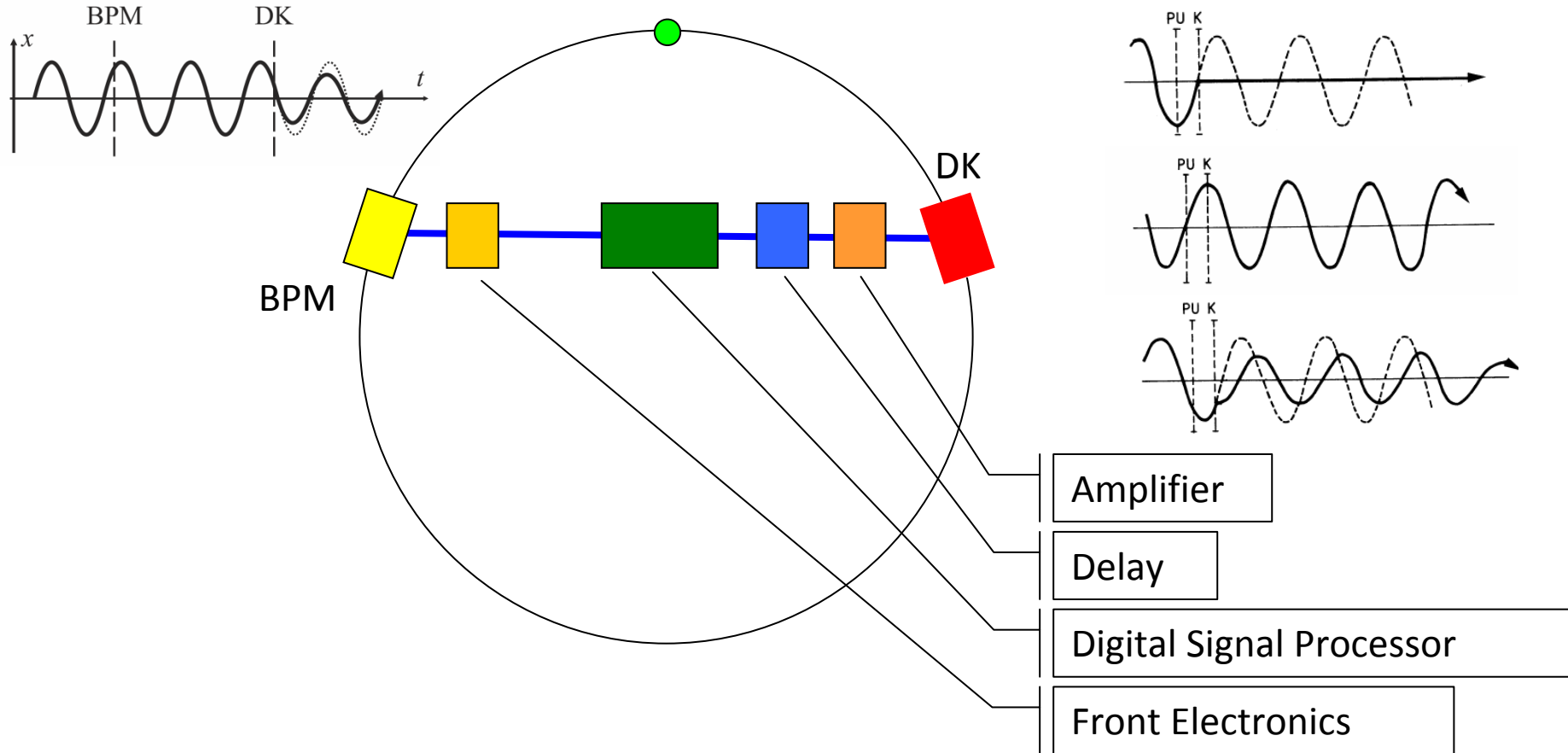


TRANSIENT BEAM RESPONSE IN SYNCHROTRONS WITH A DIGITAL TRANSVERSE FEEDBACK SYSTEM

V.M. Zhabitsky
JINR, Dubna, Russia

Transverse Feedback System in Synchrotrons



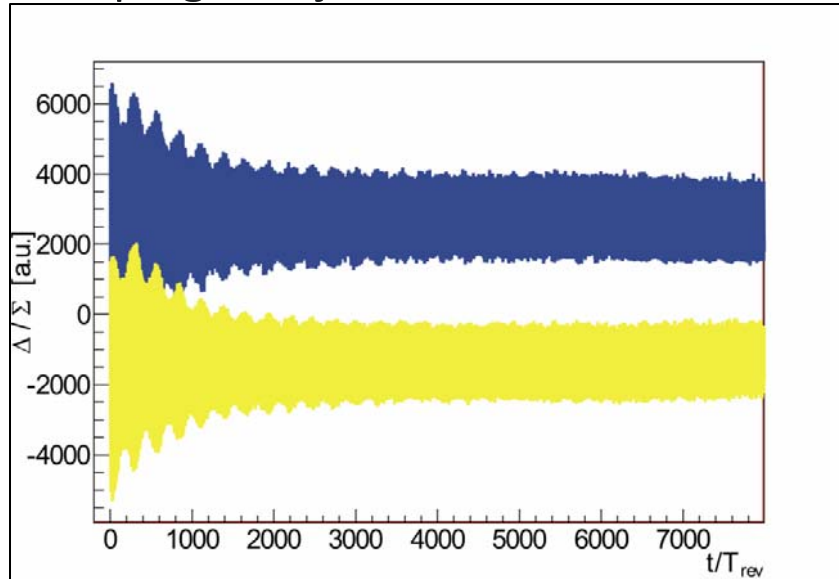
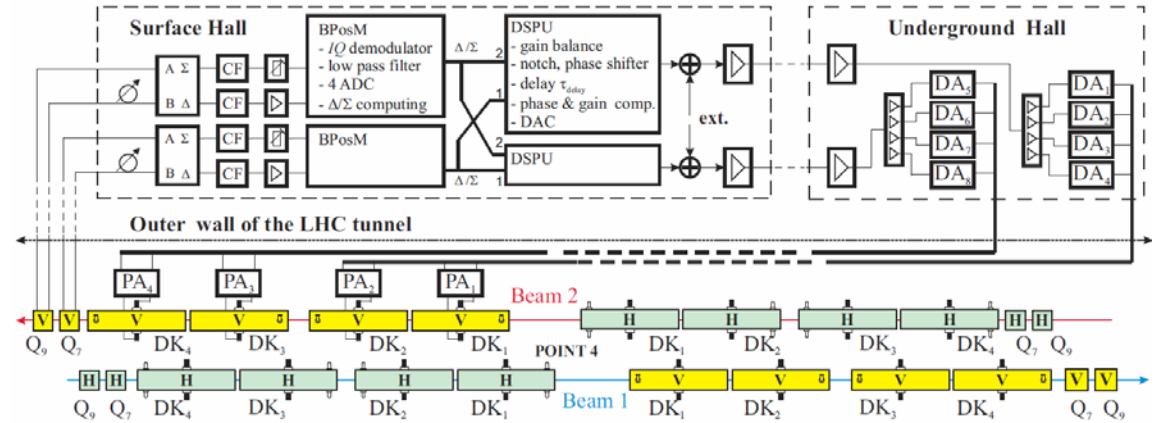
$$g = \frac{\langle \Delta x'[n, s_K] \rangle \sqrt{\hat{\beta}_P \hat{\beta}_K}}{\langle x[n, s_P] \rangle} = \frac{2 T_{\text{rev}}}{\tau_d} \Big|_{\text{opt}}$$

LHC Damper: $g = 0.05 \rightarrow \tau_d = 40 T_{\text{rev}}$

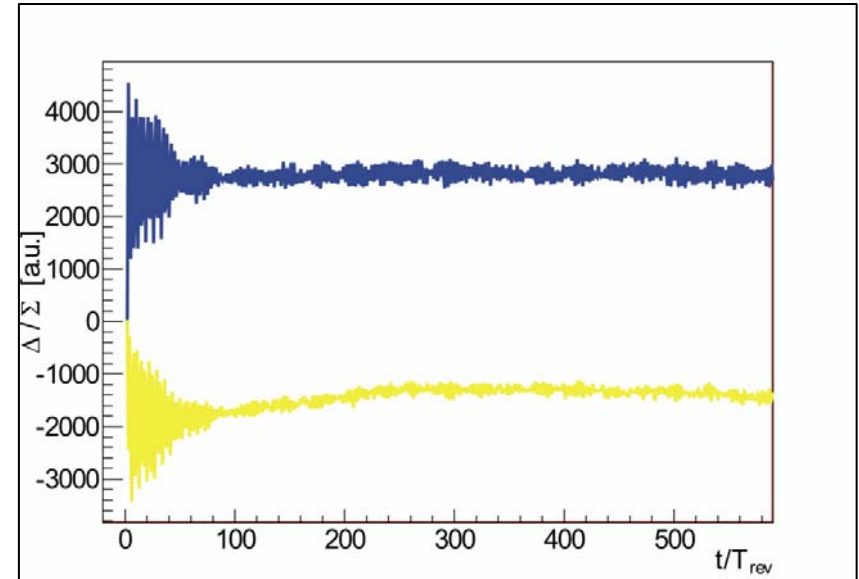
LHC Damper (CERN, JINR)



