VACUUM STUDIES FOR THE EIC HADRON RING

Time Dependent Residual Gas Density Beam Lifetime and Emittance Growth due to Interactions with Residual Gas

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Electron Ion Collider – EIC at BNL

BROOKHAVEN



EIC Hadron Ring (HR) cold and warm sections



(I) Beamline vacuum level and beam quality

- What is the highest allowable vacuum level above which beam quality is compromised?
- Which beam parameter (emittance, intensity, etc.) sets this threshold?

Work these questions in parallel to evaluation of the expected residual gas composition and pressure in the warm and cold sections of the EIC hadron rings vacuum chamber.

 \Rightarrow Goal of this work: prepare, debug calculation file to evaluate beam lifetime, emittance growth due to residual gas when we know expected residual gas composition and pressure

Beam loss due to interaction with residual gas

41 – 110 GeV/u ¹⁹⁷Au⁷⁹⁺ beams

GOAL	Aim for hundreds of hours of beam lifetime ($ au_{gas}$ ~ 100 h)
ASSUMPTIONS	$W^{(w)} \sim 0.1; W^{(c)} \sim 0.9$ (EIC hadron lattice)
RESULTS	$p^{(w)} \leq$ 9e-11 Torr; $p^{(c)} \leq$ 3e-11 Torr
$\Rightarrow au_{gas} \geq 110$ h for	• EIC 41-110 GeV ¹⁹⁷ Au ⁷⁹⁺ beams and studied gas compositions



MOST RESTRICTIVE SCENARIO DUE TO LARGER CROSS-SECTION OF GOLD NUCLEUS

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Emittance growth due to residual gas

41 GeV proton beams, vertical emittance

Using Rhoades-Brown

GOALAnalogously to beam lifetime case, aim for hundreds of hoursASSUMPTIONS $W^{(w)} \sim 0.1; W^{(c)} \sim 0.9$ (EIC hadron lattice) $p^{(w)} \leq 1.5e-11$ Torr; $p^{(c)} \leq 7e-12$ Torr

 $\Rightarrow \tau_{gas} \ge$ 110 h for EIC 41 GeV proton beams and studied gas compositions



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Lower energy beams suffer faster emittance growth.

Outlook (I)

VACUUM LEVEL REQUIREMENT FOR THE EIC HADRON RING

Taking most stringent residual gas composition (with presence of C, O):

							<u>Beam loss</u>		<u>Emittance growth</u> ☆	
Species Z	7	A	E	β	βγ	$\mathcal{E}_{n,x}, \ \mathcal{E}_{n,y}$	p ^(w)	p ^(c)	p ^(w)	p ^(c)
	2		(GeV)	(m)		(mm mrad)	(Torr)	(Torr)	(Torr)	(Torr)
proton 1	1	1	275	20.24	294	5.2 / <u>0.47</u>	3e-9	9e-10	6e-11	4.5e-11
	Ι		41		45	1.9/ <u>0.45</u>			<u>1.5e-11</u>	<u>7e-12</u>
gold	70	197	110	20.34	118	5.1/ <u>0.7</u>	<u>9e-11</u>	<u>3e-11</u>	*5e-10	*1.5e-10
	79		41		45	3.0/ <u>0.3</u>			*	*

*Stochastic cooling available. *Computed for vertical plane.

- Beam lifetime is almost independent on beam energy but does greatly decrease with the atomic weight A of the beam species. The most demanding vacuum level conditions are found for gold beams.
- Emittance growth is more severe for low energy beams. The EIC will be equipped with stochastic cooling and thus, we focus on proton beams that, independently of stochastic cooling or not, set the most demanding vacuum level to mitigate emittance growth.

