



Elettra Sincrotrone Trieste

Advanced Concepts for Seeded FELs

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Paul Scherrer Institut, Villigen PSI

What is seeding?

Single-pass Free Electron Lasers (FELs) can amplify the noise present in the electron beam density profile obtaining powerful SASE radiation.

Using instead an “external” electro-magnetic field (a seed) it is possible to initiate the FEL process in a deterministic way.

The properties of the output FEL radiation are inherited and determined by the seed itself.

The “external” seed can be:
High order harmonics generated in gas (HHG)

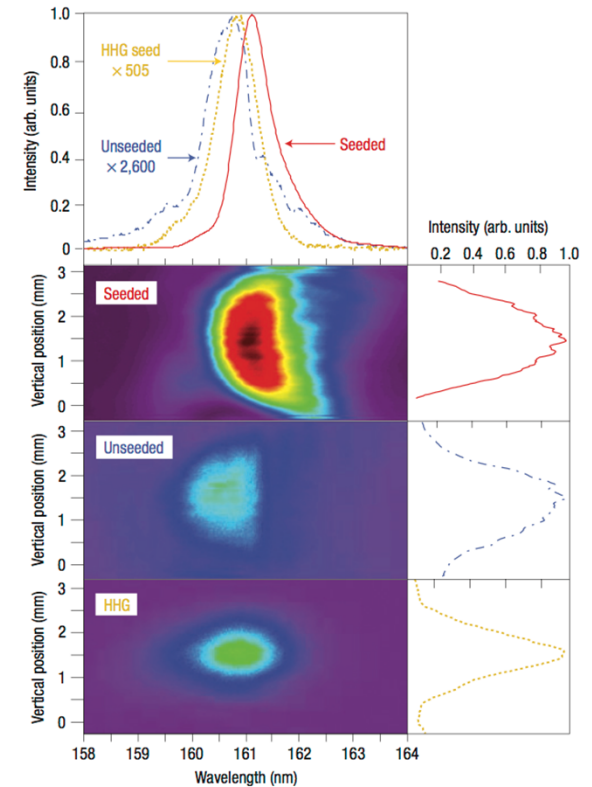


Nature Physics **4**, 296 - 300 (2008)

Published online: 9 March 2008 | doi:10.1038/nphys889

Injection of harmonics generated in gas in a free-electron laser providing intense and coherent extreme-ultraviolet light

G. Lambert^{1,2,3}, T. Hara^{2,4}, D. Garzella¹, T. Tanikawa², M. Labat^{1,3}, B. Carre¹, H. Kitamura^{2,4}, T. Shintake^{2,4}, M. Bougeard¹, S. Inoue⁴, Y. Tanaka^{2,4}, P. Salieres¹, H. Merdji¹, O. Chubar³, O. Gobert¹, K. Tahara² & M.-E. Couprie³

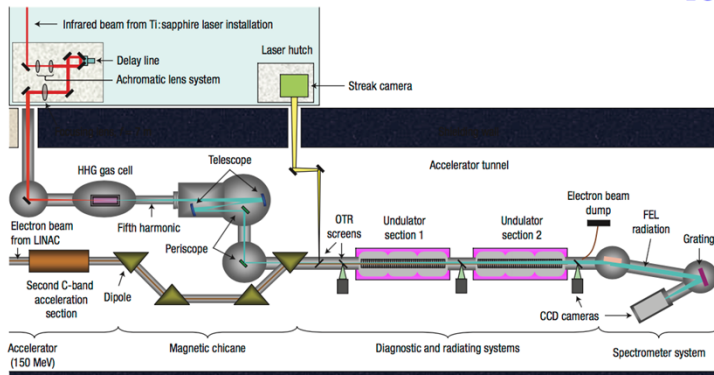


TUPA04

Proceedings of FEL2011, Shanghai, China

sFLASH - PRESENT STATUS AND COMMISSIONING RESULTS*

S. Ackermann, A. Azima, J. Bödewadt, F. Curbis, M. Drescher, E. Hass, M. Mittenzwey, T. Maltezopoulos, V. Miltchev[†], J. Rönsch-Schulenburg, J. Rossbach, R. Tarkeshian, Hamburg University, Hamburg, Germany
H. Delsim-Hashemi, K. Honkavaara, T. Laarmann, H. Schlarb, S. Schreiber, M. Tischer, DESY, Hamburg, Germany
R. Ischebeck, PSI, Villigen, Switzerland



The “external” seed can be: Another FEL Monochromatised SASE (self-seeding) or multiple stages

Journal of Modern Optics

Vol. 58, No. 16, 20 September 2011, 1391–1403

A novel self-seeding scheme for hard X-ray FELs

Gianluca Geloni^{a*}, Vitali Kocharyan^b and Evgeni Saldin^b

^aEuropean XFEL GmbH, Hamburg, Germany; ^bDeutsches Elektronen-Synchrotron (DESY), Hamburg, Germany

(Received 17 December 2010; final version received 29 April 2011)

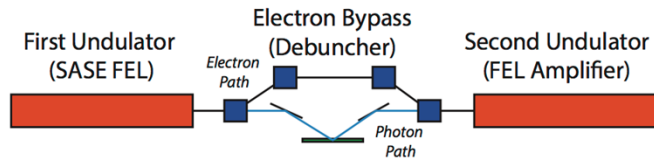
nature
photonics

ARTICLES

PUBLISHED ONLINE: 12 AUGUST 2012 | DOI: 10.1038/NPHOTON.2012.180

Demonstration of self-seeding in a hard-X-ray free-electron laser

J. Amann¹, W. Berg², V. Blank³, F.-J. Decker¹, Y. Ding¹, P. Emma^{4*}, Y. Feng¹, J. Frisch¹, D. Fritz¹, J. Hastings¹, Z. Huang¹, J. Krzywinski¹, R. Lindberg², H. Loos¹, A. Lutman¹, H.-D. Nuhn¹, D. Ratner¹, J. Rzepiela¹, D. Shu², Yu. Shvyd'ko², S. Spampinati¹, S. Stoupin², S. Terentyev³, E. Trakhtenberg², D. Walz¹, J. Welch¹, J. Wu¹, A. Zholents² and D. Zhu¹



PRL 114, 054801 (2015)

PHYSICAL REVIEW LETTERS

week ending
6 FEBRUARY 2015

Experimental Demonstration of a Soft X-Ray Self-Seeded Free-Electron Laser

D. Ratner,^{1,*} R. Abela,² J. Amann,¹ C. Behrens,¹ D. Bohler,¹ G. Bouchard,¹ C. Bostedt,¹ M. Boyes,¹ K. Chow,³ D. Cocco,¹ F. J. Decker,¹ Y. Ding,¹ C. Eckman,¹ P. Emma,¹ D. Fairley,¹ Y. Feng,¹ C. Field,¹ U. Flechsig,² G. Gassner,¹ J. Hastings,¹ P. Heimann,¹ Z. Huang,¹ N. Kelez,¹ J. Krzywinski,¹ H. Loos,¹ A. Lutman,¹ A. Marinelli,¹ G. Marcus,¹ T. Maxwell,¹ P. Montanez,¹ S. Moeller,¹ D. Morton,¹ H. D. Nuhn,¹ N. Rodes,³ W. Schlott,¹ S. Serkez,⁴ T. Stevens,³ J. Turner,¹ D. Walz,¹ J. Welch,¹ and J. Wu¹

¹SLAC National Accelerator Laboratory, Menlo Park, California 94720, USA

²Paul Scherrer Institut, CH-5232 Villigen PSI, Switzerland

³Lawrence Berkeley National Laboratory (LBNL), 1 Cyclotron Road, Berkeley, California 94720, USA

⁴Deutsches Elektronen-Synchrotron (DESY), Notkestrasse 85, D-22607 Hamburg, Germany

(Received 3 October 2014; published 6 February 2015)

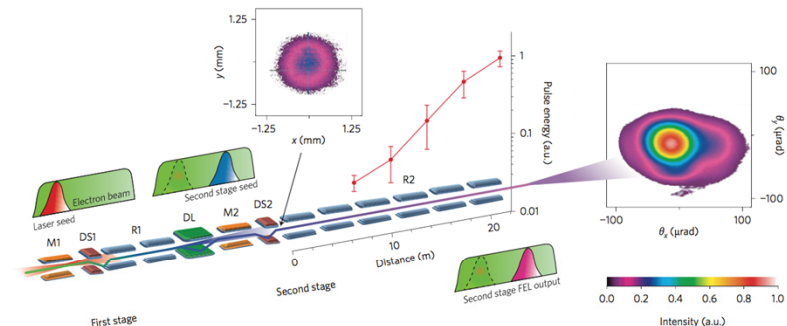
nature
photonics

ARTICLES

PUBLISHED ONLINE: 20 OCTOBER 2013 | DOI: 10.1038/NPHOTON.2013.277

Two-stage seeded soft-X-ray free-electron laser

E. Allaria¹, D. Castronovo¹, P. Cinquegrana¹, P. Craievich^{1†}, M. Dal Forno^{1,2}, M. B. Danailov¹, G. D'Auria¹, A. Demidovich¹, G. De Nino^{1,3}, S. Di Mitri¹, B. Diviacco¹, W. M. Fawley^{1*}, M. Ferianis¹, E. Ferrari¹, L. Froehlich¹, G. Gaio¹, D. Gauthier^{1,3}, L. Giannessi^{1,4*}, R. Ivanov¹, B. Mahieu^{1,5}, N. Mahne¹, I. Nikolov¹, F. Parmigiani^{1,2}, G. Penco¹, L. Raimondi¹, C. Scafuri¹, C. Serpico¹, P. Sigalotti¹, S. Spampinati^{1,3}, C. Spezzani¹, M. Svandrlík¹, C. Svetina^{1,2}, M. Tovo¹, M. Veronese¹, D. Zangrando¹ and M. Zangrando^{1,6}



The “external” seed can be*:

A conventional laser

PHYSICAL REVIEW A

VOLUME 44, NUMBER 8

15 OCTOBER 1991

Generation of intense uv radiation by subharmonically seeded single-pass free-electron lasers

L. H. Yu

National Synchrotron Light Source, Brookhaven National Laboratory, Upton, New York 11973

(Received 1 April 1991)

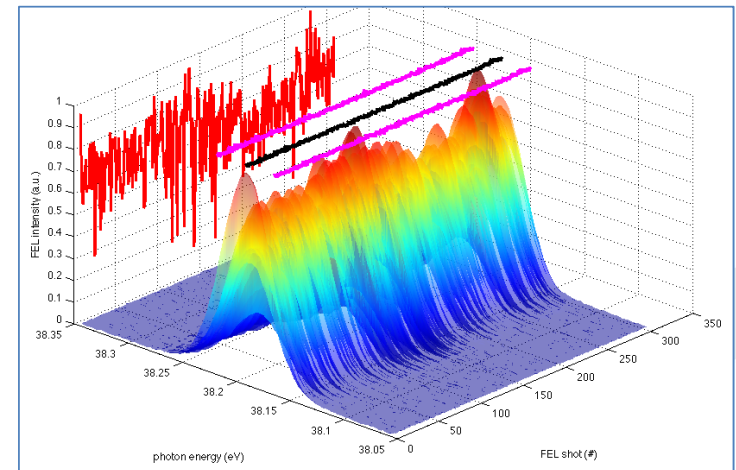
nature
photonics

ARTICLES

PUBLISHED ONLINE: 23 SEPTEMBER 2012 | DOI: 10.1038/NPHOTON.2012.233

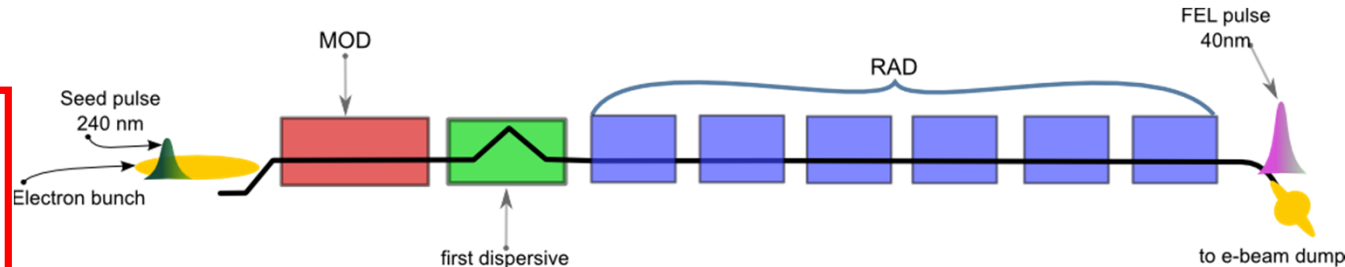
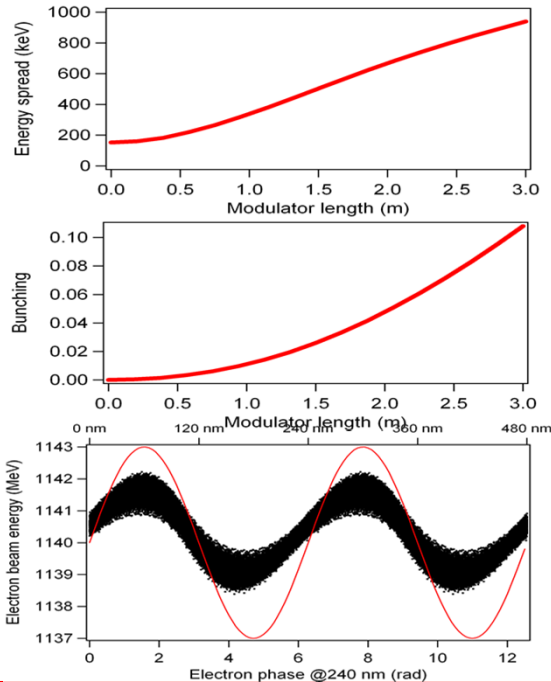
Highly coherent and stable pulses from the FERMI seeded free-electron laser in the extreme ultraviolet

E. Allaria *et al.**



High Gain Harmonic Generation (HGHHG)

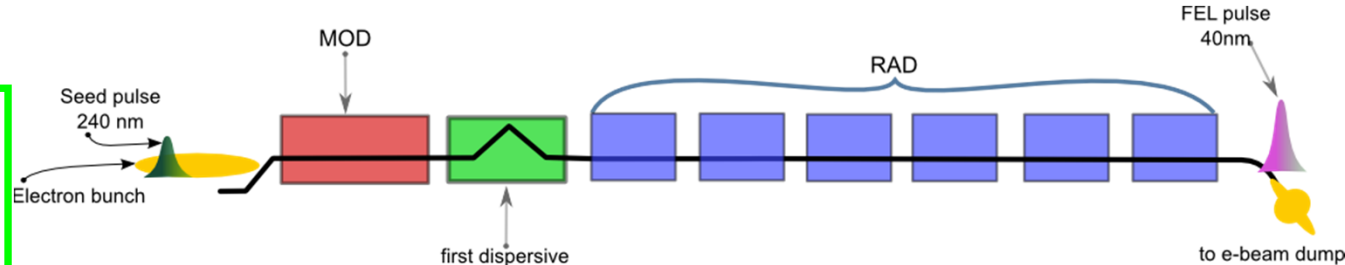
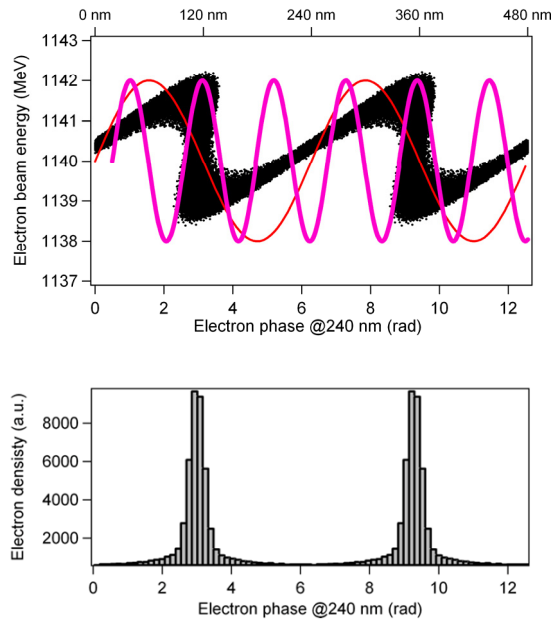
Modulator



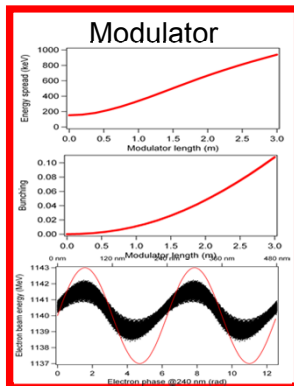
- An external laser interacts with the electron beam in the modulator inducing a periodic modulation of the electron beam energy.
- In a dispersive region, energy modulation is converted into density modulation with significant harmonic content. Beam current modulation inherits coherence properties from the input laser.
- Coherent bunching produces coherent emission in the radiator at the desired harmonic that is then amplified.

