

TRANSFORMATION of BEAMS in PLASMA LENS and INVESTIGATION of Z-PINCH DYNAMICS



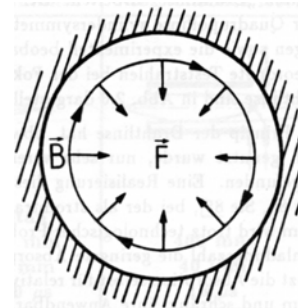
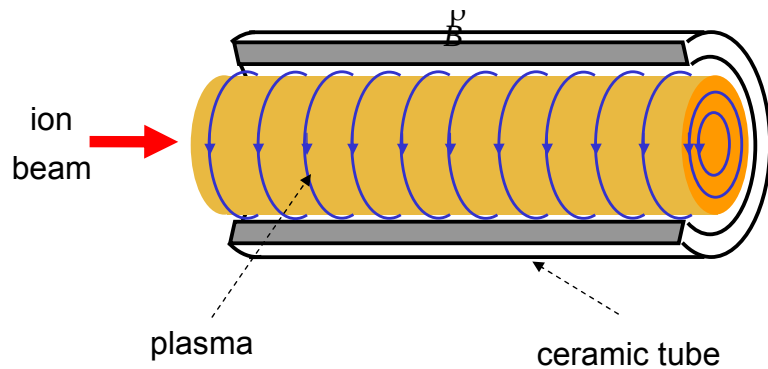
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Yu.Novozhilov, P.Sasorov, S.Savin, V.Yanenko*

Abstract. The plasma lens can carry out not only sharp focusing of ions beam. Plasma lens provides formation of hollow beams of ions. Application of the several plasma lenses allow to get a conic and a cylindrical beams. The plasma lens can be used for obtaining a beams with homogeneous spatial distribution.

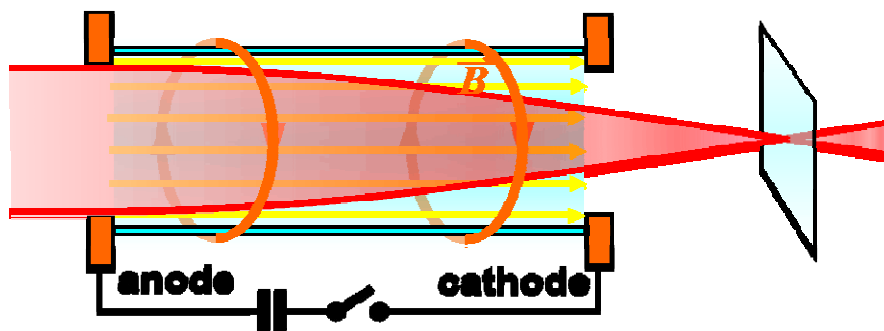
By measuring the distribution of ions in the beam at the input and output lenses can determine the current distribution in the pinch through the numerical calculation.

Ion focusing in a plasma lens



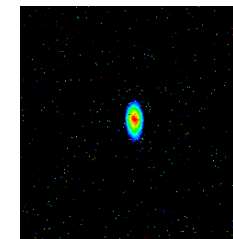
$$F = Ze v_z B_\phi(r)$$

$$B = \frac{\mu_0}{2} j r$$

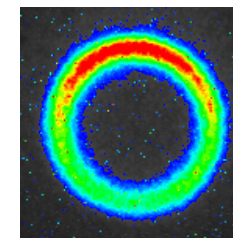


linear B-field

nonlinear B-field



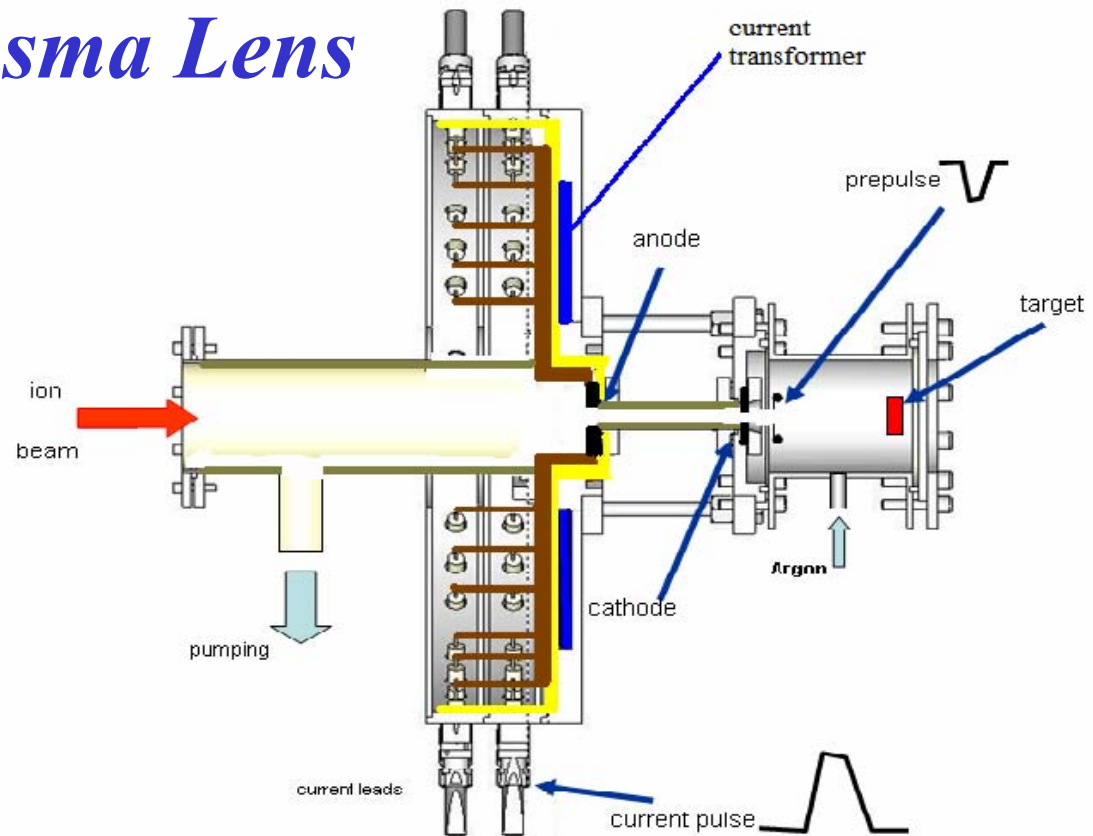
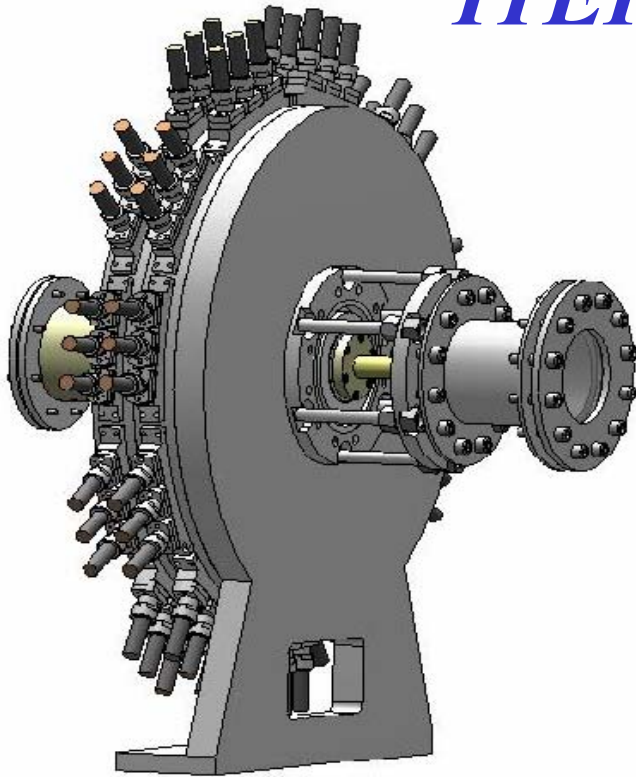
$$d_{min} \sim \epsilon l^{-1/2} .$$