Single Bunch Operation:
Damping of Injection Errors

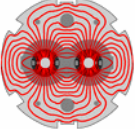


Damper OFF

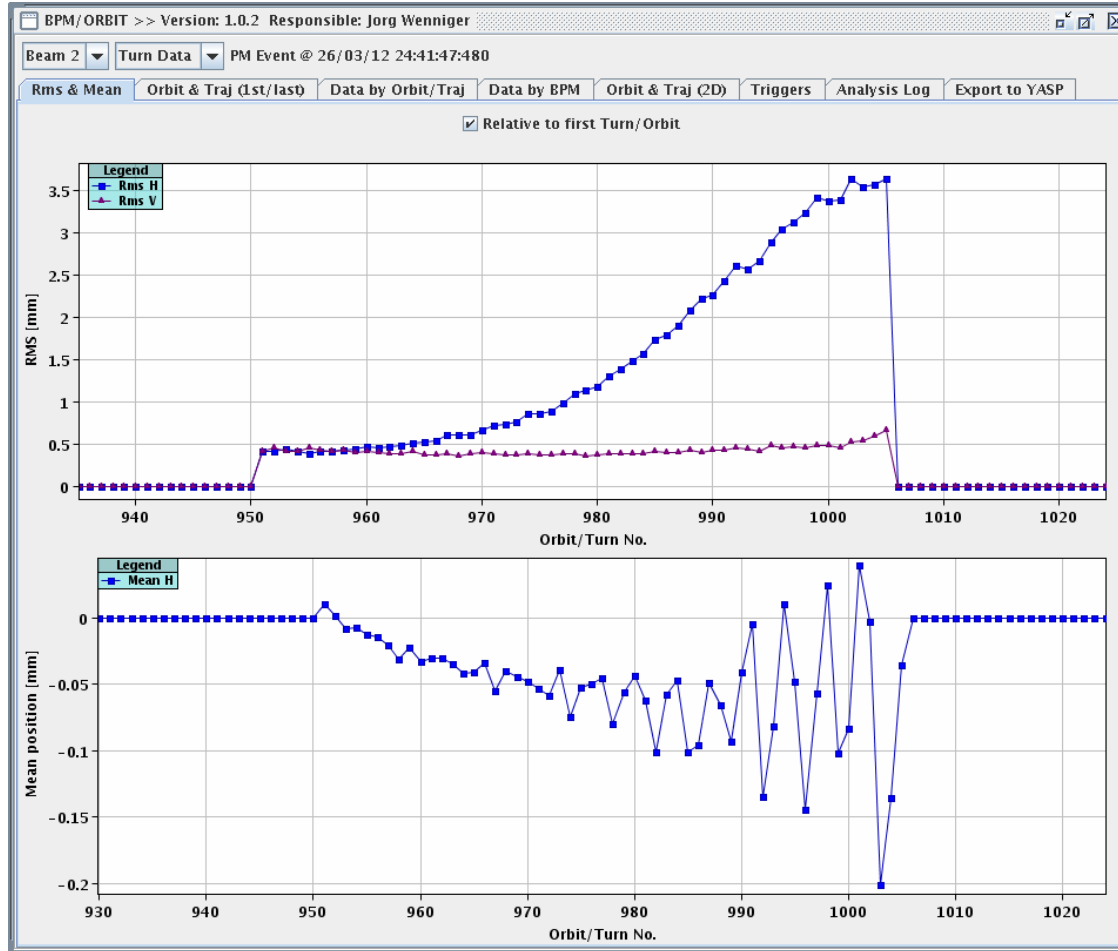


Damper ON

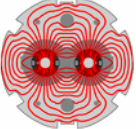
V.Zhabitsky, W.H. Fle et al. RuPAC 2010



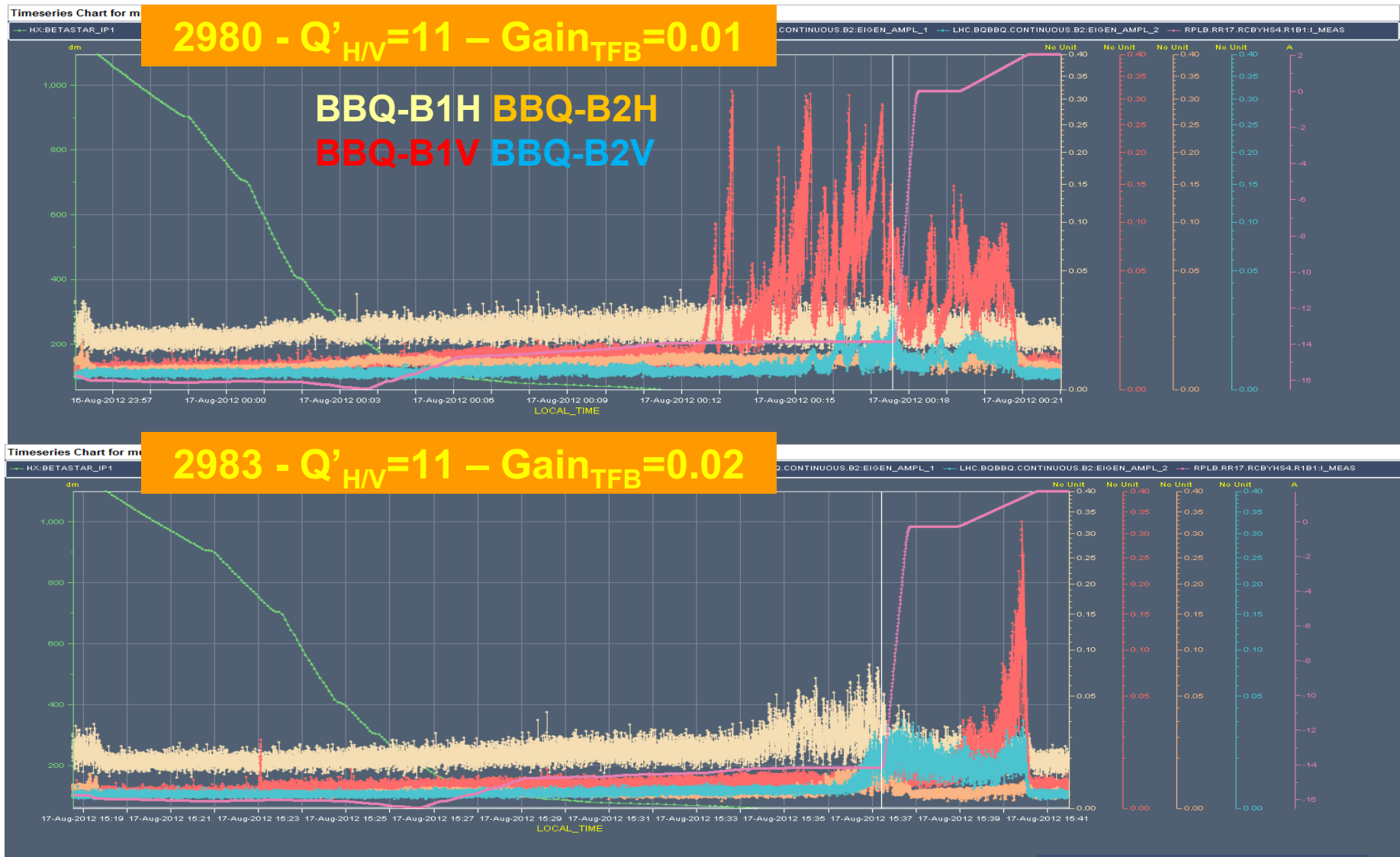
- 24:00 – 01:00 ADT fast losses test – one pilot at 450 GeV
- Achieved very fast losses (could be used for UFO studies)



J. Uythoven, E.B. Holzer (CERN)



Instabilities



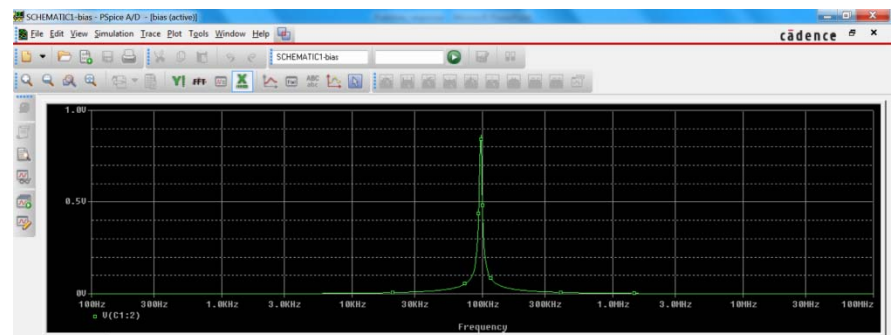
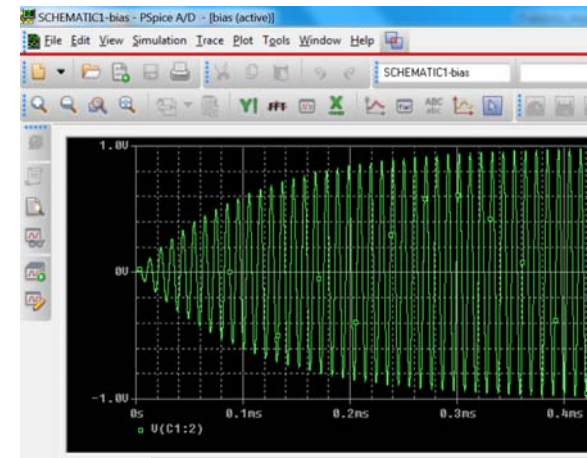
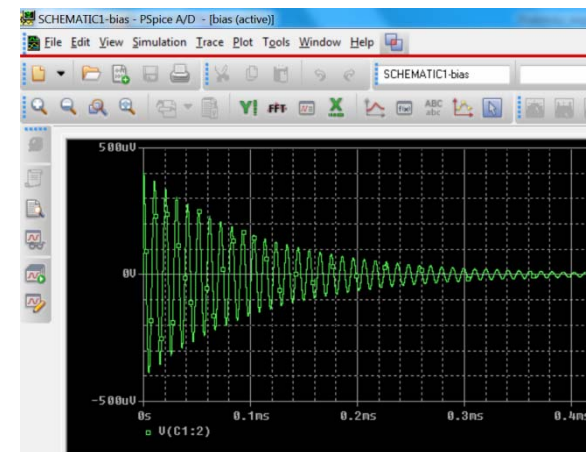
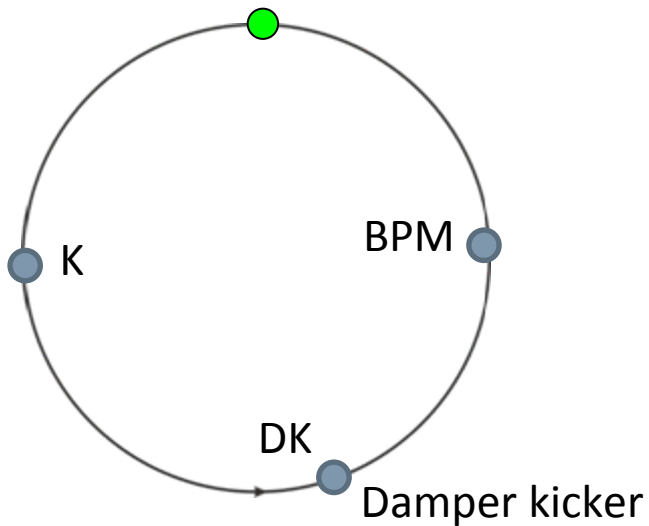
G. Arduini (CERN)

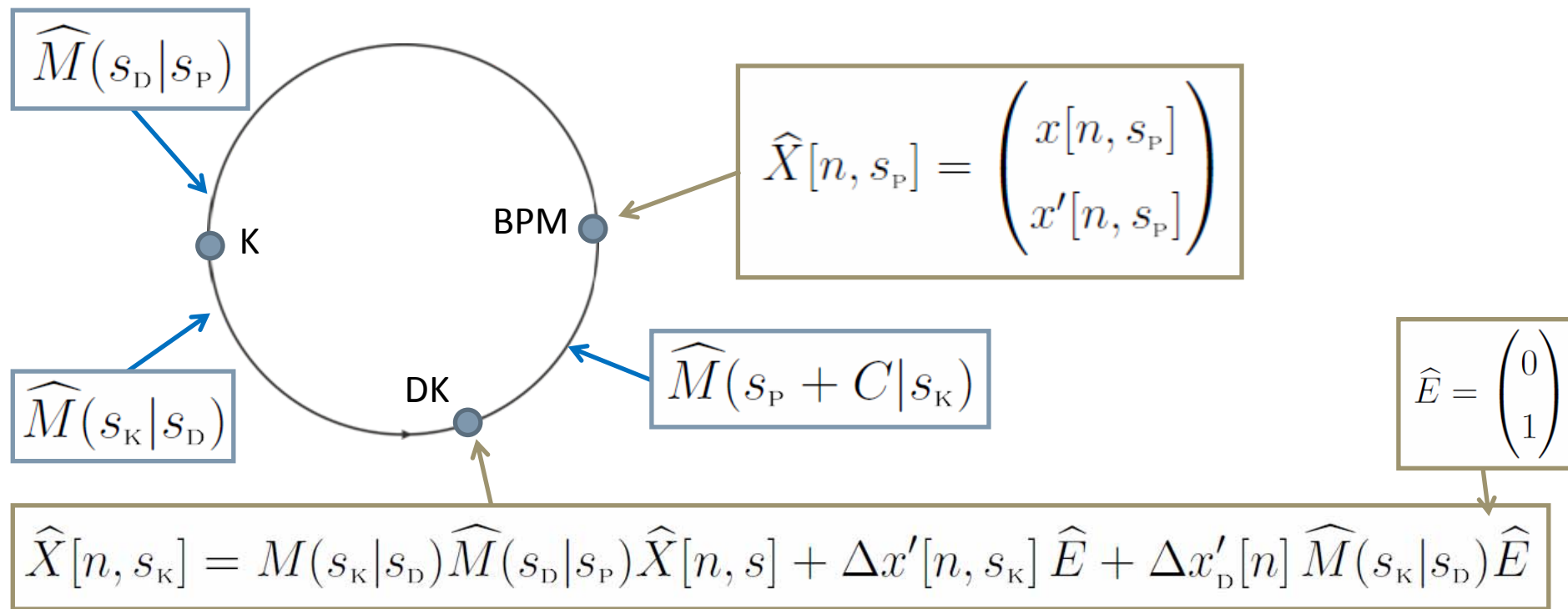
Response to impulse

Radio technical device



Synchrotrons





$$\widehat{X}[n, s_P + C] = \widehat{M}(s_P + C | s_K) \widehat{X}[n, s_K]$$

$$\widehat{X}[n + 1, s] \equiv \widehat{X}[n, s + C] = \widehat{M}(s) \widehat{X}[n, s] + \Delta x'[n, s_K] \widehat{M}_K \widehat{E} + \Delta x'_D[n] \widehat{M}_D \widehat{E}$$

$$\widehat{M}(s_P) \equiv \widehat{M}(s_P + C | s_P), \quad \widehat{M}_K \equiv \widehat{M}(s_P + C | s_K), \quad \widehat{M}_D \equiv \widehat{M}(s_P + C | s_D).$$

$$\widehat{X}[n+1, s] \equiv \widehat{X}[n, s+C] = \widehat{M}(s) \widehat{X}[n, s] + \Delta x'[n, s_K] \widehat{M}_K \widehat{E} + \Delta x'_D[n] \widehat{M}_D \widehat{E}$$

