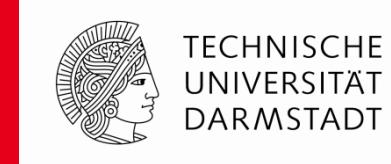


Challenges in Simulating

Beam Dynamics of

Dielectric Laser Acceleration



Uwe Niedermayer* for the ACHIP collaboration

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64289 Darmstadt, Germany

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*currently as guest at Ginzton Lab,
Stanford University,
Stanford, CA 94305, USA



ICAP'18 OCT 20 - 24, 2018
KEY WEST, FL, USA
CASA MARINA RESORT

13th International Computational
Accelerator Physics Conference - www.icap18.org

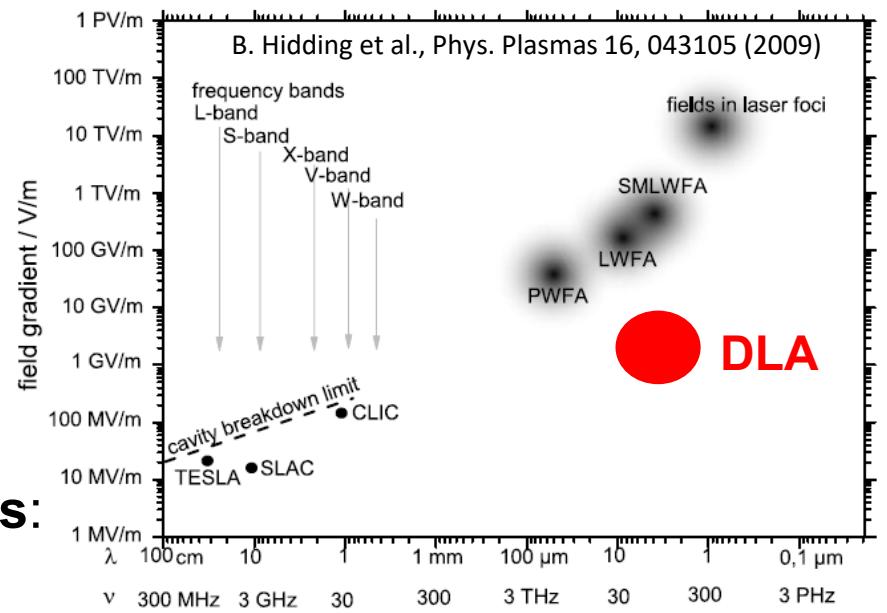
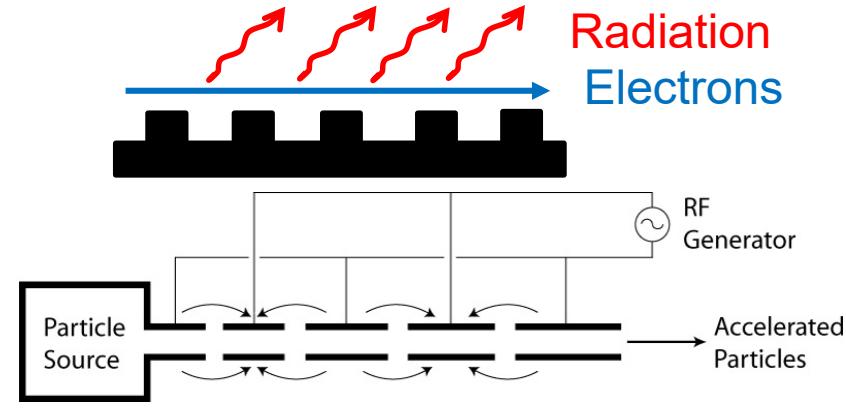
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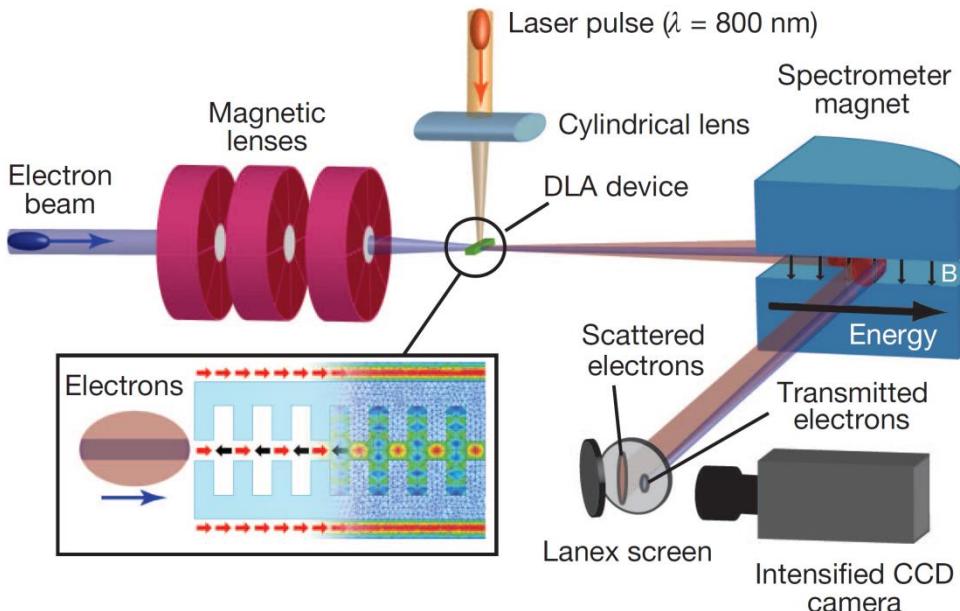
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- Current and future experiments

Dielectric Laser Accelerator (DLA) principle

- Idea is quite old
→ Inverse effects
 - Smith-Purcell (grating radiation)
 - Cherenkov (electrons superluminal in material)
- Same principle as **Wideroe-Linac** (Non- resonant)
- Dielectrics** can withstand fields up to **10GV/m**
- Gradients larger than 1GeV/m** (limited by breakdown)
- Recent technological improvements:
Laser pulses, micro-fabrication



Experiments with relativistic e-beams

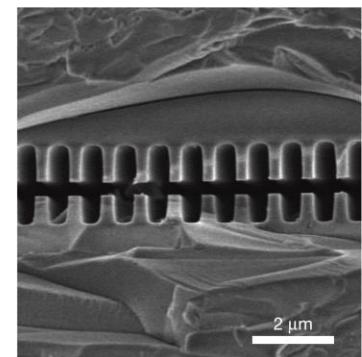
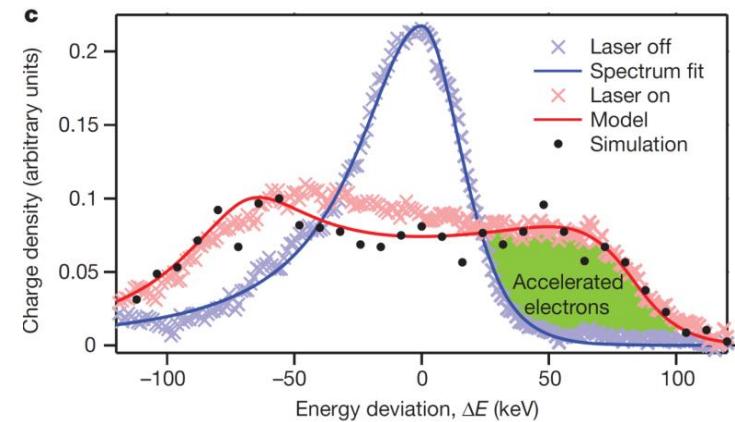