(II) EIC HSR Cryogenic Beam Screen Vacuum Assessment

- (1) Initial conditions Ideal gas law and bulk gas thermodynamics suggest cold bore arcs with a-C coating will have very low initial monolayer H2 coverage
- (2) Time Dependence: pressure (gas density) will increase over time IAW the adsorption isotherm & magnitude of gas sources as gas evolves and propagates along the length of the beamtube (sources, interconnect thermal desorption, warm-to-cold, leaks etc.)
 - a-C Carbon in limited experiments exhibits high capacity (~100x sSST/Cu) and high thermal desorption temperature
 - So far only need to consider H2. The following may require adding CO partial pressure to vacuum analysis
 - Interconnect Bellows thermal (analysis performed for range of conditions)
 - Beam stimulated desorption
 - Beam screen thermal contact (higher temperature >10K)

(1) Initial Conditions at Cooldown Recent RHIC performance





P(RT) ~1e-6 torr -> gas density start cooldown -> surface coverage -> all condensed at cooldown -> θ after cooldown

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(1) Initial Conditions: Adsorbed Gas RHIC & EIC HR Coldbore

F. Prepumping in cold sections

At high proton beam intensities an increase in the gas density in the cold sections was observed (Fig. 24). The cold sections initially relied on cryopumping, and had been evacuated before cooldown with mobile turbo pumps to about 10^{-1} Torr only in some areas. The surface density σ of gas molecules after cooldown is

$$=\frac{Pr}{2kT},$$
(20)

Apply more rigorous 1-D wave propagation method for more accurate prediction of cold bore gas density versus time.

With existing ion pumps: pre-pumping << 1 monolayer, sorption dominant

Adsorption + Condensation

where *P* and *T* are the pressure and temperature before cooldown, respectively, *r* the beam pipe radius, and *k* the Boltzmann constant. For a flat surface, a monolayer has of order 10^{19} molecules/m² [77], and a pressure of 10^{-1} Torr before cooldown will result in about 5 monolayers. Near a warm-cold transition there can be many more monolayers.

Smooth SST Surface some EIC SST BS T>4.2K & aC Surface... PV=NkBT H2 migration from warm beamline to coldbore... Q=CP

 σ

Pressure before cool down	to form on	e monolaye	er &			H2 diffusion from war	m bore to fo	orm one mo	nolayer
σ/σm for				1.E-04	torr warm				
	RHIC ARC	HSR ARC					NINC ANC	HSI AIC	
A (cm^2/cm)	22	32 R	RHIC + scr	een flats		P (Torr)	1.E-10	1.E-10	/
V (cm^3/cm)	37	37				C (I/s)	500	500	
One monolayer						Q (Torr.l/s)	5.E-08	5.E-08	
molecule/cm^2	5.E+14	5.E+16a	C σ~100>	(> SST (CE	RN ColdEx Ref)	Q (Torr.l / 7 mos)	9.E-01	9.E-01	
molecule/cm	1.E+16	2.E+18				O(moloculos/wr)	2 E±10	2 E+10	
Torr-cm^3 (PV)	4.E-01	5.E+01				Q (molecules/ yr)	5.2719	5.6719	/
Torr (one mono)	9.7E-03	1.4E+00				A (cm ²) each end	5.E+04	5.E+02	om 1 year
σ/σm=	1.0E-02	7.1E-05	¥			L (m) - Arcs each end	23.93	0.16	om 1 year
θ(molecule/cm^2)=	5.E+12	4.E+12	CERN (COLDEX >	<mark>>10x</mark>	Assuming no sorption	pump and r	no RW bear	n heating
									and the second se

Electron Ion Collider – EIC at BNL

(1) Initial Conditions: Adsorbed Gas RHIC & EIC HR Coldbore

Vapor Pressure: Temperature and Surface Coverage Dependence



Must maintain << 1 monolayer H2 coverage

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(2) 1-D Time Dependent Pressure Profile RHIC Basis Arc Interconnect



(2) General 1-D Solution Method

- Calculate Steady State Pressure Profile from the gas source (interconnect X=0) to the leading edge of the propagating wave front (Xf) based on the interconnect outgassing rate and cold bore conductance at Xf (Independent of adsorption)
- Using a modified adsorption Isotherm equation, and knowing the pressure at all X from X=0 to Xf, calculate the quantity of adsorbed gas from X= 0 to Xf
- The time to reach this pressure profile condition is the total adsorbed gas divided by the outgassing rate
- The cold bore <u>gas density</u> is taken as the average for the pressure profile, corrected for temperature.
- Analysis performed for
 - Various outgassing rates (Interconnect design and temperature)
 - Various Adsorption Isotherm conditions (a-C temperature)

