

Novel emittance measurement combining foil focusing and pepper-pot techniques

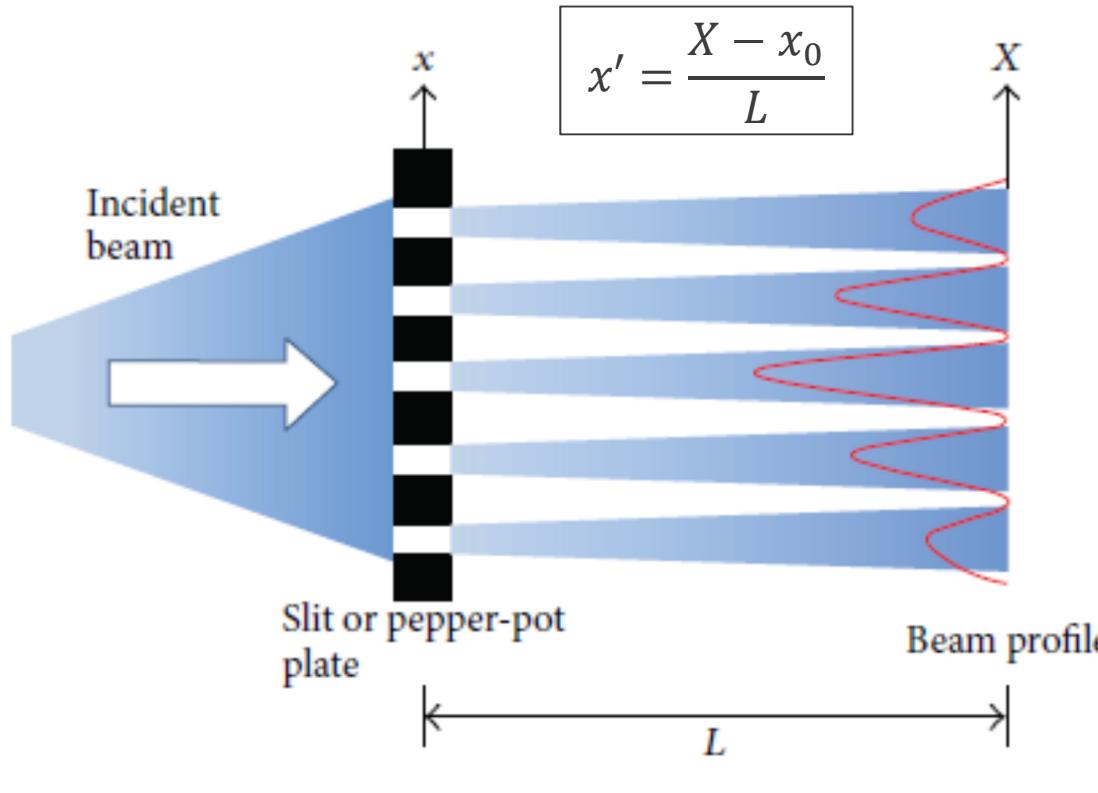


NAPAC 2019
Lansing, MI
September 5th, 2019

Kim Schultz¹
M. E. Schulze¹, G. T. Ortiz¹,
C. A. Carlson², D. Guerrero²

¹Los Alamos National Laboratory
²National Security Technologies, LLC

Pepper-pot emittance measurements

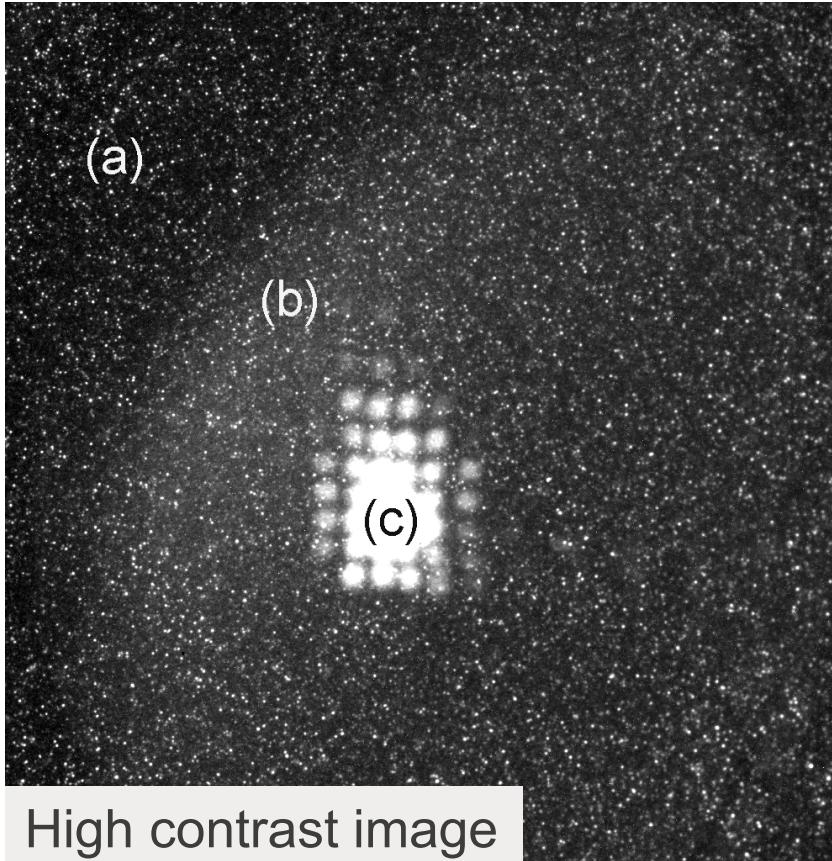


- A pepper-pot mask separates the beam into individual beamlets.
 - Typically a range-thick mask to reduce background counts.
- The beamlets are imaged after a drift.
 - Divergence is determined from the drifted centroid location.
 - RMS beamlet size is a measure of the beam's spread in divergence.
- Emittance is calculated from the beamlet distribution

$$\varepsilon_x = 4\gamma \sqrt{\langle x^2 \rangle \langle x'^2 \rangle - \langle xx' \rangle^2}$$

Use Multiple Coulomb scattering to reduce background at imaging screen instead of a range-thick mask.

- No vignetting introduced by a thin pepper-pot mask.
- Average scattering angle is 156 mrad for 254- μm thick Mo.
- Sufficient signal-background ratio to resolve spots.



