Diffusion measurement from transverse echoes Yuan Shen Li, Carleton College Advisor: Tanaji Sen, Fermilab

Introduction

- Measuring and managing diffusion is crucial in modern "intensity frontier" machines, where nonlinear phenomena, e.g. intrabeam scattering and space charge effects, can significantly increase emittance over time.
- Traditional methods to measure diffusion, e.g. beam scraping, take up to hours to complete. The transverse echo technique will require minutes or less.
- The echo displays high amplitude sensitivity to small phase space perturbations, making it an ideal tool to probe weak diffusion.
- Simultaneously, we need amplitude-boosting techniques to counter strong diffusion (e.g. space charge effects), so that the echo signal remains measurable. In this study, we develop theory and simulation to:

Diffusion

- Linear diffusion model simulated by dipole noise.
- Echo amplitude becomes attenuated with diffusion.
- $\frac{\partial \psi}{\partial t} = \frac{\partial}{\partial J} \left(D(J) \frac{\partial \psi}{\partial J} \right), \text{ where } D(J) = D_0 + D_1 \left(\frac{J}{J_0} \right)$ $D_1 = \pi \epsilon_0 \frac{\mathrm{d}\epsilon}{\mathrm{d}t} \qquad \tau_{\max} = \left(\frac{16}{3}\omega'^2 D_1\right)^{-1/3}$
- We directly measure diffusion coefficient by tracking emittance increase over a large number of turns. Results agree excellently with theory.
- Simulation results also demonstrate predicted relationship between echo amplitude and relevant parameters (below).



- Explore the behavior of transverse echoes under diffusion.
- Investigate pulsed quadrupoles as a method to boost echo amplitude.
- Provide recommendations for the planned beam echo measurement system in the future IOTA storage ring at Fermilab.

Echo: Theory and Simulation

Theory

- The transverse echo is a recoherence of the beam distribution, following phase decoherence due to nonlinear ring elements (e.g. octupoles).
- It shows up on the BPM as an oscillation of the beam centroid, some time after an initial disturbance (e.g. dipole kick). $a_0 = \theta q \sqrt{\beta \beta_k} \omega' J_0 \tau$

• Typical echo sequence:

- At t = 0, apply one-turn dipole kick θ . \bullet
- At t = τ , apply one-turn quadrupole kick *q*. \bullet
- Near time 2τ , the echo signal appears on the BPM.
- The amplitude of the echo is dependent on ring parameters. It is also extremely sensitive to diffusion. (Refer to equations above.)
- Key assumptions:
 - Both dipole and quad kicks are weak (compared to beam spread). The timing of quad kick τ is much greater than decoherence time.

Pulsed quadrupoles

- Based on gradient echoes in NMR.
- A single quad kick introduces a small, position-dependent ΔJ to the particle distribution. With linear detuning, this leads to particles "clumping" together in phase space at time 2τ .
- Pulsed kicks apply a sequence of small ΔJ 's



Simulation

- Simulation written in C, with analysis performed in *Mathematica*.
- Machine parameters based on 2005 RHIC experiment.
- Simulation options include adjustable ring elements, variable starting distribution, variable diffusion model, pulsed quadrupoles and injection oscillation.
- Simulation results agree well with theory.





Symbol	Value
N_{part}	20000 to 50000
ϵ_0	$1.5 imes 10^{-7}$
σ_x	3×10^{-3}
$ u_0$	0.245
$eta_{m{k}},eta_{m{k}},eta_{m{q}}$	10, 10, 10
μ	10^{-4} to 10^{-3}
heta	10^{-4} to 10^{-3}
q	0.01 to 0.05
τ (or τ_q)	400 to 1000
	Symbol N_{part} ϵ_0 σ_x ν_0 β, β_k, β_q μ θ q τ (or τ_q)

- Above: Typical simulation parameters, based on 2005 RHIC Above left: Evolution of simulated beam centroid over time, *Left:* Phase space portraits showing the dipole kick and phase decoherence, followed by the quadrupole kick and
- theory prediction in red. Saturation starts to set in at high

- that amplify each other, resulting in a tighter "clump" in phase space.
- Optimal sequence highly dependent on fractional tune. We investigated several possible sequences.
- Maximum echo amplification close to 100% (up to saturation point).



Conclusions and Further Work

Key findings: Consistent measurement of diffusion coefficient based on τ_{max} ; echo amplitude boost by up to 100% using pulsed quads; optimal sequence depends on fractional tune; pulsed sequence of single polarity can be just as effective.

START

 $v_0 = 0.245$

Schematic of ring used in simulation.

where $\alpha_1 = 1 + \frac{2}{3} \left(\frac{D_1 \tau_q}{J_0^2} \right) (\omega' J_0 \tau)^2$

 $a_{\rm difn} = \frac{\theta q \sqrt{\beta \beta_k} \omega' J_0 \tau_q}{\alpha^3}$

- Some further questions:
 - What is the optimum pulse sequence for a given fractional tune?
 - Echo amplitude saturation observed empirically at A \approx 0.4. How do we explain it? Is it possible to surpass this limit?
 - How will echo dynamics change in 2D? Any coupling effects?

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