(2) 1-D Solution Method Development

$\mathbf{Qt} = \mathbf{k_2} \int_0^{xf} \boldsymbol{\Theta}(x) dx$	The total gas pumped by the magnet cold bore from the source (x=0) to the gas wave front (Xf) via integration of the Adsorption Isotherm Relation				
$Q_t = k_2 \int_0^{x_f} e^{-B(RTln(P/P_0))^2} dx$	Applying the well tested analytic expression; P. A. Redhead, J. P. Hobson, & E. Y. Kornelsen, The Physical Basis of Ultrahigh Vacuum (Chapman and Hall, London, I 968).				
$k_2 = \pi \sigma_m DrT/N_o$	σ_m = monolayer coverage at Temperature T, r surface roughness factor, N _o is Avogadro's Number				
$\Theta = k_1 P^m$ $Qt = k_1 k_2 \int_0^{xf} [P(x)]^{\prime\prime\prime} dx$	The adsorption isotherm relation, power law approximation for integration purposes				
$P(x) = (Q/Ca) \cdot [1 + .75(X_F - X)D]$	Cold bore pressure as a function of gas load and beam tube diameter				
$O_{t} = k k \int_{0}^{xf} (O_{t}(C_{a})^{m} [1 + 75(x_{t} - x)_{t})^{m} dx$	Substituting for $P(X)$				
$Q(1 - K_1 K_2)_0 (Q/Cd) [1 + .75(XF - X)D] dx$					
$Qt = -k_1 k_2 {\binom{Q}{C_a}}^m \frac{D}{.75(m+1)} \left[1 - \left(1 + .75X_f / D \right)^{m+1} \right]$	The total gas pumped by the magnet cold bore from the source (0) to the gas wave front (Xf)				
$Qt = -k_1 k_2 (Q/C_a)^m \frac{D}{.75(m+1)} \left[1 - (1 + .75X_f/D)^{m+1} \right]$ t = Qt/q	The total gas pumped by the magnet cold bore from the source (0) to the gas wave front (Xf) The time for the wave to propagate a distance Xf is the total adsorbed gas (Qt) divided by the outgassing rate (q)				
$Qt = -k_1 k_2 (Q/c_a)^m \frac{D}{.75(m+1)} \Big[1 - (1 + .75X_f/D)^{m+1} \Big]$ t = Qt/q $\frac{dx_f}{dt} = \frac{C_a^m Q^{1-m}}{k_1 k_2} (1 + .75X_f/D)^{-m}$	The total gas pumped by the magnet cold bore from the source (0) to the gas wave front (Xf) The time for the wave to propagate a distance Xf is the total adsorbed gas (Qt) divided by the outgassing rate (q) Differentiated and manipulated to yield wave front propagation velocity				

(2) EIC Time Dependent Pressure Profile



(2) Adsorption Isotherm Development

The relation between the amount of gas adsorbed at a point x i.e., a(x) in molecules/cm² and the pressure P(x) at that point is given by the adsorption isotherm. Adapting the well tested analytic expression¹: **InO = - B(RT In PIP₀)²** where $\Theta = \sigma/\sigma_m$ is the ratio of surface coverage to full monolayer coverage in molecules/cm² and P_0 being the vapor pressure of the adsorbed species at the pipe temperature T. The values for H2 on smooth bare SST and those estimated for a-C coated SST are given

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FIG. 4. Hydrogen adsorption isotherm at 4.2 K.

Modeling an appropriate adsorption isotherm equation

shown in Fig. 2. Signs of surface saturation were recorded with a coverage of $> 4 \cdot 10^{17} \text{ H}_2/\text{cm}^2$ and $> 2 \cdot 10^{17} \text{ H}_2/\text{cm}^2$ respectively, *i.e.* two orders of magnitude higher than the monolayer capacity of metallic surfaces like Cu or SS. This result points to a porous surface morphology of a-C.



Figure 2: a-C H₂ adsorption isotherms at 6.5 K and 11.1 K.