60 MeV SLAC NLCTA beam:

- E. Peralta et al. *Nature* 2013, **300 MeV/m**
- K. Wootton et al.: increased to **690 MeV/m**
(*Optics Letters* 2016)

6 MeV UCLA Pegasus beam:

- D. Cesar et al. *Nat Comm Phys.* 2018,
850 MeV/m with pulse front tilt



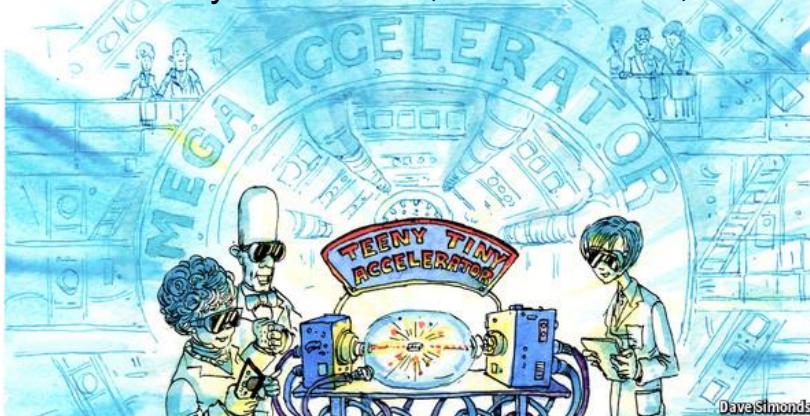
Accelerator on a Chip Intl. Program (ACHIP)

- funded by the Moore Foundation



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“Small really is beautiful”, *the economist*, 2013



“Make a chip that provides



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2015



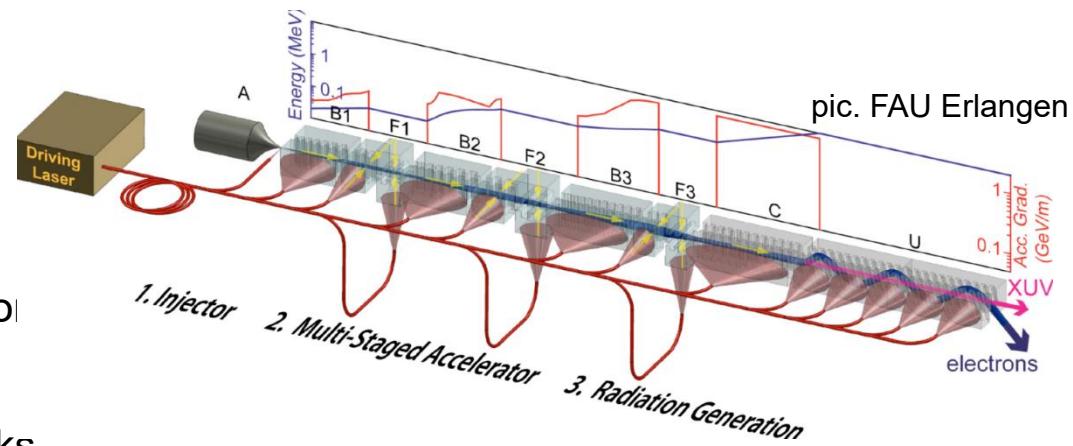
2017



Working groups:

...to work towards our dream laser accelerator

- Injectors
 - Electron sources
 - Sub-relativistic DLA
- Relativistic Acceleration
 - Large scale integration (Accelerator)
- Lasers and Laser Coupling
 - On-chip laser delivery tree-networks
- Simulations and Beam Dynamics
 - Beam dynamics schemes design
 - Large and small scale computation
 - Intensity effects
- Radiation Generation and Applications
- Integration (fit everything in a shoe-box)



Goal #1:
 $60\text{keV} \rightarrow 1\text{MeV}$
 on-chip accelerator



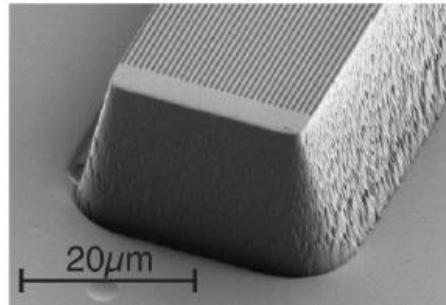
Sub-relativistic accelerators

FAU Erlangen:

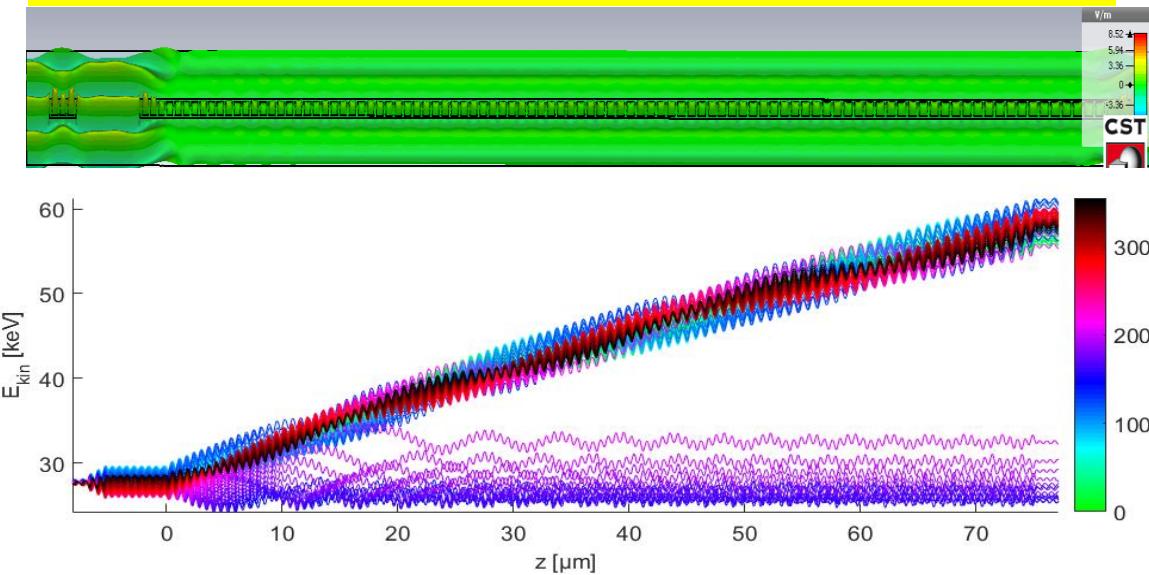
28 keV electrons
($v/c=0.32$)