High Gain Harmonic Generation (HGHHG)

Dispersive section

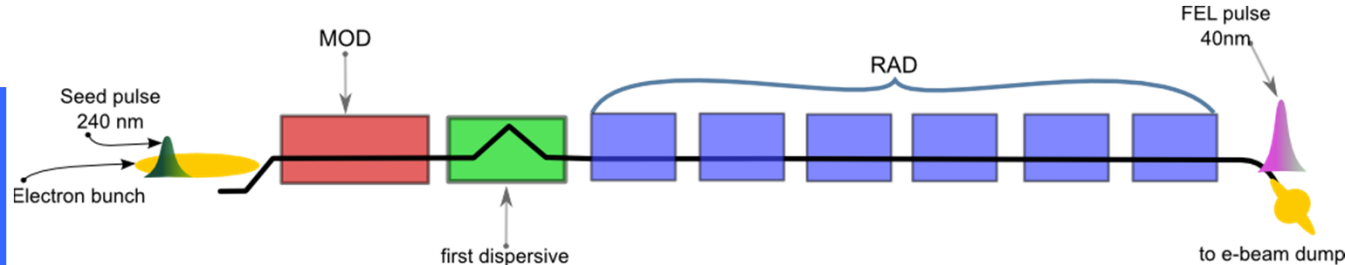
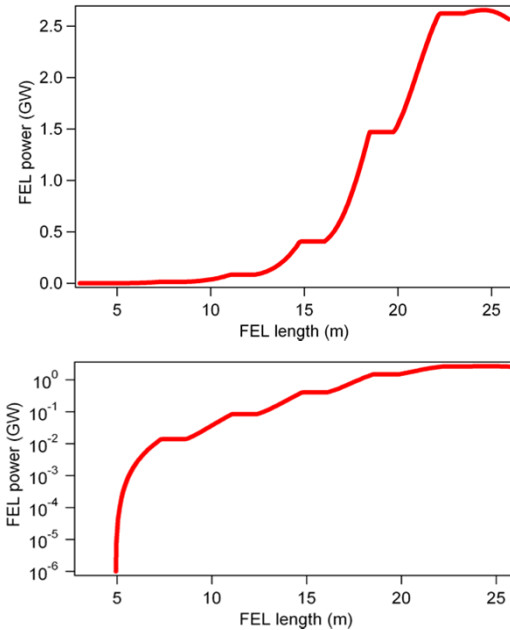


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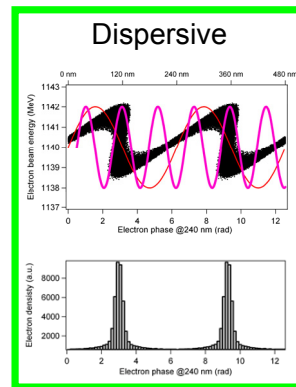
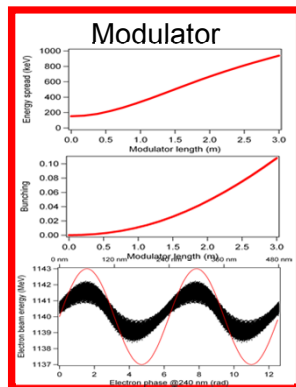


High Gain Harmonic Generation (HGHHG)

Radiator



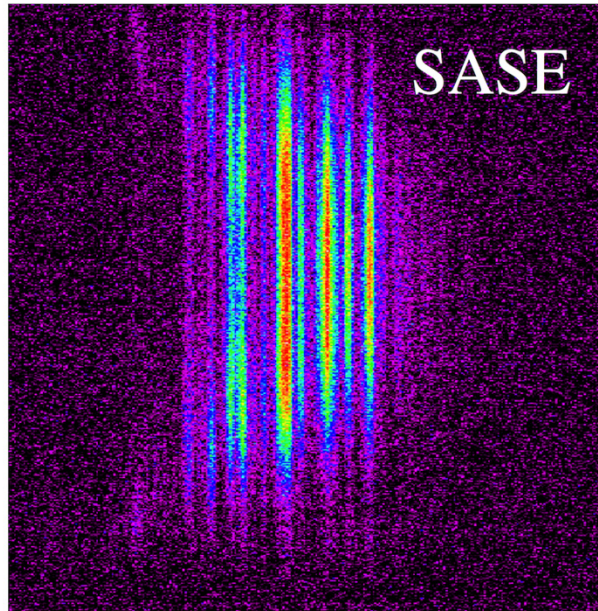
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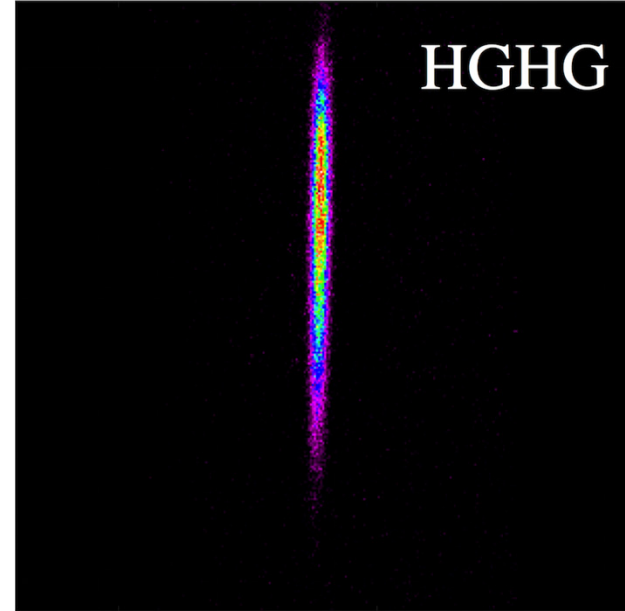
- Coherent bunching produces coherent emission in the radiator at the desired harmonic that is then amplified.

Self-Amplified Spontaneous Emission

Comparison between SASE and seeded spectra



32 λ (nm) 32.5 33



32 λ (nm) 32.5 33

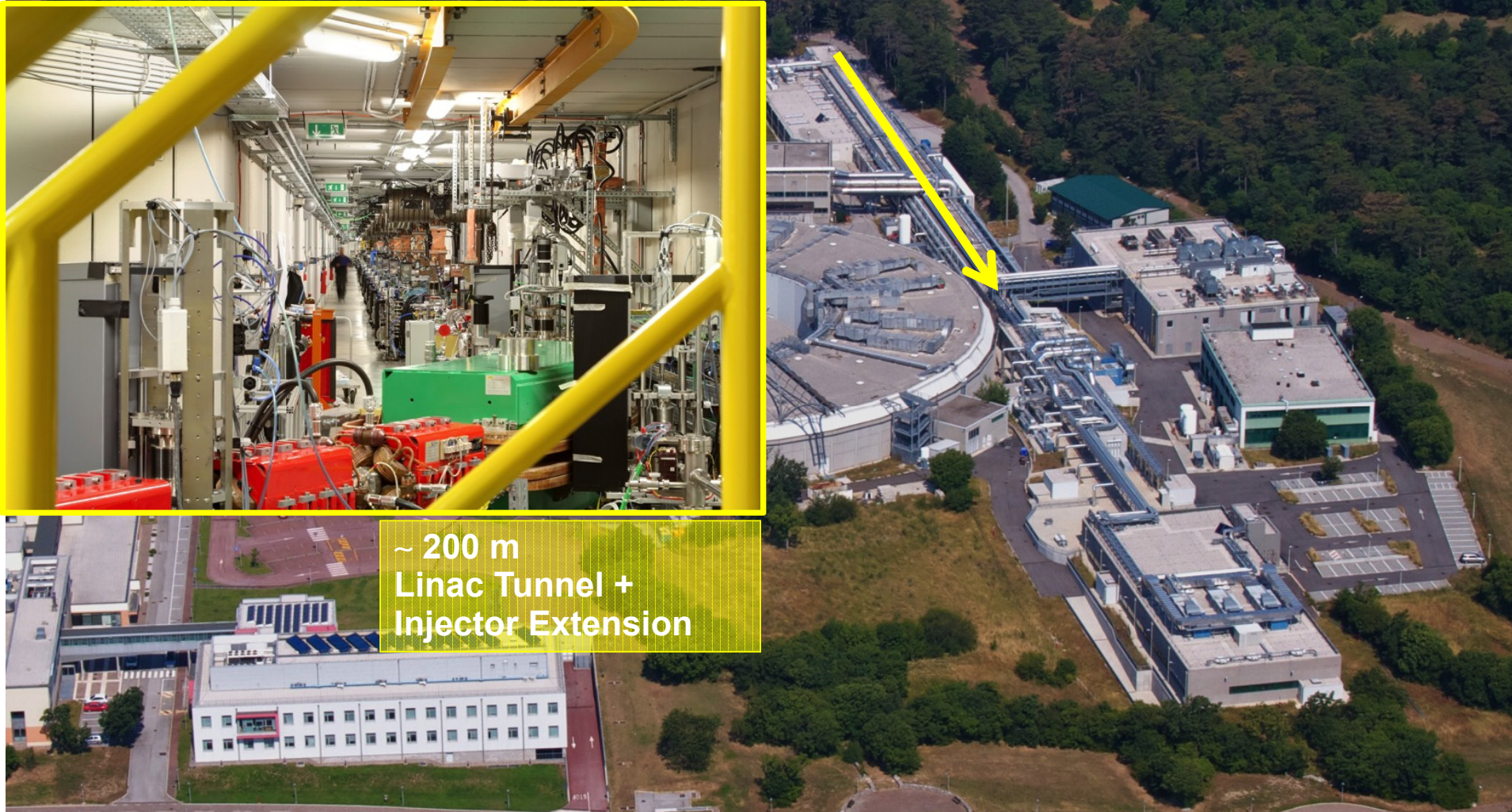
That's why you need a seed!

FERMI Free Electron Laser



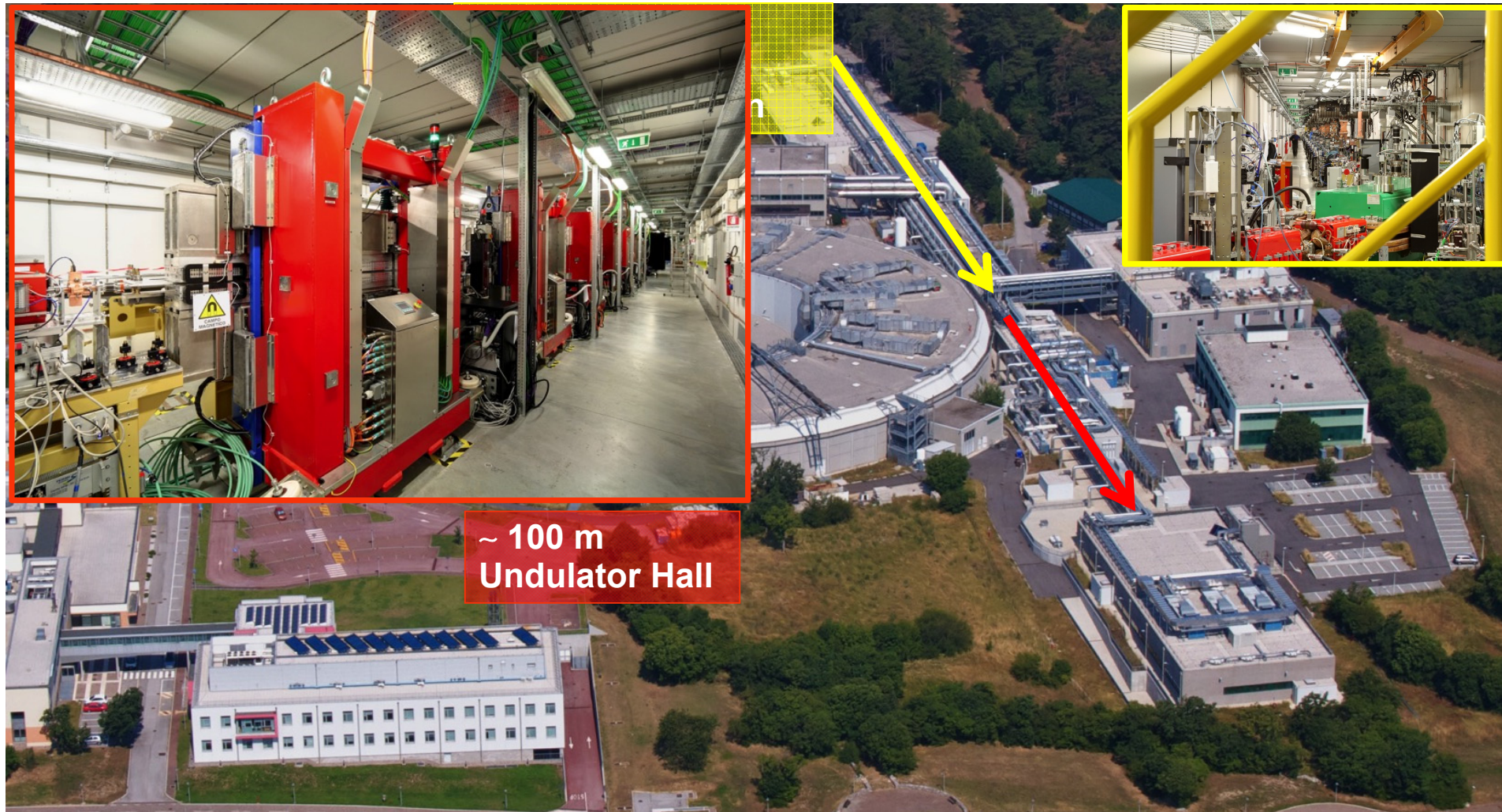
Sponsored by: Italian Minister of University and Research (MIUR), Regione Auton. Friuli Venezia Giulia, European Investment Bank (EIB), European Research Council (ERC), European Commission (EC)

