclotron



Slot-coupler structures,
manufactured at IKP FZJ



Vacuum chamber for pick-up



Vacuum chamber for kicker



Experimental set-up:

V2 MICROSOFT-9B0451

Подсистема измерения тока пучка (Bergoz) - 100kHz

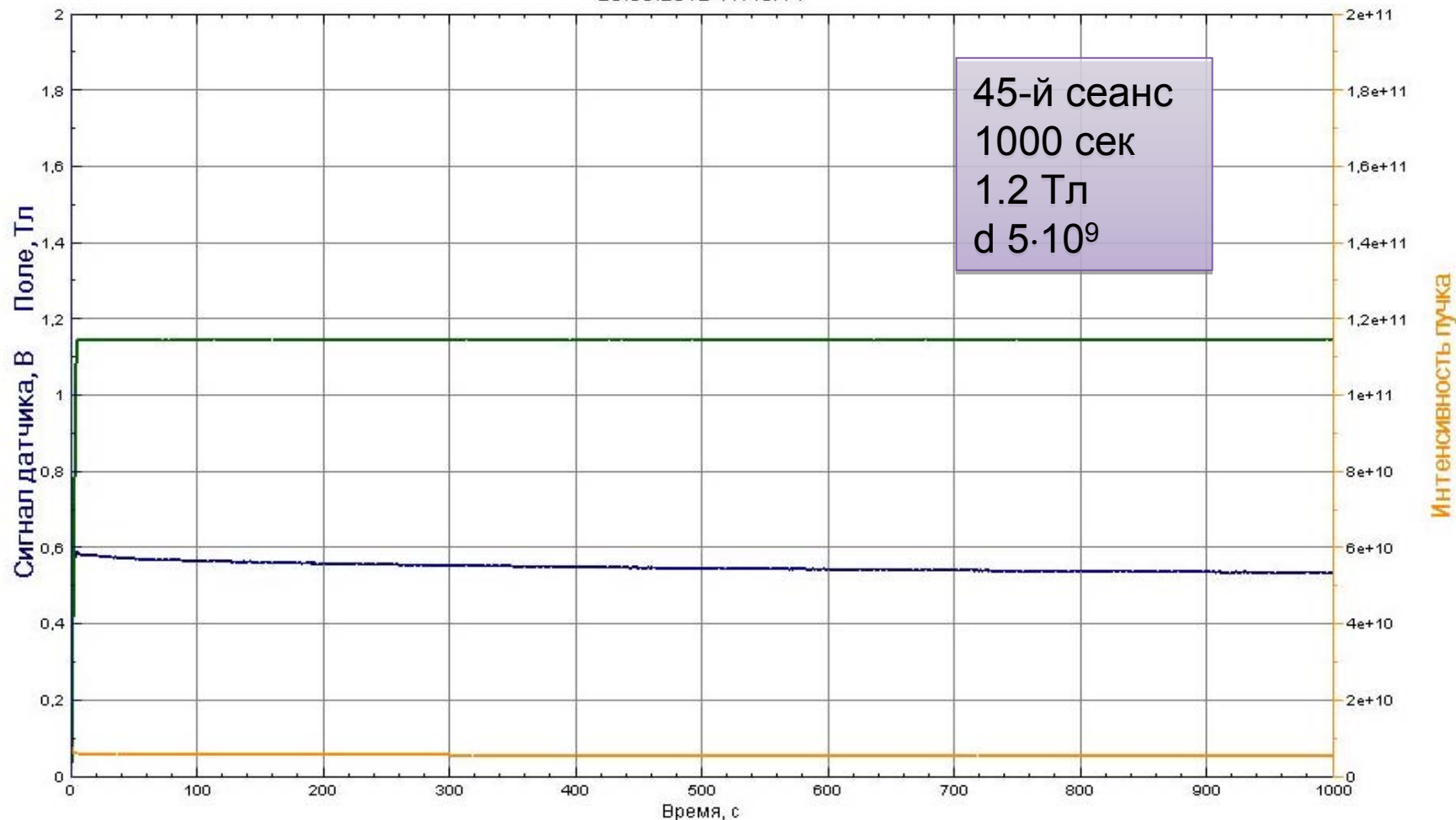
Файл Работа Помощь

Измеренный сигнал ДУ и протокол

25.03.2012 17:40:14

10 В

0.01 В



Выключить отображение!

