

TWO DIMENSIONAL ASPECTS OF REGENERATIVE BBU IN TWO-PASS RECIRCULATING ACCELERATORS *

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Abstract

In this paper, I present the formula, describing a threshold of the regenerative multi-pass Beam Breakup (BBU) for a single dipole higher order mode with arbitrary polarization in a two-pass accelerator with a general-form, 4x4 recirculation matrix. Then, I present a mathematical relation between transfer matrices between cavities of the accelerating structure and recirculation matrices for each cavity, which must be satisfied in order to successfully suppress the BBU by reflection or rotation in several cavities.

INTRODUCTION

Regenerative multi-pass BBU arises from interaction of an intense recirculated beam with Higher-Order Modes (HOM) of the accelerating structures. In the JLab FEL Upgrade [1], the regenerative, multi-pass BBU was observed at a beam current of approximately 5 mA. Because BBU imposes an immediate threat to operations of the machine, suppression of regenerative BBU is a task of paramount importance.

As shown by R. E. Rand and T. I. Smith in [2], the threshold of regenerative BBU strongly depends on a form of the 4x4 recirculation matrix. In their paper, the authors investigated the possibility of suppressing the BBU by means of reflection or 90°-rotation of the beam displacement on the second pass. In the Summer of 2004, an optical reflector was installed in the recirculation pass of the JLab FEL Upgrade. The insertion consists of 5 skew-quadrupoles and reflects the beam displacement respective to the $y = x$ plane [3]. The first experience with the reflector showed an increase of the BBU threshold for offensive modes by a factor of 5 [4].

In this paper, a formula for the BBU threshold for a single dipole HOM with arbitrary polarization in a two-pass machine with a general-form, 4x4 recirculation matrix is derived. The paper also contains formulas describing evolution of the HOM voltage above and below the threshold. To suppress BBU by optical transformations in a machine consisting of several or many accelerating cavities, the recirculation matrix for each cavity has to be of a form that provides effective BBU suppression. To preserve the form of the recirculation matrix from cavity to cavity, transfer matrices between the cavities have to satisfy a condition that is presented at the end of this paper.

BBU THRESHOLD IN TWO-DIMENSIONAL SYSTEMS AND HOM VOLTAGE BEHAVIOR ABOVE AND BELOW THE THRESHOLD

Single-Mode Approximation

Simulation results and experimental data presented elsewhere [5] suggest that dipole TM HOMs with a quality factor of $10^5 - 10^6$ will limit the maximum recirculated current to the order of tens to hundreds of milliamps. Assuming a quality factor of the order of $10^5 - 10^6$ and a typical HOM frequency of the order of 2 GHz, one concludes that a typical bandwidth of dangerous HOMs is of the order of 1 kHz to 10 kHz.

Imperfections and asymmetrical features of cavities such as couplers violate two-dimensional symmetry of cavity geometry and split frequencies of degenerate dipole modes. A typical separation between two orthogonal polarizations of the same dipole mode in a superconducting RF cavity is of the order of several hundred kHz to several MHz. Besides, imperfections cause the frequency of HOMs to vary from cavity to cavity. The variation of the HOM frequency from cavity to cavity is also of the order of several hundred kHz to several MHz.

The probability of overlapping of HOM frequencies is small if the number of cavities is significantly smaller than the ratio of the spread of HOM frequencies to the typical bandwidth of the modes. For example, assuming an HOM bandwidth of 10 kHz and a frequency spread of a few MHz, one can consider modes separately if the number of cavities is of the order of a few tens.

Two-Dimensional Formula for the BBU Threshold

Steering and misalignment errors lead to a constant offset of the beam centroid relative to the cavity axis on the first pass. On the second pass, the transverse beam displacement consists of two terms: a constant offset due to steering/misalignment errors and an oscillatory term induced by the HOM voltage on the first pass. The constant offset induces the HOM voltage, which is always limited in amplitude. The oscillatory term, on the contrary, constitutes a feedback that can cause regenerative beam breakup and, in theory, an infinite growth of the HOM voltage.

In calculating the threshold of the regenerative BBU, we treat the voltage induced by the oscillatory term on the second pass as a perturbation and assume that the HOM voltage is a harmonic function of time. A variation of the energy stored in a dipole TM HOM produced by a point-like bunch, ΔU , is given, to the first order in the bunch charge,

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by [6]

$$\Delta U = -q \frac{V_a}{a} \cos(\phi) (x \cos(\alpha) + y \sin(\alpha)), \quad (1)$$

where V_a is the accelerating HOM voltage at the radius of the beam pipe a induced by previous bunches, ϕ is the phase of the point-like bunch with respect to the peak of the HOM electric field, q is the bunch charge, (x, y) is the bunch displacement vector, and α is the mode polarization angle in the xy coordinate system.