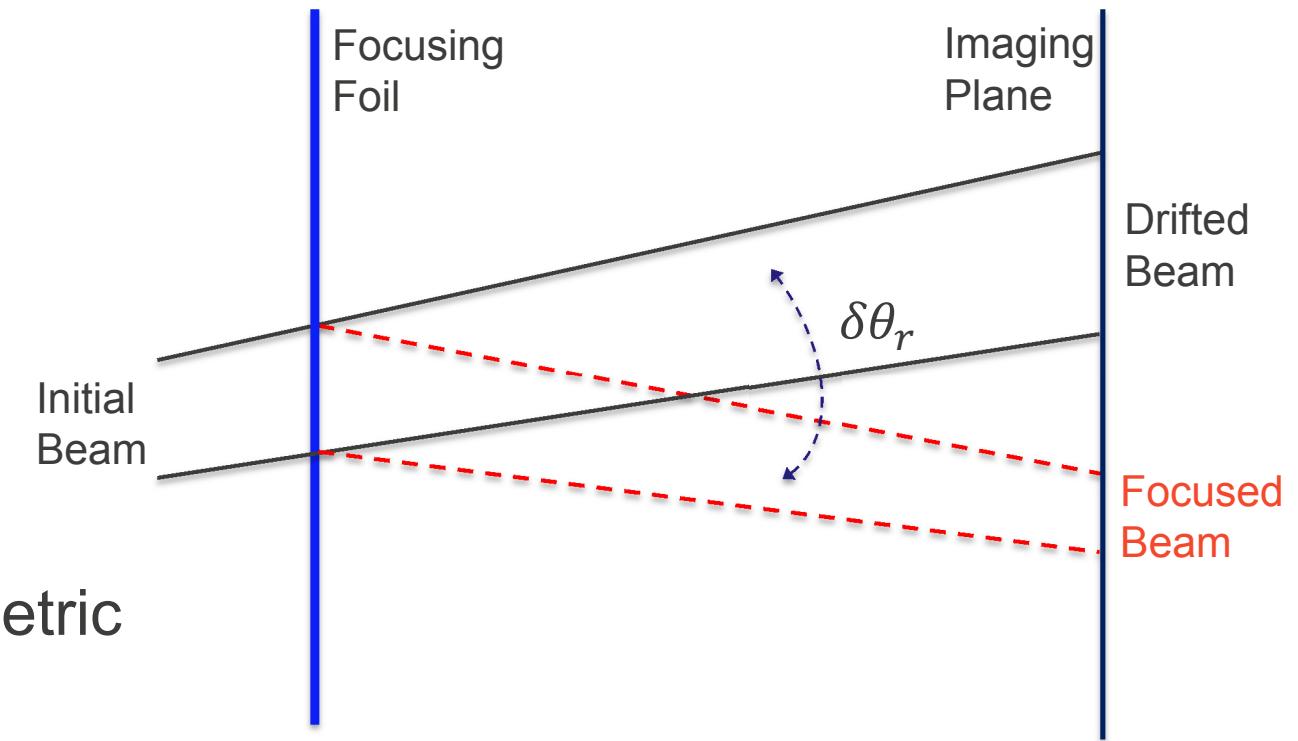
- (a) Background (no OTR target)
- (b) Background due to scattered beam on target
- (c) Foil focused pepper-pot image

Thin foils focus IREBs.

- A relativistic beam experiences a transverse ‘kick’ when passing through a thin, conducting, grounded foil.
- For an axisymmetric, uniform beam, this kick* is

$$\frac{\delta p_r}{p_z} = \delta\theta_r = -16 \frac{I_b}{I_A} \frac{b}{a} \sum_{n=1}^{\infty} \frac{J_1\left(\chi_{0n}a\right)J_1\left(\chi_{0n}r\right)}{\chi_{0n}^3 J_1(\chi_{0n})^2}$$

- I_b is the beam current (1.65 kA)
- $I_a = 17.05\beta\gamma$ is the Alven current in kA
- a is the beam radius
- b is the beam pipe radius (10 cm)
- χ_{0n} is the nth root of the Zero-order Bessel function $J_0(x)$
- Can be extended to any axisymmetric beam distribution.**



*Adler, R. J. *Part. Accel.* **12**, 39–44 (1982).

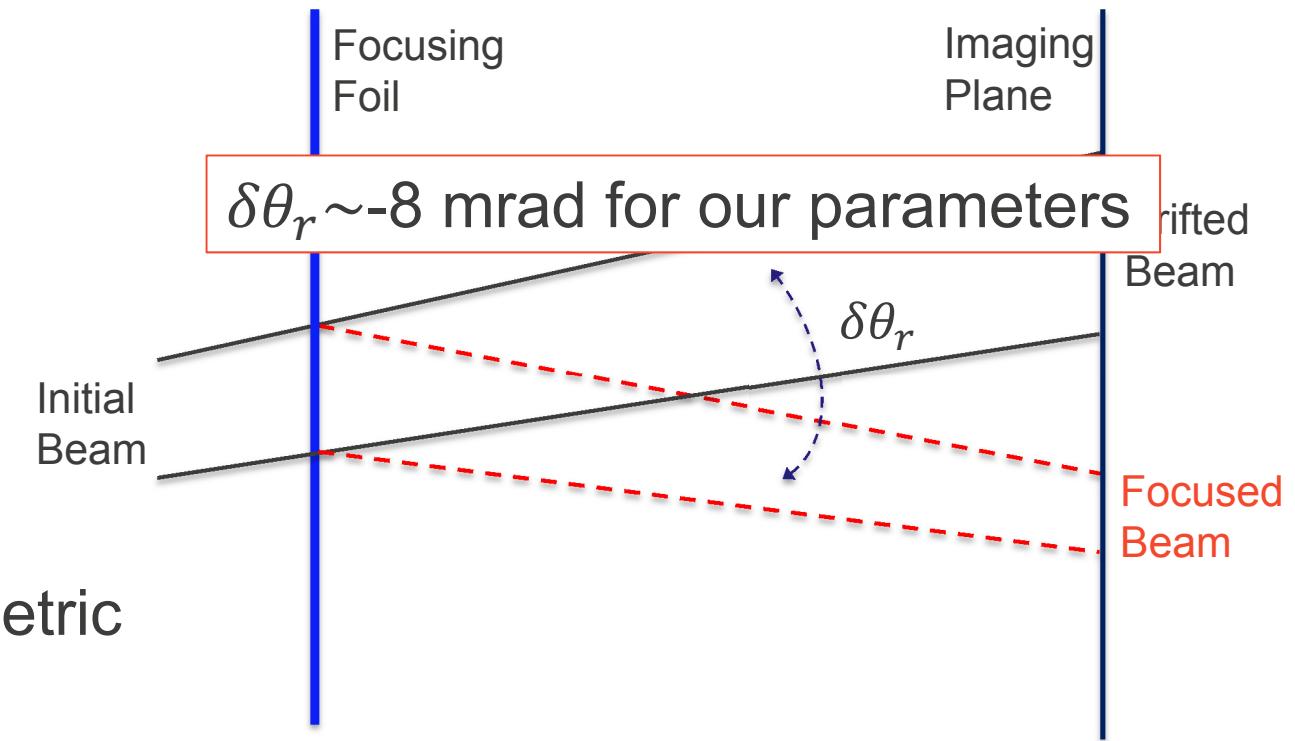
R. F. Fernsler, R. F. Hubbard, and S. P. Slinker, *J. Appl. Phys.*, **68, 5985–5994 (1990).

Thin foils focus IREBs.

- A relativistic beam experiences a transverse ‘kick’ when passing through a thin, conducting, grounded foil.
- For an axisymmetric, uniform beam, this kick* is

$$\frac{\delta p_r}{p_z} = \delta\theta_r = -16 \frac{I_b}{I_A} \frac{b}{a} \sum_{n=1}^{\infty} \frac{J_1\left(\chi_{0n}a\right)J_1\left(\chi_{0n}r\right)}{\chi_{0n}^3 J_1(\chi_{0n})^2}$$

- I_b is the beam current (1.65 kA)
- $I_a = 17.05\beta\gamma$ is the Alven current in kA
- a is the beam radius
- b is the beam pipe radius (10 cm)
- χ_{0n} is the nth root of the Zero-order Bessel function $J_0(x)$
- Can be extended to any axisymmetric beam distribution.**