FERMI Free Electron Laser



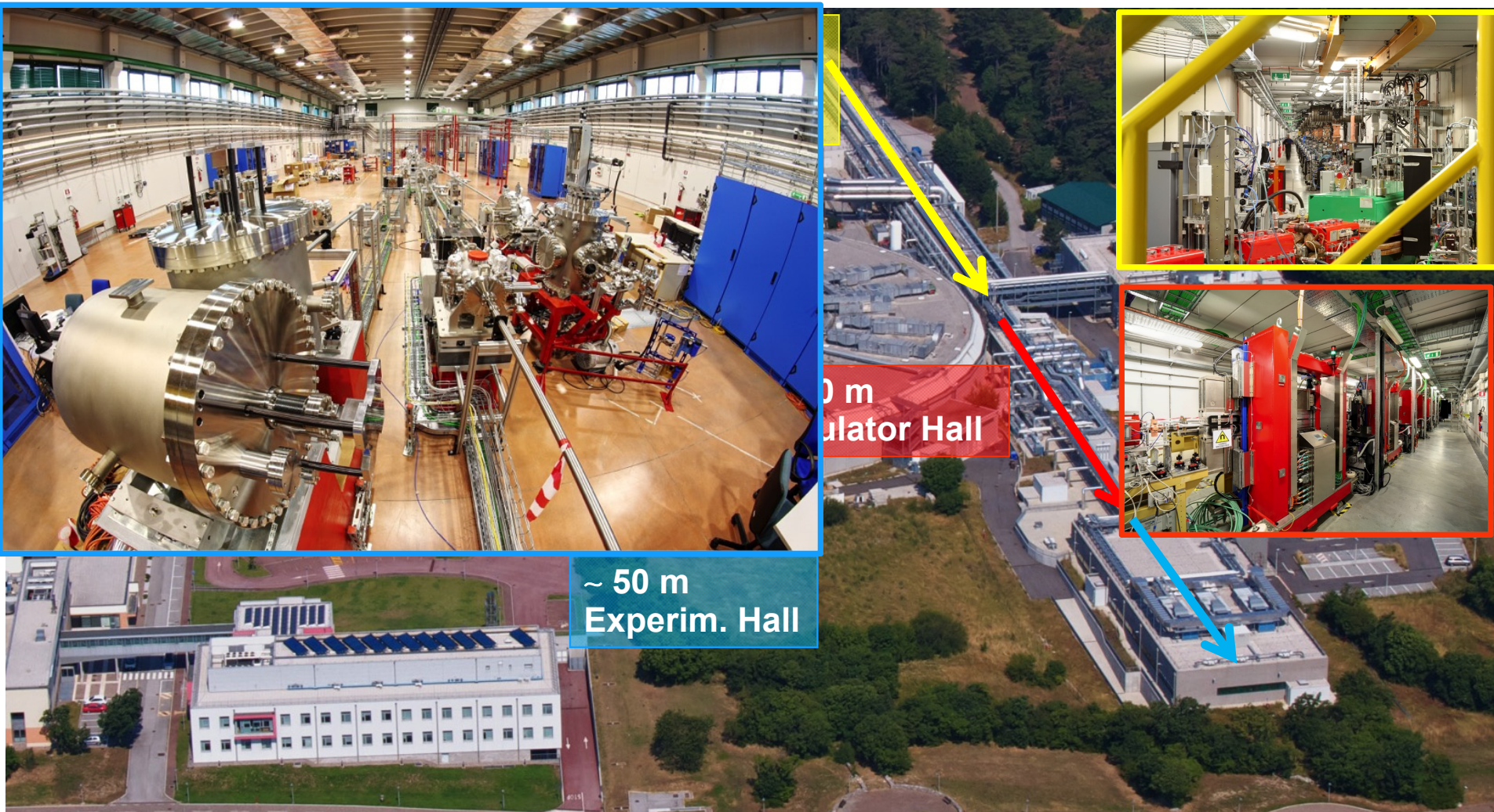
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FERMI Free Electron Laser

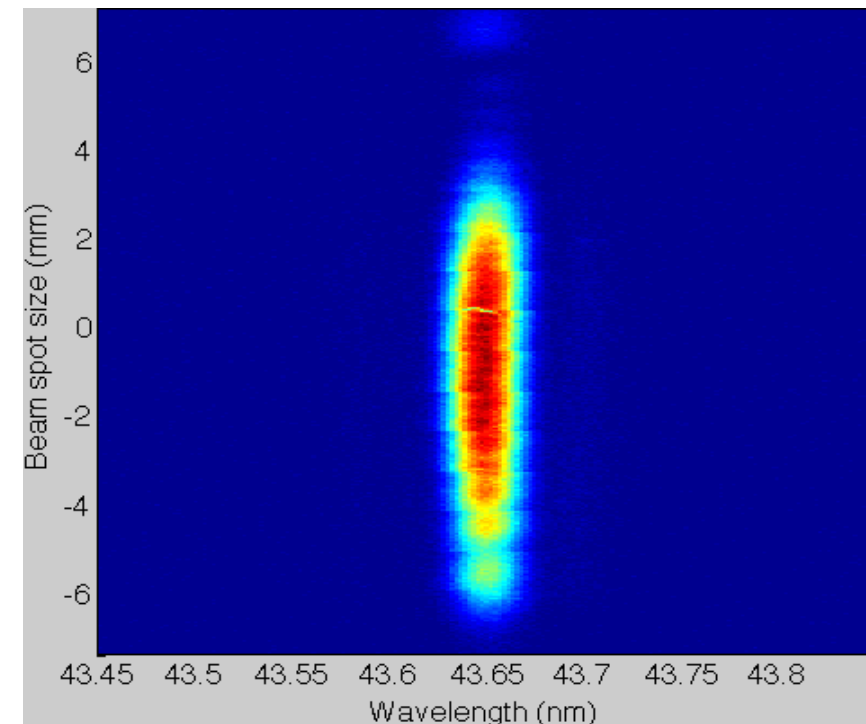


Sponsored by: Italian Minister of University and Research (MIUR), Regione Auton. Friuli Venezia Giulia, European Investment Bank (EIB), European Research Council (ERC), European Commission (EC)

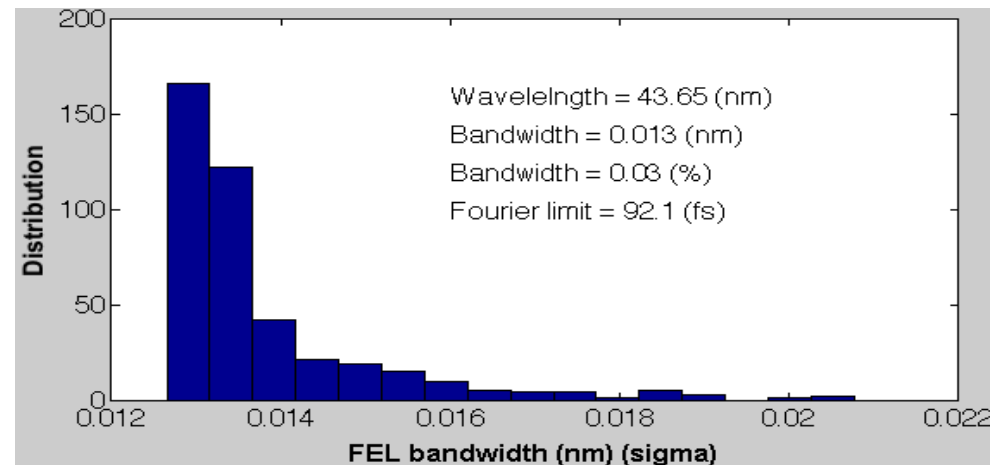
FEL bandwidth control

FEL pulses inherit the **coherence** from the **external seed** laser.

The use of **seed pulses** close to the **Fourier limit**, allow the production of **FEL pulses** that are also **close** to the **Fourier limit**.



- Typical seed pulses are **120 fs** long with **0.9 nm** spectral width.
- **FEL pulses** are characterized by **single mode spectra** with **narrow linewidth**.



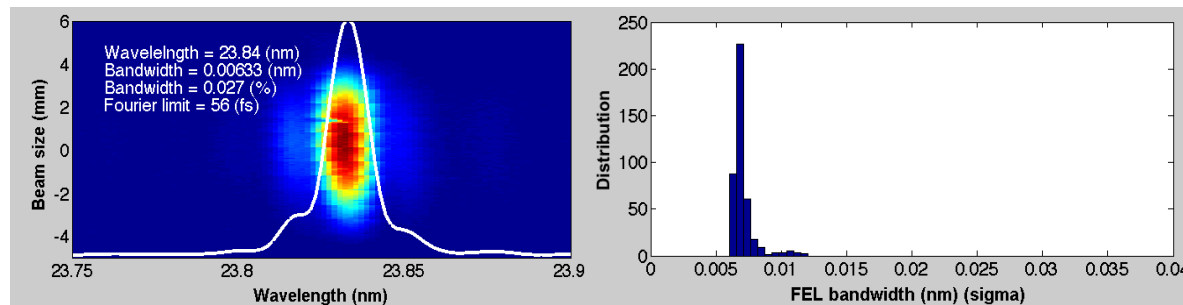
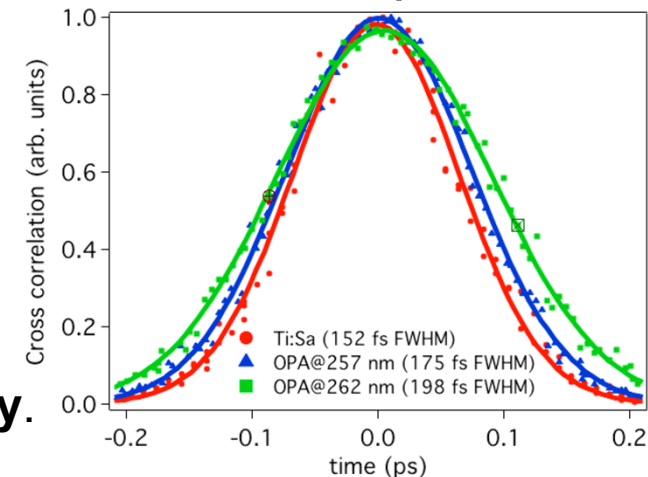
- **FEL spectra** fit **Gaussian** curves and both **wavelength** and **bandwidth** are **stable**.
- Calculated Fourier-limited pulse length matches theoretical predictions suggesting a **very high** longitudinal **coherence**.

Seed laser chirp

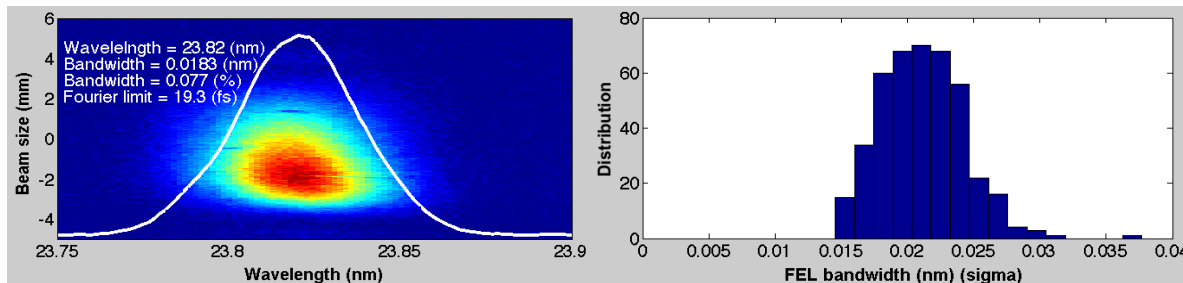
Fine **tuning** of the **FEL wavelength** requires **changes** of the **seed laser wavelength** and it implies the use of an Optical Parametric Amplifier (**OPA**) laser.

Due to his complexity and to the effects of the multilayers mirrors, with **OPA** we have **less control** on the seed laser **chirp**. At some specific wavelengths **seed pulses** can be **longer** with stronger **chirp**.

This has a **significant impact** in the FEL **spectral quality**.



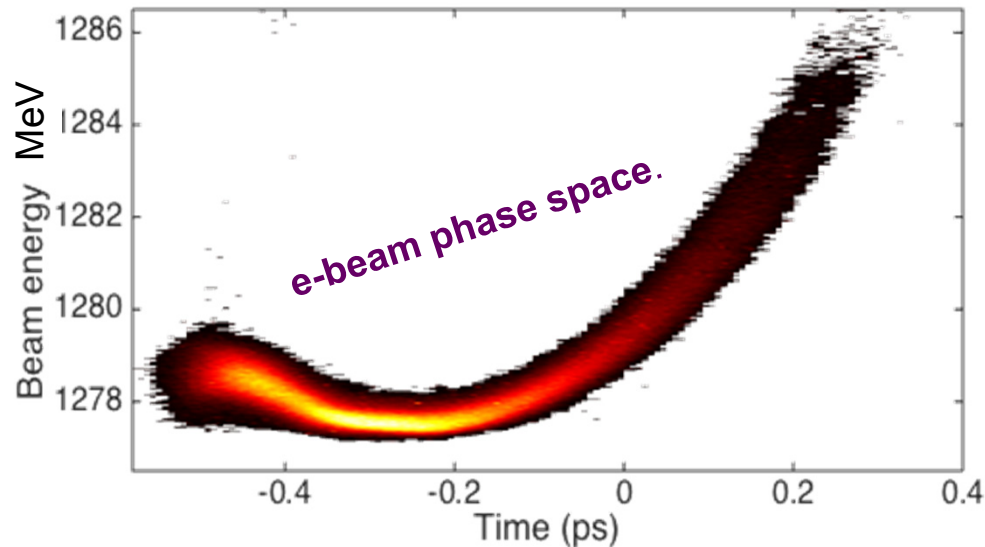
In some cases, **by switching** the seed laser from **Ti:Sa** to **OPA**, the FEL **bandwidth** can **increase by a factor ~3**.



If coherence and chirp are preserved, **larger bandwidth** can be used for compressing the pulse (**CPA**).