Damper kicker:

$$\Delta x'[n, s_K] = S_K V_{\text{out}}[n] = S_K K_{\text{out}} K_{\text{in}} \sum_{m=0}^{N_F} h[m] V_{\text{in}}[n - \hat{q} - m] u[n - \hat{q} - m]$$

$$\tau_{\text{delay}} = \tau_{\text{PK}} + \hat{q} T_{\text{rev}}$$

$$V_{\text{in}}[n] = (x[n, s_P] + \delta x_P) S_P u[n]$$

Driving force:

$$\Delta x'_D[n] \sqrt{\hat{\beta}_D \hat{\beta}_P} \equiv V_D = a_D \delta[n - n_D]$$

$$V_D = a_D \sin(2\pi(n - n_D)Q_D + \phi_D) (u[n - n_D] - u[n - n_D - N_D])$$

$$\mathbf{y}(z) = \mathcal{Z}\{y[n]\} \equiv \sum_{n=0}^{\infty} y[n] z^{-n}, \quad y[n] = 0 \quad \forall n < 0,$$

$$y[n] = \mathcal{Z}^{-1}\{\mathbf{y}(z)\} \equiv \frac{1}{2\pi j} \oint_{\Gamma} \mathbf{y}(z) z^{-1} dz = \sum_k \text{Res} [\mathbf{y}(z) z^{n-1}; z_k]$$

$$\widehat{\mathbf{X}}(z, s) = \frac{z\widehat{I} - \widehat{\mathbf{M}}^{-1} \det \widehat{\mathbf{M}}}{\det (z\widehat{I} - \widehat{\mathbf{M}})} \left(\mathcal{Z}\{\Delta x'_d[n]\} \widehat{M}_d \widehat{E} + z\widehat{X}[0, s] + \frac{g z^{-\hat{q}} \mathbf{K}(z) \delta x_p}{(1 - z^{-1})(\hat{\beta}_k \hat{\beta}_p)^{1/2} K_0} \widehat{M}_k \widehat{E} \right)$$

$$\mathbf{K}(z) = \mathbf{K}_{\text{out}} H(z) \mathbf{K}_{\text{in}}$$

$$H(z) = \mathcal{Z}\{h[n]\}$$

$$g = (\hat{\beta}_k \hat{\beta}_p)^{1/2} K_0 S_k S_p$$

It can be done if $|z_k| < 1$ and

$$\lim_{n \rightarrow \infty} \widehat{X}[n, s] = \lim_{z \rightarrow 1} (z - 1) \widehat{\mathbf{X}}(z, s) = 0.$$

Consequently $\mathbf{K}(z = 1) = 0$.

$$H_{\text{NF}}(z) = (1 - z^{-1}).$$

$$\widehat{\mathbf{M}} \equiv \widehat{\mathbf{M}}(z, s) = \widehat{M}(s) + \frac{g z^{-\hat{q}} \mathbf{K}(z)}{(\hat{\beta}_k \hat{\beta}_p)^{1/2} K_0} \widehat{M}_k \widehat{T} \widehat{M}(s_p | s), \quad \widehat{T} \equiv \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$$

$$z_k^2 - \left(2 \cos(2\pi Q) + \frac{g z_k^{-\hat{q}} \mathbf{K}(z_k)}{K_0} \sin(2\pi Q - \psi_{\text{PK}}) \right) z_k + 1 - \frac{g z_k^{-\hat{q}} \mathbf{K}(z_k)}{K_0} \sin \psi_{\text{PK}} = 0$$

$$g = 0 \quad \Rightarrow \quad z_{1,2}^{(0)} = \exp(\pm j 2\pi Q)$$

TFS: Damping Parameters

$$x[n, s_K] = \sum_k A_k \sqrt{\hat{\beta}_K} e^{-n\alpha_k + j(2\pi n\{Q_k\} + \Phi_k)}$$

$$\alpha_k \equiv -\ln |z_k| = \frac{T_{\text{rev}}}{\tau_k}, \quad \{Q_k\} = \frac{1}{2\pi} \arg z_k$$

If $g \ll 1$ and $\mathbf{K}_{\text{in}}(\omega)\mathbf{K}_{\text{out}}(\omega)$ depends weakly on frequency then:

$$\alpha_m = \frac{g |\mathbf{K}(\omega_m)|}{2K_0} \sin \Psi_{\text{PK}},$$

$$\omega_m = 2\pi(Q + m)f_{\text{rev}}$$

$$\{Q_m\} = \{Q\} - \frac{g |\mathbf{K}(\omega_m)|}{4\pi K_0} \cos \Psi_{\text{PK}},$$

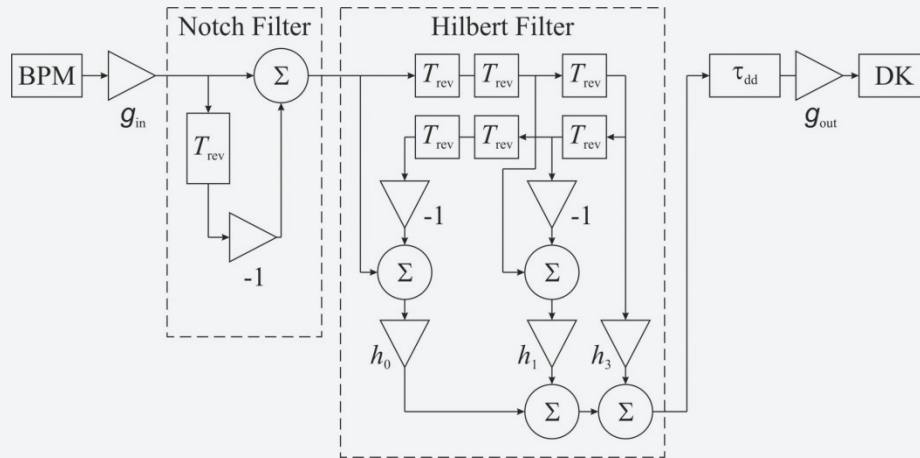
$$-0.5 < \{Q\} \leq 0.5.$$

$$\Psi_{\text{PK}} = \psi_{\text{PK}} + 2\pi\hat{q}Q - \arg \mathbf{K}(\omega_m),$$

$$\mathbf{K}(\omega_m) = \mathbf{K}_{\text{in}}(\omega_m) \mathbf{K}_{\text{out}}(\omega_m) H(z = e^{j2\pi Q}),$$

$$|K_0| = |\mathbf{K}(\omega_{\text{min}})|, \quad K_0 \sin \Psi_{\text{PK}}(\omega_{\text{min}}) > 0,$$

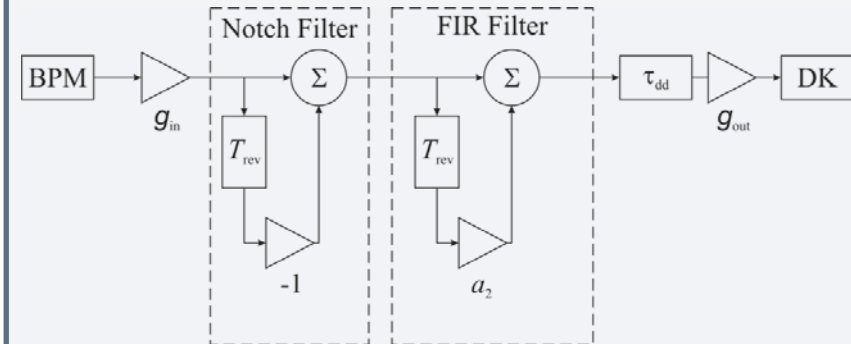
TFS with Notch and Hilbert Filters



$$H_1(z) \equiv H_N(z) H_{HF}(z) = (1 - z^{-1}) (h_0 z^{-3} + h_1 z^{-2} (1 - z^{-2}) + h_3 (1 - z^{-6}))$$

$$h_0 = \cos(\Delta\varphi), \quad h_1 = \frac{2}{\pi} \sin(\Delta\varphi), \quad h_3 = \frac{2}{3\pi} \sin(\Delta\varphi)$$

TFS with Notch and FIR-1 Filters



$$H_2(z) \equiv H_N(z) H_{FIR}(z) = (1 - z^{-1}) (1 + a_2 z^{-1})$$

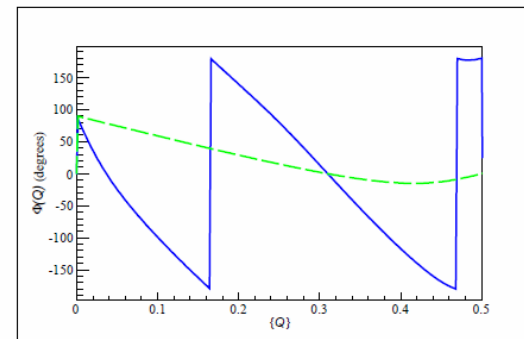
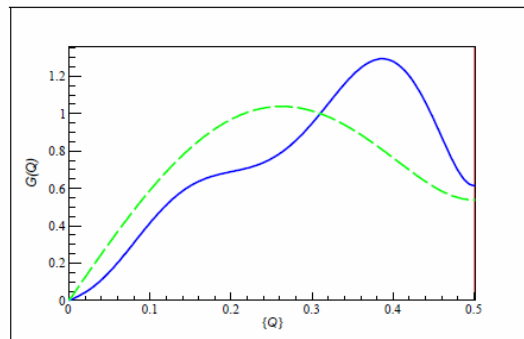
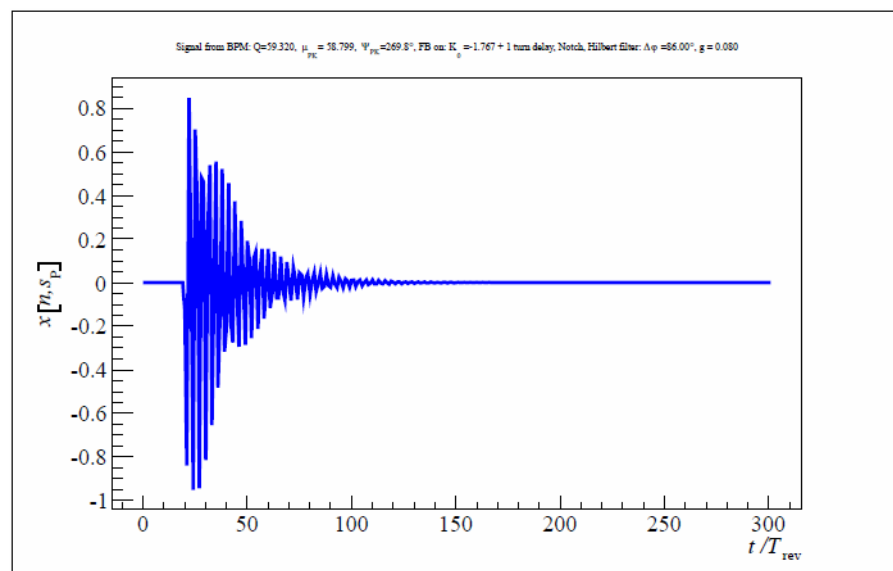
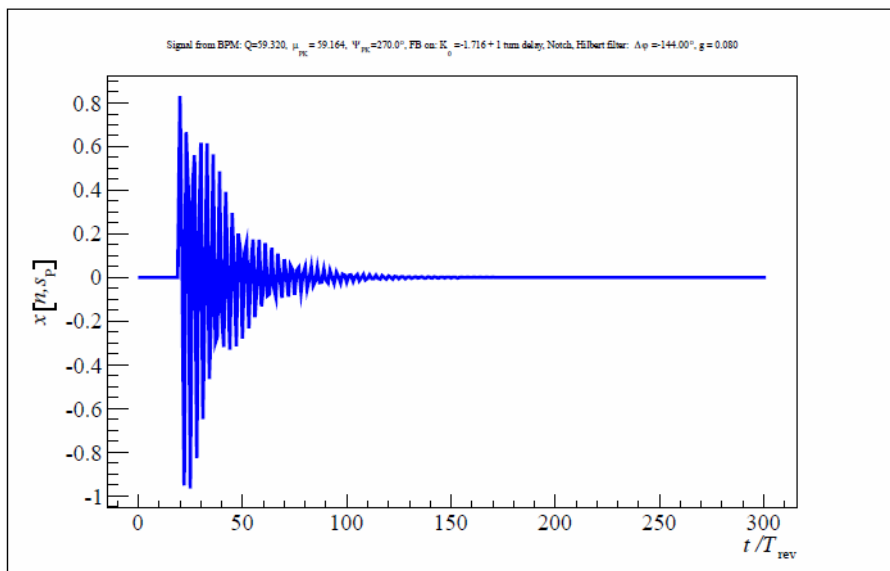
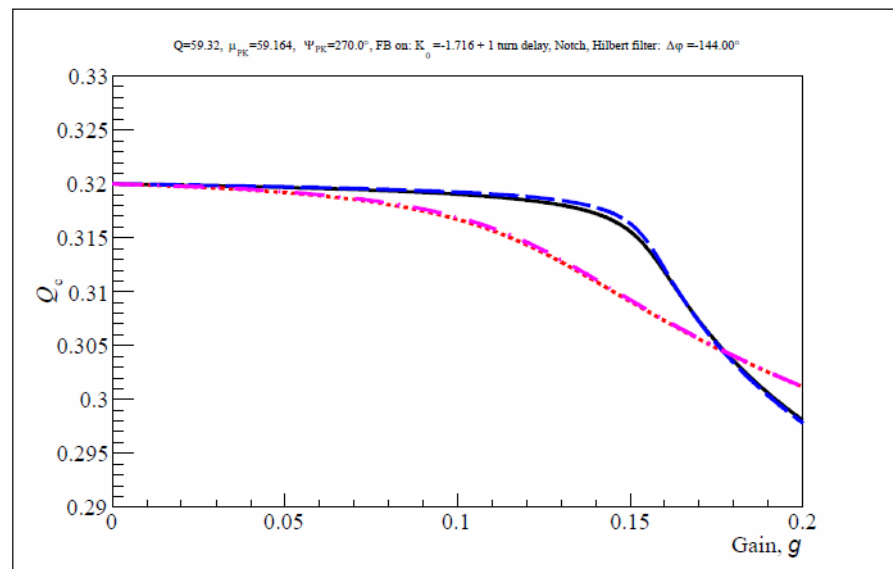
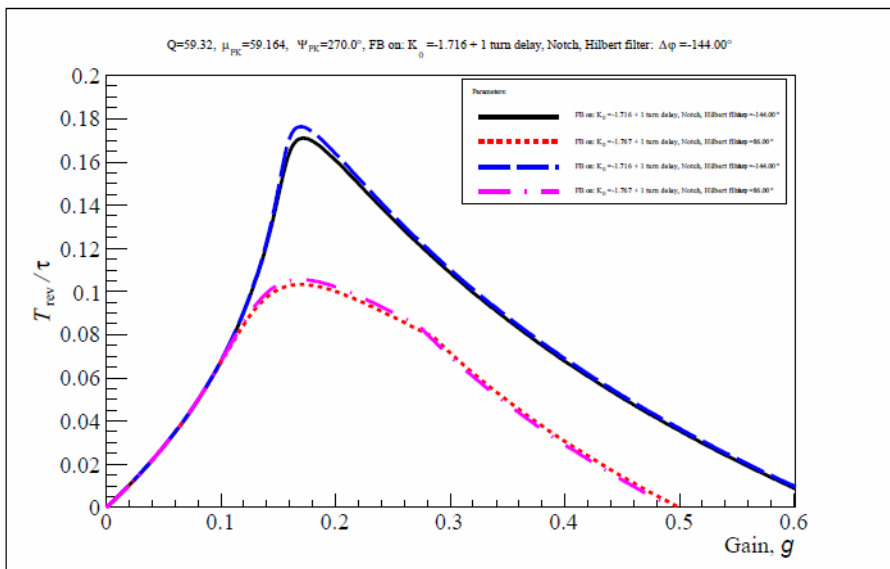