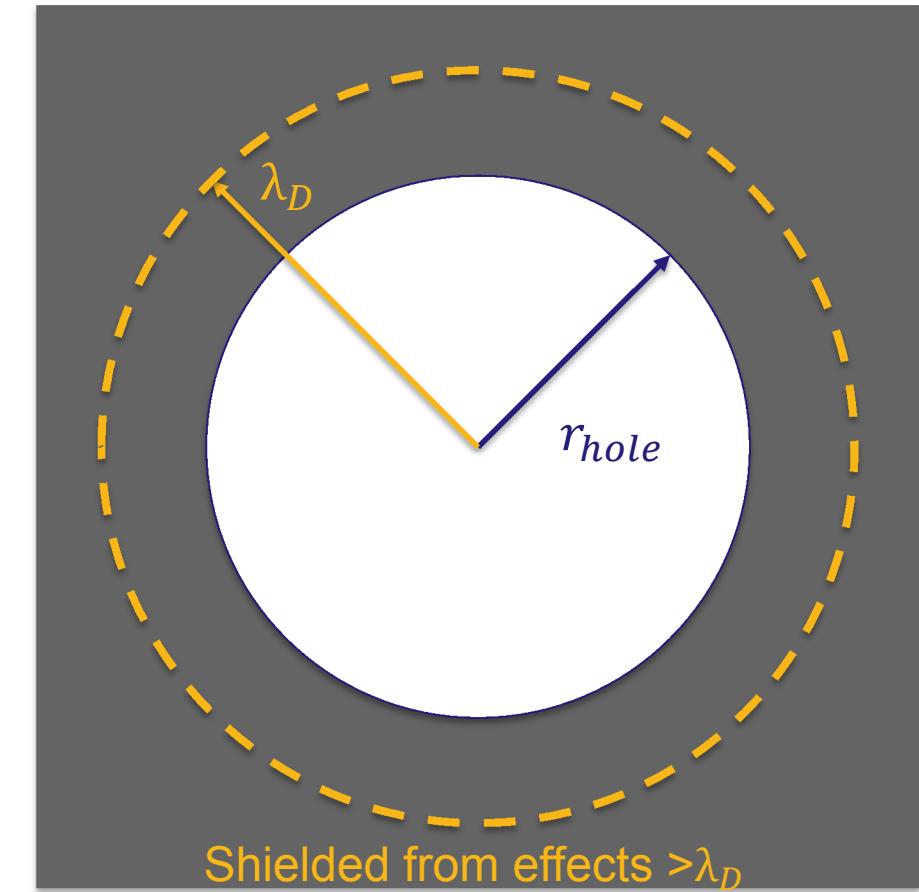


*Adler, R. J. *Part. Accel.* **12**, 39–44 (1982).

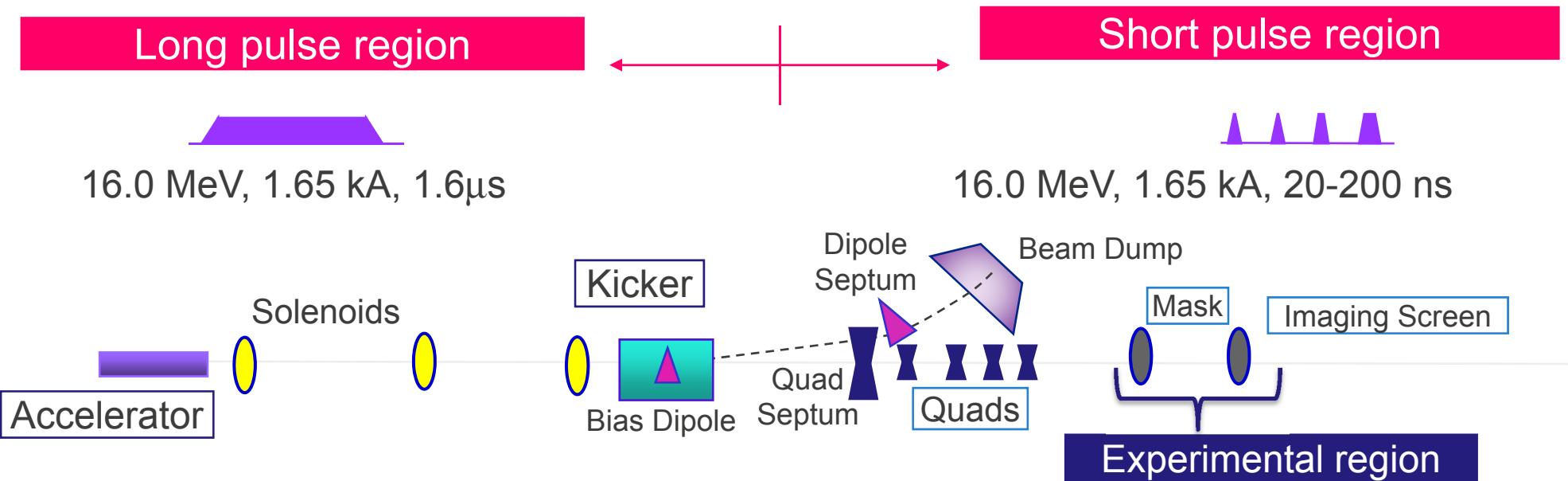
R. F. Fernsler, R. F. Hubbard, and S. P. Slinker, *J. Appl. Phys.*, **68, 5985–5994 (1990).

Debye length and foil focusing

- Beam passing through foil (not holes) experiences foil focusing and scattering deflections.
 - Scattering is dominant force.
 - $156 \gg -8$ mrad
- Beamlets **only** experience foil focusing as long as the hole radius is smaller than the Debye length.
 - Hole radius .75 mm
 - Debye length $\lambda_D = \frac{\varepsilon_n}{4} \sqrt{\frac{I_A}{I_b}} \sim 2\text{mm}$ for DARHT Axis-II



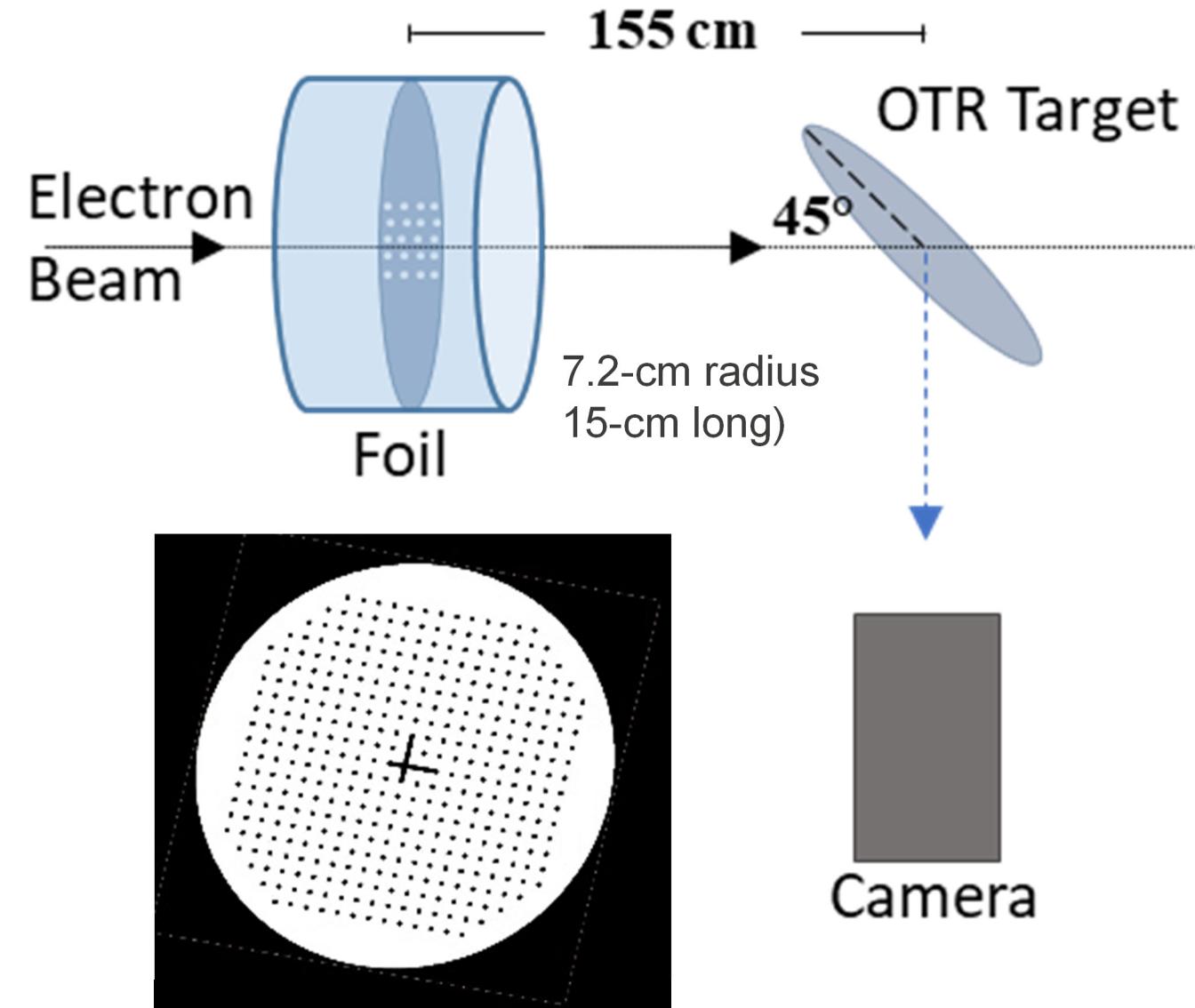
Dual-Axis Radiographic Hydrotest Facility: DARHT Axis-II