FERMI electron beam

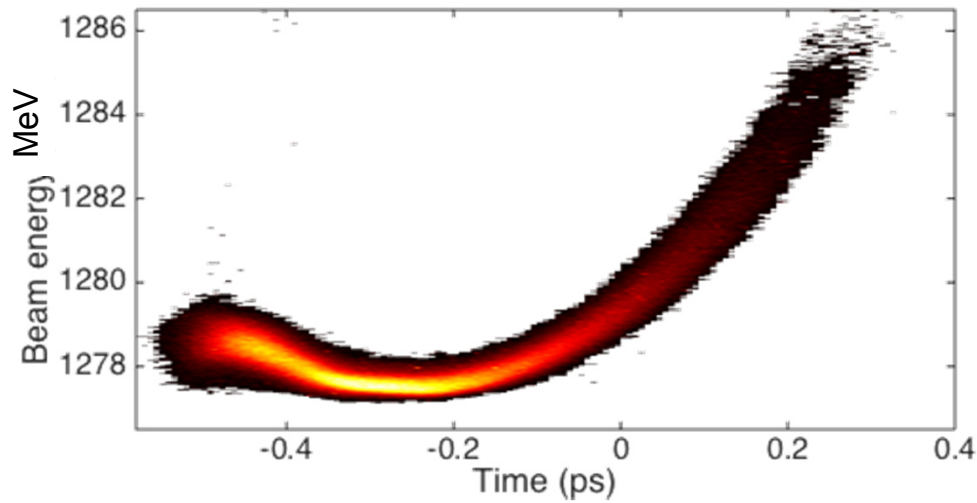
- Current **spikes** are **not suitable** for seeded FELs.
- **Low energy spread** and **flat longitudinal phase space** are required for seeding.
- Electron beam **optimization** is **different** than for a **SASE** FEL.
- Moderate compression is used.
- Phase space nonlinearities may counteract the benefits of the seed.
- Only electrons interacting with the laser participate to the FEL emission.



Charge	700	pC
Peak current	~800	A
Energy	1 – 1.5	GeV
Energy spread	~150	keV
Energy chirp	~3	MeV
Emittance	1	mm mrad
Size (rms)	~100	μm

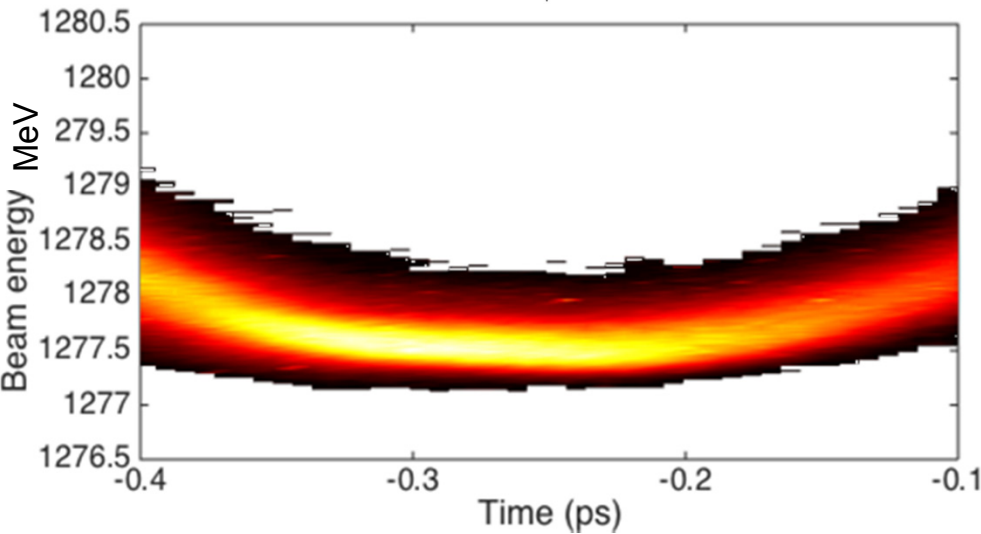
Electron beam parameters

Seeded electron beam



Head and **tail** parts of the beam
can **deteriorate** FEL properties.

Seeded electron beam



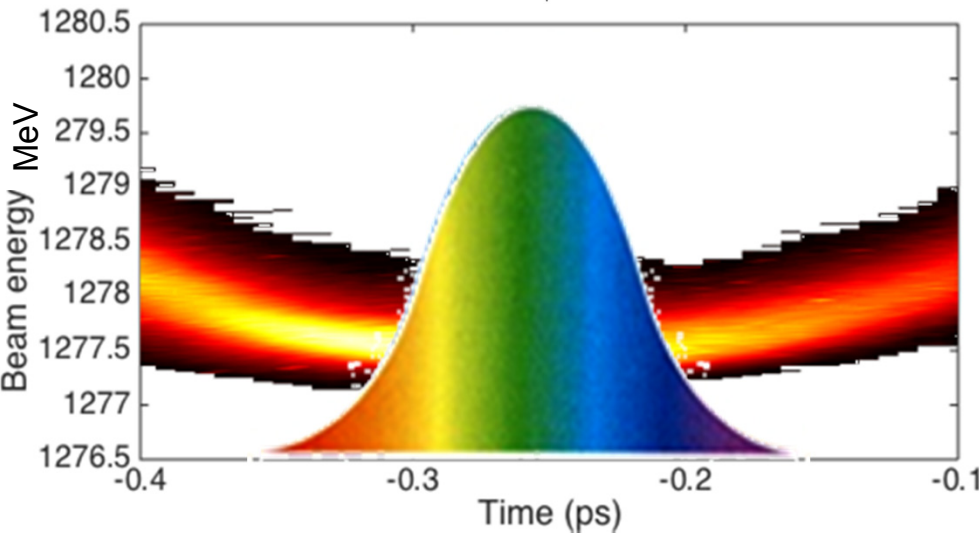
Head and **tail** parts of the beam can **deteriorate** FEL properties.

Seed laser pulse is **shorter** than the electron beam.

This allow to properly select the **best** part of the **electron** for **FEL**.

Because the final radiator is short the **unseeded** part does **not produce** any significant **radiation**.

Seeded electron beam



Head and **tail** parts of the beam can **deteriorate** FEL properties.

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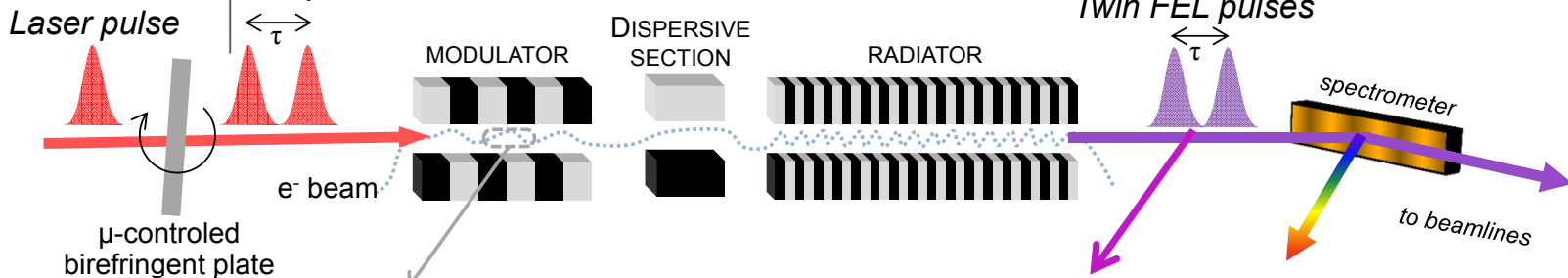
Because the final radiator is short the **unseeded** part does **not** produce any significant **radiation**.

If the **electron beam** is **not uniform** the FEL properties may depend on seed timing.

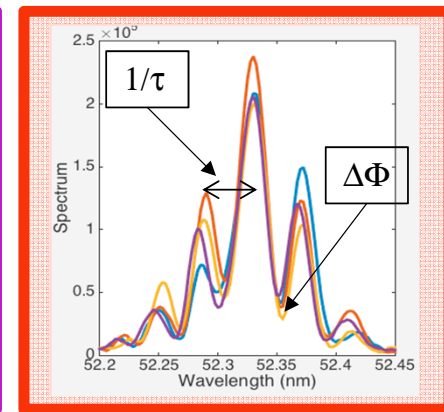
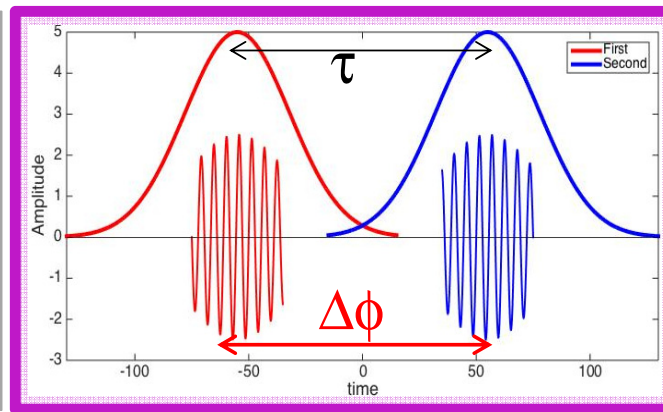
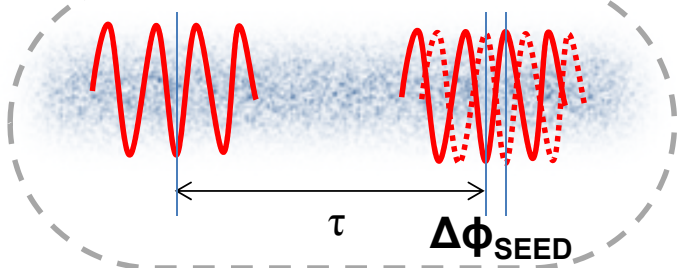


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Phase-locked FEL pulses

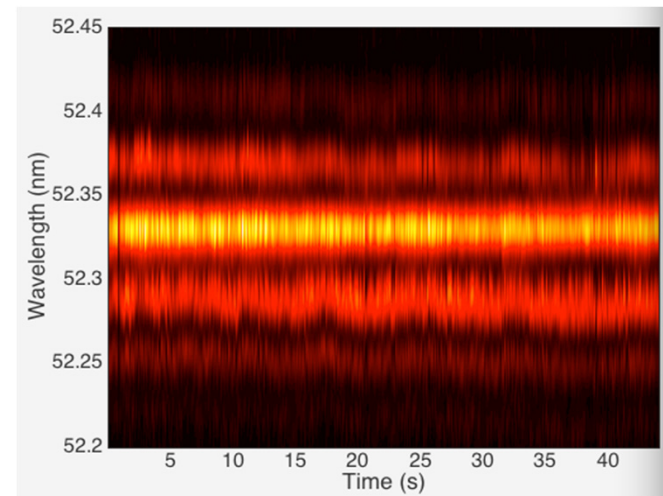


Twin seed-e- beam
interaction



With two seed lasers one can control and change the relative time between two FEL pulses. For coherent pulses, a fine tuning allows to control the relative phase between the two FEL pulses.

Interference between two coherent and phase-locked pulses is evident in the spectral domain.



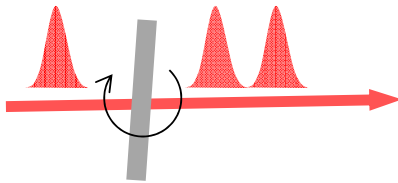
FEL phase-locking control

An accurate **control** of the relative **phase** between the **seed pulses** with steps

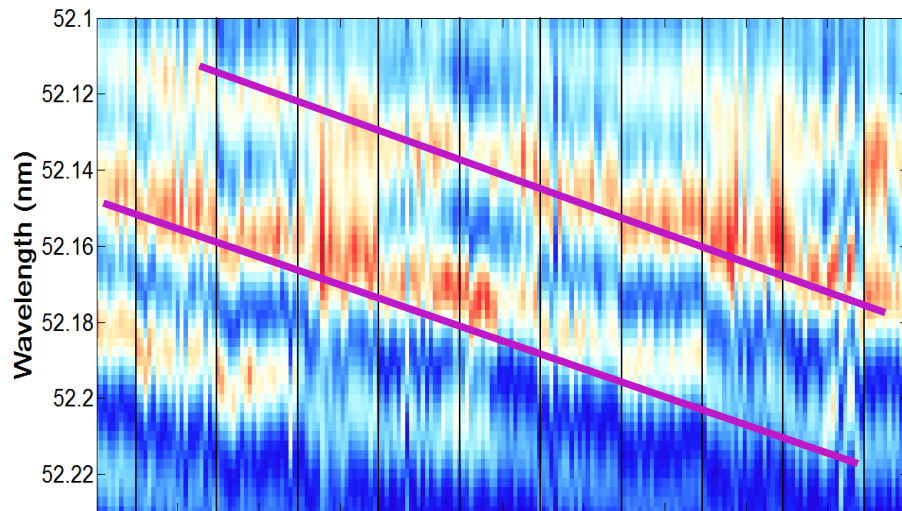
$$\Delta\phi_{\text{seed}} \sim 12^\circ @ 260\text{nm}$$

Allow to change the **phase** between the **FEL pulses** by

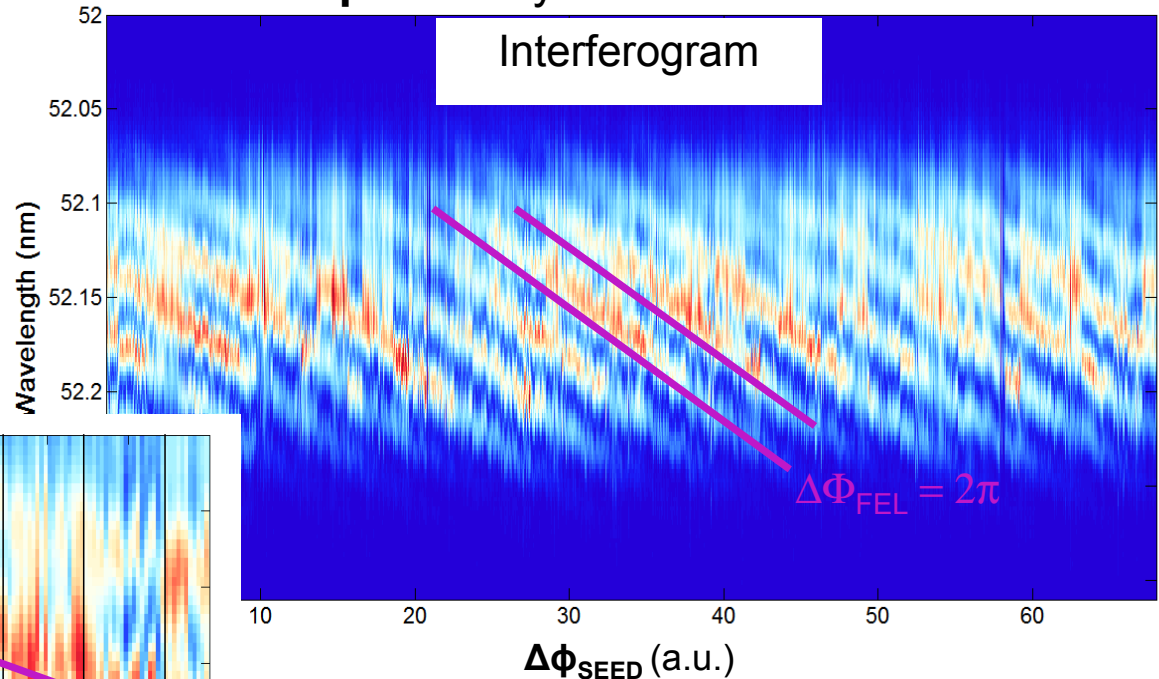
$$\Delta\phi_{\text{FEL}} \sim 60^\circ @ 52\text{nm}$$



This correspond to a **control** with ~ 30 as resolution.