Advantages of plasma lens :

- focusing power is proportional to the magnetic field strength;
- there is no limit for the magnetic field magnitude connected with saturation;
- charge neutralization of ion beam into the plasma lens ;

ITEP Plasma Lens

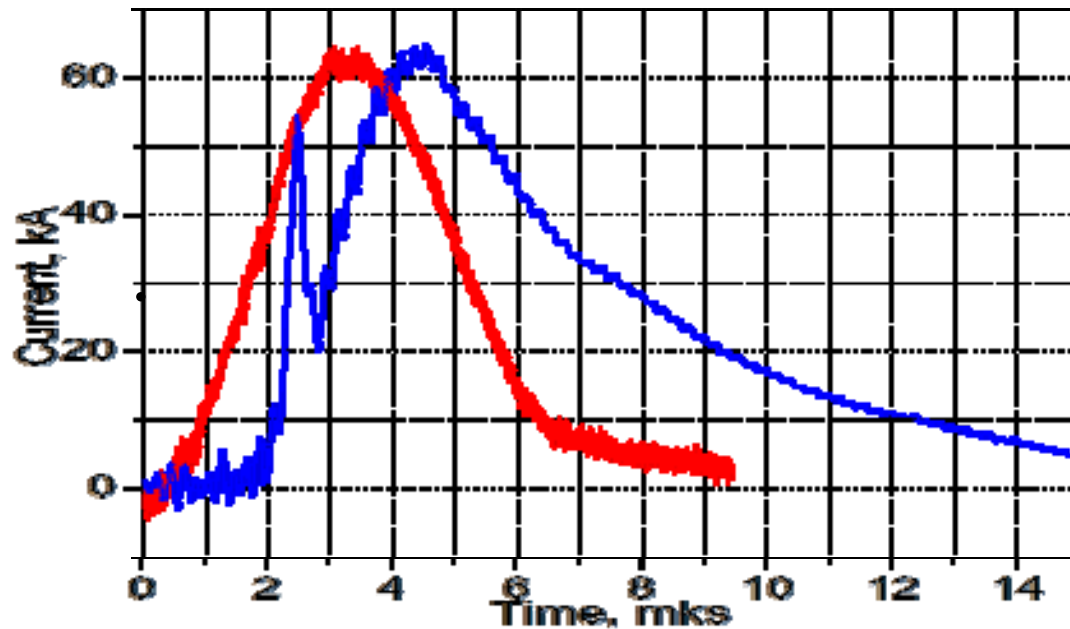


Features of discharge	current generator
Switch (Short pulse $T = 5 \mu s$) Discharge current pulse duration Max discharge current	Thyatron TDI1-150/25 $T = 5 \mu s$ at $C = 25 \mu F$ $I = 200 \text{ kA}$
Switch (Long pulse $T = 20 \mu s$) Discharge current pulse duration Max. discharge current	Thyatron TDI1-200k/25H $T = 20 \mu s$ at $C = 160 \mu F$ $I = 400 \text{ kA}$

The plasma discharge current 120 kA



Discharge

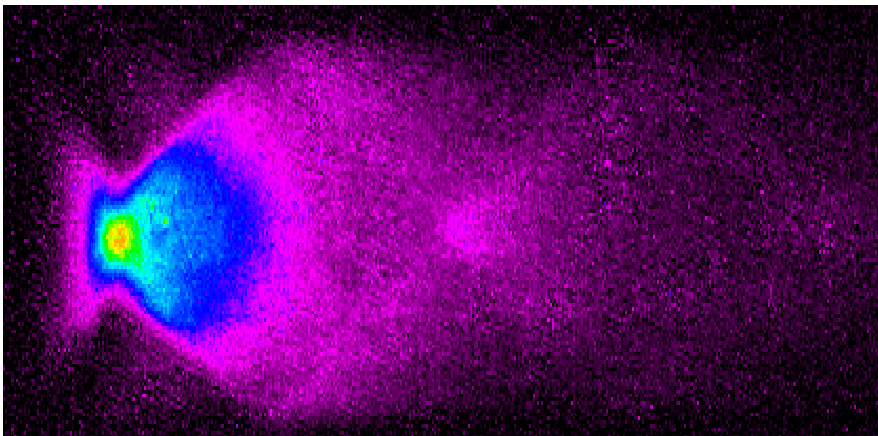


$$T \sim 17 \text{ eV (200 000 } ^\circ\text{K)}$$

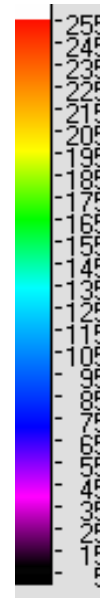
$$P \sim 400 \text{ bar}$$

$$n_e \sim 1-2 \cdot 10^{19} \text{ 1/cm}^3$$

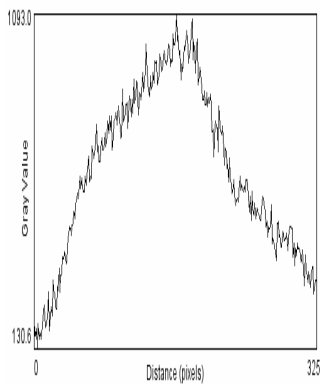
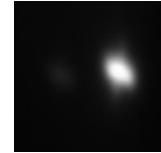
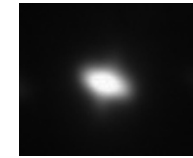
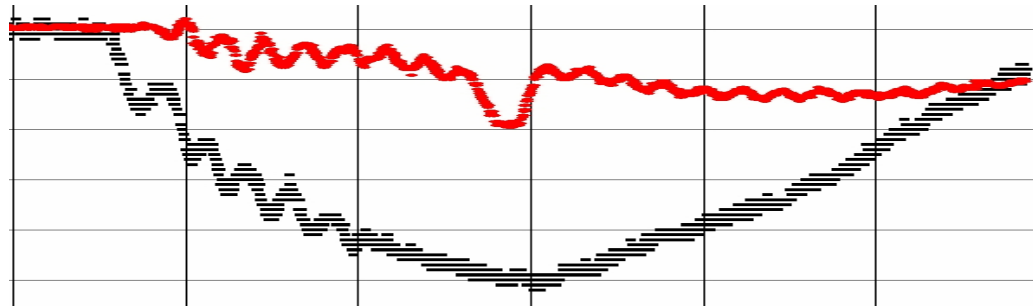
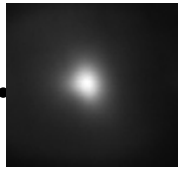
$$Z_i^{\text{max}} \sim 7$$