The energy deposited by the bunch in the HOM on the first and second passes can be written as

$$\Delta U_1 = -q \frac{V_a}{a} \cos(\phi) (x_1 \cos(\alpha) + y_1 \sin(\alpha)), \quad (2)$$

$$\Delta U_2 = -q \frac{V_a}{a} \cos(\phi + \omega T_r) (x_2 \cos(\alpha) + y_2 \sin(\alpha)), \quad (3)$$

where T_r is the recirculation time. It is assumed in (2-3) that the voltage amplitude V_a does not change significantly during recirculation. Exactly at threshold, V_a is constant. Coordinates of the recirculated bunch, x_2 and y_2 , can be expressed via the bunch coordinates before deflection and the HOM accelerating voltage on the first pass as

$$x_2 = m_{11}x_1 + m_{12}x'_1 + m_{13}y_1 + m_{14}y'_1 - \frac{qV_a}{\omega a p} \sin(\phi) (m_{12} \cos(\alpha) + m_{14} \sin(\alpha)), \quad (4)$$

$$y_2 = m_{31}x_1 + m_{32}x'_1 + m_{33}y_1 + m_{34}y'_1 - \frac{qV_a}{\omega a p} \sin(\phi) (m_{32} \cos(\alpha) + m_{34} \sin(\alpha)), \quad (5)$$

where p is the beam momentum.

The energy balance equation for the HOM stored energy is

$$\dot{U} = \dot{U}_{\text{beam}} - P_c = \langle \Delta U_1 + \Delta U_2 \rangle \cdot f_b - P_c. \quad (6)$$

The averaging is done respectively to the phase of the HOM, ϕ , taken at moments when bunches pass through the cavity on the first pass. Ohmic losses in the cavity can be expressed as [6]

$$P_c = \frac{V_a^2}{\left(\frac{\omega}{c}\right)^2 a^2 \left(\frac{R_d}{Q}\right) Q}. \quad (7)$$

Terms proportional to $\cos(\phi)$, $\sin(\phi)$, $\cos(\phi + \omega T_r)$, and $\sin(\phi + \omega T_r)$ yield zero after averaging if x_1, x'_1 and y_2, y'_2 in (4-5) are steering/misalignment errors. If the HOM frequency is not equal to a harmonic of the bunch repetition rate, terms proportional to $\cos(\phi) \cdot \sin(\phi)$ also yield zero and the average value of the $\sin^2(\phi)$ is equal to 1/2. Thus, Equation (6) can be rewritten as

$$\begin{aligned} \dot{U} &= I_b \frac{qV_a^2}{\omega p a^2} m_{12}^* \langle \sin(\phi) \cos(\phi + \omega T_r) \rangle \\ &\quad - \frac{V_a^2}{(\omega/c)^2 a^2 (R_d/Q) Q} \\ &= -I_b \frac{qV_a^2}{\omega p a^2} m_{12}^* \frac{\sin(\omega T_r)}{2} - \frac{V_a^2}{(\omega/c)^2 a^2 (R_d/Q) Q}, \end{aligned} \quad (8)$$

where

$$m_{12}^* = m_{12} \cos^2(\alpha) + (m_{14} + m_{32}) \sin(\alpha) \cos(\alpha) + m_{34} \sin^2(\alpha). \quad (9)$$

At threshold, \dot{U} is equal to zero. Thus, the threshold current is given by

$$I_{\text{th}} \frac{m_{12}^* q}{2p \omega} \sin(\omega T_r) + \frac{1}{(\omega/c)^2 (R_d/Q) Q} = 0, \quad (10)$$

which yields the threshold current as

$$I_{\text{th}} = - \frac{2pc}{q \frac{\omega}{c} \left(\frac{R_d}{Q}\right) Q m_{12}^* \sin(\omega T_r)}. \quad (11)$$

For positive values of the product $m_{12}^* \sin(\omega T_r)$, Equation (11) yields a negative threshold, implying a stable beam. However, the beam can be unstable at extremely high values of the beam current even if $m_{12}^* \sin(\omega T_r) > 0$. This discrepancy is caused by the assumption that the voltage induced by the beam on the second pass is a small perturbation to the HOM voltage, which fails at high beam intensities. In one-dimensional cases, similar dependence of the threshold on $m_{12} \sin(\omega T_r)$ was predicted analytically and observed in simulations in early works by G. A. Krafft, J. J. Bisognano, and S. Laubach [7] and later by G. H. Hoffstaetter and I. V. Bazarov [8].