Zoom of the central part

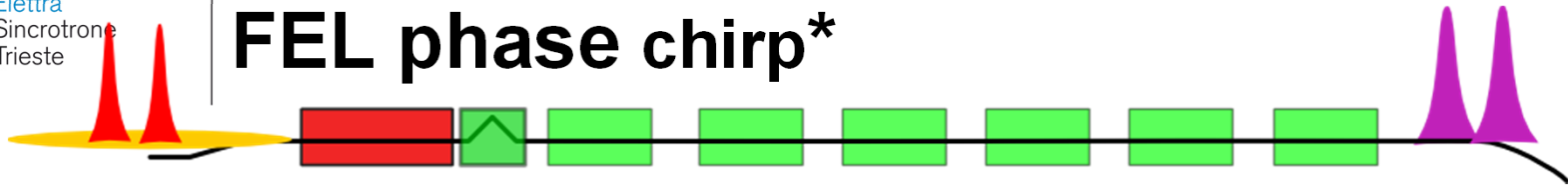


Residual fringe instability is associated to shifts of the central wavelength of the spectral envelope and has strong correlation with the electron beam bunch compression monitor.



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Measurements and control of the FEL phase chirp*

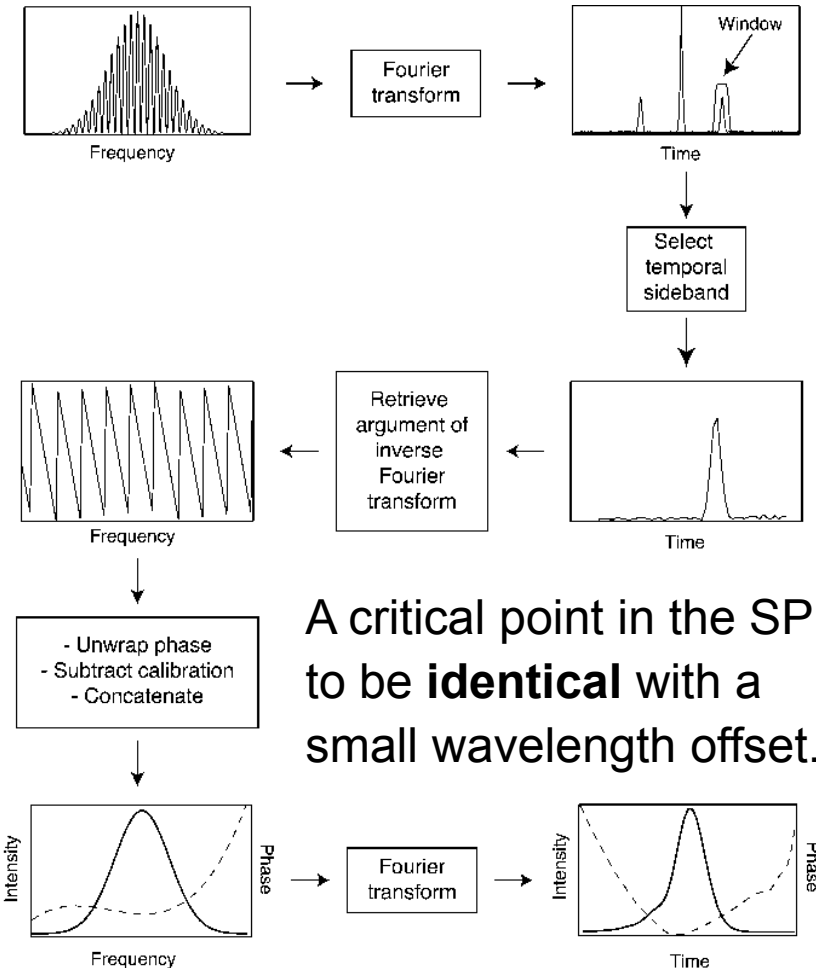


The phase locked FEL pulses are suitable for **SPIDER**[§] measurement (amplitude and phase)

[§]“Spectral Phase Interferometry for Direct Electric-field Reconstruction ...”. OPTICS LETTERS 23, 792 1998

From a **single measurement of the interference spectrogram** of two identical pulses (with slightly different wavelength) the **SPIDER** algorithm **reconstructs the intensity and phase profiles** in the temporal and spectral domain.

A critical point in the SPIDER reconstruction is that the **two pulses** have to be **identical** with a small wavelength offset.

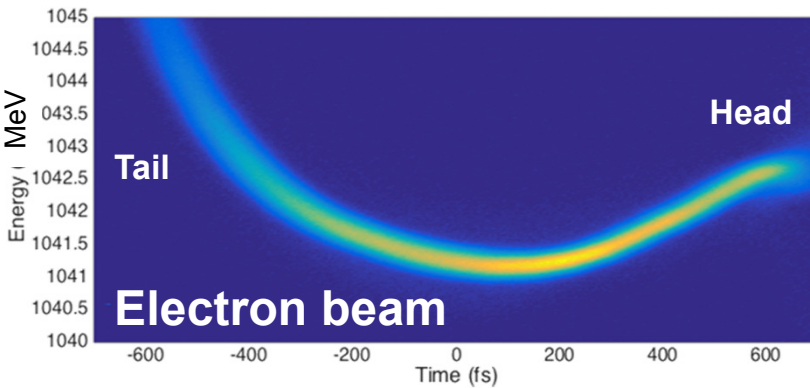
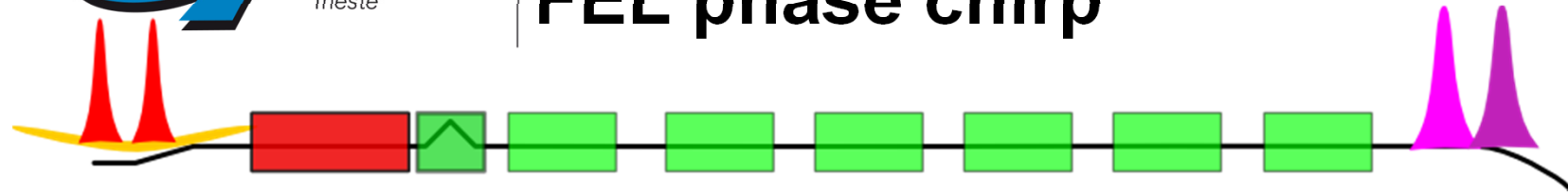


FEL wavelength	52.3 nm
seed wavelength/harmonic number	261.5/5
seed pulse duration	126 fs/180 fs (FWHM)
Peak current	≈ 600 A
temporal delay between seed pulse replicas (τ)	231 fs
spectral shear between pulse replicas (Ω)	1.7×10^{13} Hz

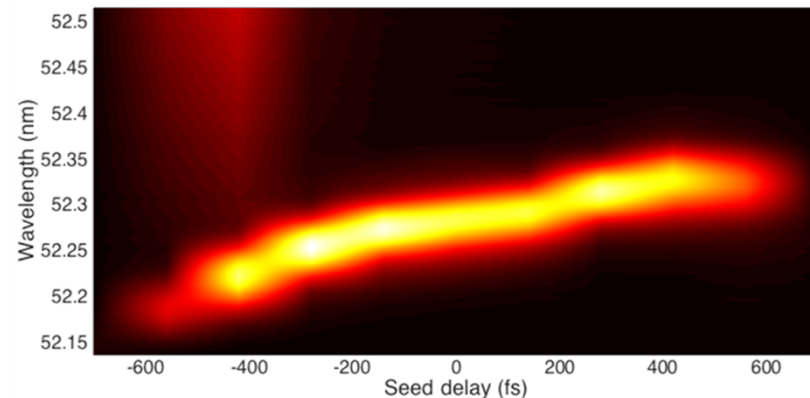


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Measurements and control of the FEL phase chirp

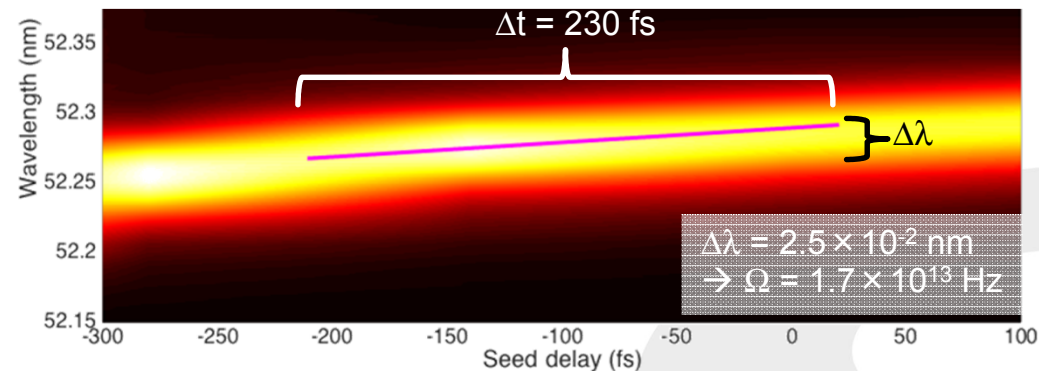


A simple **solution** at FERMI is to seed with **two identical pulses** and to allow the **electron beam phase-space** to slightly **detune** the FEL **wavelengths**.

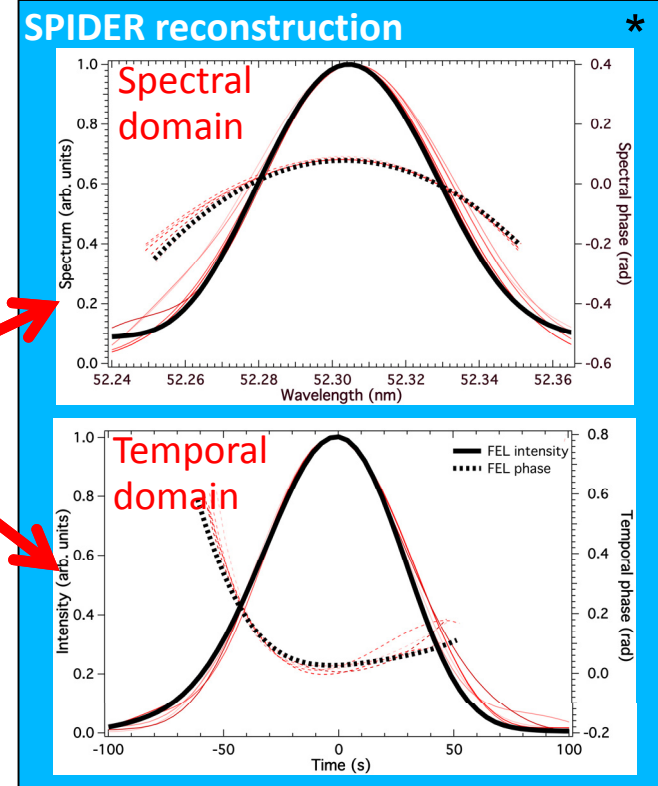
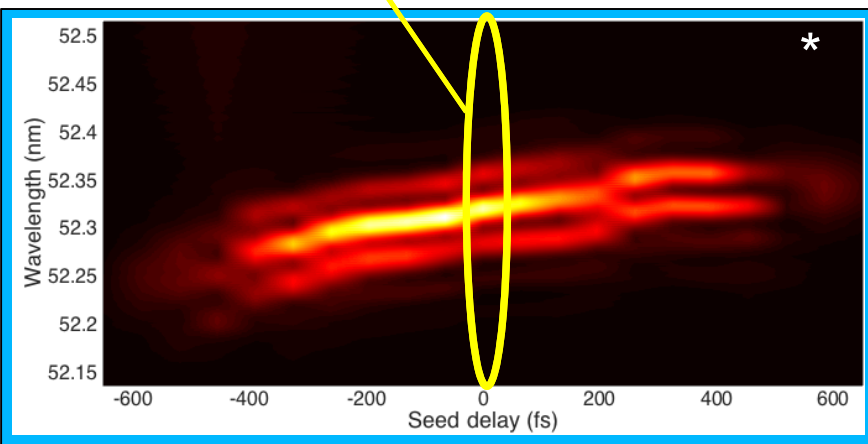
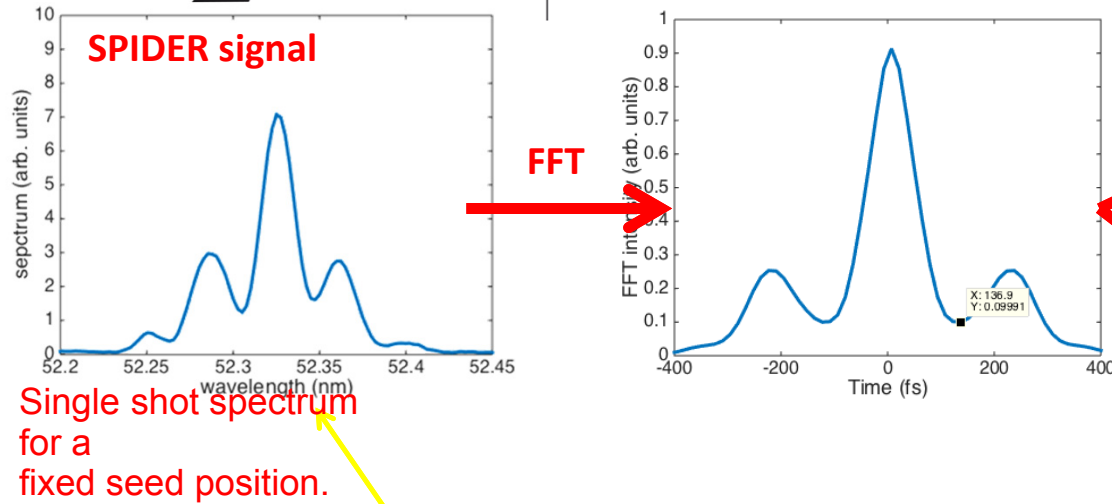


The frequency separation Ω is **used** in the **SPIDER** reconstruction.

Spectral measurements at different seed laser **delay** give the **wavelength separation**.



SPIDER reconstruction



When the **two identical seed lasers** are used clear interference **fringes** are visible.