(2) 1-D Solution Results

INTERCONNECT THERMAL					
Т	(К)	50	20	10	5
q(H2)	(torr-l/(s-cm^2))	1E-11	1E-12	1E-13	1E-14
Q = q(Interconnect Surface Area)	(torr-l/s)	5.18E-09	5.18E-10	5.18E-11	5.18E-12
Pressure Profile func	tion of outgassing r	ate and beam tub	e conductance		
PRESSURE PROFILE $P(x) = (Q/Ca) \cdot [1+.75($	<i>XF</i> – X)D]	Results when H2 wa	ve has propagated	d to dipole midpoir	nt
Pf at X=Xf=5m	(torr)	2.11E-11	2.11E-12	2.11E-13	2.11E-14
P0 at X=0 when Xf=5m	(torr)	5.47E-08	5.47E-09	5.47E-10	5.47E-11
Pavg = (Pf+PO)/2	(torr)	2.74E-08	2.74E-09	2.74E-10	2.74E-11
Pavg => target	(torr)	1.82E-10	1.82E-10	1.82E-10	1.82E-10
ρavg => target	(H2/cm^3)	2.00E+06	2.00E+06	2.00E+06	2.00E+06
<pre>pavg => PV=NkBT => (N/V)=P/kB*T (Xf=5m)</pre>	(H2/cm^3)	3.01E+08	3.01E+07	3.01E+06	3.01E+05
Xf at target ρ (Q dependent only)	(cm)	40.0	128.2	407.0	500.0
ISOTHERM PARAMETERS		H-W (base)	H-W (a-C)	11K	6.5K
r Propagation time a function of		1	100	1	1
σm the adsorption isotherm	(H2/cm^2)	2E+15	2E+15	2E+17	5E+17
Ot = k k (Q/L)	$m D \begin{bmatrix} 1 \\ 1 \end{bmatrix}$	$(75 \times (0)^{m+1}]$			
WAVE PROPAGATION $QI = R_1 R_2 (7c_a)$	$\frac{1}{.75(m+1)} \begin{bmatrix} 1 - (1) \end{bmatrix}$	$+.75\Lambda_f/D$	k₂ = πσmDr	RT/N₀	t = Qad/Q
Qad (H-W, RHIC baseline, r=1)	(torr-l)	2.73E-04	2.69E-04	2.65E-04	2.61E-04
Time (t) for Xf=5m	(yr)	0.002	0.016	0.162	1.598
time when Xf at target ρ	(yr)	0.0001	0.004	0.13	1.60
Qad (H-W: r=100)	(torr-l)	2.05E-02	2.02E-02	1.99E-02	1.96E-02
Time (t) for Xf=5m	(yr)	0.13	1.23	12.16	119.86
time when Xf at target ρ	(yr)	0.01	0.31	9.89	119.86
Qad (σCOLDEX: 11K, r=1)	(torr-l)	2.73E-02	2.69E-02	2.65E-02	2.61E-02
Time (t) for Xf=5m	(yr)	0.17	1.65	16.22	159.82
time when Xf at target ρ	(yr)	0.01	0.42	13.18	159.82
Qad (σCOLDEX, 6.5K, r=1)	(torr-l)	6.82E-02	6.73E-02	6.63E-02	6.53F-02
Time (t) for Xf=5m	(yr)	0.42	4.11	40.54	399.55
time when Xf at target ρ	(yr)	0.03	1.05	32.96	399.55

(2) 1-D Method Parameter Assumptions

• Interconnect:

- <u>Temperature</u>: Unknown pending final thermal analysis, Assumed range between 4 and 50K
- <u>Outgassing rate</u>: estimated as a function of interconnect temperature, validation needed as there is no available rates at the cryogenic temperature range considered. Dependent on specific material treatment and conditions

Beam tube/Beam screen

- <u>Temperature</u>: 6.5 and 11K based on available COLDEX adsorption isotherm data
- <u>Cryogenic Outgassing rate</u>: assumed 0 as in the case of RHIC model at 4K. A value may be needed if screen is higher temperature. Needs validation

- <u>Adsorption Capacity:</u> Range of values applied. Developed EIC a-C coating needs validation
 - Assumed COLDEX values at 6.5K and 11K
 - Initial condition θ very small ~ 0 for wave propagation analysis. Room temperature a-C coating characteristics required to validate.