25 MeV/m gradient
Single grating

Breuer et. al.
PRL 2013



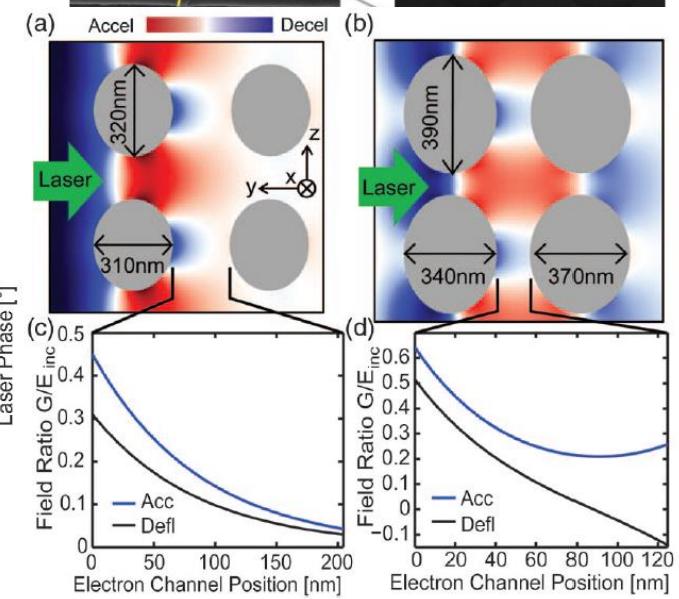
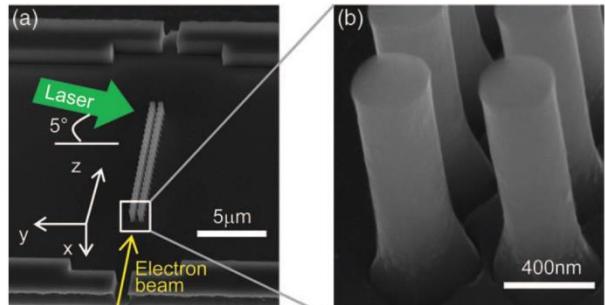
Inherent problem: dephasing due to velocity increase
Solution: chirped grating



U Niedermayer et al. J. Phys.: Conf. Ser. 874 012041 (2017).

22 October 2018 | TU Darmstadt | Fachbereich 18 | Institut Theorie Elektromagnetischer Felder | Uwe Niedermayer | 7

Stanford: 370MeV/m @ ~90keV
K.Leedle et al. Optics Letters, 2015



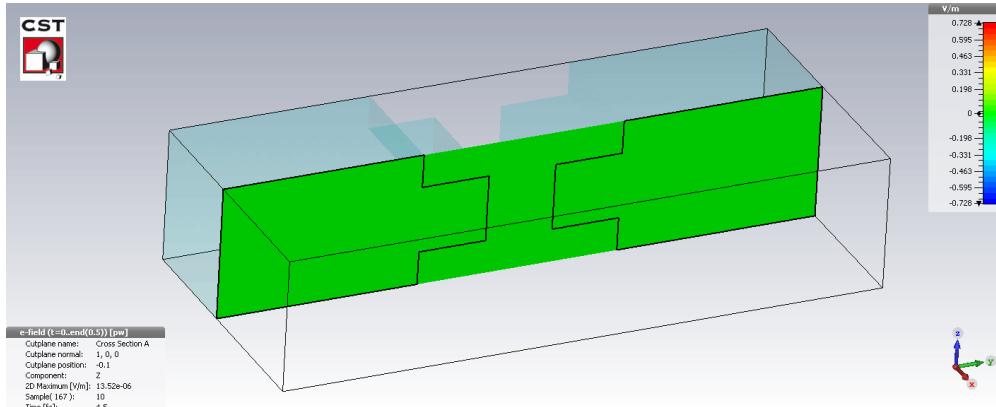
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Fields in a single DLA cell



Simulation not very challenging

- FD / FI TD
- FD FD
- FE FD

$$\begin{aligned}\Delta W(x, y; s) &= q \int_{-\lambda_g/2}^{\lambda_g/2} E_z(x, y, z; t = (z + s)/v) dz \\ &= q \int_{-\lambda_g/2}^{\lambda_g/2} \Re\{\underline{E}_z(x, y, z) e^{i\omega(z+s)/v}\} dz\end{aligned}$$

$$\lambda_g = m\beta\lambda_0$$

$$= q\lambda_g \Re \left\{ e^{2\pi i \frac{s}{\beta\lambda_0}} \underline{e}_m(x, y) \right\}$$