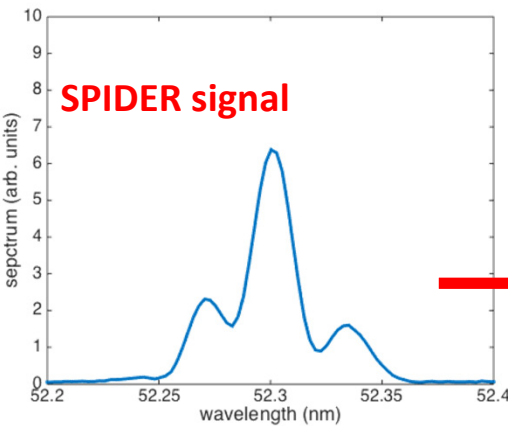
Fringes shape is constant when **scanning** the **seed** lasers over a good range of the beam (**~400 fs**).

Under the assumption that the two FEL pulses have similar characteristics the SPIDER algorithm allows reconstruction of the temporal pulse profile (positive chirp).

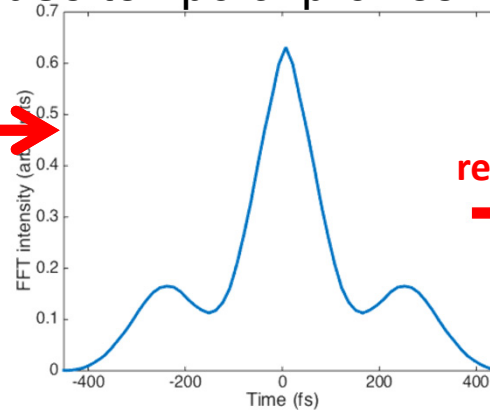
A pulse length of 70 fs is measured.

FEL phase chirp control

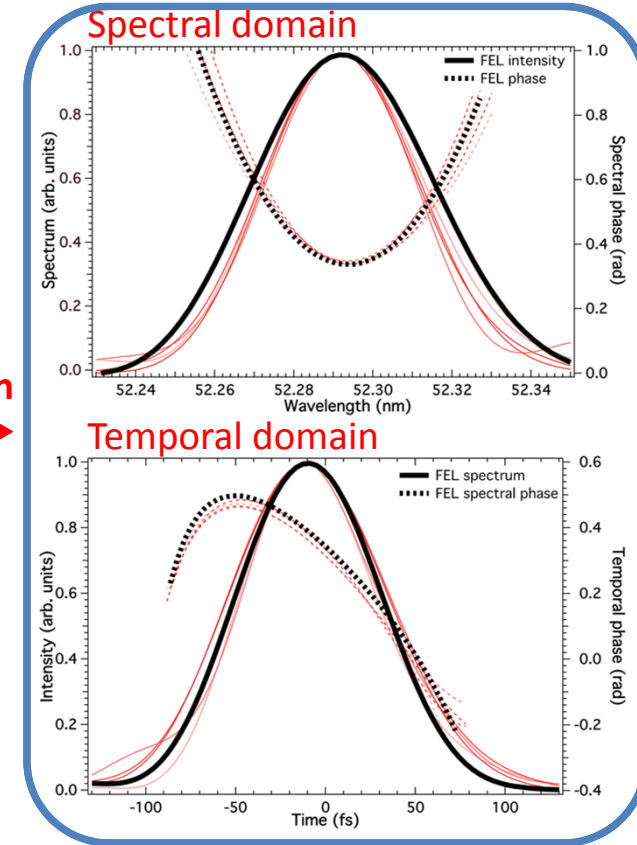
When changing the seed laser chirp from positive to negative, the shape of the interference fringes changes significantly.



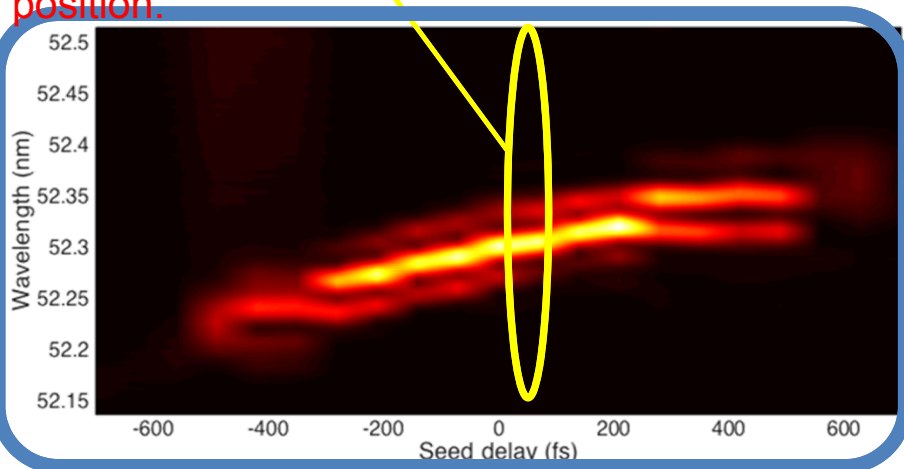
Evident difference in the reconstructed intensity and phase temporal profiles.



SPIDER reconstruction

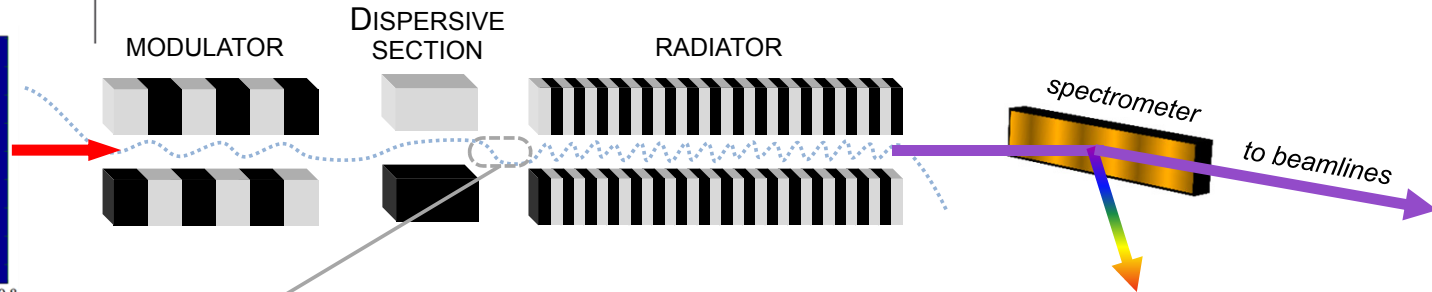
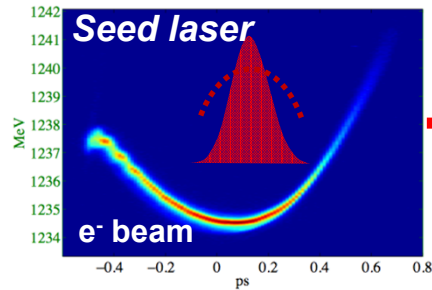


Single shot spectrum for a fixed seed position.

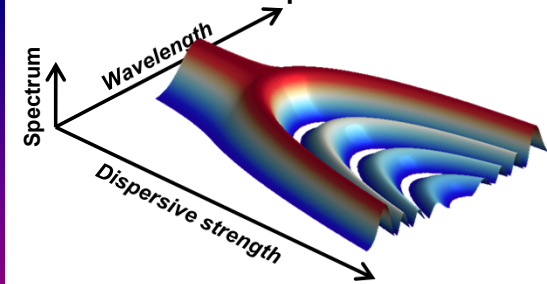


Experiments with a **controlled chirp** along the **FEL pulse** may become **possible**.

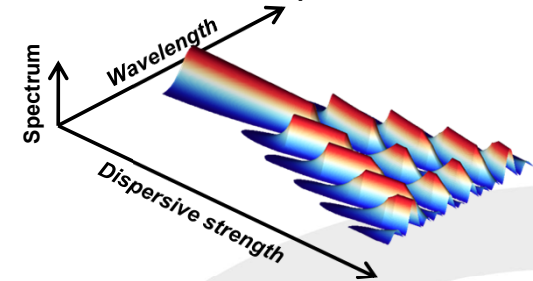
Spectro-temporal shaping (theory)



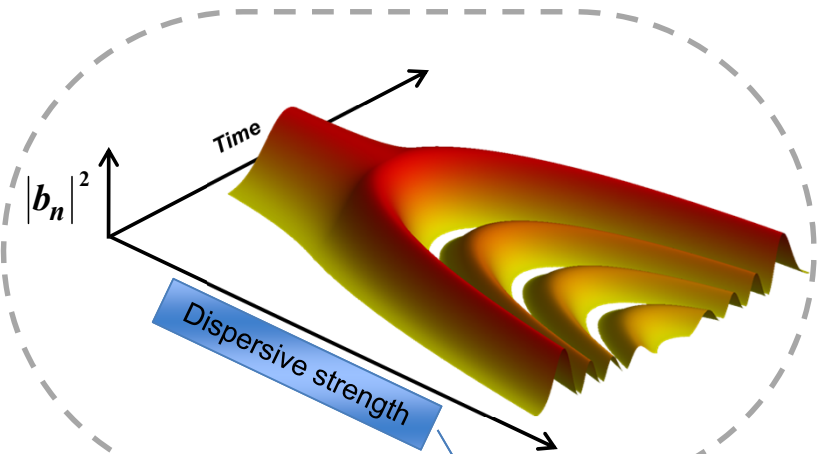
Spectral profile of the FEL pulse



Spectral responses to the chirps



Temporal profile of the bunching

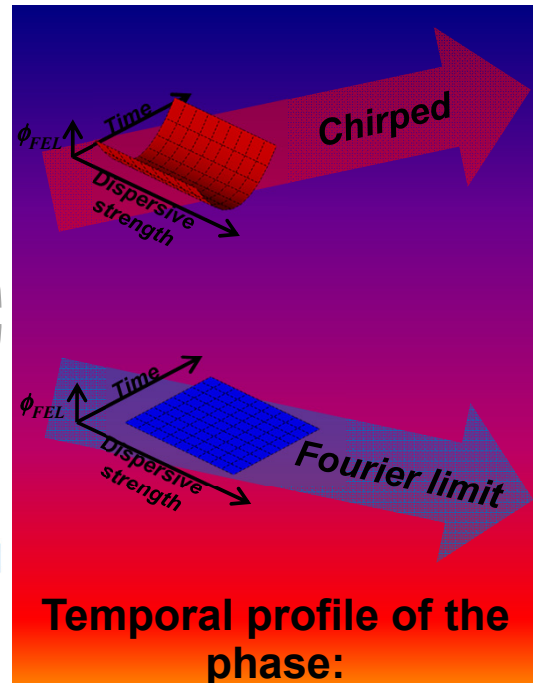


$$b_n(t) = e^{-n^2 B^2 / 2} J_n[-n B A(t)] e^{in[\phi_s(t) + \phi_e(t)]}$$

seed envelope

seed slowly varying phase

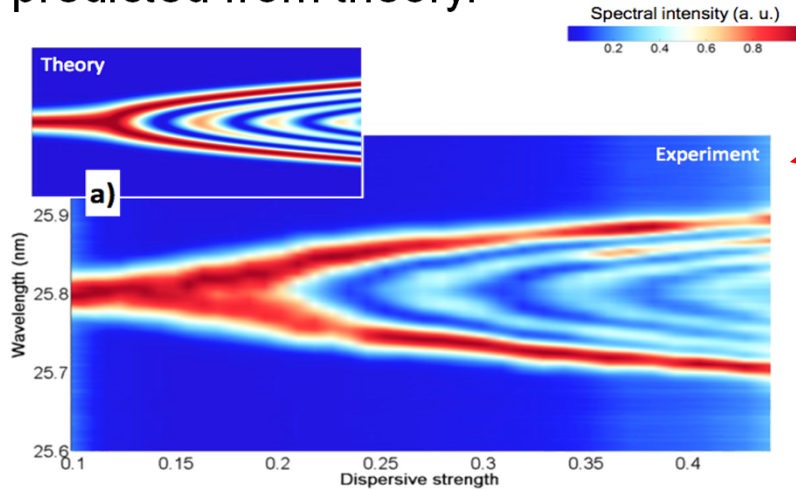
e- beam phase from time-dependent energy profile



contribution of the bunching + FEL amplification

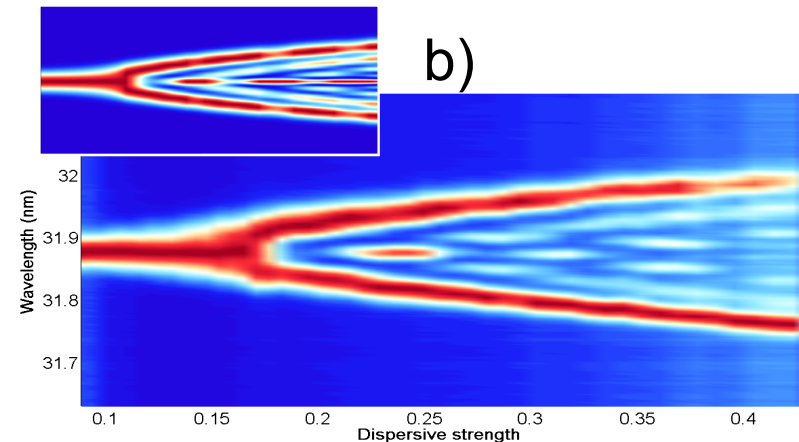
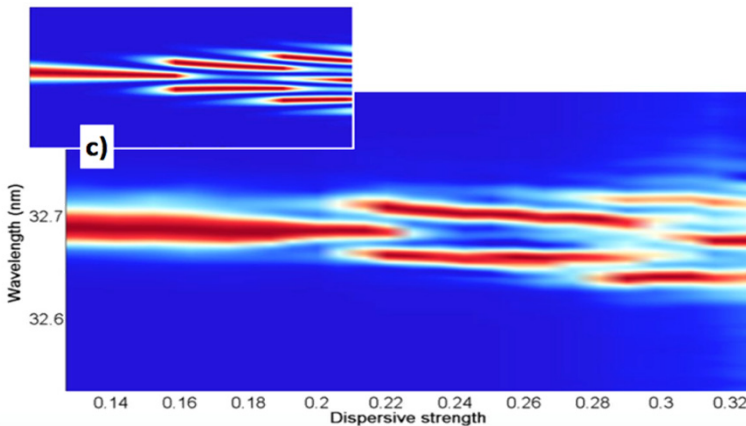
Spectro-temporal shaping (results)

Properties of spectra modulation are strongly dependent on the FEL chirp and can be predicted from theory.



With a significant **positive chirp** the **temporal-spectral** mapping allows to **see the bunching modulation**.

Reducing the **positive chirp** changes the spectral properties.