Figure 2: The magnitude $G(Q)$ and phase response $\Phi(Q)$ graphics for the notch and Hilbert filters (solid curve) and for the notch filter and the FIR filter of the first order (dashed curve)

TFS & Beam response to δ -kick



TFS & Beam response on δ -kick

W. Höfle et al.
CERN ATS/Note/2011/131
2011-12-12

Beam	Plane	Gain setting (dB)	Damping time [turns]					
			Q7+	Q7-	Q9+	Q9-	ave	StDev
1	hor.	-49	1188	729	1247	1172	1084	239
1	hor.	-43	1096	613	969	1066	936	222
1	hor.	-37	884	564	476	873	699	210
1	hor.	-31	370	383	902	414	517	257
1	hor.	-25	692	278	594	288	463	212
1	hor.	-19	294	203	280	200	244	50
1	hor.	-13	214	290	214	285	251	42
1	hor.	-7	83	89	78	89	85	5
1	hor.	-1	43	39	36	44	41	5
1	vert.	-49	887	800	1049	999	934	112
1	vert.	-25	558	496	592	594	560	46
1	vert.	-19	460	408	499	467	458	38
1	vert.	-13	301	165	341	317	281	79
1	vert.	-7	86	89	—	—	88	2
1	vert.	-1	47	39	48	—	44	5
2	hor.	-49	1735	1758	1828	1677	1750	62
2	hor.	-25	398	422	395	419	408	14
2	hor.	-19	251	328	574	314	367	142
2	hor.	-13	224	188	400	238	262	94
2	hor.	-7	89	83	87	91	87	4
2	hor.	-1	45	41	31	56	43	10
2	vert.	-49	405	360	917	1224	726	417
2	vert.	-25	347	363	484	411	401	62
2	vert.	-19	172	243	295	170	220	60
2	vert.	-13	70	90	77	96	83	12
2	vert.	-7	43	42	46	46	44	2
2	vert.	-1	22	20	20	20	20	1

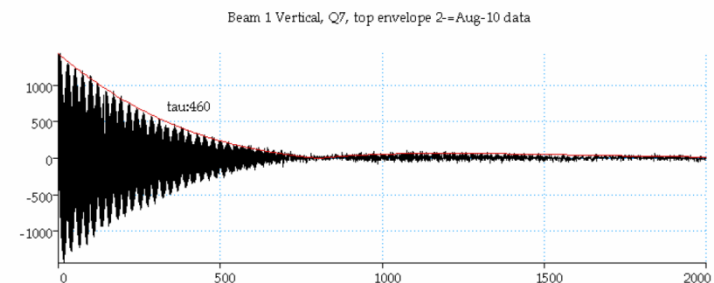
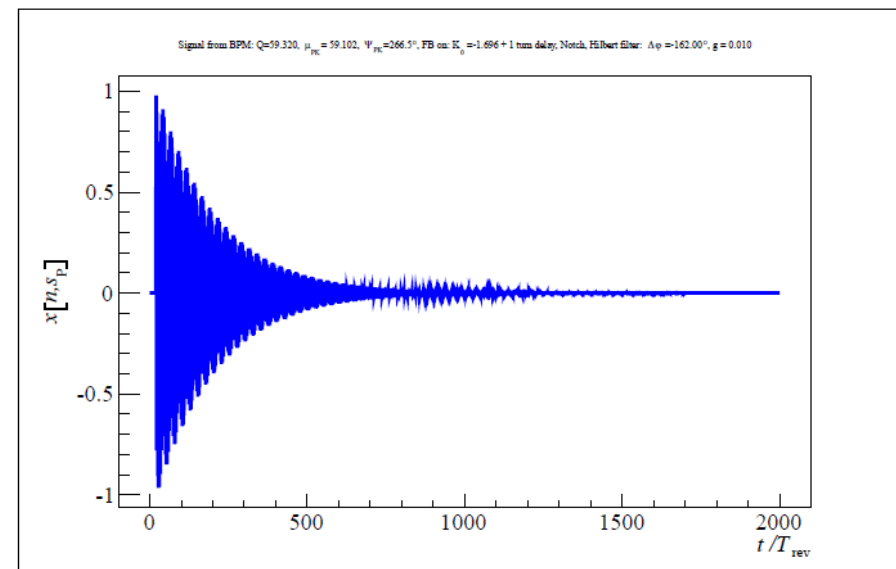
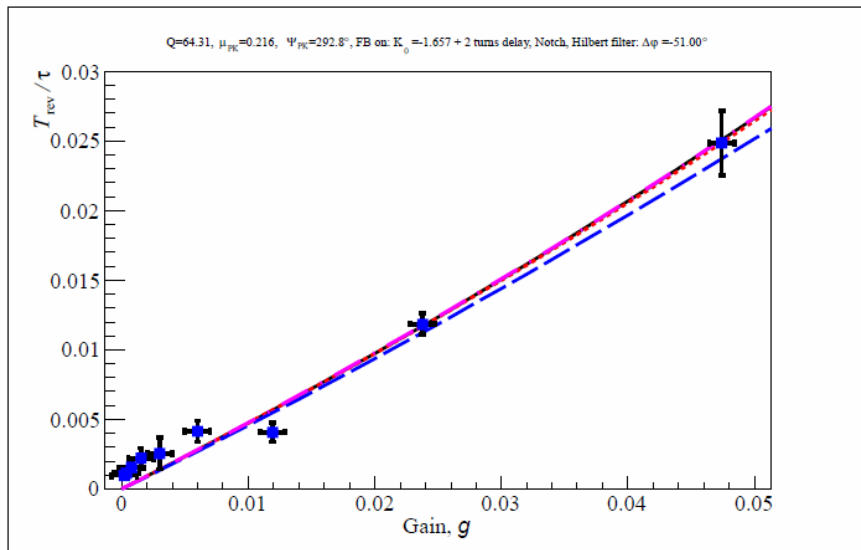
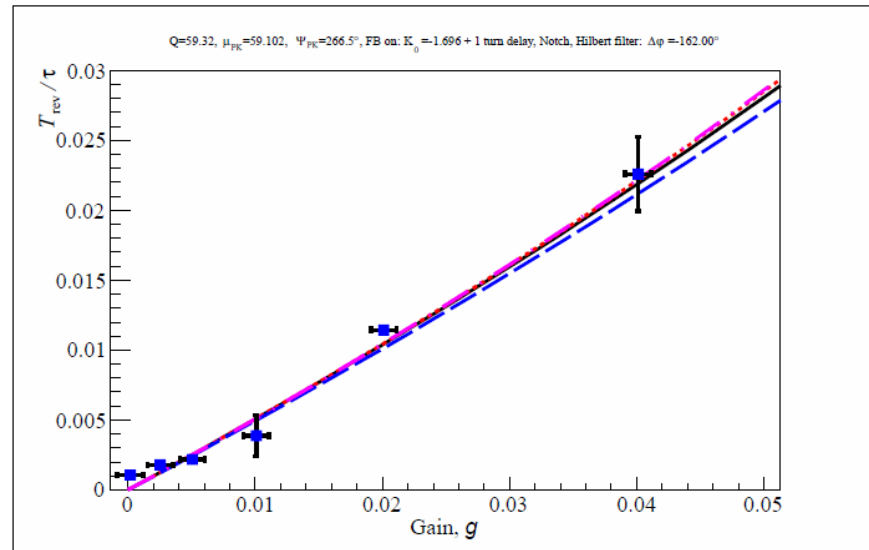


Figure 1.2: Envelope fit example for Beam 1, Q7 data, -19 dB gain setting.

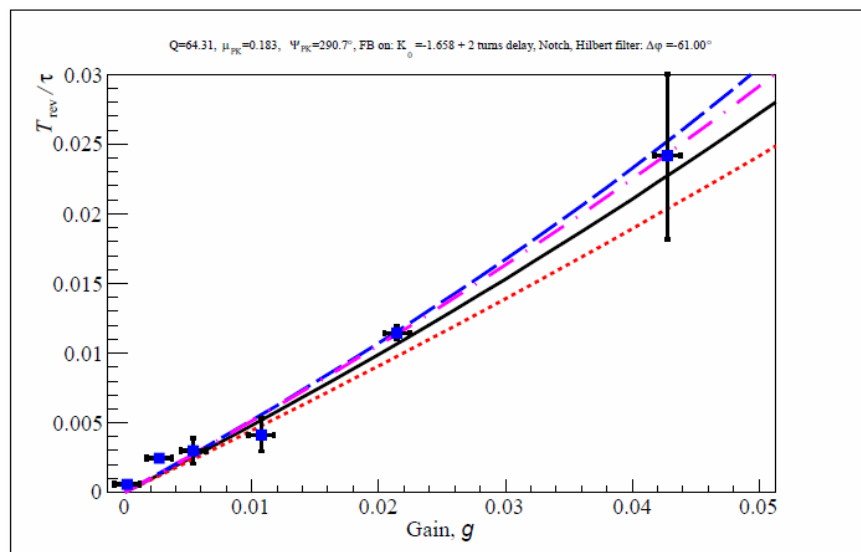




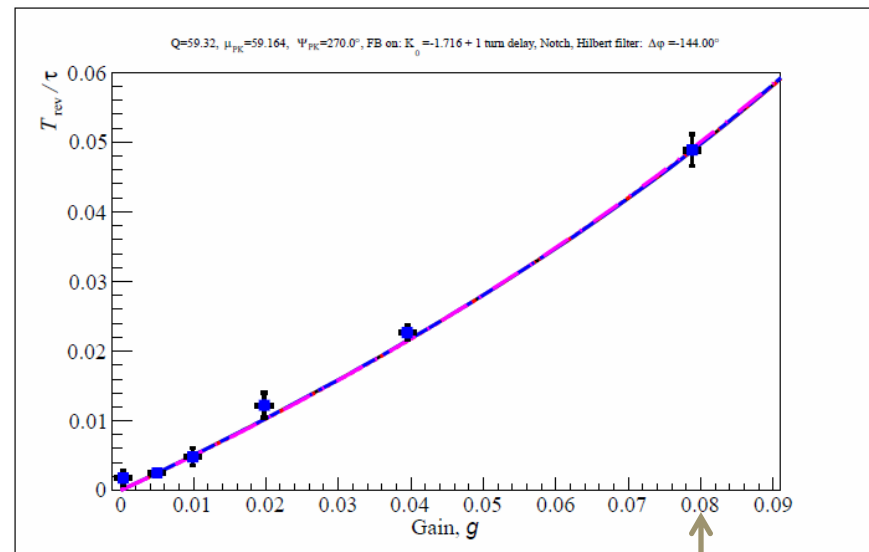
Beam 1, hor.



Beam 1, ver.