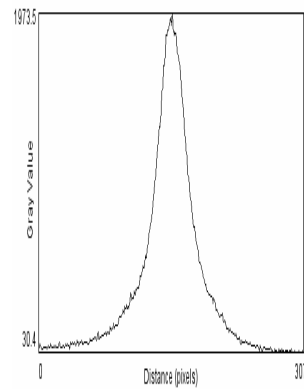
Streak



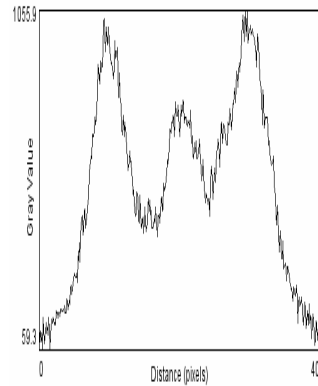
Light output from scintillator and the density distribution of ion Fe^{+26} at $T \mu s$ after beginning of discharge



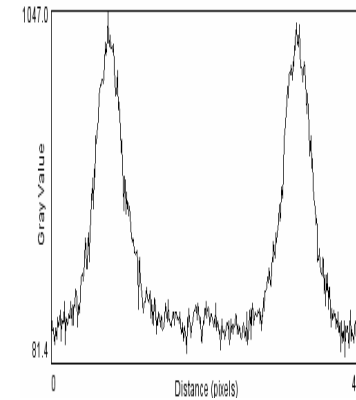
$T = 0 \mu s$



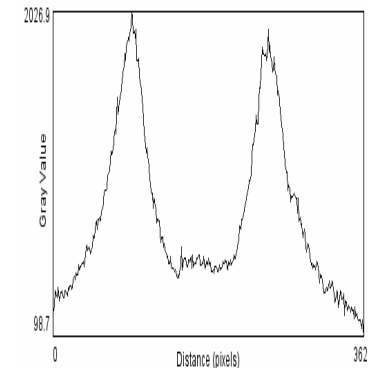
$T = 0.5 \mu s$



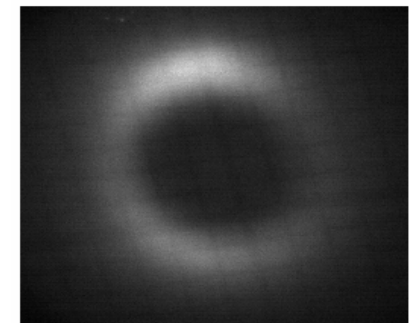
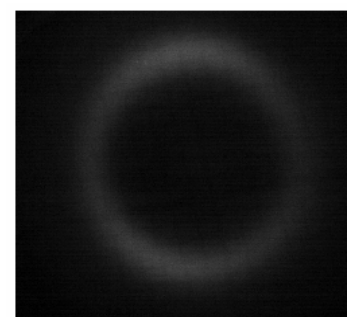
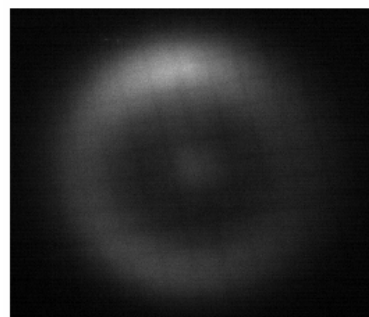
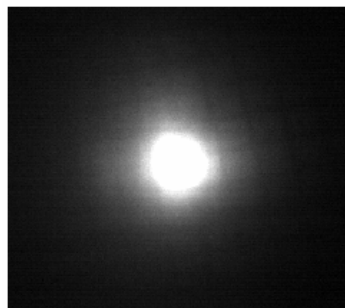
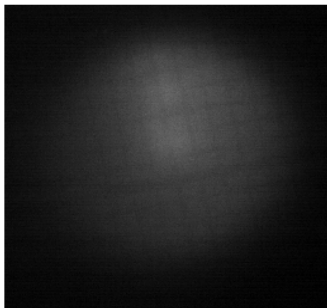
$T = 1.5 \mu s$



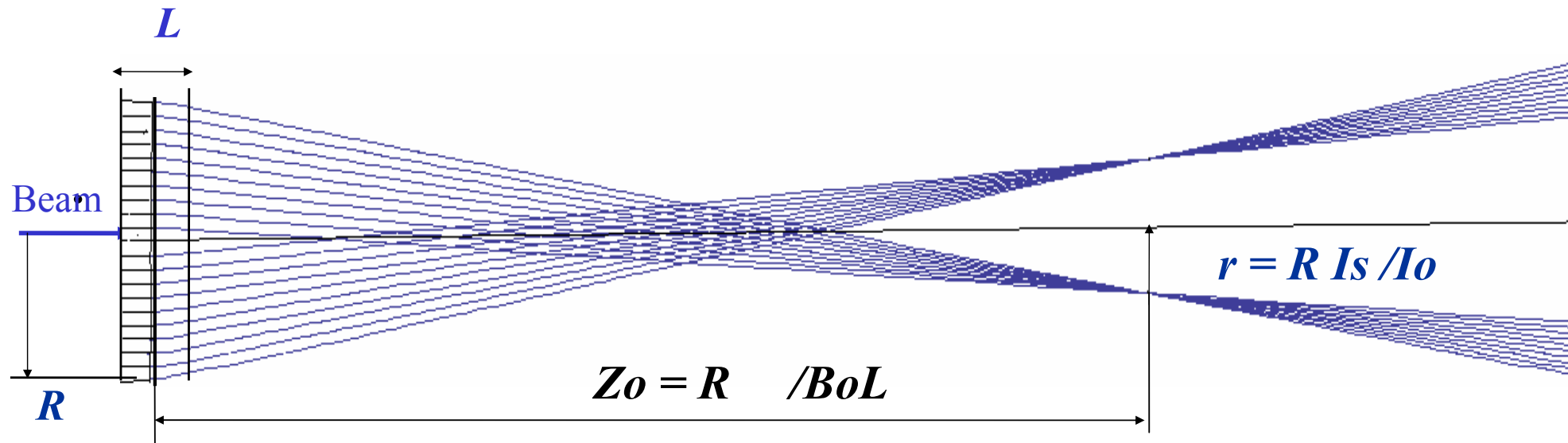
$T = 1.7 \mu s$



$T = 2.3 \mu s$

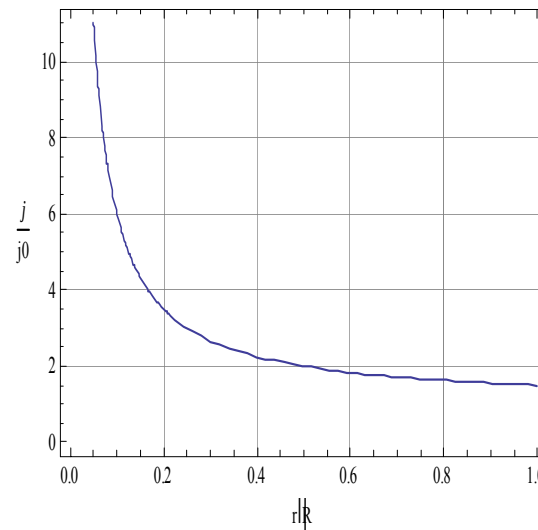


Forming ideal hollow beams (thin lens)