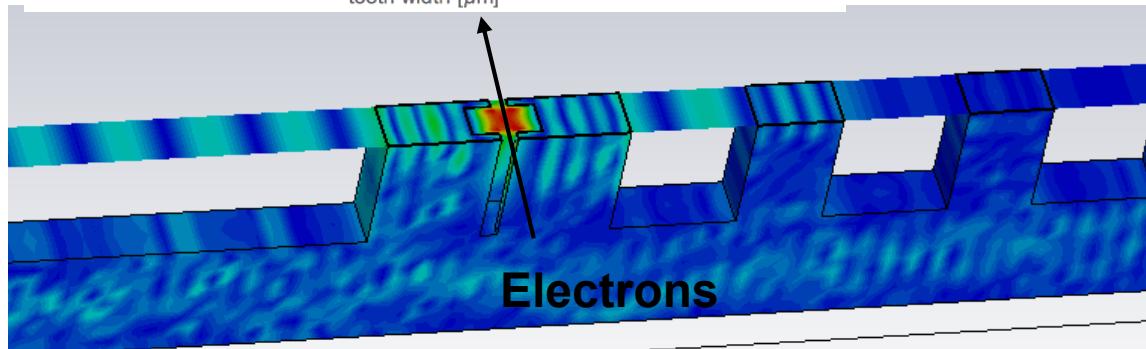
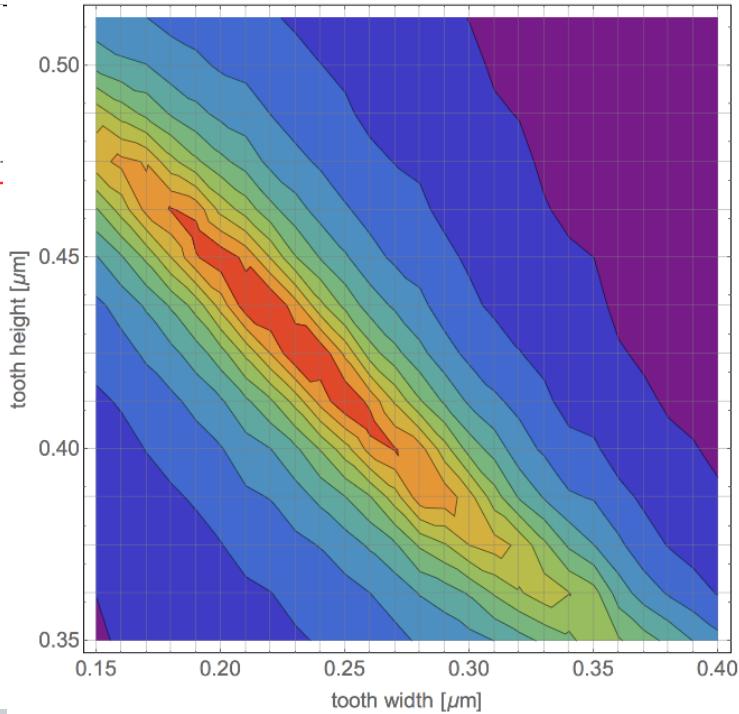
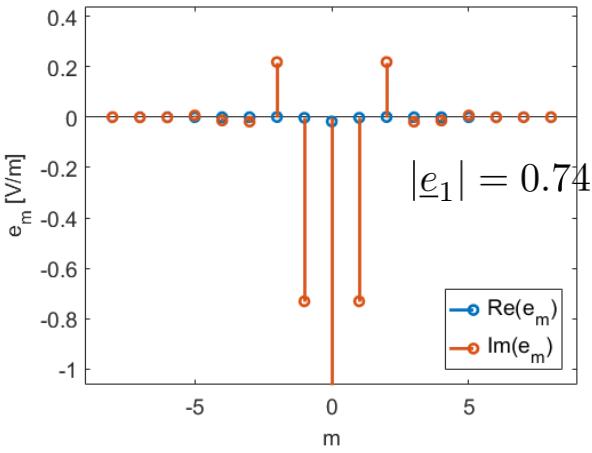
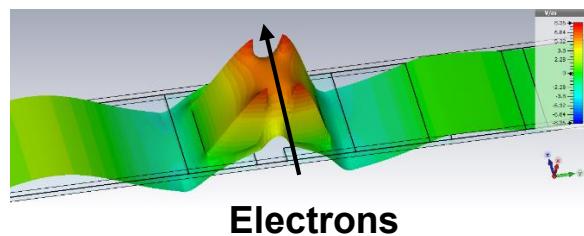
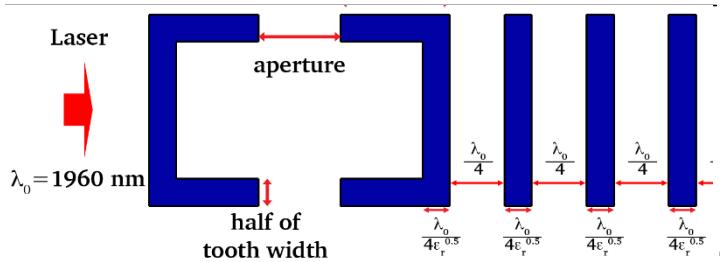
$$\underline{E}_z(x, y, z) = \sum_m \underline{e}_m e^{-im\frac{2\pi}{\lambda_g}z}$$

$$\Delta W_{\max} = q\lambda_g |\underline{e}_m|$$

Single DLA cell optimization process



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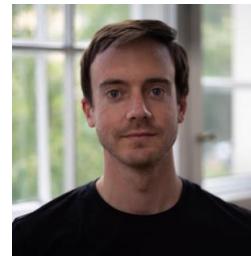
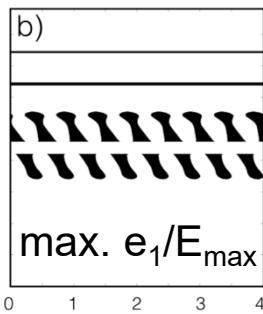
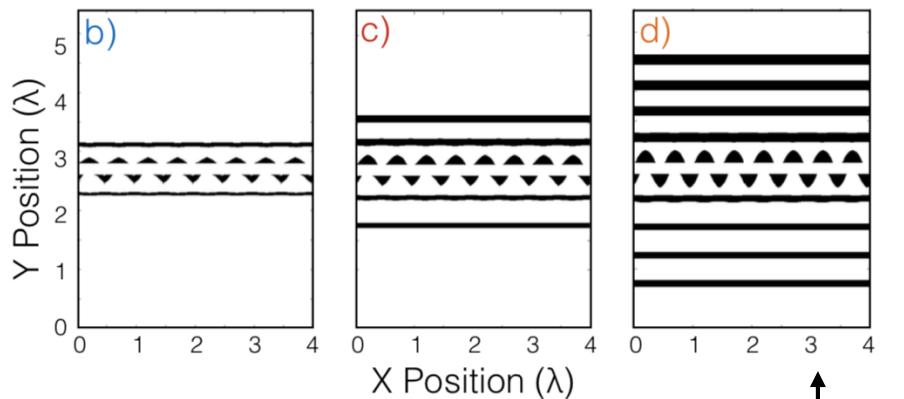
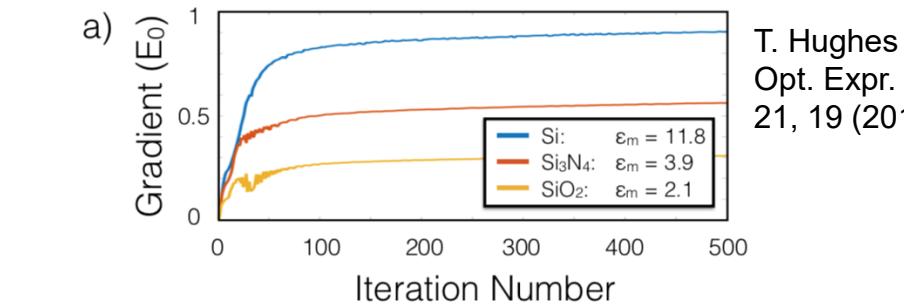
T. Egenolf,
TU Darmstadt

Advanced Optimization Schemes

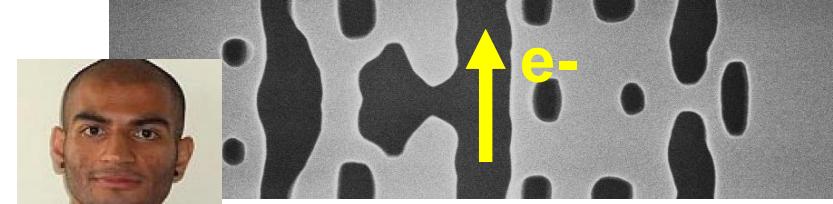
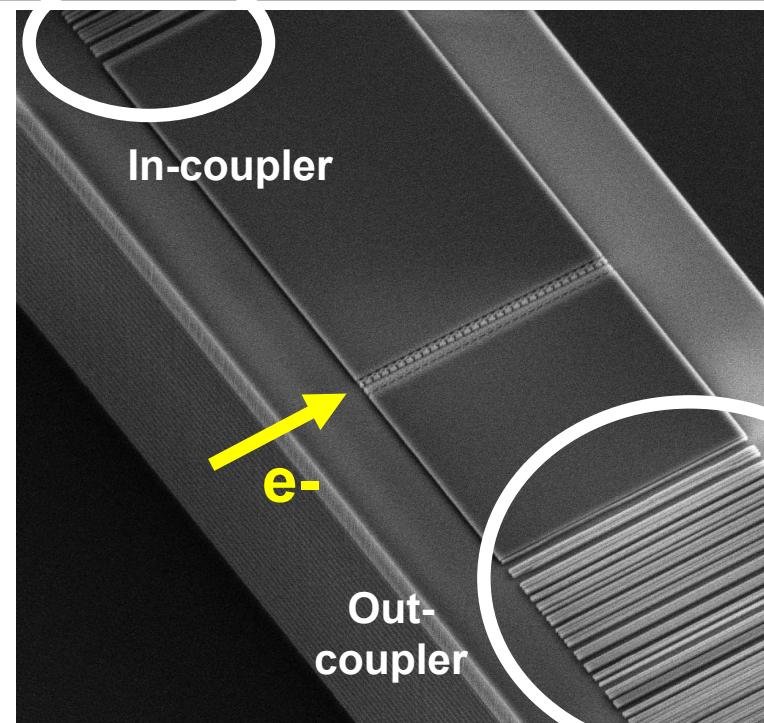
Inverse Design by Adjoint Variable Method (AVM)