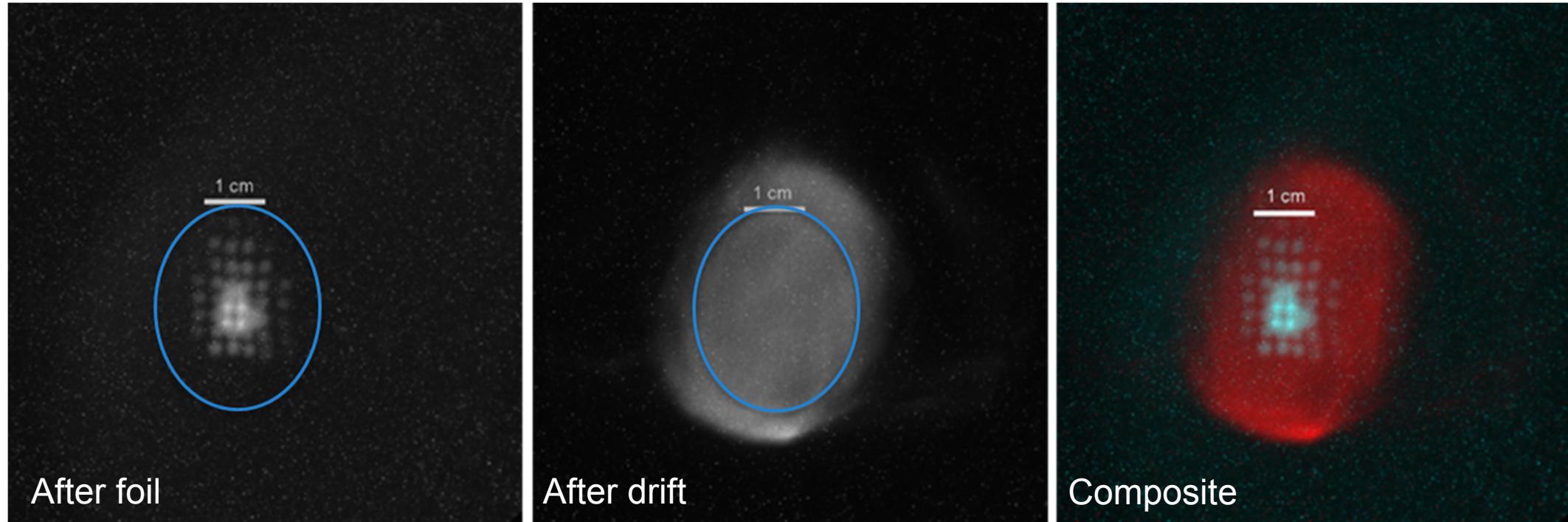
- Completed 2008
- Linear induction accelerator
- 16-MeV electron kinetic energy
- Programmable pulse widths

Experimental setup

- DARHT Axis-II Beam
 - 16 MeV
 - 1.65 kA
 - 40-ns pulse kicked from the 1.6- μ s flattop
- Mo pepper-pot/focusing mask
 - 254 μ m thick
 - 1.5 mm diameter holes
 - 5 mm center-to-center spacing
- NB: Mask fiducial length (17 mm) is much greater than the Debye length (2 mm)

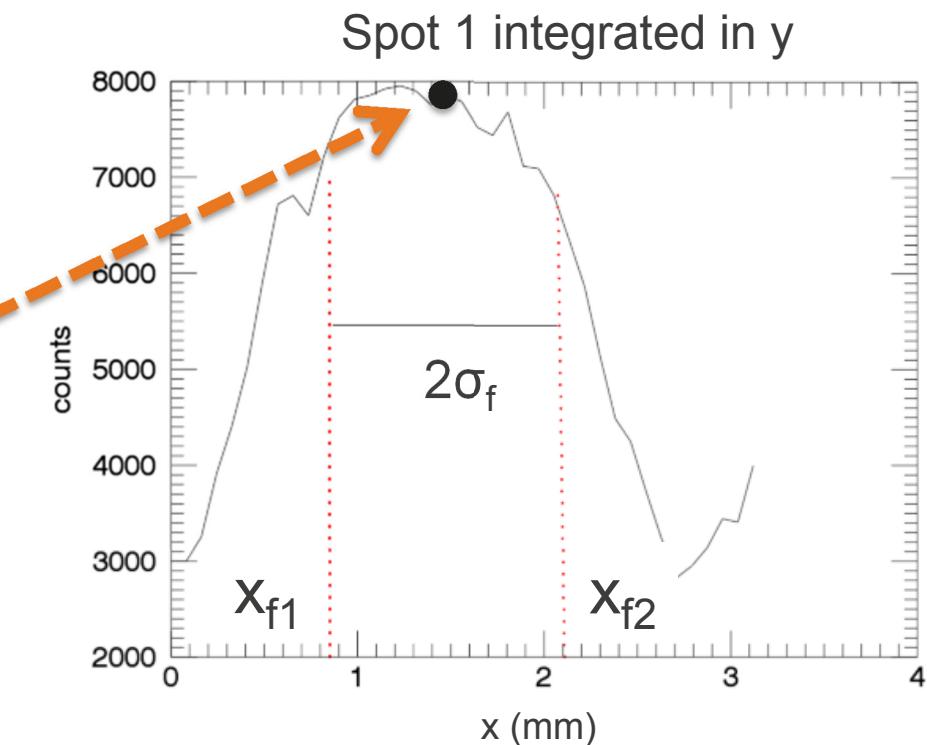
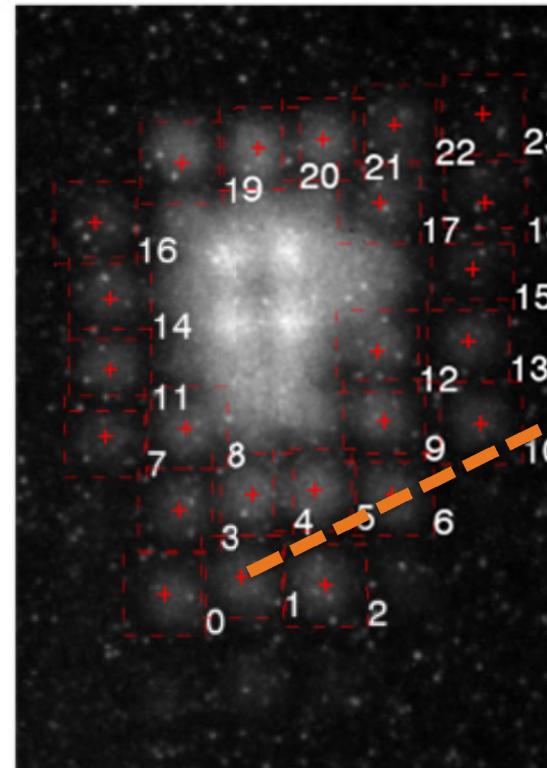
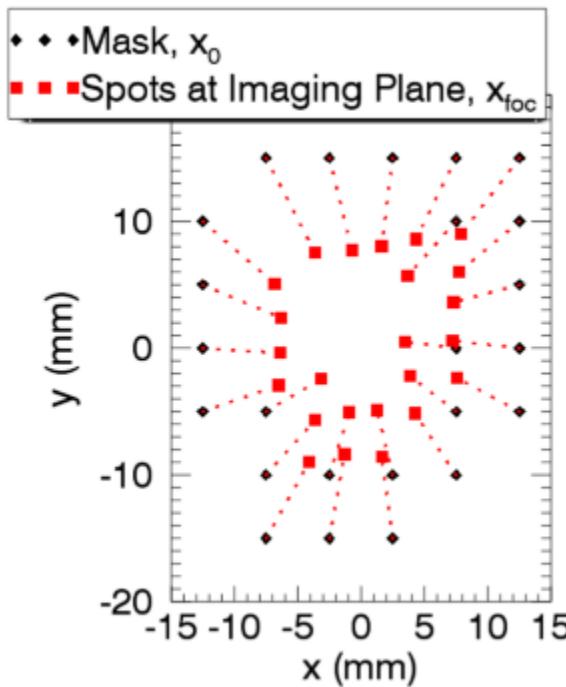