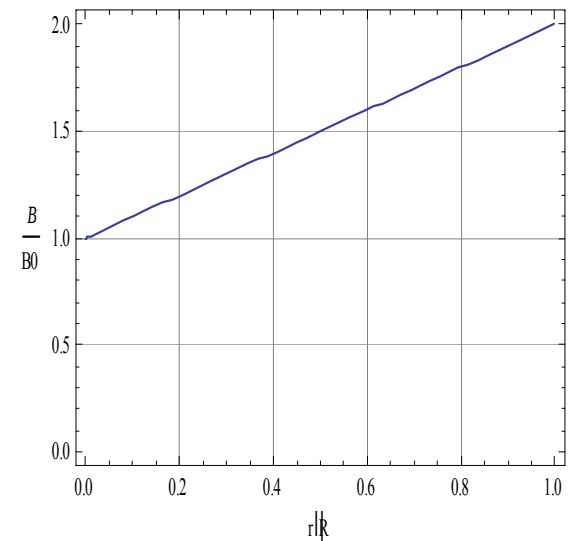


–the beam rigidity, B_0 –field strength due to uniform discharge current I_0 of radius R .

The paraxial beam is converted to a ring, if the distribution of the discharge current density is a superposition of the homogeneous and the singular distribution

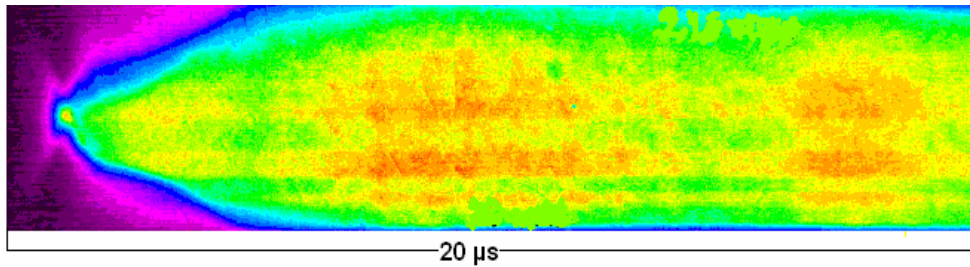


$$j = I_0 / \pi R^2 + I_s / 2\pi R r$$

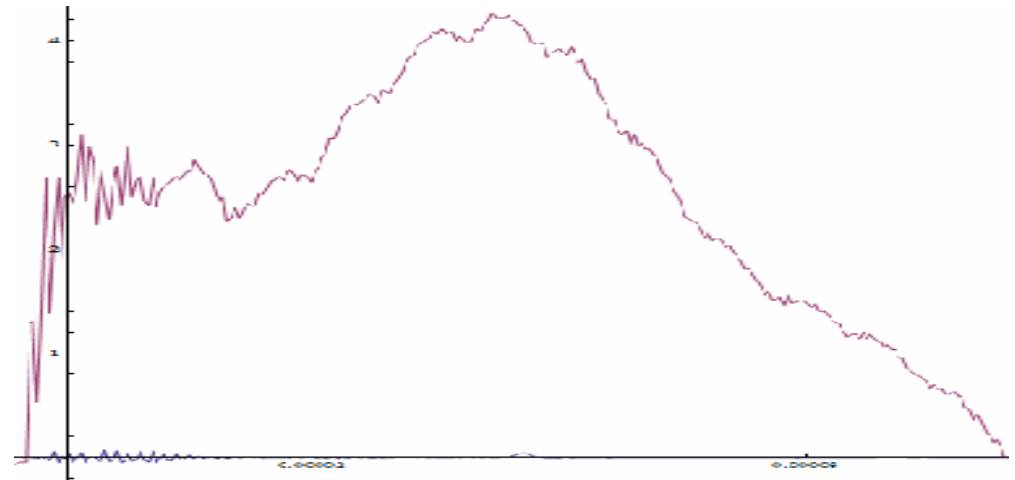


$$B = B_0 (r/R + I_s / I_0).$$

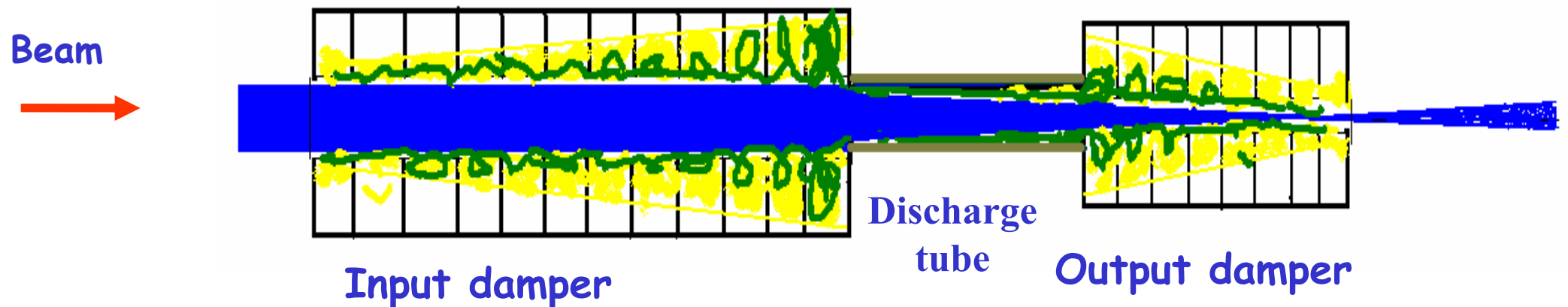
Features of the long discharge



The luminosity of the plasma



The discharge current



FORMATION OF THE HOMOGENEOUS BEAM - Theoretical



Equilibrium quasiBennett distribution of the discharge current

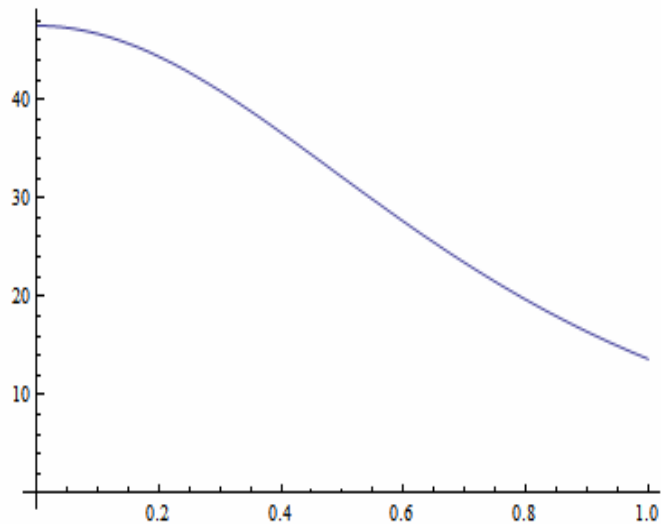
$$j = I(1 + \tilde{A}) / \pi R^2 (1 + \tilde{A} (r/R)^2)^2$$