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T Hughes,
N. Sapra,
Ginzton Lab,
Stanford



N. Sapra et al.
<https://arxiv.org/abs/1808.07630>

Bottom-Up design of DLA structures: DLAtrack6D



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- One kick per grating cell (**numerically lightweight**)
- **Symplectic code** → No artificial emittance increase
- Kicks by resonant Fourier coefficient

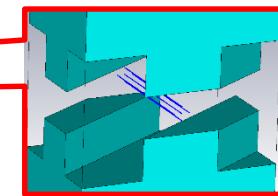
$$\underline{e}_1(x, y) = \underline{e}_1(0, 0) \cosh(ik_y y) e^{ik_x x}$$

- Transverse kick by Panofsky-Wenzel theorem

$$\Delta \vec{p}_{\perp}(x, y; s) = -\nabla_{\perp} \int ds \Delta p_{\parallel}(x, y; s) = \frac{q\lambda_g \lambda_0}{2\pi c} \Im \left\{ e^{2\pi i \frac{s}{\lambda_g}} \nabla_{\perp} \underline{e}_1 \right\}$$

- Can be applied to laterally coupled structures

- Subrelativistic structures / Relativistic structures
- Tilted grating structures
- Alternating phase / Spatial Harmonic focusing structures



U. Niedermayer et al. Phys. Rev. Accel. Beams **20**, 111302 (2017)

DLAtrack6D

→One kick per grating period



$$x' = \frac{p_x}{p_{z0}} , \quad \Delta x' = \frac{\Delta p_x(x,y,\varphi)}{p_{z0}} , \quad y' = \frac{p_y}{p_{z0}} , \quad \Delta y' = \frac{\Delta p_y(x,y,\varphi)}{p_{z0}}$$

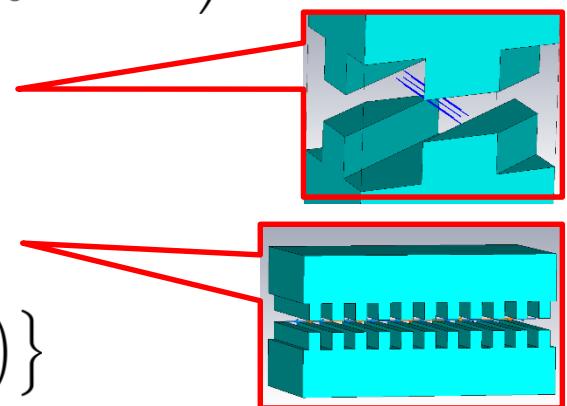
$$\varphi = 2\pi \frac{s}{\lambda_{gz}} , \quad \delta = \frac{W-W_0}{W_0} , \quad \Delta \delta = \frac{\Delta W(x,y,\varphi) - \Delta W(0,0,\varphi_s)}{W_0}$$

$$\begin{pmatrix} x \\ x' \\ y \\ y' \\ \varphi \\ \delta \end{pmatrix}^{(n+1)} = \begin{pmatrix} x \\ Ax' + \Delta x'(x,y,\varphi) \\ y \\ Ay' + \Delta y'(x,y,\varphi) \\ \varphi \\ \delta + \Delta \delta(x,y,\varphi; \varphi_s) \end{pmatrix}^{(n)} + \begin{pmatrix} \lambda_{gz} x'(x,y,\varphi) \\ 0 \\ \lambda_{gz} y'(x,y,\varphi) \\ 0 \\ -\frac{2\pi}{\beta^2 \gamma^2} \delta(x,y,\varphi) \\ 0 \end{pmatrix}^{(n+1)}$$

$$\Delta x' = -\frac{q\lambda_0}{p_{z0}c} \tan(\alpha) \cosh(ik_y y) \Re \left\{ \underline{e}_1 e^{i\varphi + i\frac{2\pi x}{\lambda_{gx}}} \right\}$$

$$\Delta y' = \frac{-ik_y \lambda_0^2 q \beta}{2\pi p_{z0} c} \sinh(ik_y y) \Im \left\{ \underline{e}_1 e^{i\varphi + i\frac{2\pi x}{\lambda_{gx}}} \right\}$$

$$\Delta \delta = \frac{q\lambda_{gz}}{\gamma m_e c^2} \Re \left\{ \underline{e}_1 \left(\cosh(ik_y y) e^{i\varphi + i\frac{2\pi x}{\lambda_{gx}}} - e^{i\varphi_s} \right) \right\}$$



U. Niedermayer et al. Phys. Rev. Accel. Beams **20**, 111302 (2017)

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Stabilization and scalability of DLA

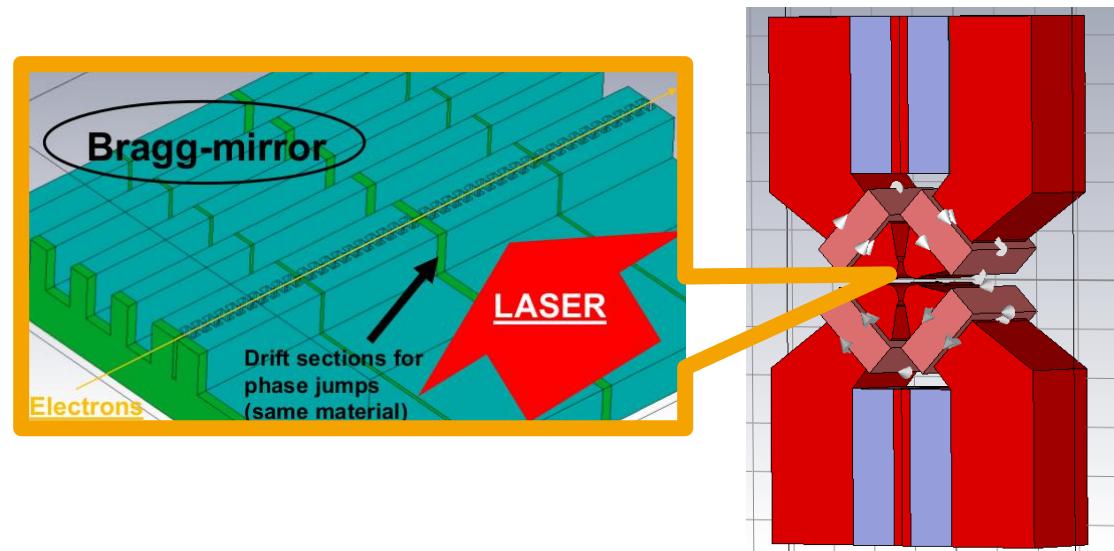
- Magnetic focusing is too weak to counteract acceleration defocusing (see e.g. Ody et al. NIM A 865, 75-83 (2017))
- Spatial harmonic focusing in travelling wave structures
 - see B. Naranjo et al., PRL 109, 1 (2012) → Galaxie DLA
 - Shows only stability, not what emittance fits into the tiny aperture
- Focusing in laterally coupled (standing wave) structures
 - Programmable Spatial Light Modulator (SLM)
 - Small drift sections → Alternating Phase Focusing (APF)