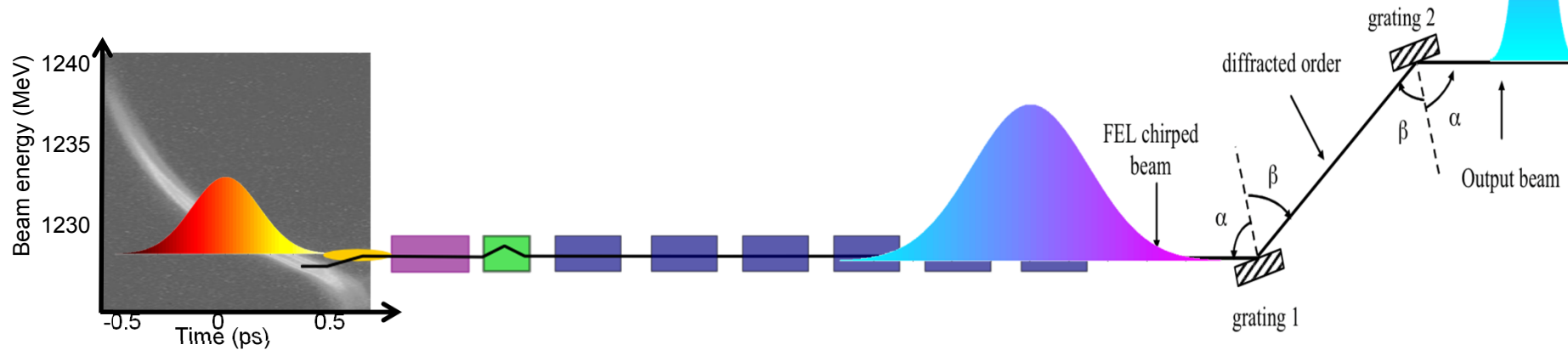


With slightly **negative** chirp of the seed (compensating the e-beam chirp) the FEL response is completely different and one can identify **signatures of Fourier limit FEL pulses**.

Chirped pulsed amplification (CPA)

Another way to **control** the FEL pulse length is **CPA**.

A **coherent light** pulse characterized by a **linear dependence** of the **wavelength** along the pulse (**chirp**) can be **compressed** with dispersive elements.



- A **chirped e-beam** is used in combination with a **chirped seed laser** to create **FEL** pulses with **time-wavelength dependence** and higher intensity as the number of electrons participating to the FEL process increases.
- **Chirped FEL** pulses are sent to a **compressor** based on **double grating**.

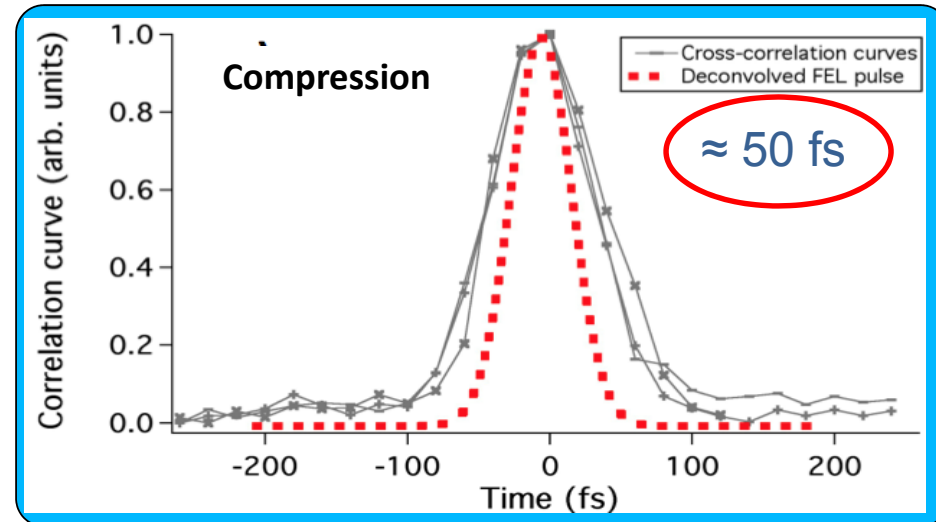
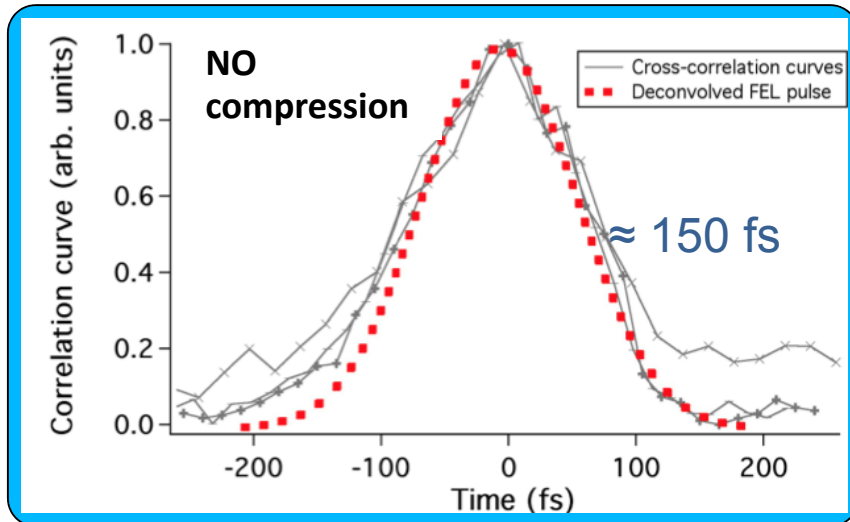
CPA in FEL, first demonstration

The standard FERMI **seed laser** has been **stretched** to **290 fs** (FWHM).

With the **nominal** seed laser **bandwidth** (0.9 nm FWHM) this seed has a **significant chirp** and is about **a factor 3** far from the **Fourier limit**.

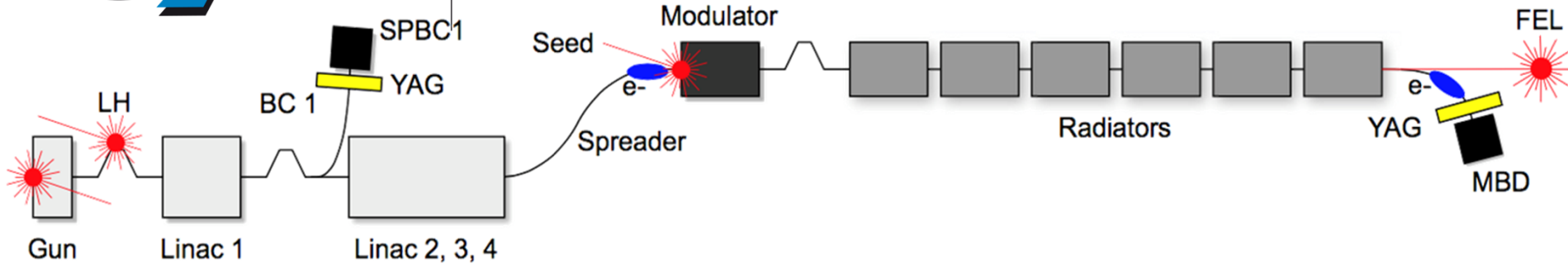
FERMI FEL-1 has been operated at **37 nm** corresponding to one of the **good wavelength** for the **available compressor**.

FEL **pulse length** has been **measured** for different **settings** of the optical **compressor**.



The demonstrated capability of **compressing** FEL pulses is an **indication** of the **high degree** of coherence of the FERMI FEL pulses and open the way to new possibilities for very short pulses.

Two frequencies seeding



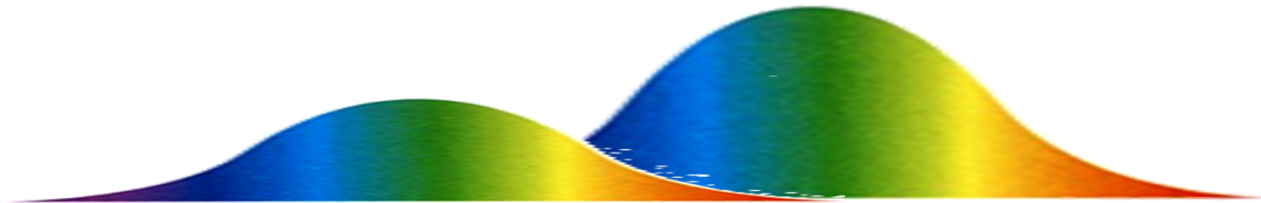
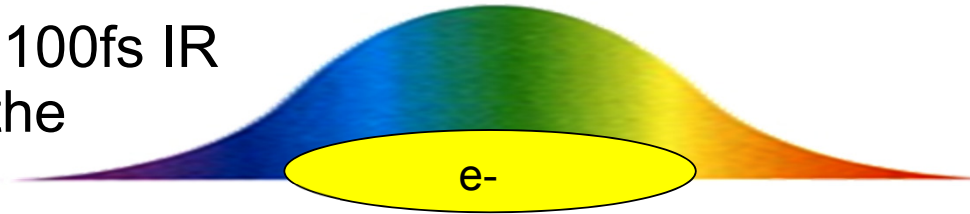
Instead of seeding with two wavelengths at the end of the accelerator we can seed the beam at different locations along the accelerator.

Most FELs already have a second seeding section in the low energy part (laser heater), but induced modulation is washed out.

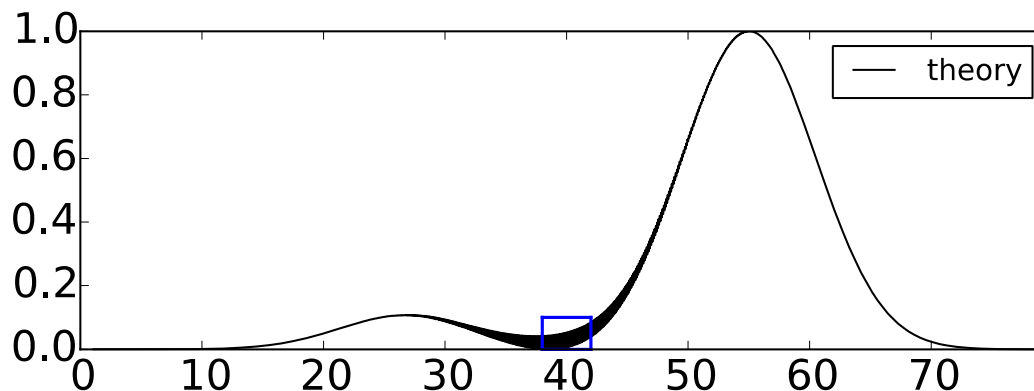
Using the chirped pulse beating technique one can seed in LH a modulation wavelength that survives or is amplified in the accelerator and is then interfering with the final seeding.

Laser heater beating

LH laser pulses are produced from 100fs IR pulses that are stretched to match the electron beam pulse length ~ 10 ps.



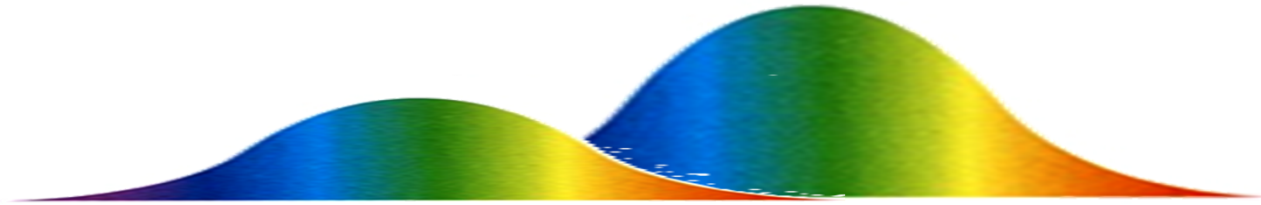
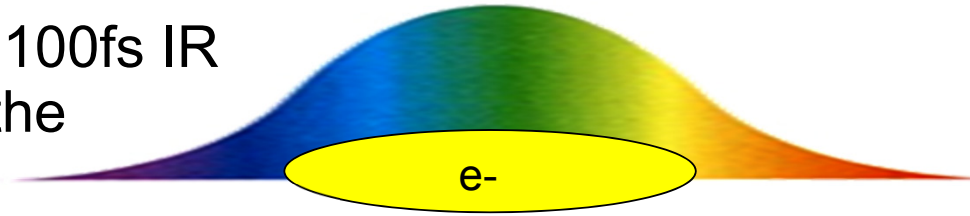
From the **superposition** of **two** laser **pulses** with a small **delay** we can **create** a **beating** of the **laser amplitude** due to the **superposition** of **two** fields with slightly **different** wavelength.



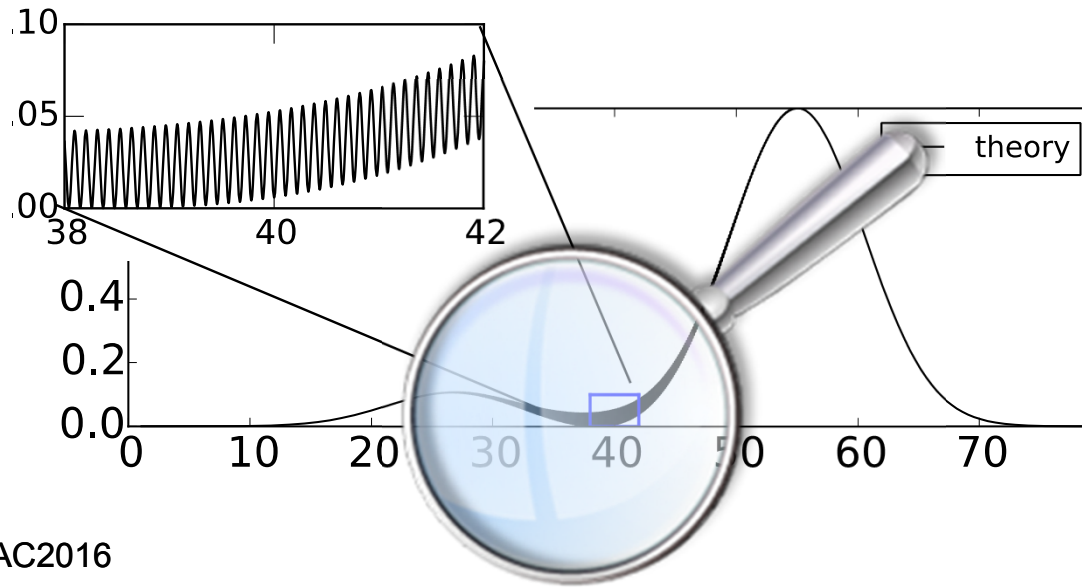
In our configuration, **beating** wavelength is $\sim 30 \mu\text{m}$ and can be **controlled** by acting on the laser **chirp** or pulse **separation**.