☒ Сигнал датчика ☒ Поле ☒ Интенсивность

Длительность цикла (сек):

1000

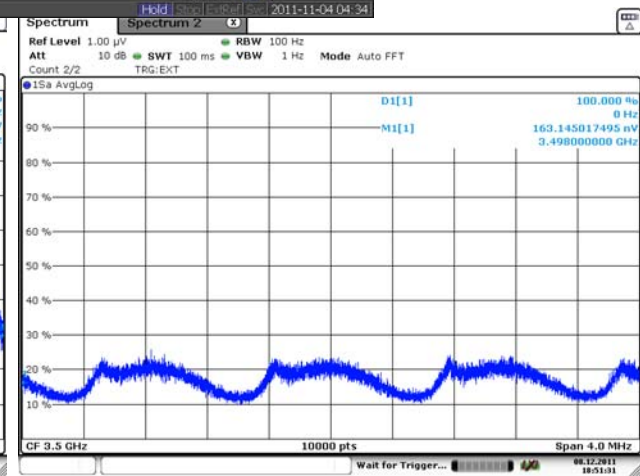
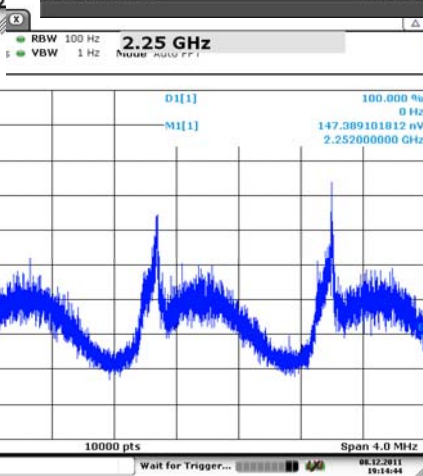
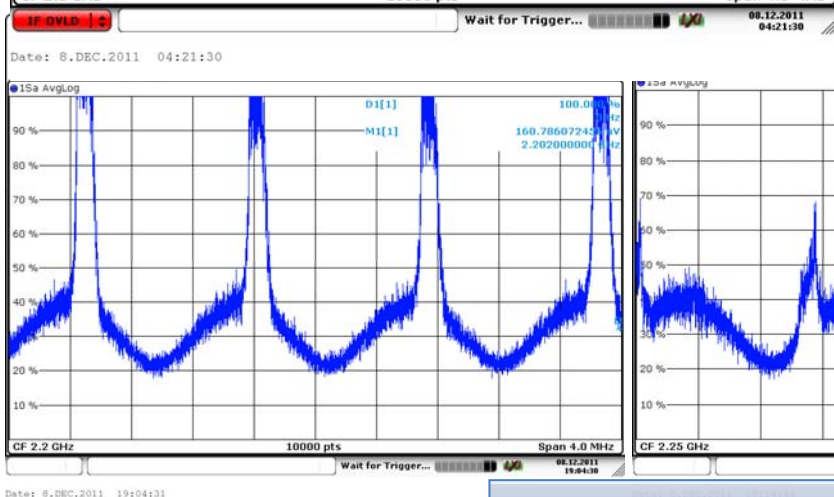
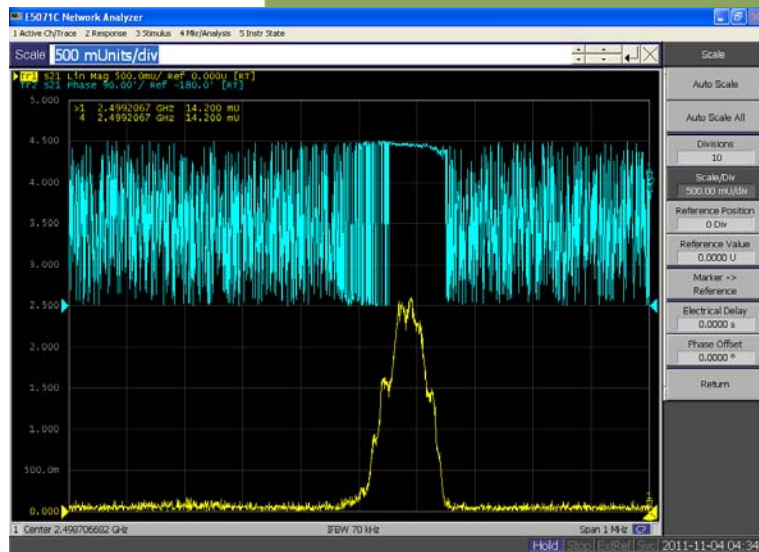
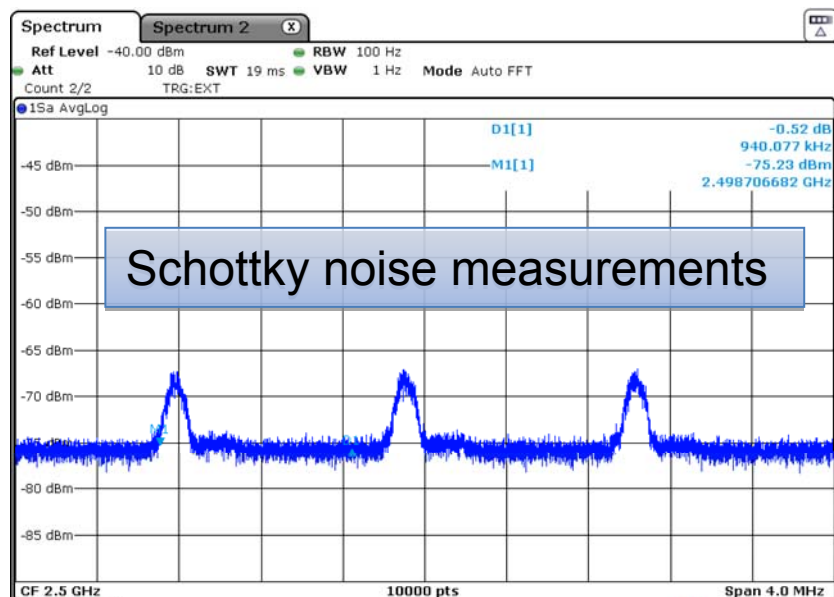
Принять