U. Niedermayer et al. accepted for publication in PRL. (see arXiv 1806.07287)

Alternating Phase Focusing (APF)



- Originally invented for conventional RF-ion-linacs in the 1950', but outperformed by RFQ
- “Earnshaw’s theorem“ forbids simultaneous focusing in all directions
- We cannot make 3D structures on the μ -chip scale
- APF lattice is 2D: Possible to fabricate by lithographic methods



Constant focusing in z
Alternating in x and y

Constant focusing in x
Alternating in z and y

Hamiltonian (dual pillar structure)



Read off directly from symplectic tracking scheme:

$$\dot{y} = \frac{p_y}{m_e \gamma}$$

$$\dot{p}_y = -qe_1 \frac{\lambda_{gz}}{2\pi} \frac{\omega}{\beta\gamma c} \sinh\left(\frac{\omega y}{\beta\gamma c}\right) \sin\left(\frac{2\pi s}{\lambda_{gz}}\right)$$

$$\dot{s} = \frac{\Delta p_z}{m_e \gamma^3}$$

$$\dot{\Delta p_z} = qe_1 \left[\cosh\left(\frac{\omega y}{\beta\gamma c}\right) \cos\left(\frac{2\pi s}{\lambda_{gz}}\right) - \cos \varphi_s \right].$$

$$H = \frac{1}{2m_e \gamma} (p_x^2 + p_y^2 + (\Delta p_z/\gamma)^2) + V$$

Panofsky-Wenzel: $\vec{F} = -\nabla' V$

$$V = qe_1 \left[\frac{\lambda_{gz}}{2\pi} \cosh\left(\frac{\omega y}{\beta\gamma c}\right) \sin\left(\frac{2\pi s}{\lambda_{gz}}\right) - s \cos \varphi_s \right]$$

U. Niedermayer et al. Phys. Rev. Accel. Beams **20**, 111302 (2017)

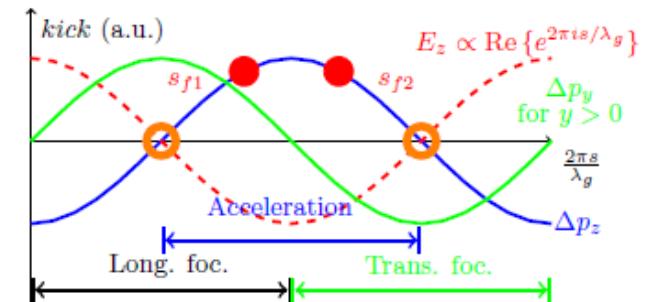
Design of an Alternating Phase Focusing Accelerator Chip



- Time dependent potential

$$V = q \Im \left\{ e_1 \left[\frac{\lambda_g}{2\pi} \cosh \left(\frac{\omega y}{\beta \gamma c} \right) e^{2\pi i s / \lambda_g} - i s e^{i \varphi_s} \right] \right\}$$

$$\varphi_s^{(n)} = \varphi_0 - \arg(e_1)^{(n)} \quad \varphi_0 = \arccos \left[\frac{m_e c^2}{q \lambda_0} \frac{\gamma^3}{|e_1|} \frac{\Delta \lambda_g}{\lambda_0} \right]$$



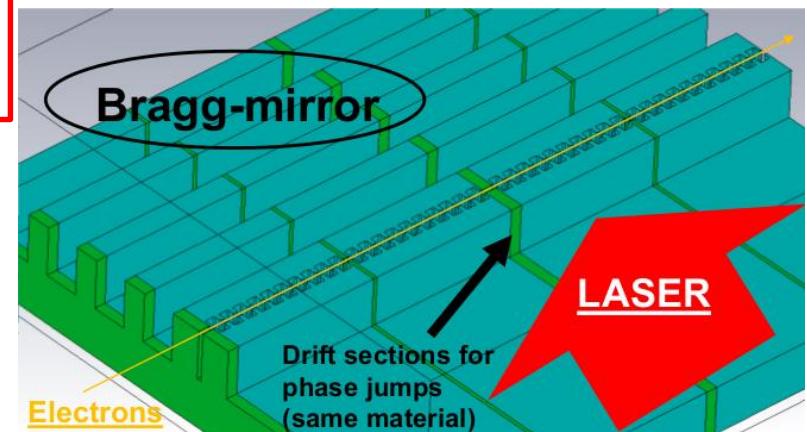
$$V(x, y, s = s_{f1} + \Delta s) = -V(x, y, s = s_{f2} + \Delta s)$$

$$= \frac{q |e_1| \lambda_g}{2\pi} \left[\frac{1}{2} \left(\frac{\omega y}{\beta \gamma c} \right)^2 - \frac{1}{2} \left(\frac{2\pi}{\lambda_g} \Delta s \right)^2 \right] \sin(\varphi_0)$$

$$y'' + K y = 0 \quad K = \frac{|q| \omega e_1}{(m_e \beta^3 \gamma^3 c^3)} \sin(\varphi_s)$$