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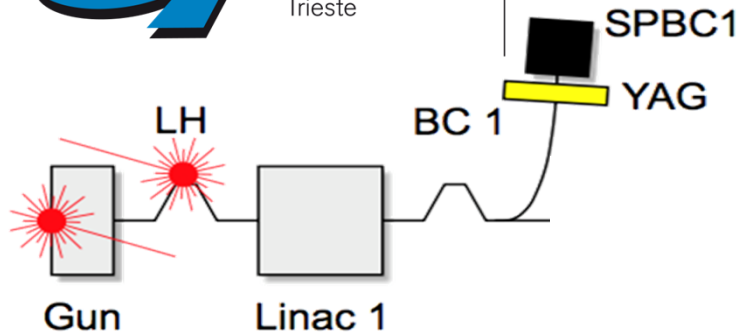


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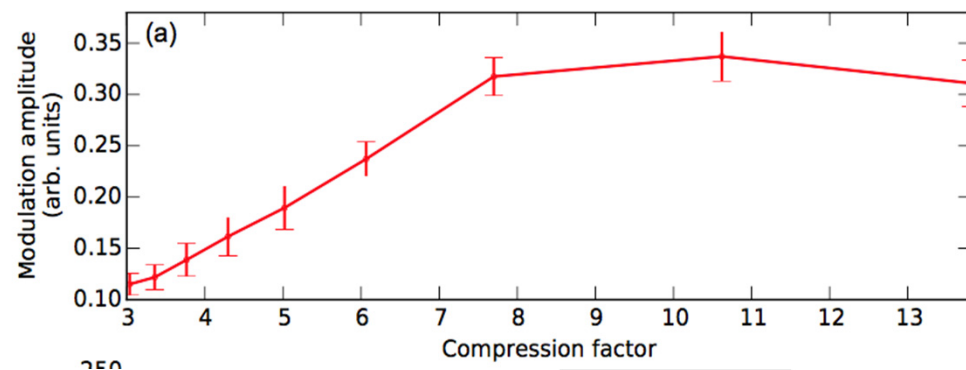
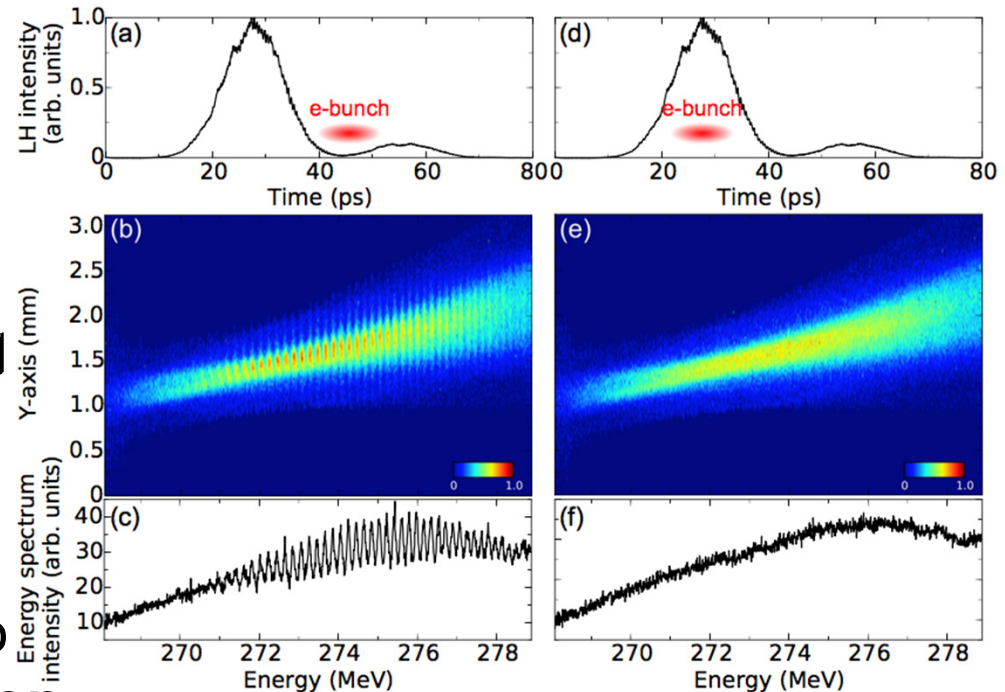
Beating induced modulation to e-beam



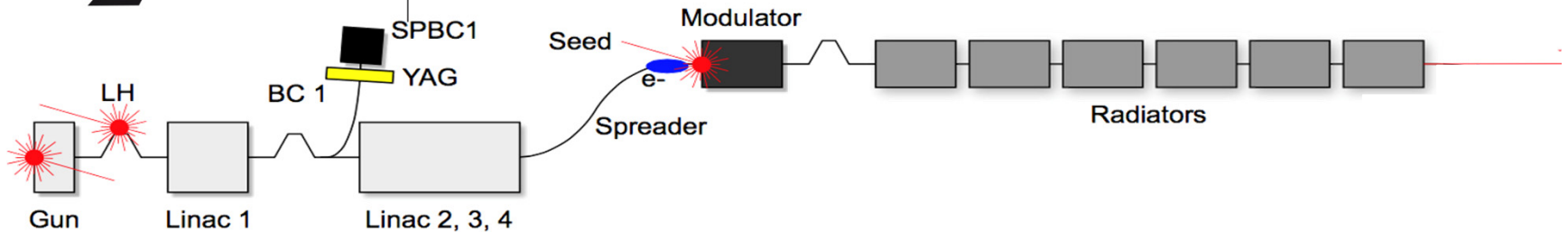
Placing the e-beam in the **beating** region creates **energy spread modulation**.

In the accelerator energy spread modulation can be **converted** into **energy and/or density modulation**.

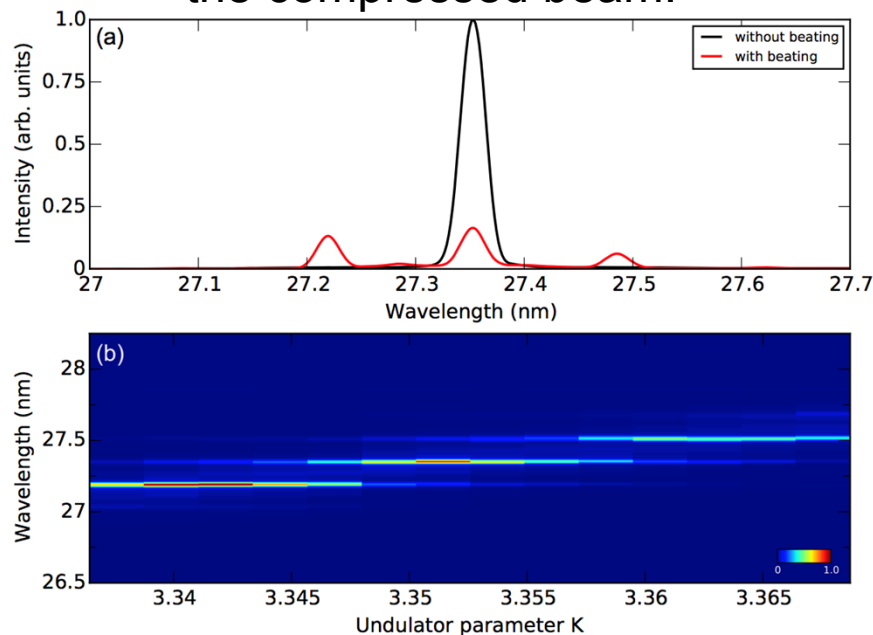
Modulation amplitude depends on the **compression** suggesting that **μB instability** is playing a role and amplifies the initial modulation.



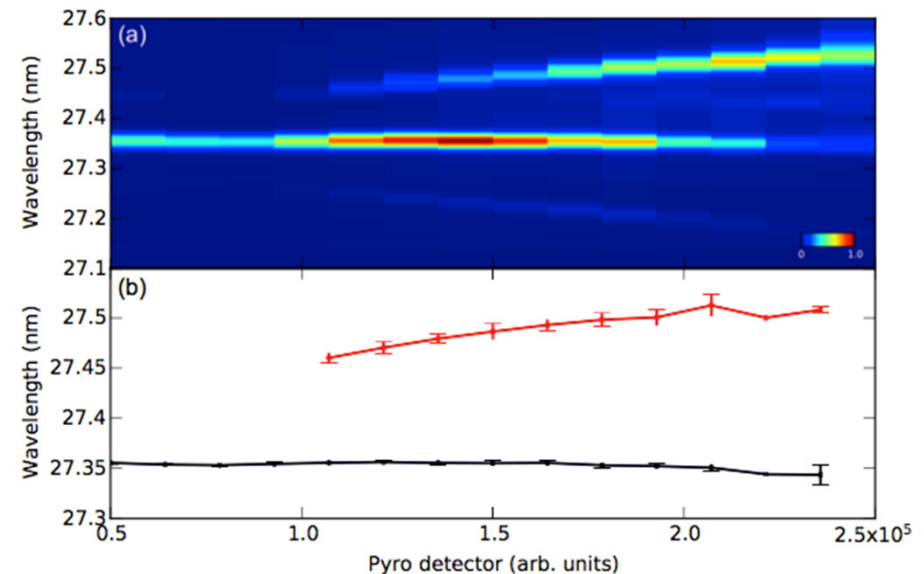
Beating induced modulation to e-beam



Also the seeded **FEL spectra** show **signatures** of coherent **modulation** in the **e-beam**. Sidebands distance matches the estimated periodicity for the compressed beam.



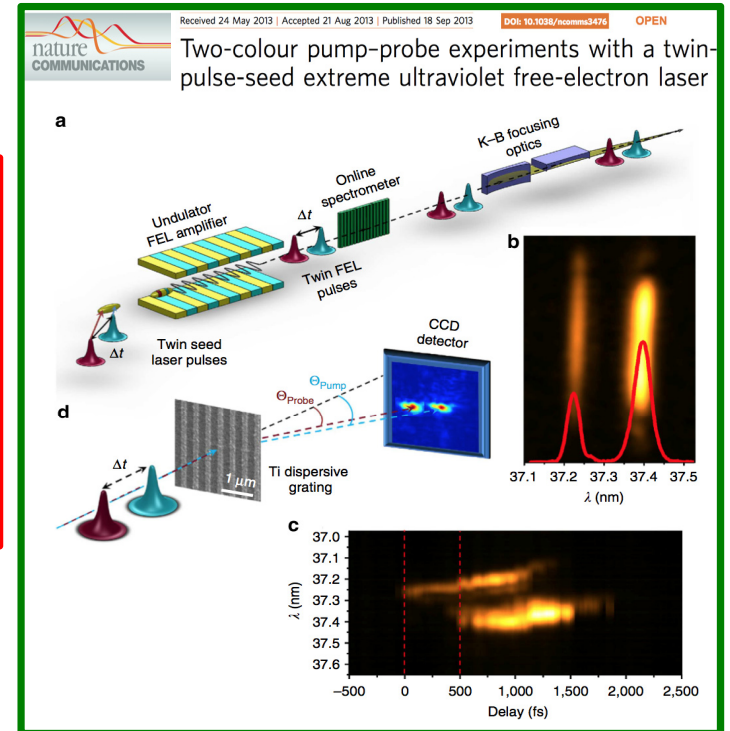
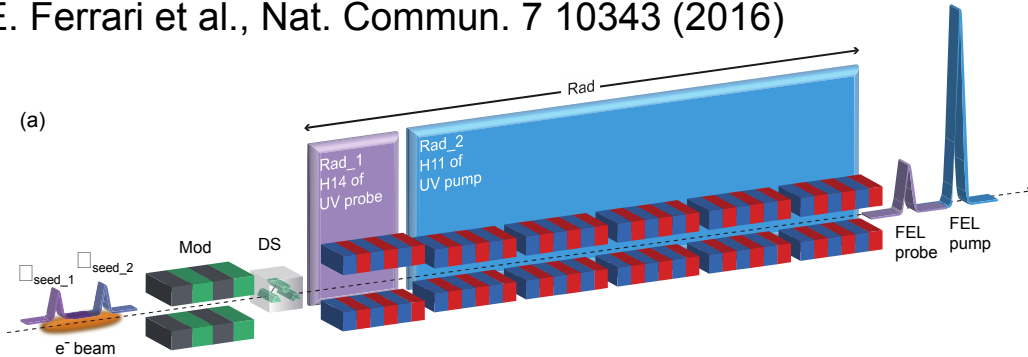
The sideband central wavelength changes with the compression (i.e., the periodicity of the modulation):
Tunable FEL source without tunable seed laser



Two color seeded FEL

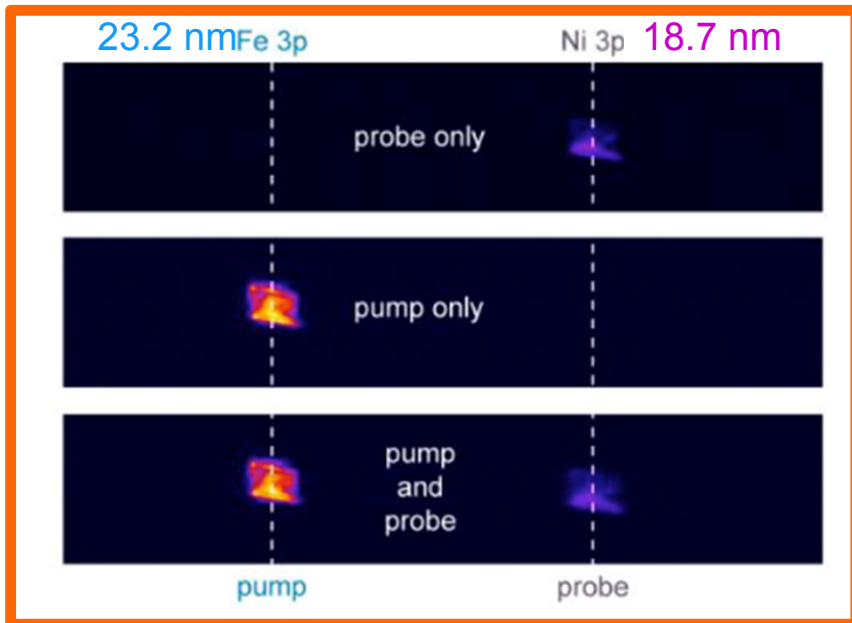
Various **schemes** for **two color FEL** pulses have been **implemented** and provided to the users.

E. Ferrari et al., Nat. Commun. 7 10343 (2016)



FEL pump - FEL probe experiments

- Different **diffraction angles** to have distinguishable signals
- **Selectively excite a resonance** in an atom and resonantly probe another atom bound to the excited one



Conclusions

- ✓ The seeded FEL radiation inherits the coherence properties of the seed laser, mediated by the electron beam.
- ✓ The use of the external seed allows control of critical FEL properties:
 - Wavelength, bandwidth;
 - Pulse length and pulse separation;
 - Phase and chirp;
 - ...
- ✓ The high degree of coherence and control capabilities open the possibility of “engineer” the FEL pulses specifically for the experiment one wants to perform (it’s a laser!!!)

Acknowledgements

- ✓ FERMI commissioning, laser and PADReS teams.
- ✓ FERMI beamline scientists (LDM, DiProI, TIMEX-TIMER) and their users.





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Thank you!



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