Evidence of Foil Focusing on DARHT-II.



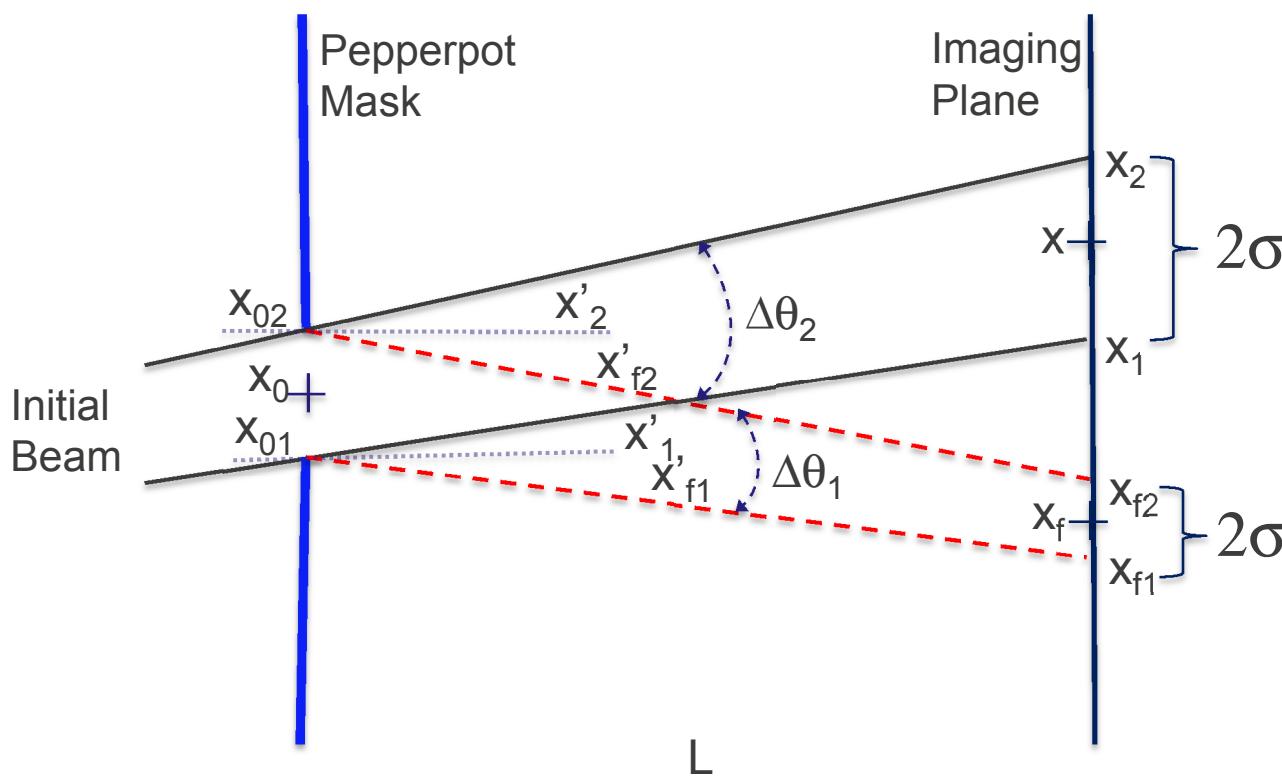
Beam radius at (mm)	X (mm)	Y (mm)	
Mask 6x7.5 holes	13	16	
After foil	6	8	Converging beam!
After drift	20	25	Diverging beam!

Spot Analysis



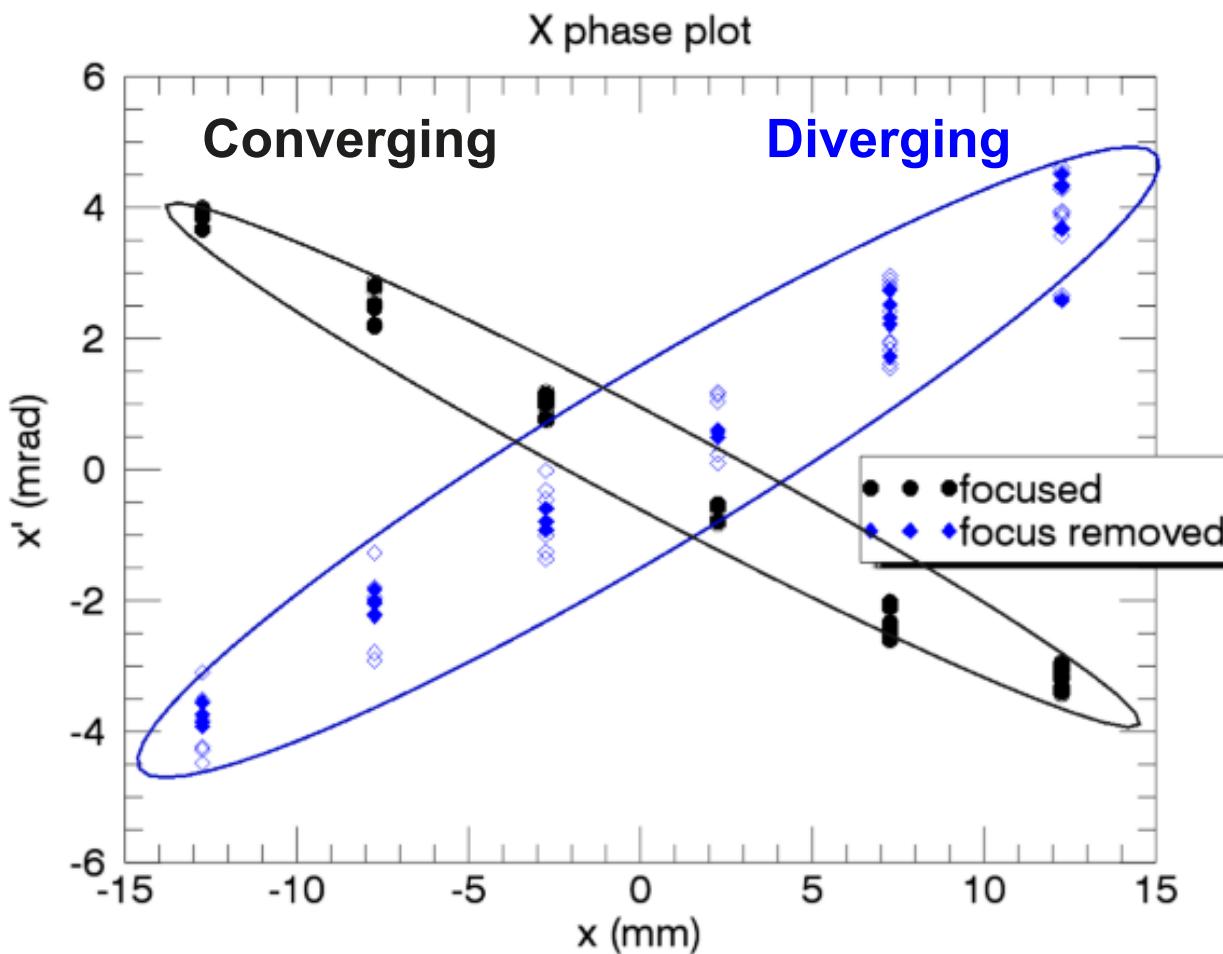
- Subtract constant background with 5-pix median smoothing
- Spot size: $\sigma_f^2 = \langle x_{beamlet}^2 \rangle = \frac{\sum counts(x-\langle x \rangle)^2}{\sum counts}$
- RMS edges of beamlet: $x_{f1}, x_{f2} = x_f \pm \sigma_f$
- $f \rightarrow$ focused beamlet pattern