Outlook II

- H2 Propagation magnitude, time and avg gas density analyzed for various interconnect and beam screen conditions
 - Pressure Profile approaches the pressure profile absent any adsorption over time.
 - Surface and adsorption isotherm properties only serve to adjust the time to establish that profile.

• a-C cryogenic vacuum properties in published literature are promising

- High H2 adsorption capacity (~100x SST)
- Higher thermal desorption temperatures compared to SST
- EIC a-C film cryogenic vacuum properties require validation
- Carbon Fiber Cryosorber (& screen cross-section design) and perforations remain a consideration in beam screen design

Initial a-C adsorption coverage depends on pressure before cooldown

- Assumed ~0 for 1-D analysis
- EIC a-C film ambient temperature vacuum properties require validation
- more ion pumps and longer pumpdown prior to cooldown remain considerations

Higher Confidence if conditions below are also achieved

Beam screen thermal analysis (thermal contact => BS Temperature ~5K)

- Interconnect thermal analysis (Temperature <10K)
- Beam stimulated desorption analysis: ion and e- dose rates

References (I)

[1] A. Chao et al., Handbook for Accelerator Science And Technology, Section 3.3.2: Beam and Luminosity Lifetime, Second Ed. (2013)

[2] M. Ferro-Luzzi, Beam-Gas Interactions, Proc. CAS 2017

[3] M. J. Rhoades-Brown and M. Harrison, *Vacuum Requirements for RHIC*, AD-RHIC-1-142 (1992)

[4] D. Trobojevic et al., *Beam Lifetime Dependence on the Beam-Gas Interactions in RHIC*, Proc. PAC'01: <u>https://accelconf.web.cern.ch/p01/PAPERS/RPAH125.PDF</u>

[5] Particle Data Group, *Hadronic Cross-Sections* (2019): <u>https://pdg.lbl.gov/2019/hadronic-xsections/hadron.html</u>

[6] J. D. Jackson, Classical Electrodynamics, Section 13.7 Elastic Scattering of Fast Particles by Atoms, Second Ed. (1975)

References (II)

[1] Time-dependent helium and hydrogen pressure profiles in a long, cryogenically cooled tube, pumped at periodic intervals J. Vac. Sci. Technol. A 11(4), Jul/Aug 1993 0734-2101/93/11(4)/1566/9/ 1993 American Vacuum Society, J.P. Hobson, National Vacuum Technologies Inc., Box 4160, Postal Station E, Ottawa, Ontario KIS 5B2, Canada, K. M. Welch, RHIC Project, Brookhaven National Laboratory, Associated Universities Inc., Upton, New York 11973 (Received 25 September 1992; accepted 7 December 1992)

[2] Measurements of the helium propagation at 4.4 *K* in a 480 m long stainless steel pipe H.C. Hseuh, RHIC Project, Brookhaven National Laboratory, Upton, New York, 11973, USA, and E. Wallen, MAX-lab, Box 118, S-22100 Lund, Sweden 1997

[3] Vacuum System Performance for the First Sextant Test of the Relativistic Heavy Ion Collider, AD/RHIC/RD-116, R. Davis, H. C. Hseuh, D. Pate, L. Smart, R. Todd, D. Weiss

[4] Hydrogen and Its Desorption in RHIC BNL-71032-2003-CP, International Workshop on Hydrogen in Materials & Vacuum. Systems JLAB, Newport News, VA, November 11-13, 2002H.C. Hseuh, *Collider-Accelerator Department Brookhaven National Laboratory Upton, NY 11973, USA*

[5] Capture Pump Technology: An Introduction, Kimo M. Welch, Pergamon Press, 1991

[6] Cryopumping and Vacuum Systems, Proceedings of the 2017 CERN–Accelerator– School course on Vacuum for Particle Accelerators, Glumslöv, (Sweden), Vincent Baglin CERN, Geneva, Switzerland