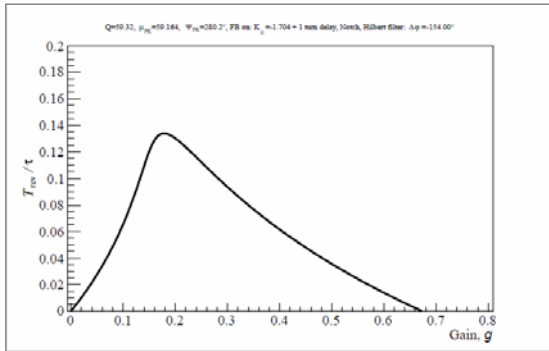
Beam 2, hor.



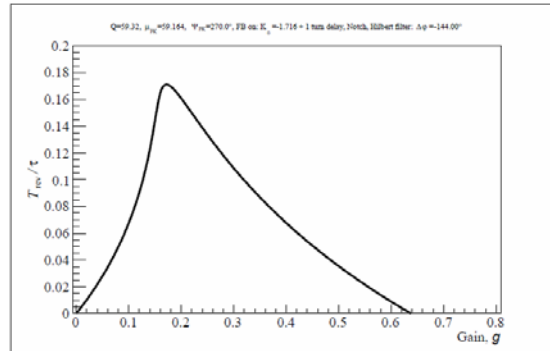
Beam 2, ver.

$a_{\max} < 0.15$ mm; $\Delta x' < 0.24$ μ rad; $E = 3.5$ TeV

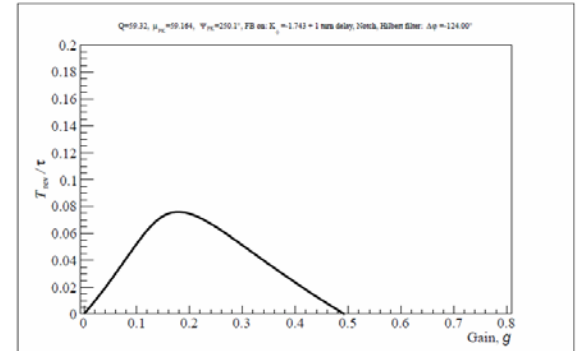
TFS: variation of parameters



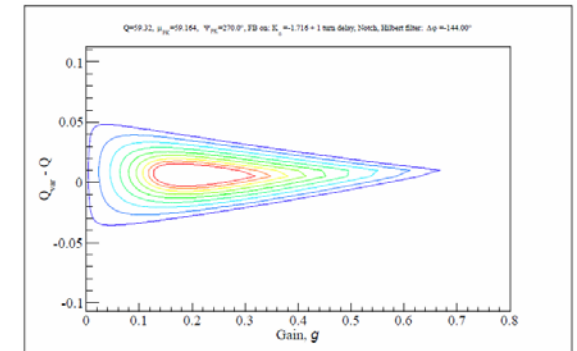
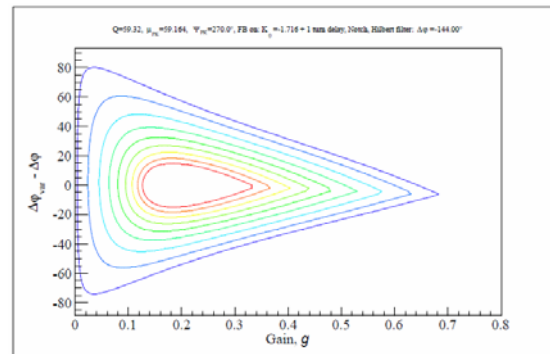
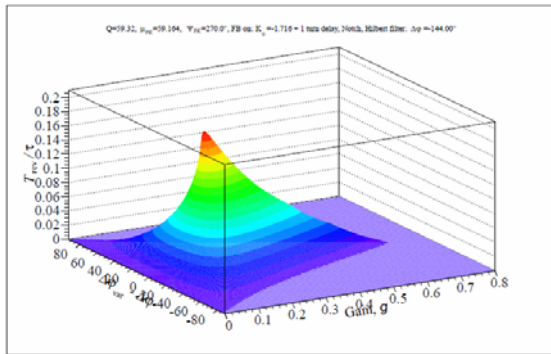
$$\Delta\varphi_{\text{var}} = \Delta\varphi_{\text{opt}} - 10$$



$$\Delta\varphi = \Delta\varphi_{\text{opt}}$$



$$\Delta\varphi_{\text{var}} = \Delta\varphi_{\text{opt}} + 20$$



$$T_{\text{rev}}/\tau \equiv \alpha(g, \Delta\varphi_{\text{var}} - \Delta\varphi) = \alpha_n;$$

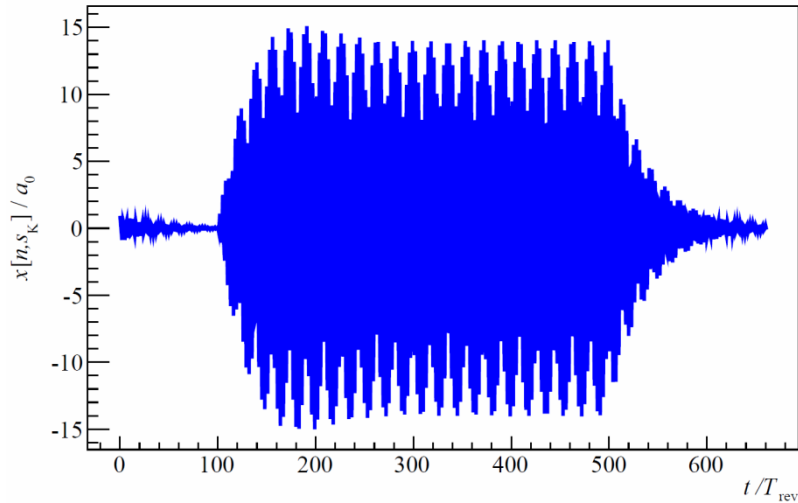
$$\alpha_0 = 0.002, \quad \alpha_n = n_c/80, \quad 1 \leq n_c \leq 8$$

$$T_{\text{rev}}/\tau \equiv \alpha(g, Q - Q_0) = \alpha_n;$$

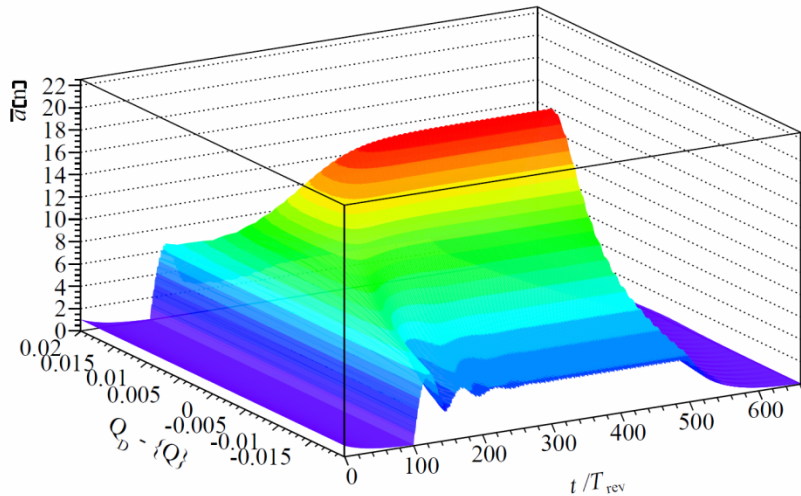
TFS & Sinusoidal driving force

Isolines: $a(t/T_{\text{rev}}, Q_D - \{Q\}) = a_n$

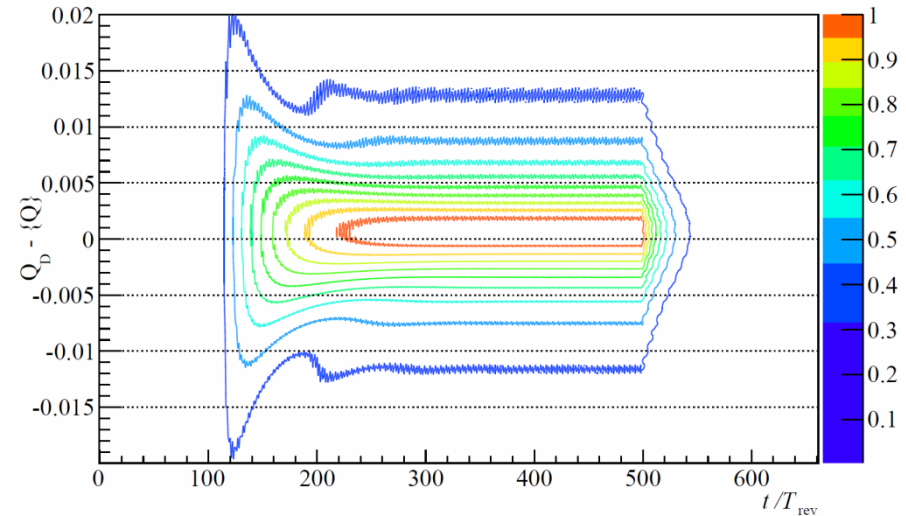
$Q=59.310$, $\mu_{\text{pk}} = 59.250$, $\Psi_{\text{pk}}=97.0^\circ$, FB on: $K_0=1.633$, FIR: $a_1=-1.00 + \text{FIR}$: $a_2=0.700$, $g = 0.050$



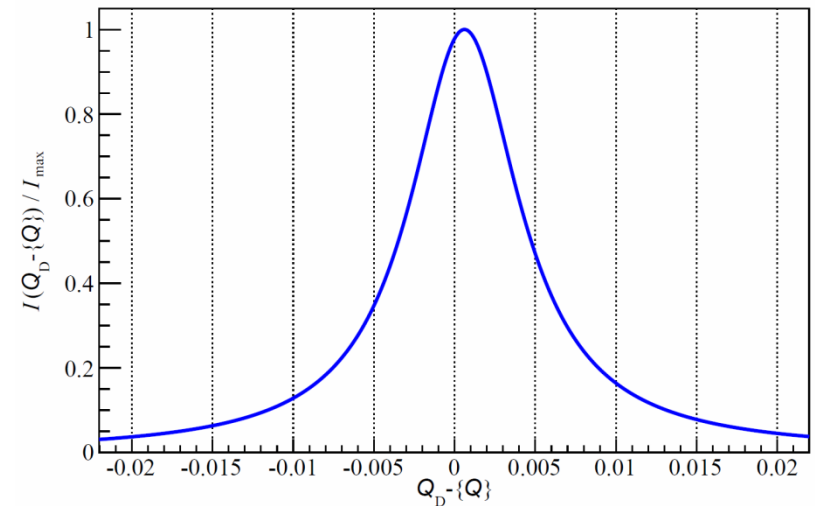
$Q=59.310$, $\mu_{\text{pk}} = 59.250$, $\Psi_{\text{pk}}=97.0^\circ$, FB on: $K_0=1.633$, FIR: $a_1=-1.00 + \text{FIR}$: $a_2=0.700$, $g = 0.050$



$Q=59.310$, $\mu_{\text{pk}} = 59.250$, $\Psi_{\text{pk}}=97.0^\circ$, FB on: $K_0=1.633$, FIR: $a_1=-1.00 + \text{FIR}$: $a_2=0.700$, $g = 0.050$



Sinusoidal driving force, $Q=59.31$, $\mu_{\text{pk}} = 59.250$, $\Psi_{\text{pk}}=97.0^\circ$, FB on: $K_0=1.633$, FIR: $a_1=-1.00 + \text{FIR}$: $a_2=0.700$, $g=0.050$



TFS: Resonance Curve

$$\frac{I(Q_D)}{I_{\max}} = \frac{1}{I_{\max}} \left| \det \left(z_D \hat{I} - \widehat{\mathbf{M}}(z_D) \right) \right|^{-2},$$