For lens current $I = 80$ kA $\tilde{A} = 0.85$

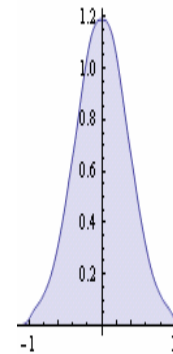
Distributions of ions of C^{+6} (300 MeV/n)

The distribution of the discharge current

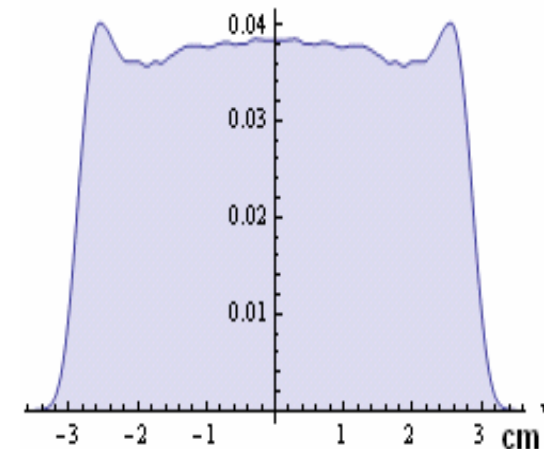
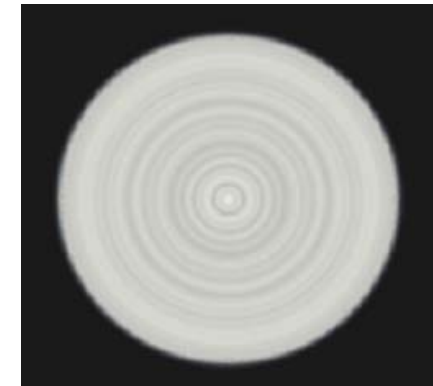
$j, \text{kA/cm}^2$



R, cm

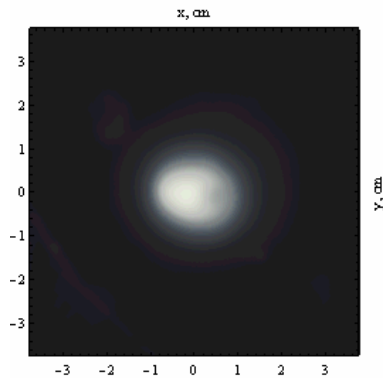


The initial Gaussian distribution.



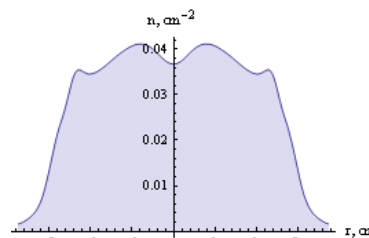
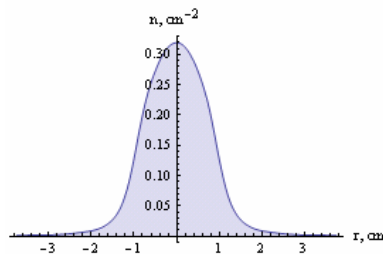
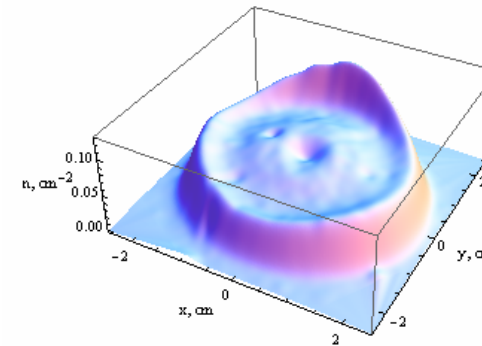
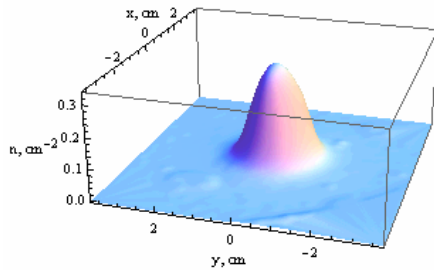
Transformed distributions at position of 150 cm behind the lens

FORMATION OF THE HOMOGENEOUS BEAM of C^{+6} (300 MeV/n) – Experimental

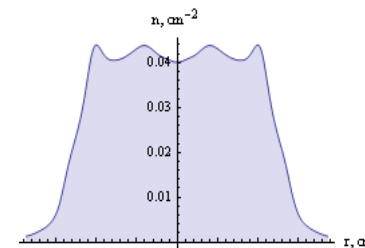


The initial distribution.

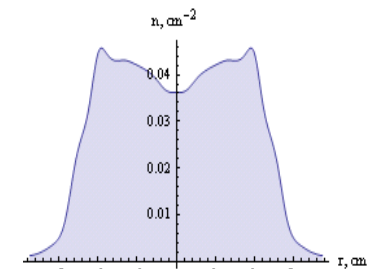
Transformed distributions at position of 110 cm behind the lens for different initial pressure of argon