Accounting for foil focusing of beamlets



- Beamlet envelopes are also foil focused.
 - $\Delta\theta_i = \theta_{r0} x_{0i}/r_{0i}$, kick in x-direction
 - $x'_i = x'_{fi} - \Delta\theta_i$, corrected divergence
 - $x_i = x_{0i} + x'_i L$, corrected beamlet edges
- Beam sizes before and after accounting for focusing
 - σ_f is 0.78 mm
 - σ is 1.3 mm

Accounting for foil focusing changes beam convergence angle and the measured emittance



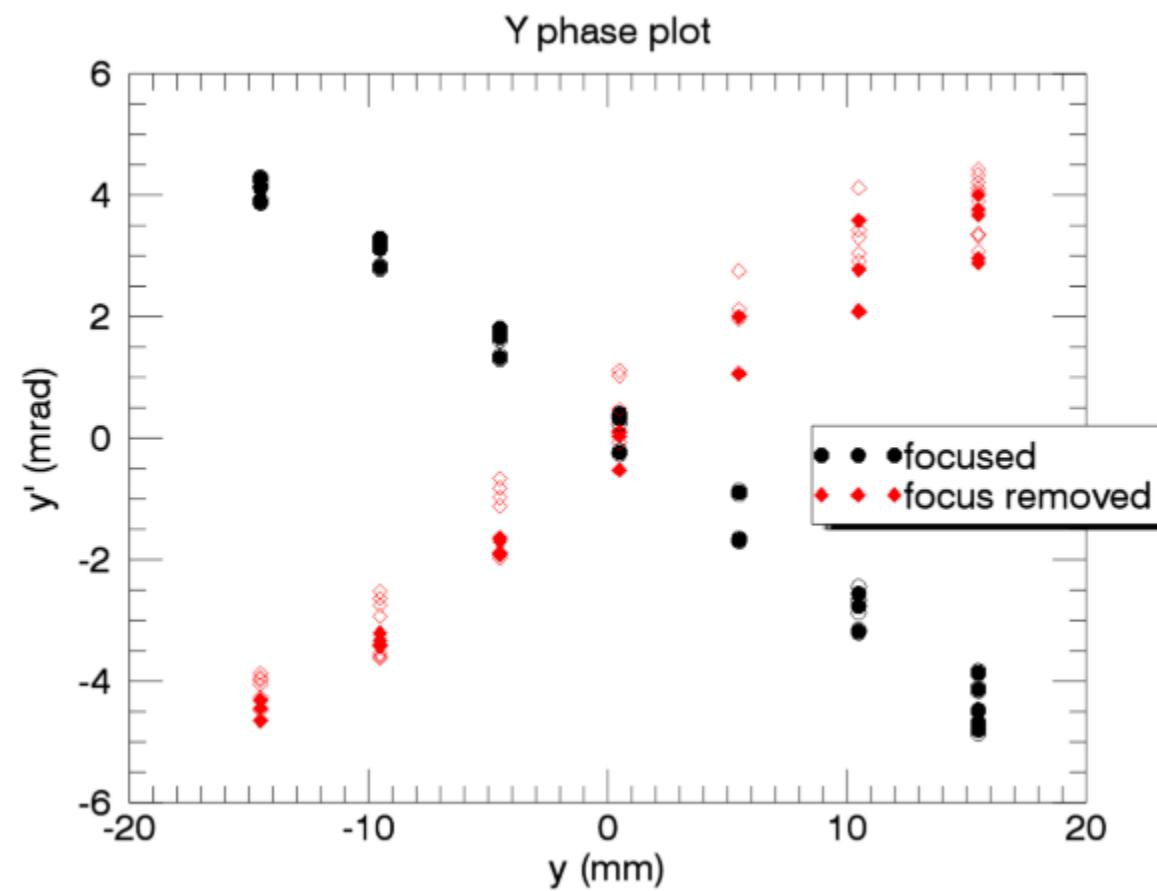
in $\pi\text{-mm-mrad}$	ε_x	ε_y
Focused	147	306
Corrected	226	415

- Analysis of another shot with a different quadrupole tune gives $\varepsilon_x = 300$ and $\varepsilon_{x,corrected} = 423$ $\pi\text{-mm-mrad}$.
- Solid-centroid locations.
- Open-edge locations.