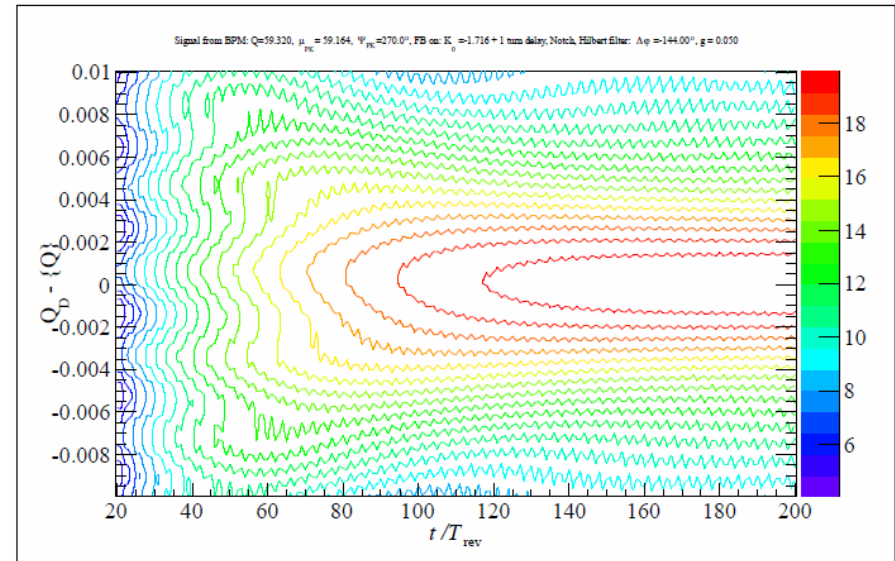
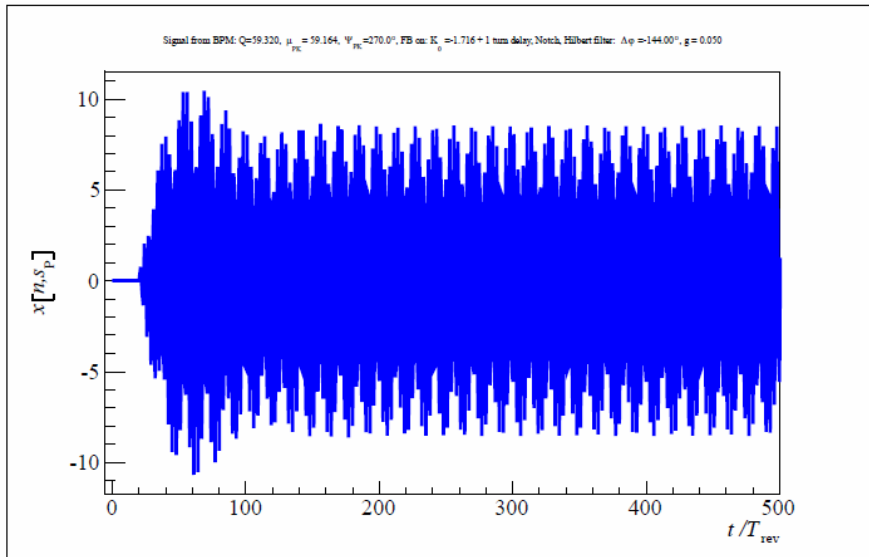
$$z_D = \exp(j 2\pi \{Q_D\}), \quad I_{\max} = I \left(Q_D^{(\max)} \right).$$

If $g \ll 1$:

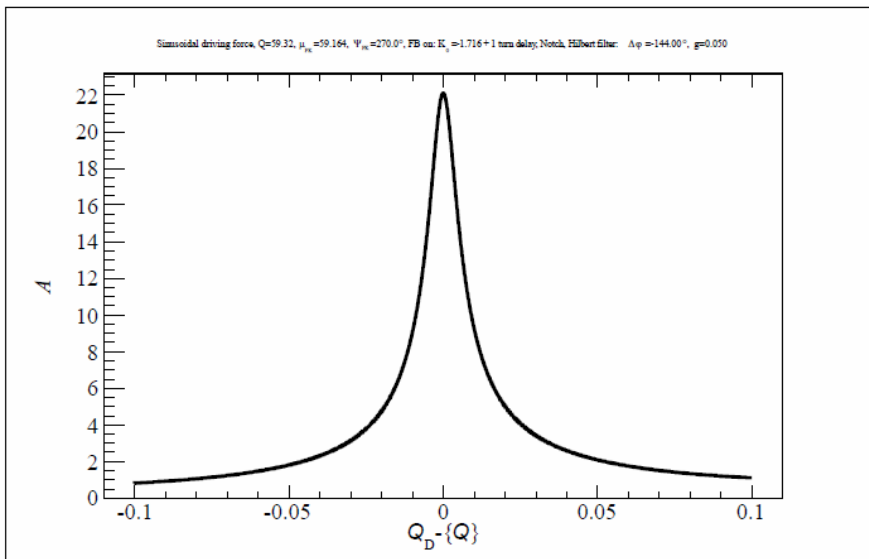
$$\frac{I(Q_D)}{I_{\max}} = \frac{\alpha_m^2}{4\pi^2(\{Q_D\} - \{Q_m\})^2 + \alpha_m^2}.$$

$$\{Q_D^{(\max)}\} = \{Q_m\} = \{Q\} - \frac{g |\mathbf{K}(\omega_m)|}{4\pi K_0} \cos \Psi_{PK}.$$

TFS & Sinusoidal driving force



$$a(t/T_{rev}, Q_D - \{Q\}) = a_n$$

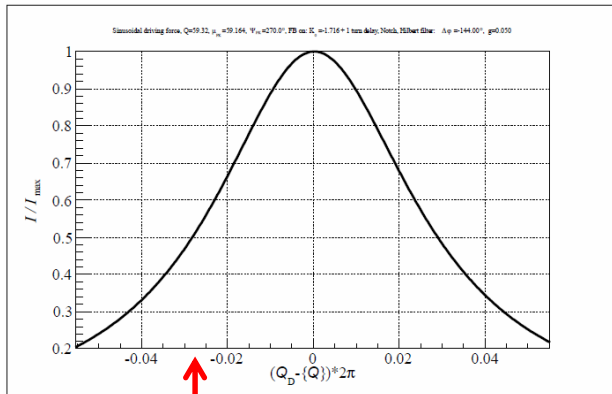
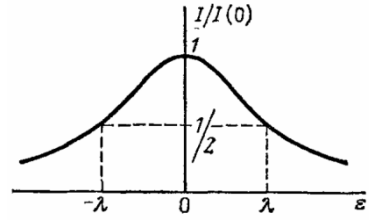


Resonance

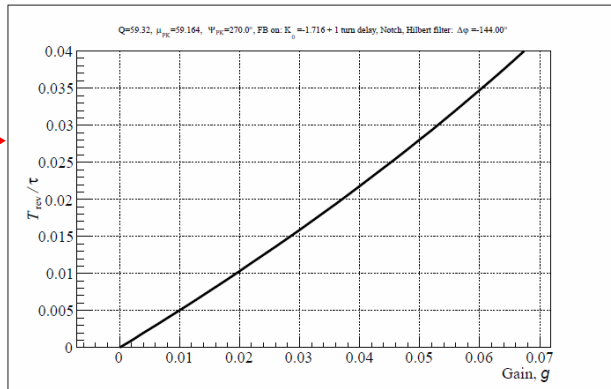
$$\ddot{x} + 2\lambda x + \omega^2 x = \frac{f_D}{m} \cos(\omega_D t)$$

Intensity of oscillations:

$$I(\epsilon = \omega_D - \omega) \propto \frac{\lambda^2}{\epsilon^2 + \lambda^2}$$



$\alpha=0.028$

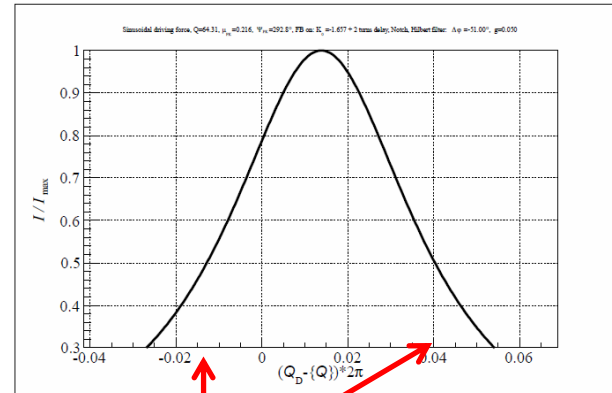


Synchrotrons:

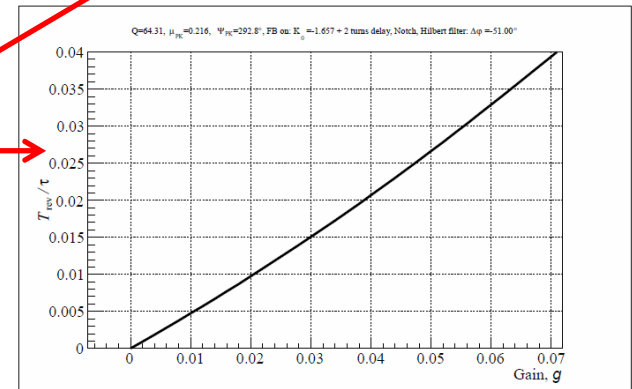
$$\omega_m = Q_m \omega_{\text{rev}} = 2\pi Q_m / T_{\text{rev}}; \quad \alpha_m = \lambda T_{\text{rev}}$$

$$I(Q_D) \propto \frac{\alpha_m^2}{4\pi^2(\{Q_D\} - \{Q_m\})^2 + \alpha_m^2}$$

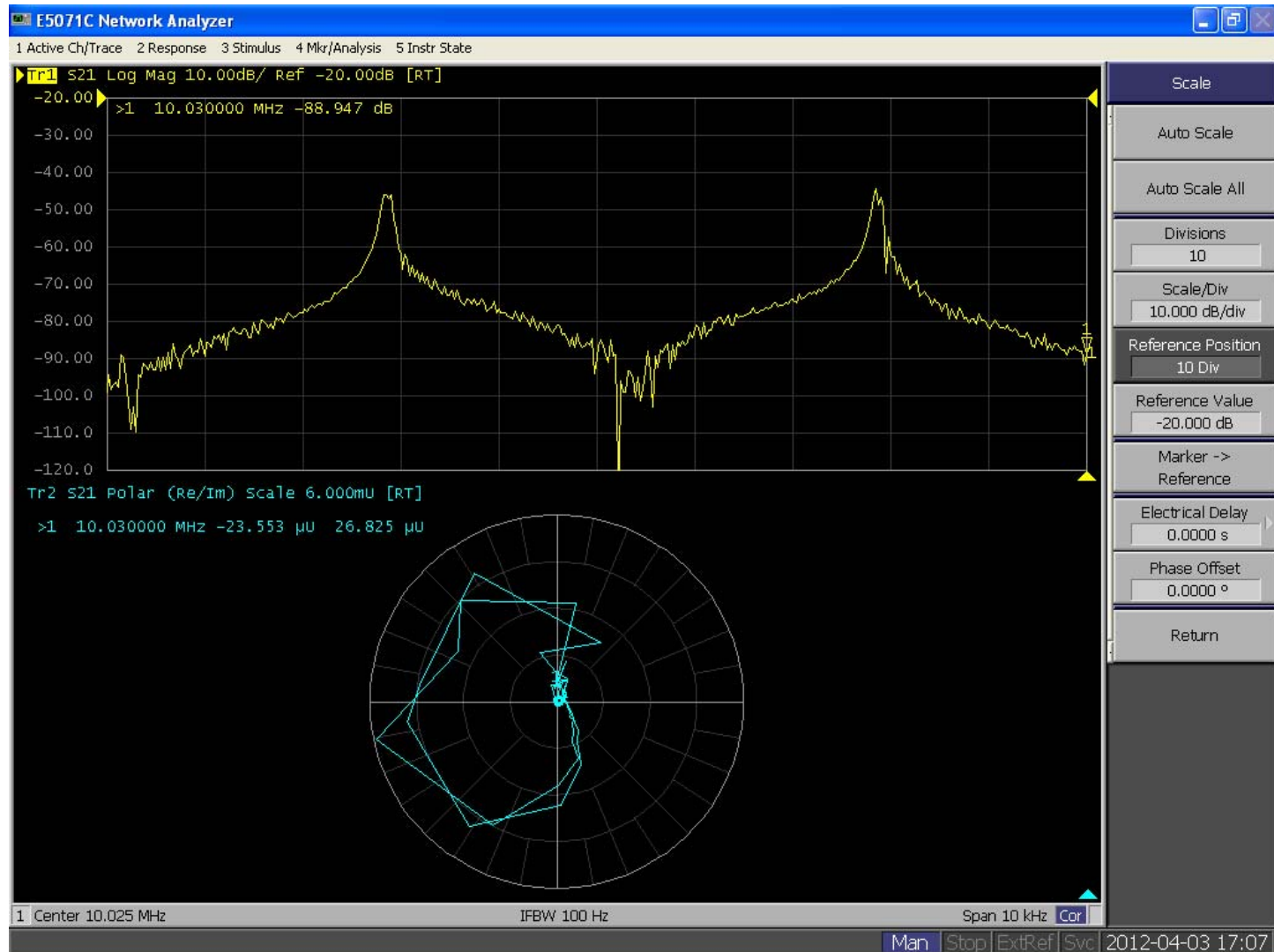
$$\{Q_m\} = \{Q\} - \frac{g |K(\omega_m)|}{4\pi K_0} \cos \Psi_{PK}$$