2 Torr

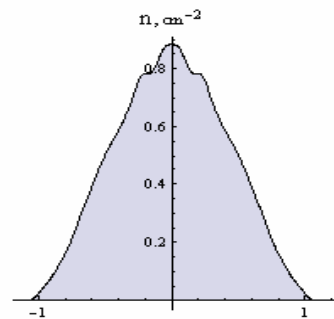
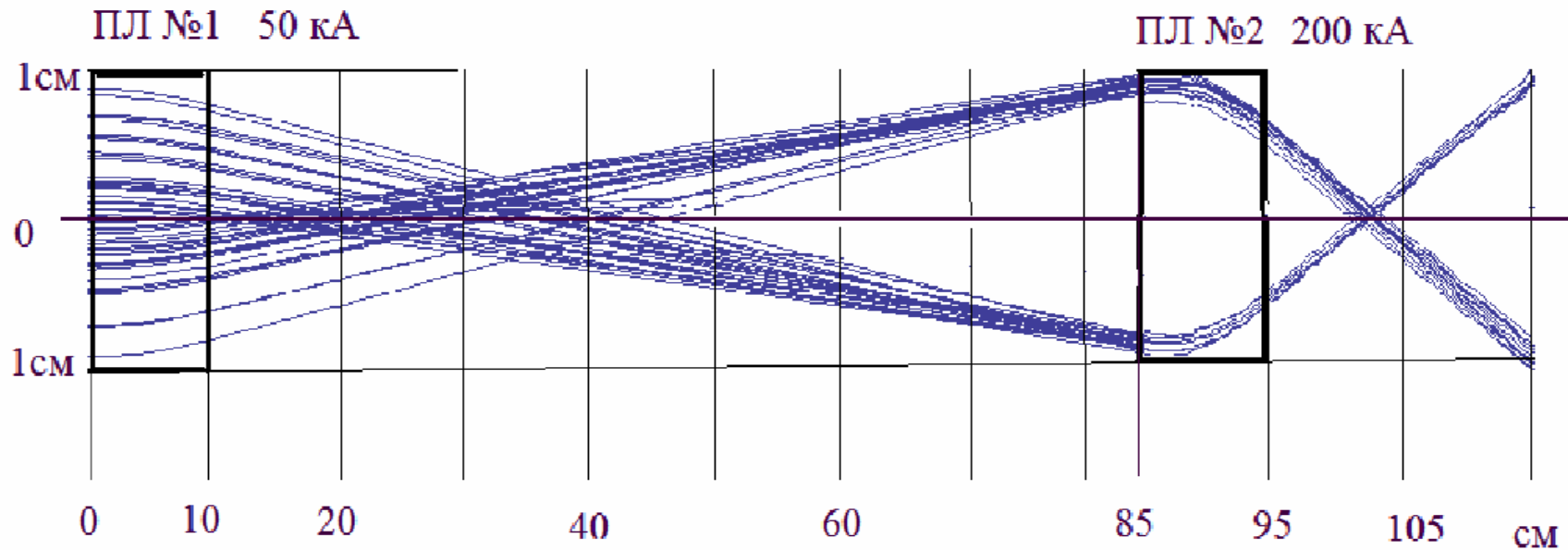


4 Torr

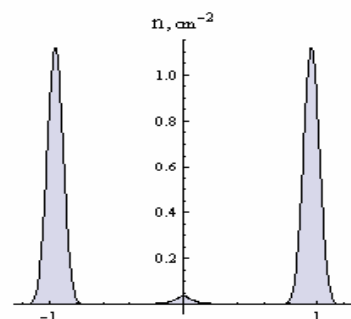


7 Torr

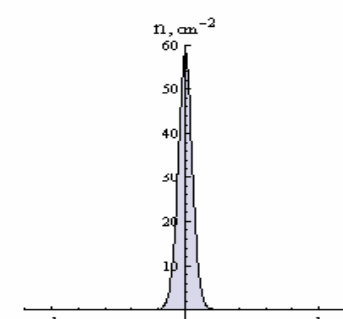
TWO-DIMENSIONAL TRANSFORMATION OF THE BEAM



Пучок на входе в ПЛ №1



Пучок на входе в ПЛ №2

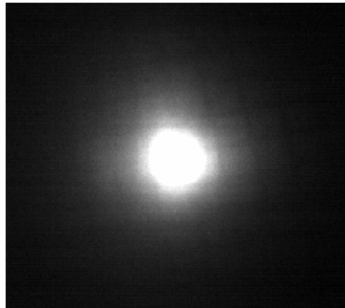


Пучок в фокусе

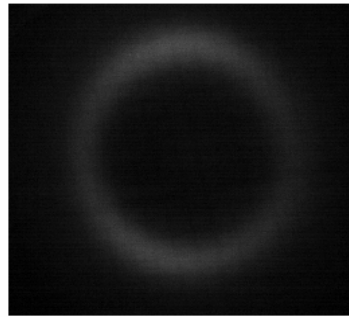
Light output from scintillator and the density distribution of ion Fe^{+26} at the time T after discharge switch



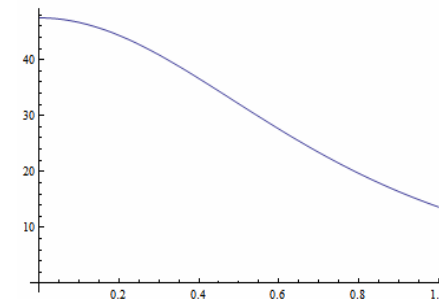
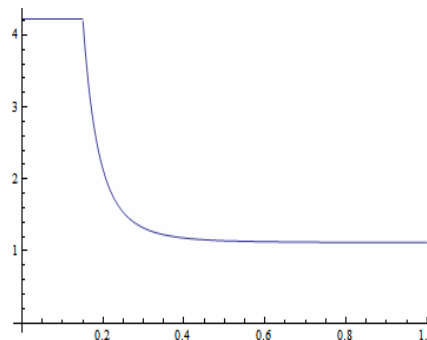
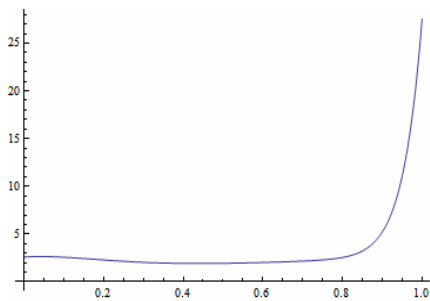
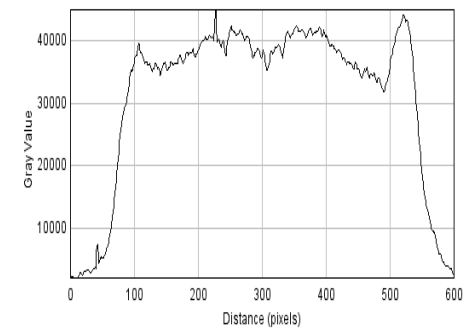
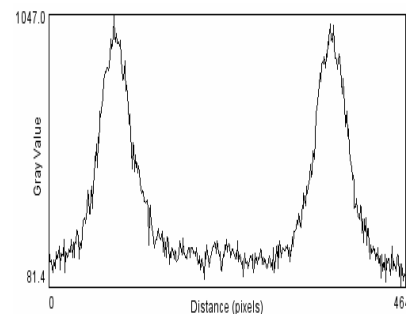
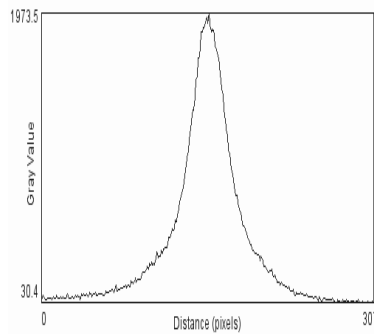
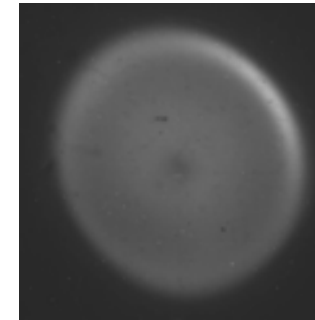
$T = 0.5 \mu s$