Evolution of the HOM Voltage Below and Above the Threshold

Equation (8) can be rewritten as

$$\frac{dU}{U} = -dt \frac{\omega}{Q} \frac{I_{\text{th}} - I}{I_{\text{th}}}. \quad (12)$$

The solution of the last equation is

$$U = U_0 \exp\left(-t \frac{\omega}{Q} \frac{I_{\text{th}} - I}{I_{\text{th}}}\right). \quad (13)$$

The HOM voltage depends on time as

$$V = V_0 \exp\left(-t \frac{\omega}{2Q} \frac{I_{\text{th}} - I}{I_{\text{th}}}\right). \quad (14)$$

Equations (13) and (14) describe evolution of the stored energy and HOM voltage below and above the threshold. Using equations (13) and (14), one can introduce the effective quality factor Q_{eff} for the beam-HOM system, given by the equation

$$Q_{\text{eff}} = Q \frac{I_{\text{th}}}{I_{\text{th}} - I}. \quad (15)$$

Equation (15) can be also rewritten in terms of the decay time as

$$\tau_{\text{eff}} = \tau \frac{I_{\text{th}}}{I_{\text{th}} - I}. \quad (16)$$

Above the threshold, the decay time given by (16) changes its sign and becomes the growth time of the instability.

Equations (13-16) are also valid for negative values of the threshold current given by (11) for $m_{12}^* \sin(\omega T_r) > 0$ as long as the beam current is not too high and the assumptions made during the derivation of (11) are valid. For negative values of the BBU threshold given by (11), Equation (15) can be rewritten as

$$Q_{\text{eff}} = Q \frac{-|I_{\text{th}}|}{-|I_{\text{th}}| - I} = Q \frac{|I_{\text{th}}|}{|I_{\text{th}}| + I}. \quad (17)$$

The last two equations show that for $m_{12}^* \sin(\omega T_r) > 0$ the effective quality factor decreases with the beam current.

SUPPRESSION OF REGENERATIVE BBU BY OPTICAL TRANSFORMATIONS IN SEVERAL CAVITIES OF A TWO-PASS ACCELERATOR

Recirculation matrices for two cavities in a two pass-machine M_1 and M_2 are related to each other via transfer matrices between the cavities on the first and the second pass, T_1 and T_2 respectively, as

$$M_2 = T_2 M_1 T_1^{-1}. \quad (18)$$

If M_1 and M_2 are of a form that provides effective suppression of the BBU, Equation (18) yields a condition on the form of the transfer matrices T_1 and T_2 , which has to be satisfied for the particular form of the recirculation matrices M_1 and M_2 . Using the fact that the transfer matrices are symplectic, one can transform the Equation (18) into

$$M_2 = -T_2 M_1 S T_1^T S. \quad (19)$$

The last equation also allows one to express the condition (18) as a set of equations for 2x2 block sub-matrices of $M_{1,2}$, $T_{1,2}$, and S .

Using (18), it is trivial to show that, if M_1 and M_2 are fully coupled 4x4 matrices exchanging the x and y betatron planes and all HOMs are bound to the horizontal and vertical planes, all uncoupled transfer matrices $T_{1,2}$ satisfy Equation (18):

$$\begin{aligned} M_2 &= T_2 M_1 T_1^{-1} \\ &= \begin{bmatrix} X_2 & 0 \\ 0 & Y_2 \end{bmatrix} \begin{bmatrix} 0 & A \\ B & 0 \end{bmatrix} \begin{bmatrix} X_1^{-1} & 0 \\ 0 & Y_1^{-1} \end{bmatrix} \\ &= \begin{bmatrix} 0 & X_2 A Y_1^{-1} \\ Y_2 B X_1^{-1} & 0 \end{bmatrix}. \end{aligned} \quad (20)$$

Thus, an optical insertion interchanging the vertical and horizontal planes installed in the recirculation path can simultaneously suppress HOMs in different cavities if the HOMs are bound to the horizontal and vertical planes.

CONCLUSIONS

Equation (11), describing the BBU threshold for a single dipole HOM with arbitrary polarization and a general-form, 4x4 recirculation matrix in a two-pass accelerator, has been derived. The equation can be used to estimate the BBU threshold in the case of several high- Q HOMs if the separation between frequencies of the HOMs exceeds the bandwidth of the modes.

Reflection or rotation produced by the recirculation optics can significantly suppress or completely eliminate the regenerative BBU. This result was first predicted by R. E. Rand and T. I. Smith in [2]. To provide BBU suppression by reflection or rotation simultaneously in several cavities of the accelerating structure, the transfer matrices between the cavities have to satisfy equation (18).

If HOMs are bound to the horizontal and vertical planes, BBU can be simultaneously suppressed in several cavities by an optical transformation with a matrix of the form

$$\begin{bmatrix} 0 & A \\ B & 0 \end{bmatrix}, \quad (21)$$

where A and B are 2x2 block matrices. An optical insertion providing the $x - y$ motion exchange can consist of solenoids and/or skew quadrupoles and can be located in the recirculation path of a two-pass machine.

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