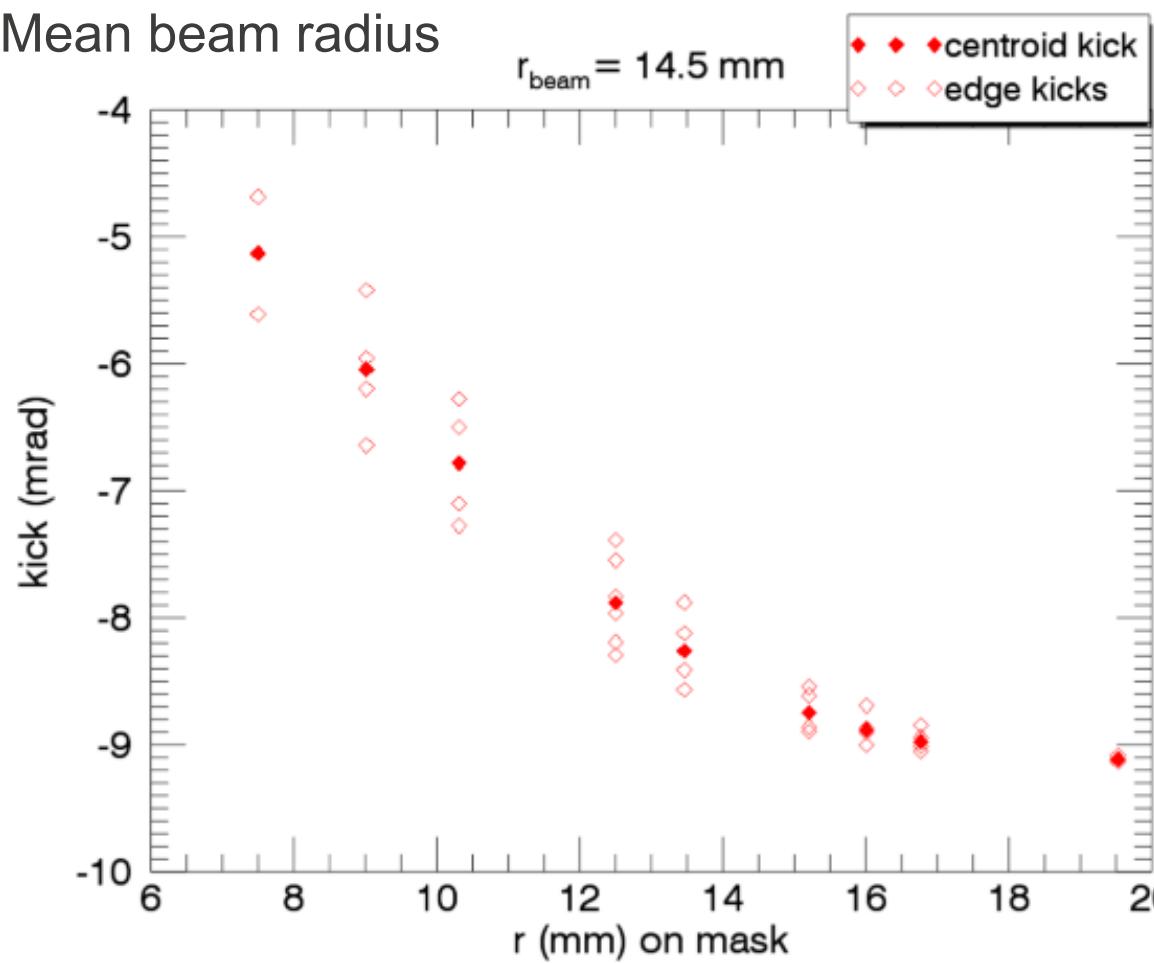
Conclusion

- Foil focusing must be taken into account when using a thin pepper-pot foil to measure a relativistic beam's emittance.
 - More important for beams that are close to pipe radius
- Future direction
 - Optimize mask
 - Automatic spot detection
 - Elliptical beams
 - Off-center beams
 - Background
 - Error analysis
 - Space charge