$T = 1.7 \mu s$



$T = 8.5 \mu s$

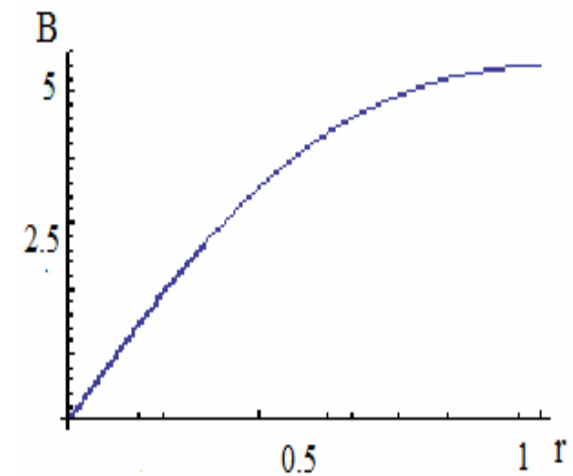
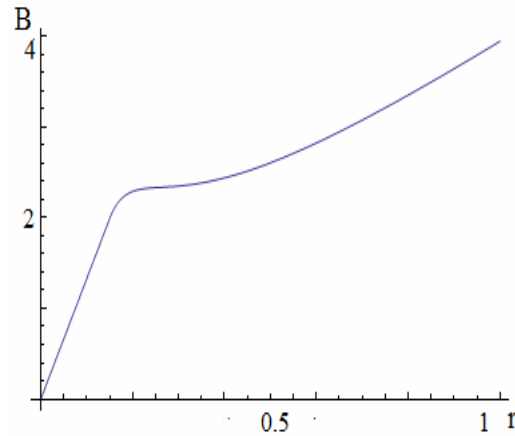
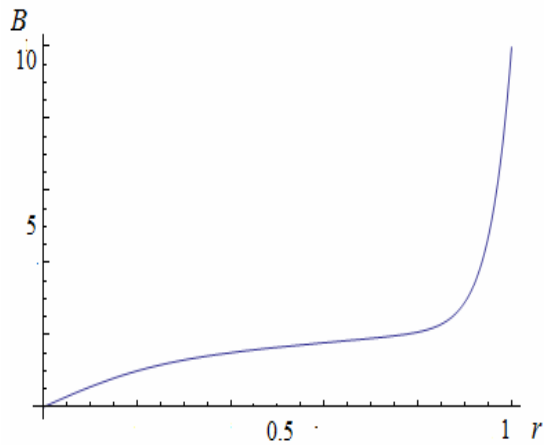


Experimental distributions current density in z -pinch of the plasma lens

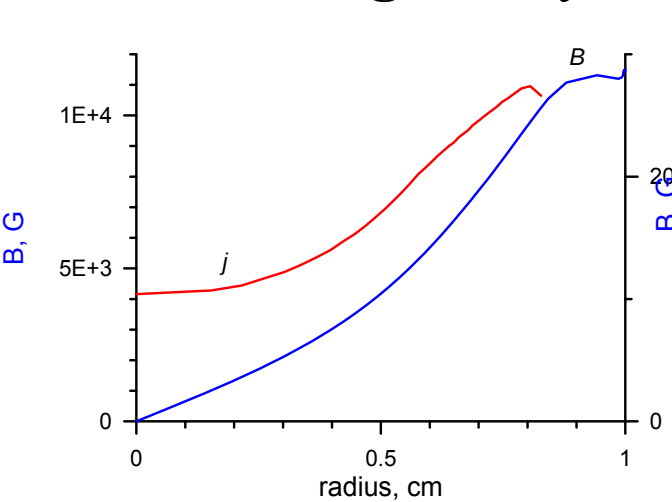
Z – pinch dynamic



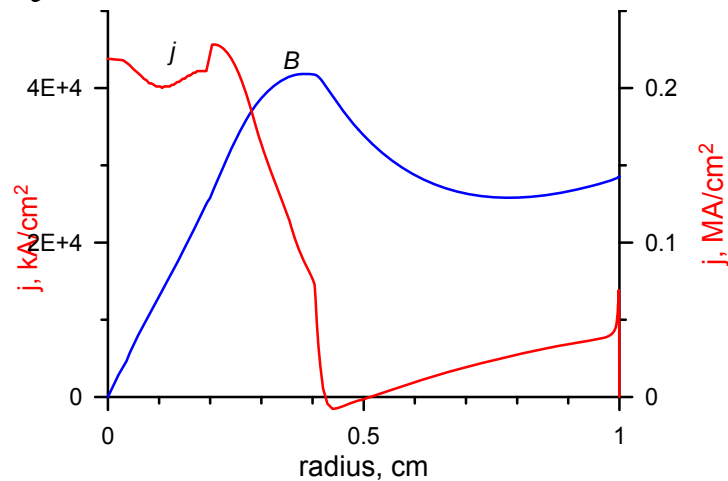
Experimental distributions self magnetic field in lens z-pinch



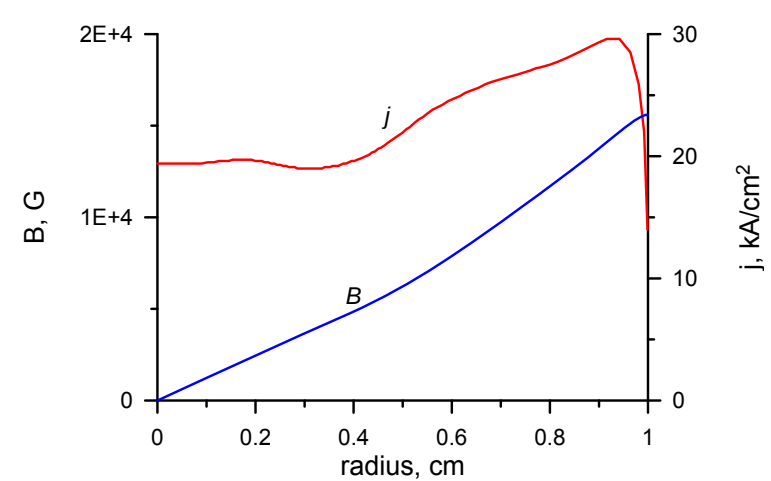
Magnetohydrodynamics simulation



$T = 0.5 \mu s$



$T = 1.6 \mu s$



$T = 8.5 \mu s$

CONCLUSION



The plasma lens can carry out not only sharp focusing of ions beam with considerable reduction of their sizes. At those stages of the plasma discharge at which the magnetic field is nonlinear, formation of other interesting configurations of beams is possible.