$$\Delta s'' - K \Delta s = 0$$

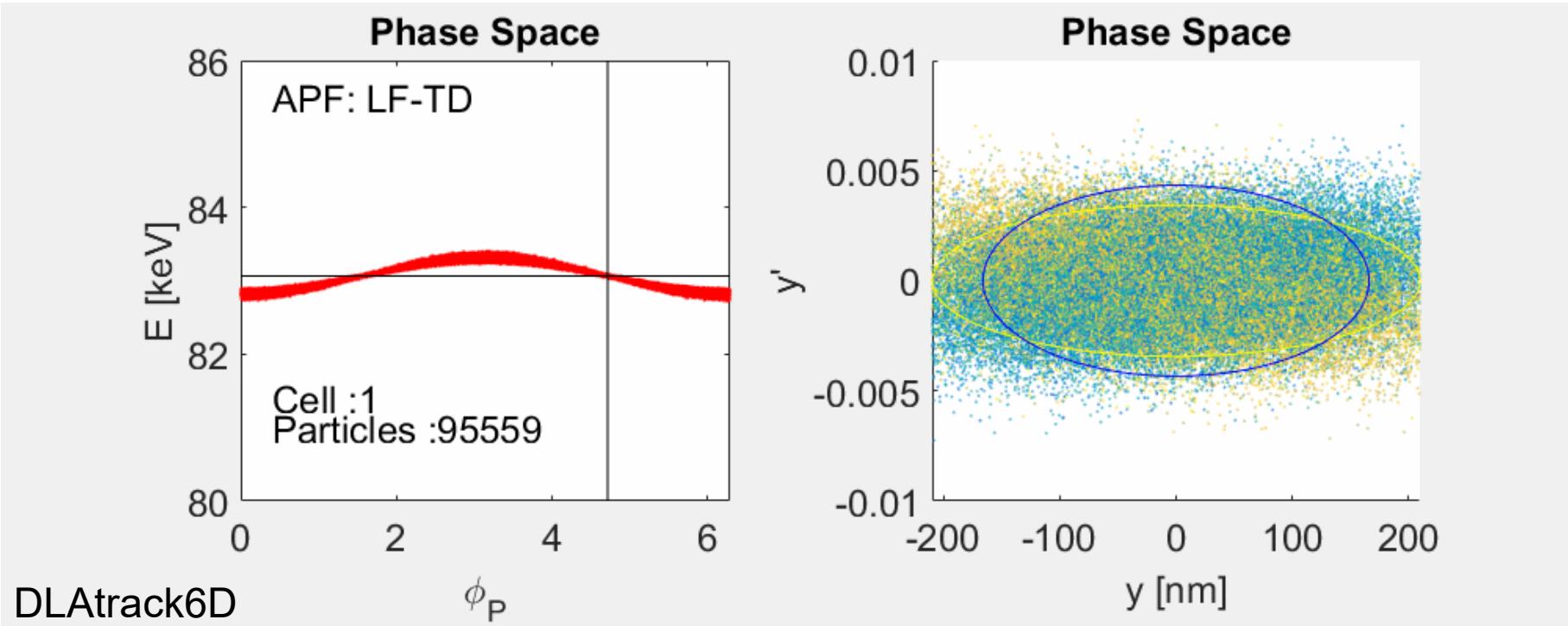
$$\arg(e_1)(P) = \begin{cases} 0, P \text{ odd} \\ 2\varphi_0, P \text{ even} \end{cases}$$



Unbunched transversely matched



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$$\varepsilon(y, y') = \hat{\gamma}y^2 + 2\hat{\alpha}yy' + \hat{\beta}y'^2$$

Longitudinal Courant-Snyder Invariant: $\varepsilon_L(\Delta s, \Delta s') = \hat{\gamma}_L \Delta s^2 + 2\hat{\alpha}_L \Delta s \Delta s' + \hat{\beta}_L \Delta s'^2$

U. Niedermayer et al. accepted for publication in PRL. (see arXiv 1806.07287)



Ramp:

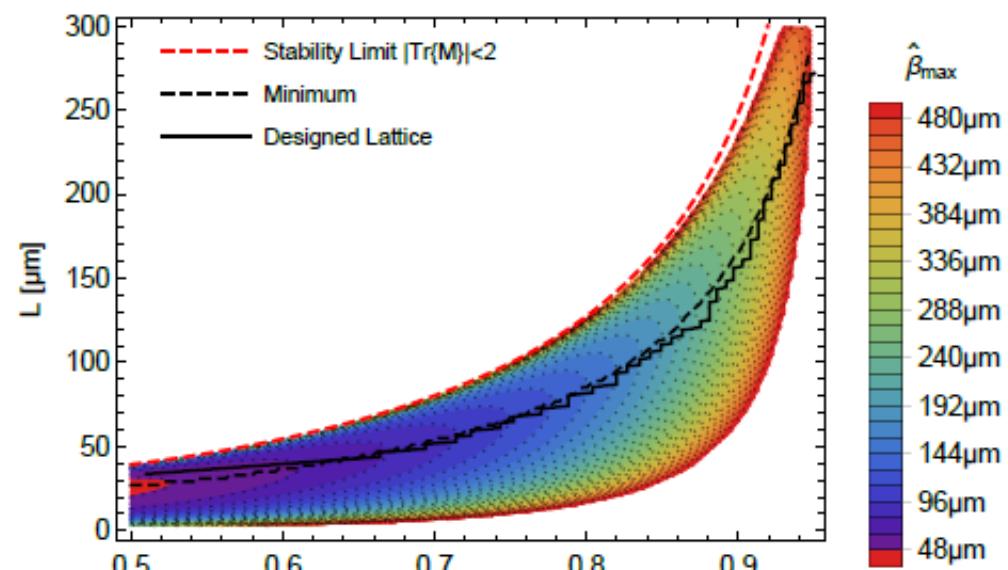
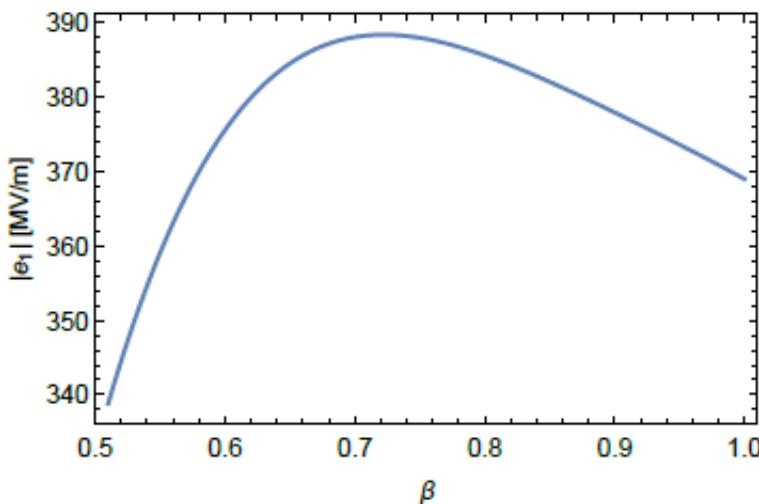
$$W_{kin}(N) = W_{kin,0} + q \sum_{n=1}^N \lambda_g^{(n)} \Re\{e_1^{(n)} e^{i\varphi_s^{(n)}}\}$$

Length chirp according to $\lambda_g = \beta \lambda_0$:

$$\frac{\lambda_g^{(n+1)} - \lambda_g^{(n)}}{\lambda_0} = \beta^{(n+1)} - \beta^{(n)} = \frac{q \lambda_0 \Re\{e_1^{(n)} e^{i\varphi_s^{(n)}}\}}{m_e c^2 \gamma^{(n)} \gamma^3}$$