$\alpha=0.027$



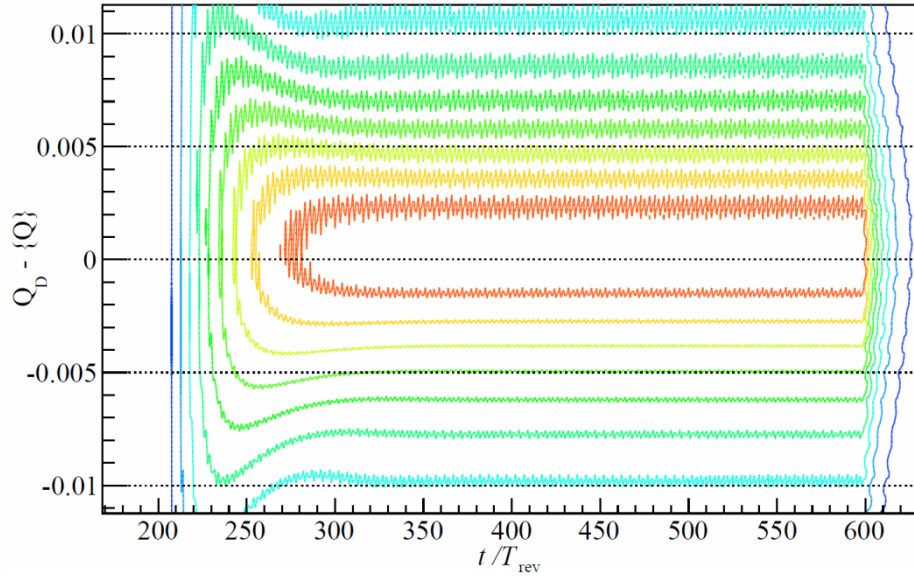
LHC Damper



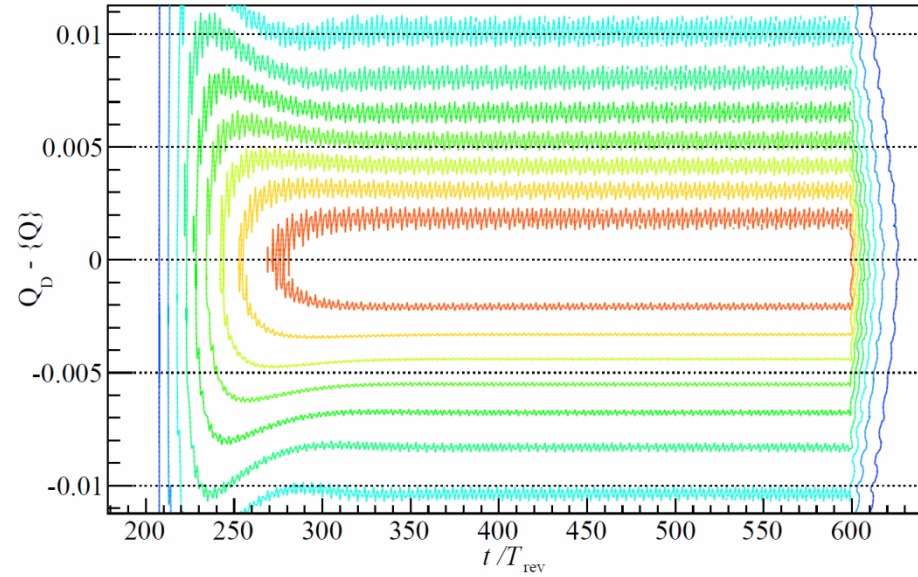
W. Höfle, D. Valuch (CERN)

TFS & Sinusoidal driving force

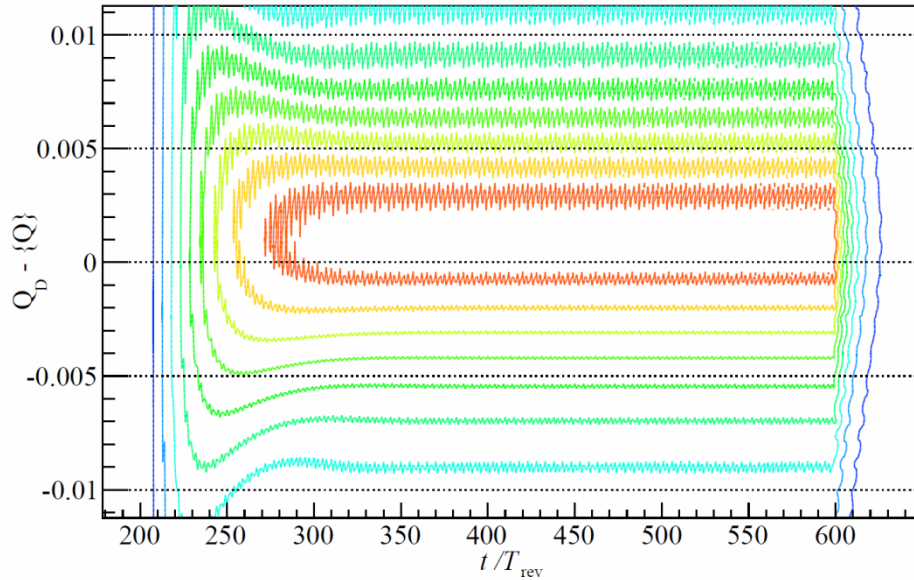
$Q=59.310$, $\mu_{PK}=59.250$, $\Psi_{PK}=92.0^\circ$, FB on: $K_0=1.589$, FIR: $a_1=-1.00 + \text{FIR}$: $a_2=0.610$, $g=0.080$



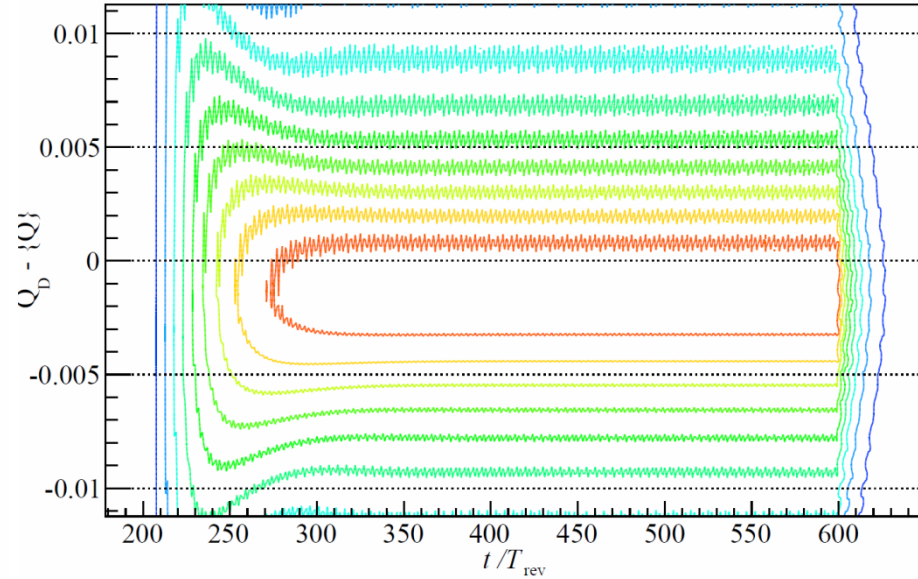
$Q=59.310$, $\mu_{PK}=59.250$, $\Psi_{PK}=87.9^\circ$, FB on: $K_0=1.564$, FIR: $a_1=-1.00 + \text{FIR}$: $a_2=0.540$, $g=0.080$



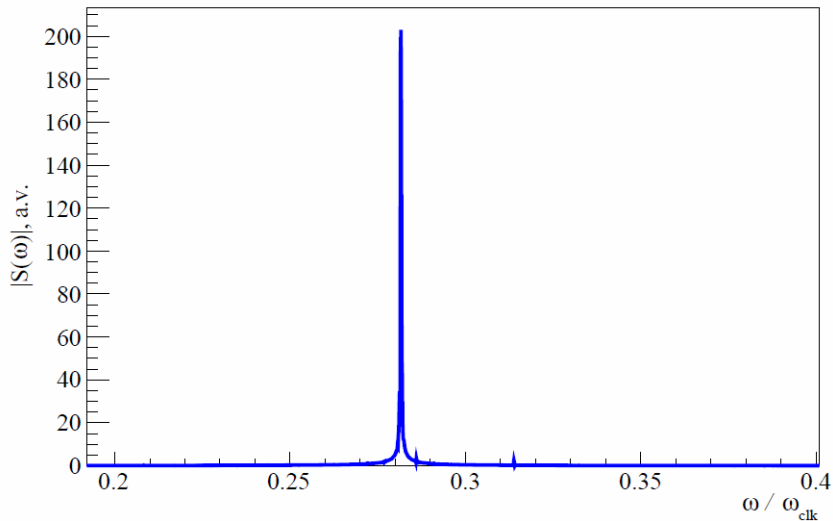
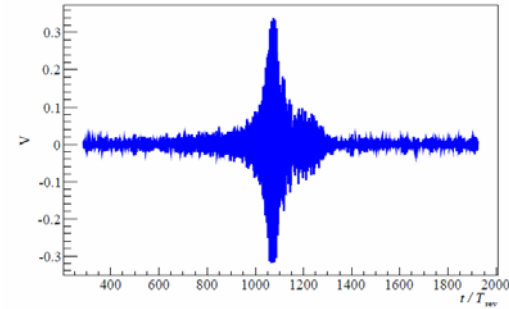
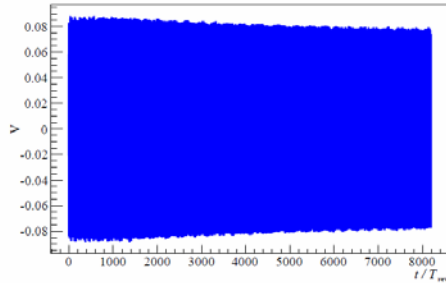
$Q=59.310$, $\mu_{PK}=59.250$, $\Psi_{PK}=97.0^\circ$, FB on: $K_0=1.633$, FIR: $a_1=-1.00 + \text{FIR}$: $a_2=0.700$, $g=0.080$



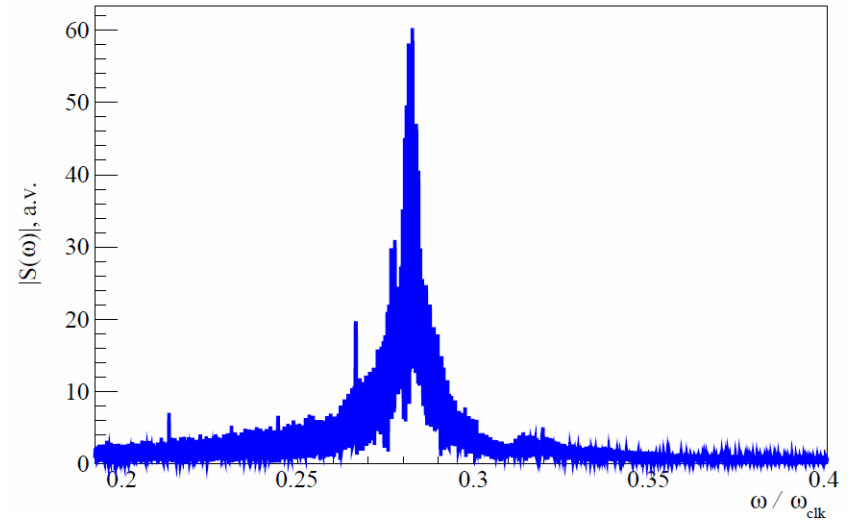
$Q=59.310$, $\mu_{PK}=59.250$, $\Psi_{PK}=79.4^\circ$, FB on: $K_0=1.539$, FIR: $a_1=-1.00 + \text{FIR}$: $a_2=0.400$, $g=0.080$



Base-Band Q (BBQ) Measurements at the LHC



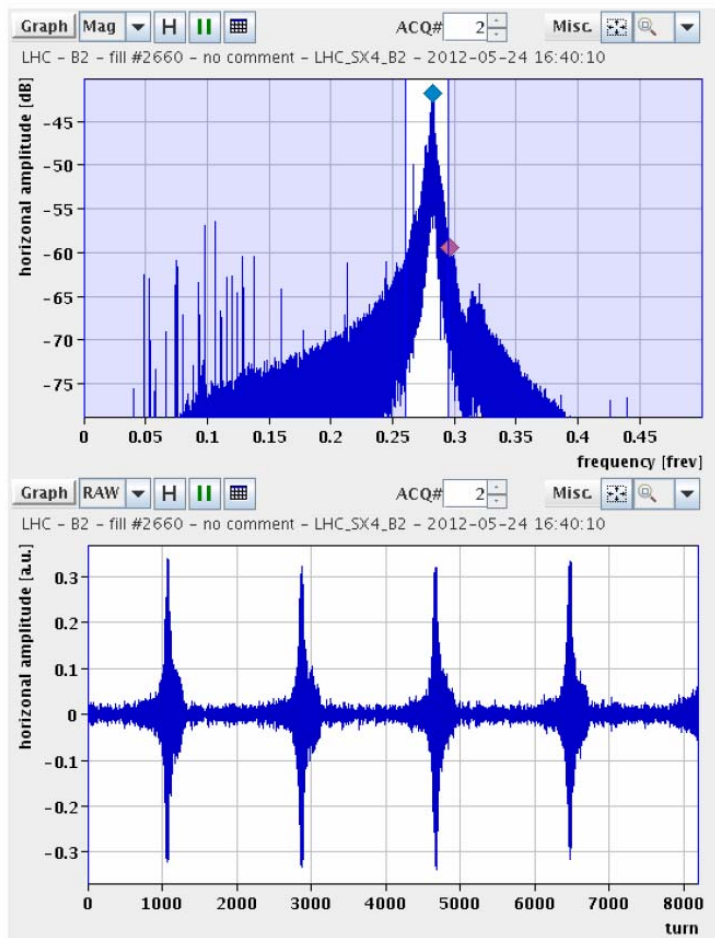
Damper OFF (oscillations after injection)