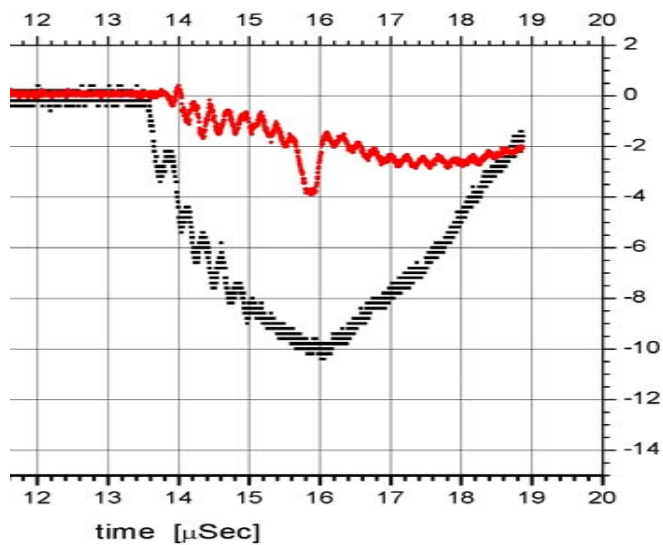
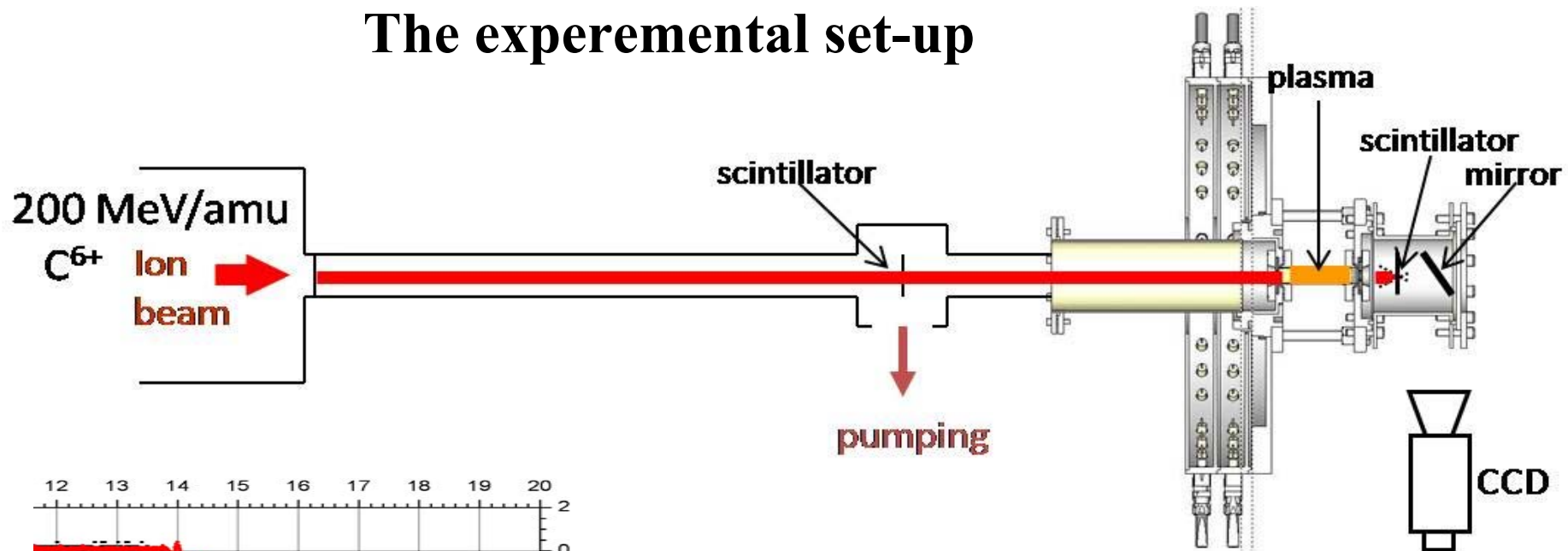
The plasma lens provides formation of hollow beams of ions in a wide range of parameters that allows to consider it as a possible variant of a terminal lens for realization of inertial thermonuclear synthesis.

The plasma lens can be used for transformation of beams with Gaussian distribution of particles density in a beams with homogeneous spatial distribution.

Application of the several plasma lenses which are in different stages of the plasma discharge, presumes to create some nontrivial spatial configurations of ions beams.

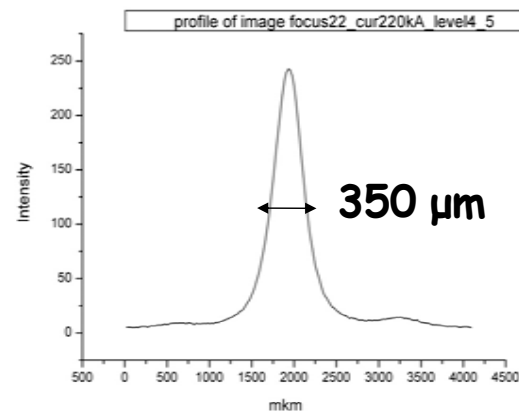
The plasma lens represents the universal tool for investigation of plasma discharges.

The experimental set-up

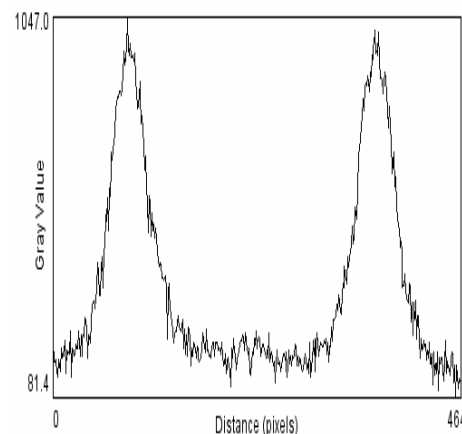
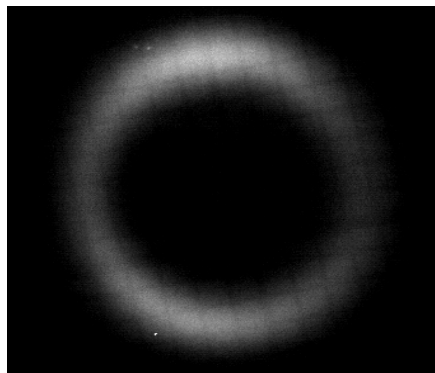


Time scanning of a discharge luminescence and a discharge and beam currents for short pulse mode.

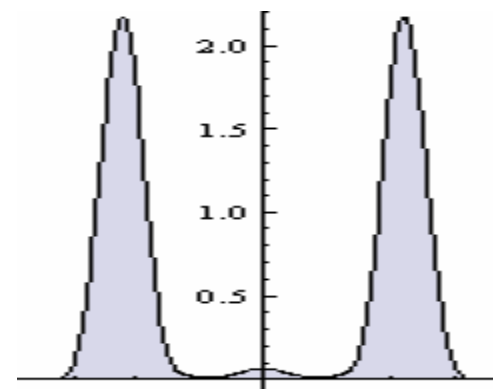
Results of focusing C^{+6} ions



The light output from a scintillator



The distribution of ion Fe^{+26} in time $1.7 \mu\text{s}$ behind discharge switch for distances 30 cm and discharge current 150 kA. The ring diameter - 9 mm .



The light output from a scintillator and the distribution of ion Fe^{+26} density calculated in the model approximation for the experimental condition.