Отменить

Experimental results:

Energy range 0.5-4.0 GeV/u
Deutrons, C6+
During 2 runs in 2011 and
2012

Beam BTF
d, 2 GeV/u



Notch filter rejects revolution frequencies



Conclusions

1. Electron cooling at Booster is required to have wide machine possibilities at beam injection from HILAC and to perform applied research
2. Stochastic cooling at NICA collider is sufficient for IBS suppression and for beam stacking also
2. Electron cooling can be used for cooling at experiment in the total energy range
3. Electron cooling can provide effective stacking at small energy only
4. At energy larger than 3 GeV/u the experiment can be provided using electron cooling as well as stochastic, or combination of both methods
5. At energy below 3 GeV/u the experiment is provided using electron cooling. Possibility to increase of luminosity at minimum energy is related with space charge dominated mode
6. Slow acceleration of the beam is presumed for the case when effective stacking (or injection) is complicated at the experiment energy

Thank you for your attention !

Many thanks to my colleagues for fruitful discussions and help:

*R.Stassen, T.Katayama, L.Thorndal, D.Mohl, A.Sidorin, N.Shurkhno,
I.Meshkov, S.Kostromin, O.Kozlov, A.Smirnov, H.Khodzhibagiyan,
V.Volkov, S.Romanov, E.Gorbachev, A.Butenko, V.Seleznev, and others*