Damper ON (chirp excitation)

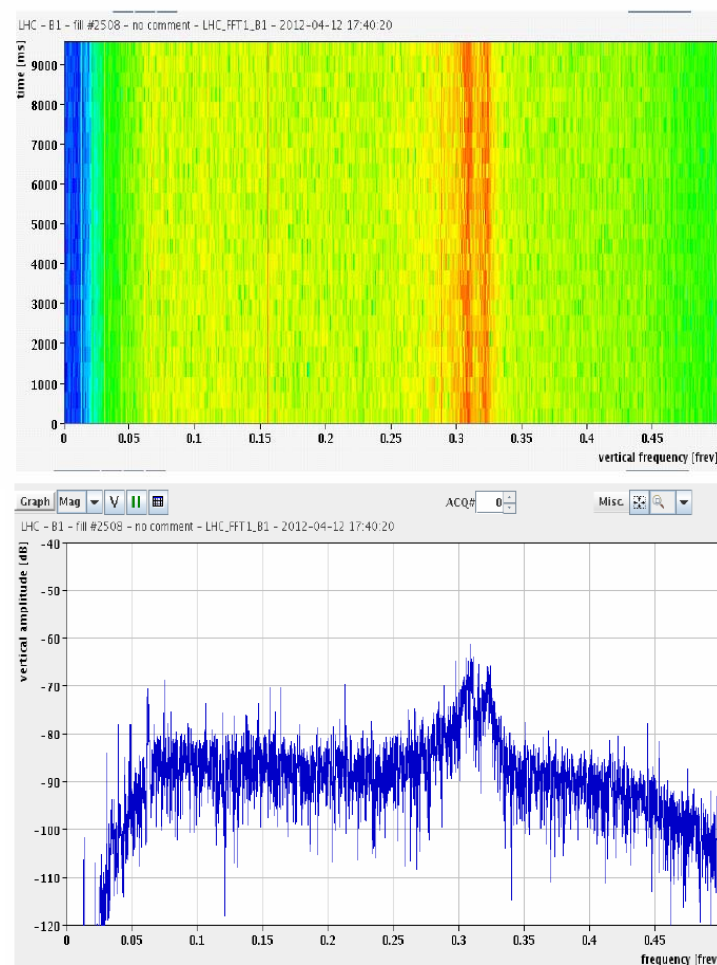
M. Gasior (CERN)

Base-Band Q (BBQ) Measurements at the LHC



Damper ON (chimp excitation)

M. Gasior (CERN)



Q-map (top) and Q-line (bottom)
from BBQ Display

CONTROLLED TRANSVERSE BLOW-UP OF HIGH-ENERGY PROTON BEAMS FOR APERTURE MEASUREMENTS AND LOSS MAPS

W. Hofle ✱, R. Assmann, S. Redaelli, R. Schmidt, D. Valuch, D. Wollmann, M. Zerlauth,
CERN, Geneva, Switzerland

Proceedings of IPAC2012, New Orleans, Louisiana, USA. Pp.4059-4061

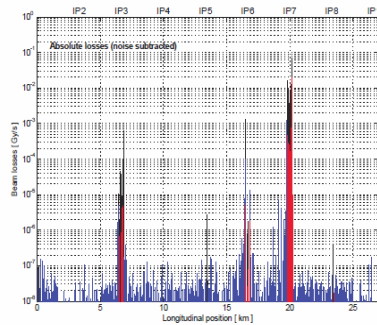


Figure 4: Loss maps obtained with the damper blow-up method with one nominal bunch in the LHC (blue: cold part of LHC, red and black: warm part).

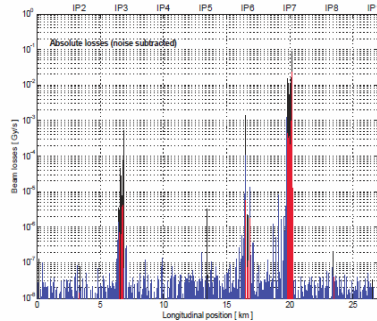
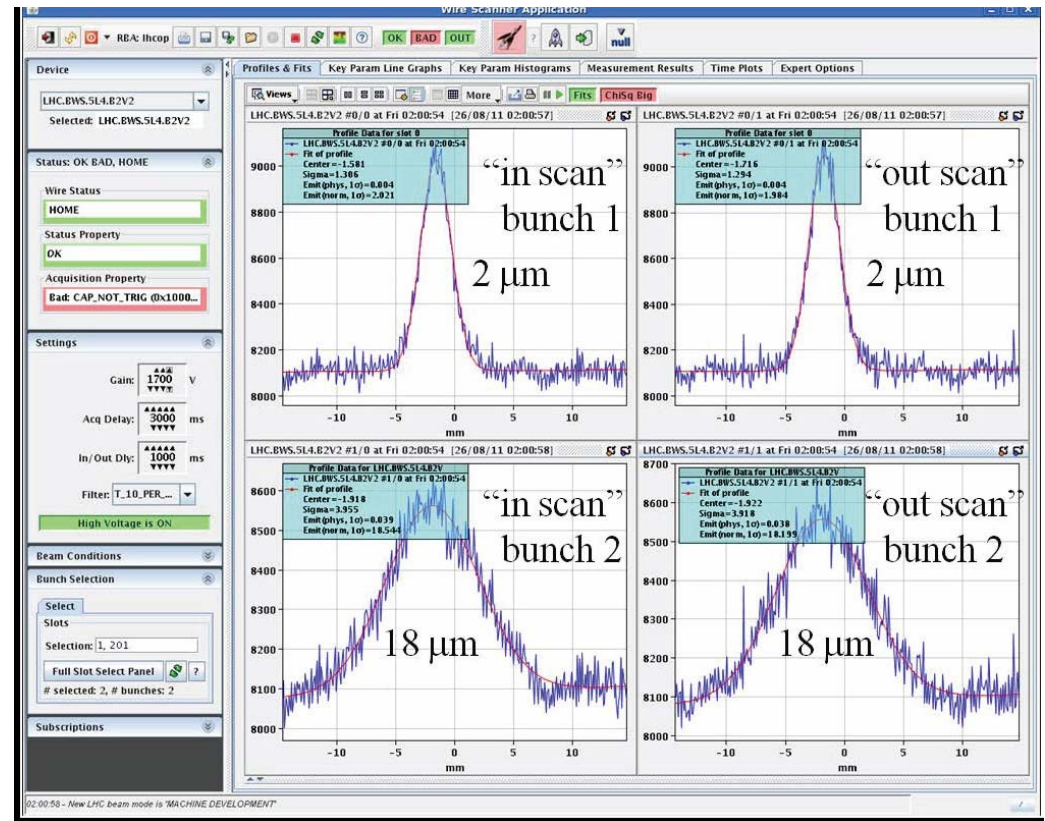
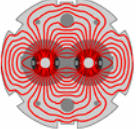


Figure 5: Loss maps obtained by crossing the third-order resonance with one nominal bunch in the LHC (blue: cold part of LHC, red and black: warm part).

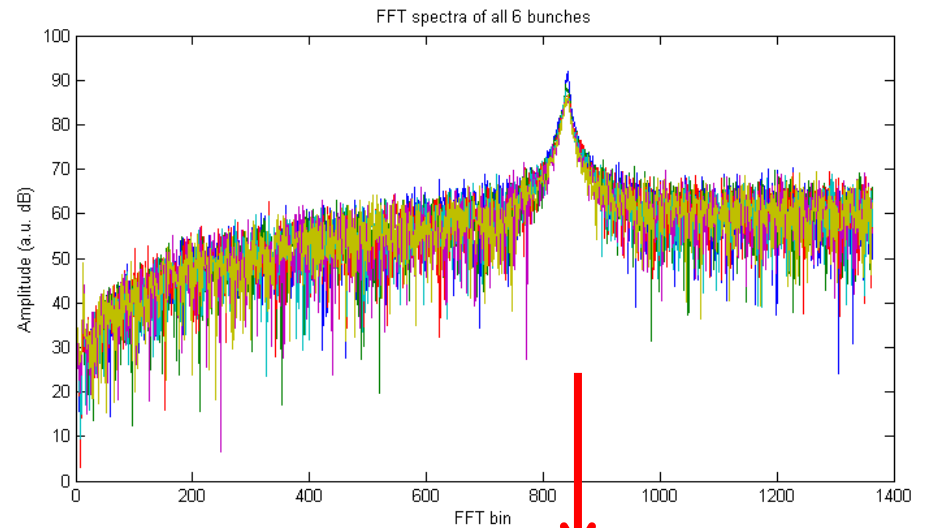
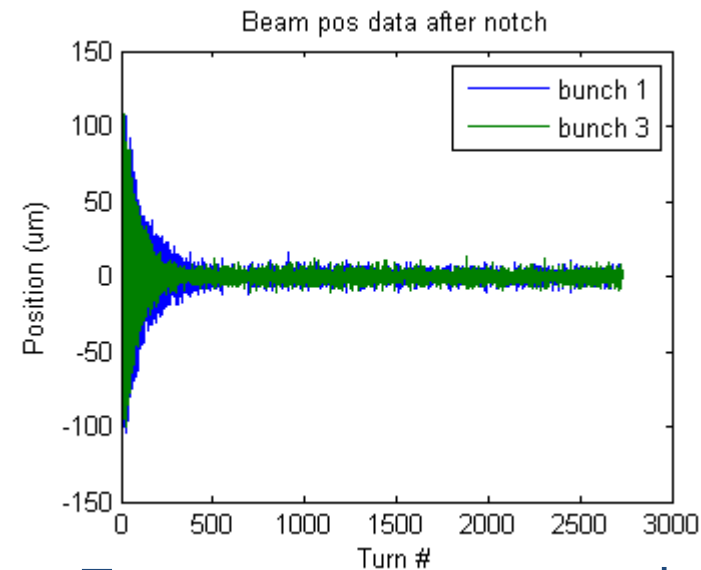
One can get the aperture in one single measurement taking 1-2 minutes only.



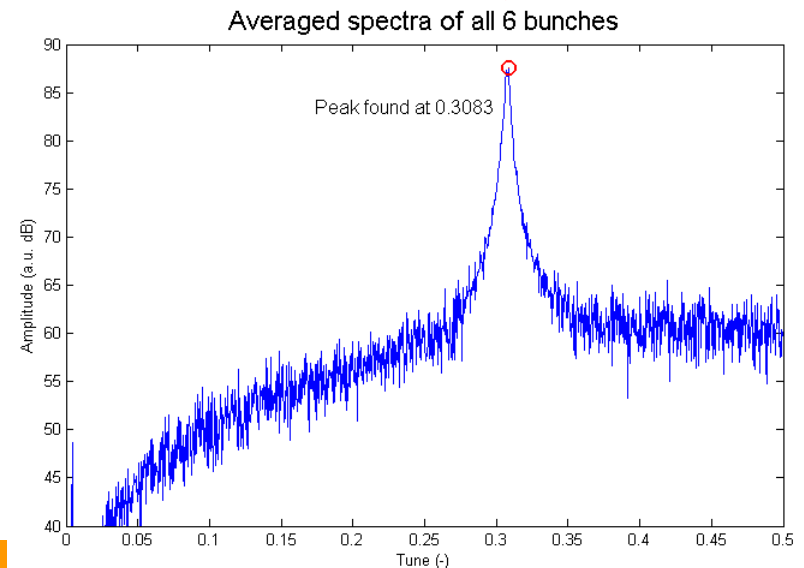
Profiles measured with the wire scanners after a blow-up targeted to bunch 2: blow-up to aperture limit of this bunch; bunch 1 emittance unchanged.



Wednesday: ADT tune measurement test



- Test tune measurement by selective excitation of 6 bunches.
- FFT of data for **each observed bunch**, average the spectra, find the peak
- **Very successful test.** Validation during a ramp required (including emittance blow-up measurement) + operational implementation



W. Höfle, D. Valuch (CERN)

Conclusion

Measurements of a beam response to the δ -impulse are the effective approach to tune the transverse feedback system in the damping mode of coherent betatron oscillations of the bunch.

Observation of a beam response to the harmonic impulse can be a good instrument for selective measurements of circulated bunches because of the dedicated resonance behavior of the detected signal in synchrotrons with a digital transverse feedback system.

Thank you for your attention!