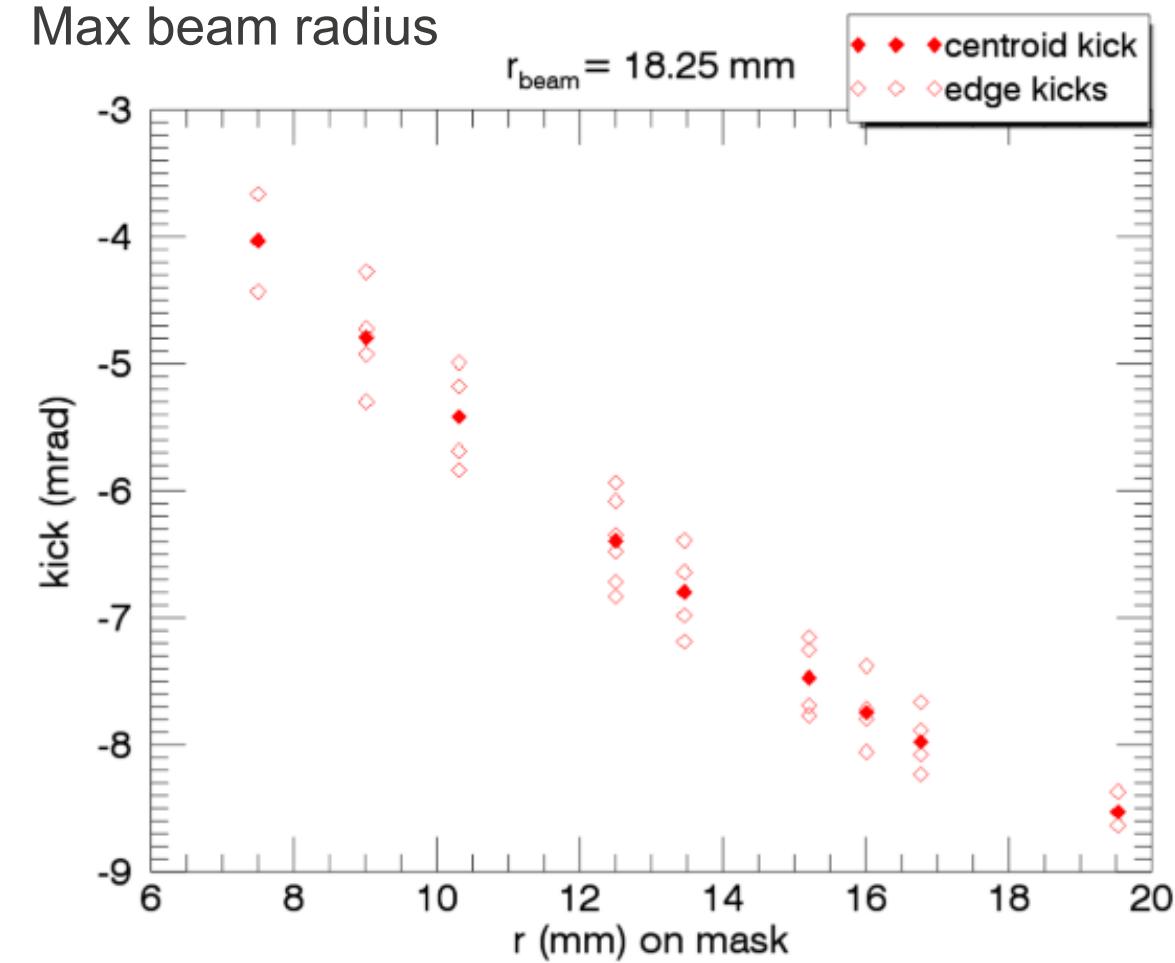
Thank you.
kimberlys@lanl.gov



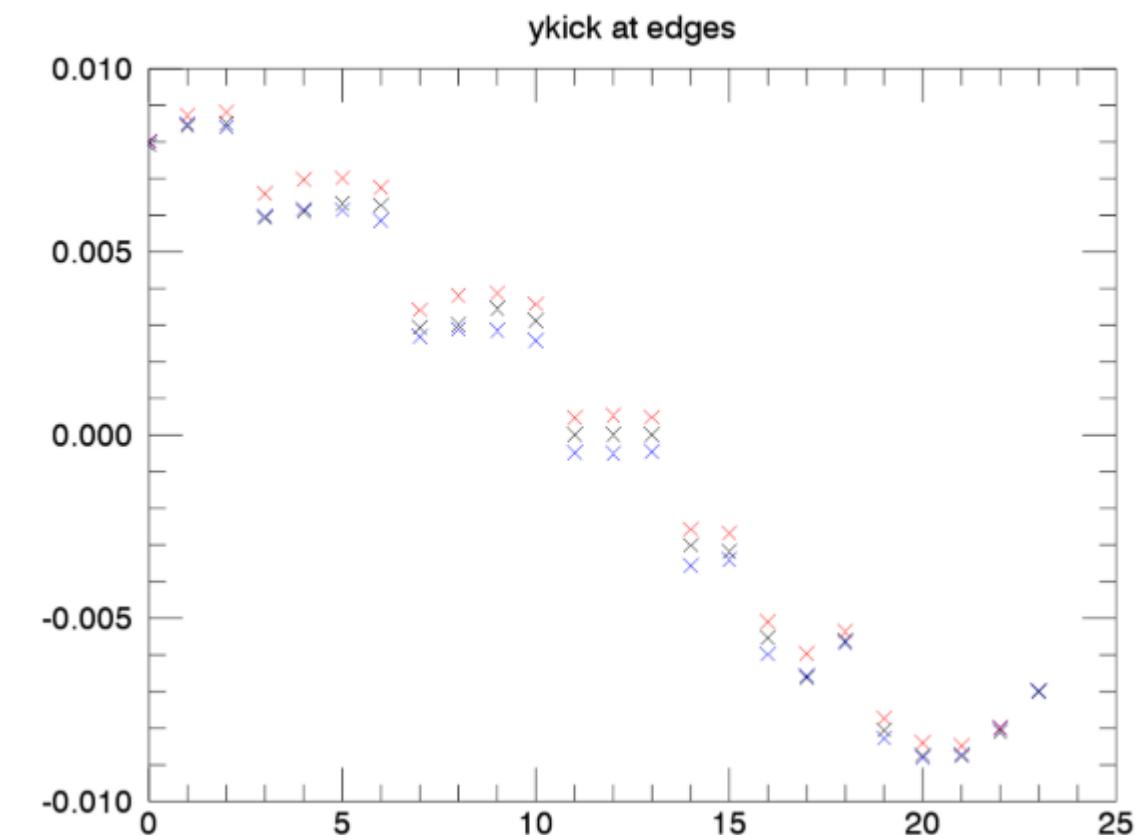
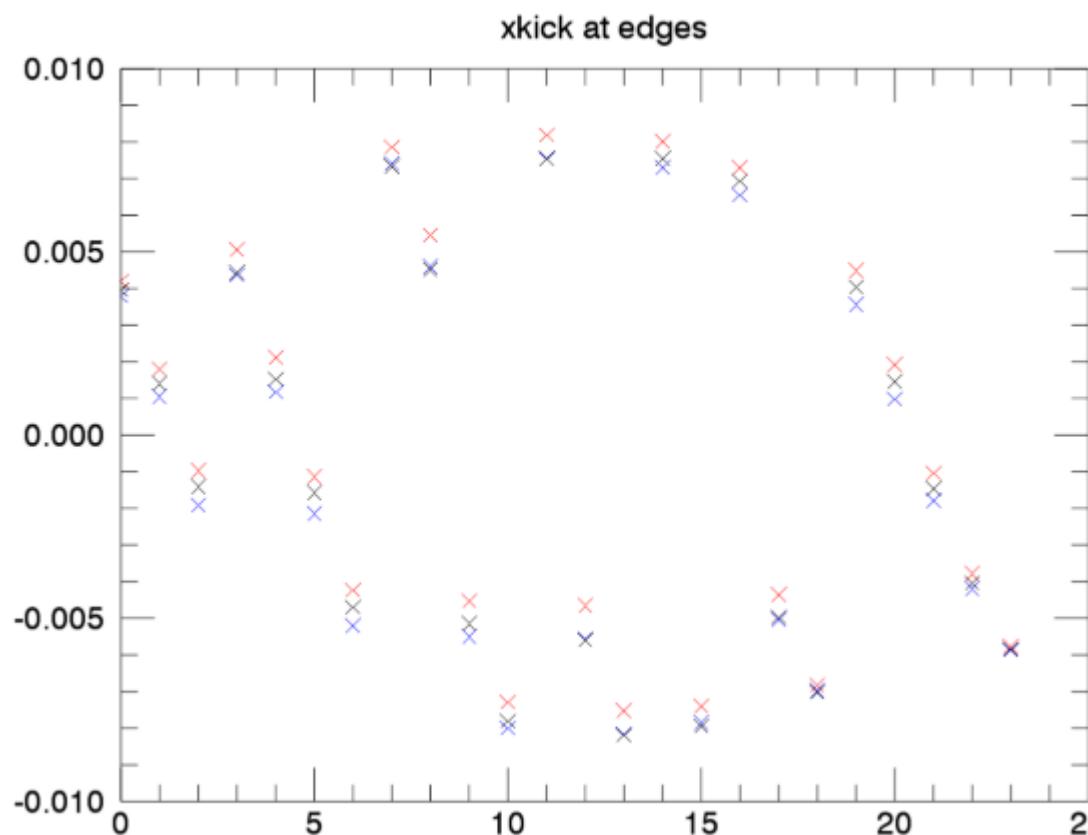
Mean beam radius

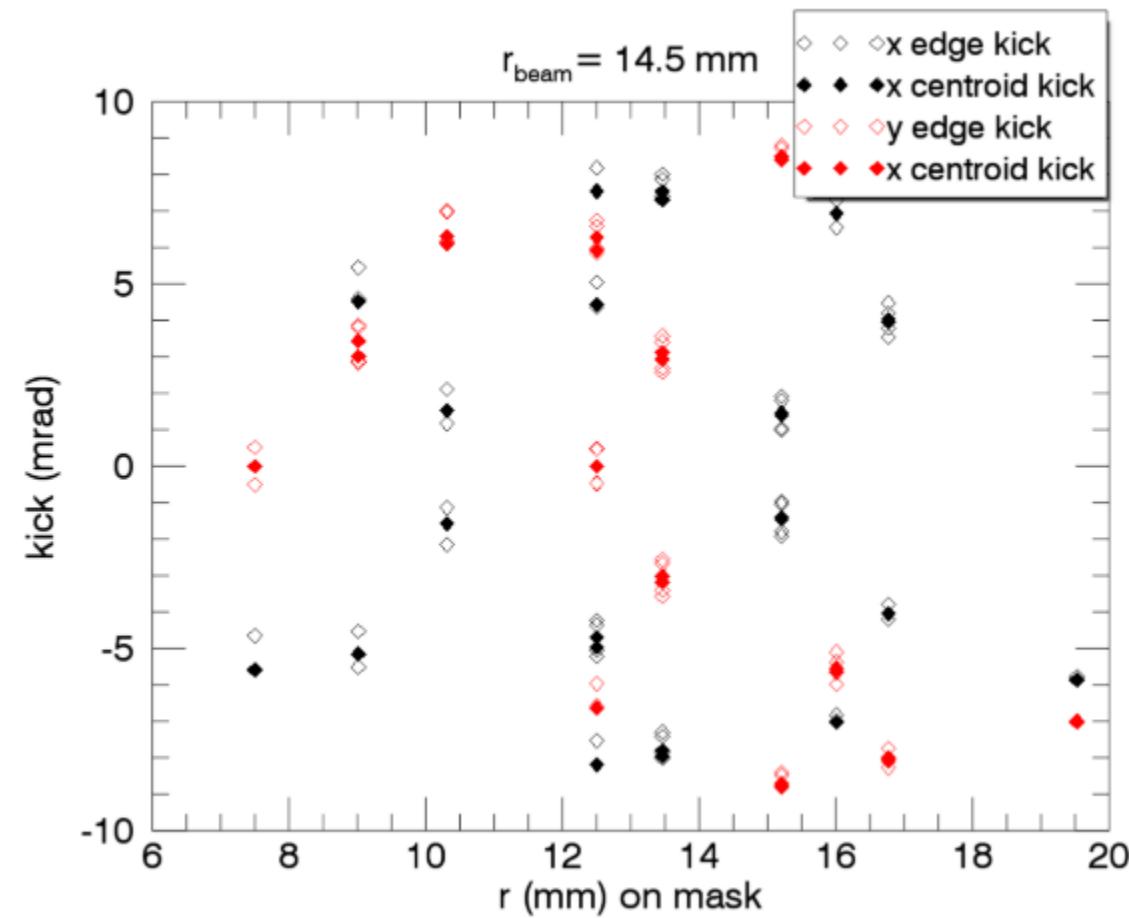


Max beam radius



Kick (in rad) as a function of spot





RMS spot values

$$\langle x \rangle = \frac{\sum_i I_i x_i}{\sum_i I_i}$$

$$\langle x' \rangle = \frac{\sum_i I_i x'_i}{\sum_i I_i}$$

$$\langle x^2 \rangle = \frac{\sum I (x - \langle x \rangle)^2}{\sum I}$$

$$\langle x'^2 \rangle = \frac{\sum I (x' - \langle x' \rangle)^2}{\sum I}$$

$$\langle xx' \rangle = \frac{\sum I (x - \langle x \rangle) (x' - \langle x' \rangle)}{\sum I}$$

$$\varepsilon_x = 4\gamma \sqrt{\langle x^2 \rangle \langle x'^2 \rangle - \langle xx' \rangle^2}$$