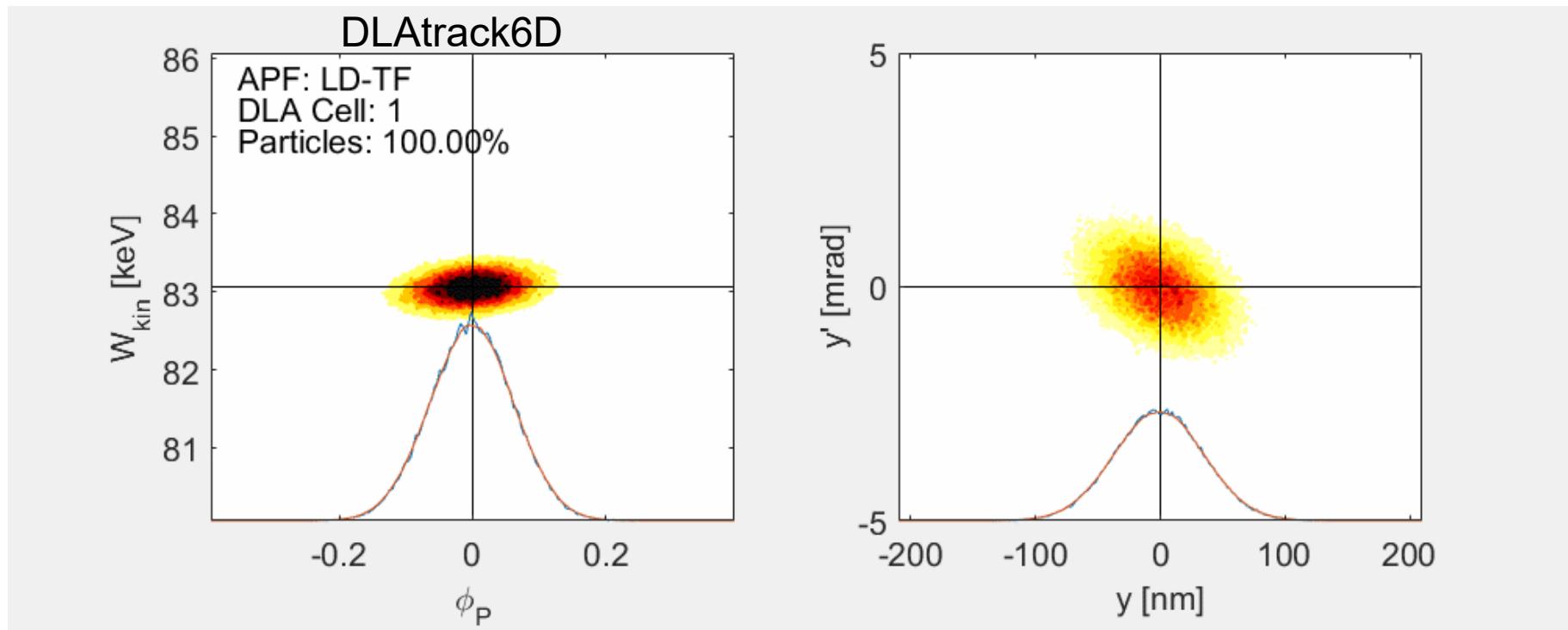
Optimization such that

$$\arg(e_1)^{(n)} \neq f(\lambda_g^{(n)})$$



U. Niedermayer et al. accepted for publication in PRL. (see arXiv 1806.07287)

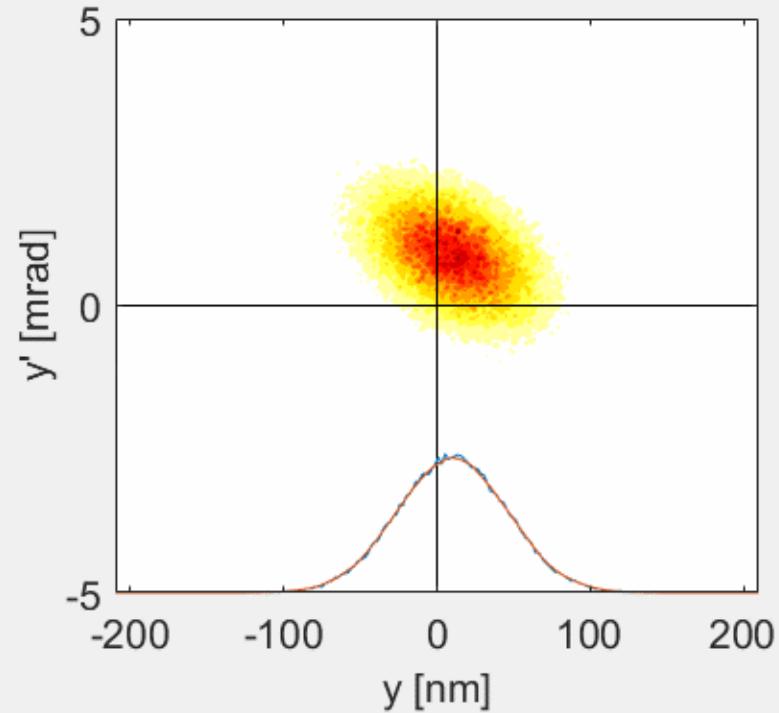
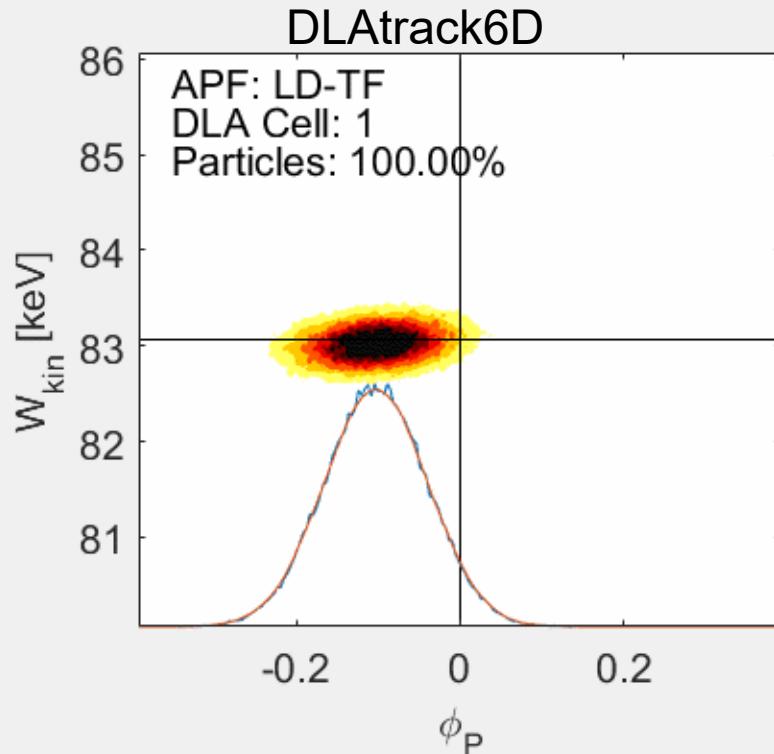
The ideal ACHIP scenario



93% transmission at $\varepsilon = 0.025$ nm

U. Niedermayer et al. accepted for publication in PRL. (see arXiv 1806.07287)

Misalignment example



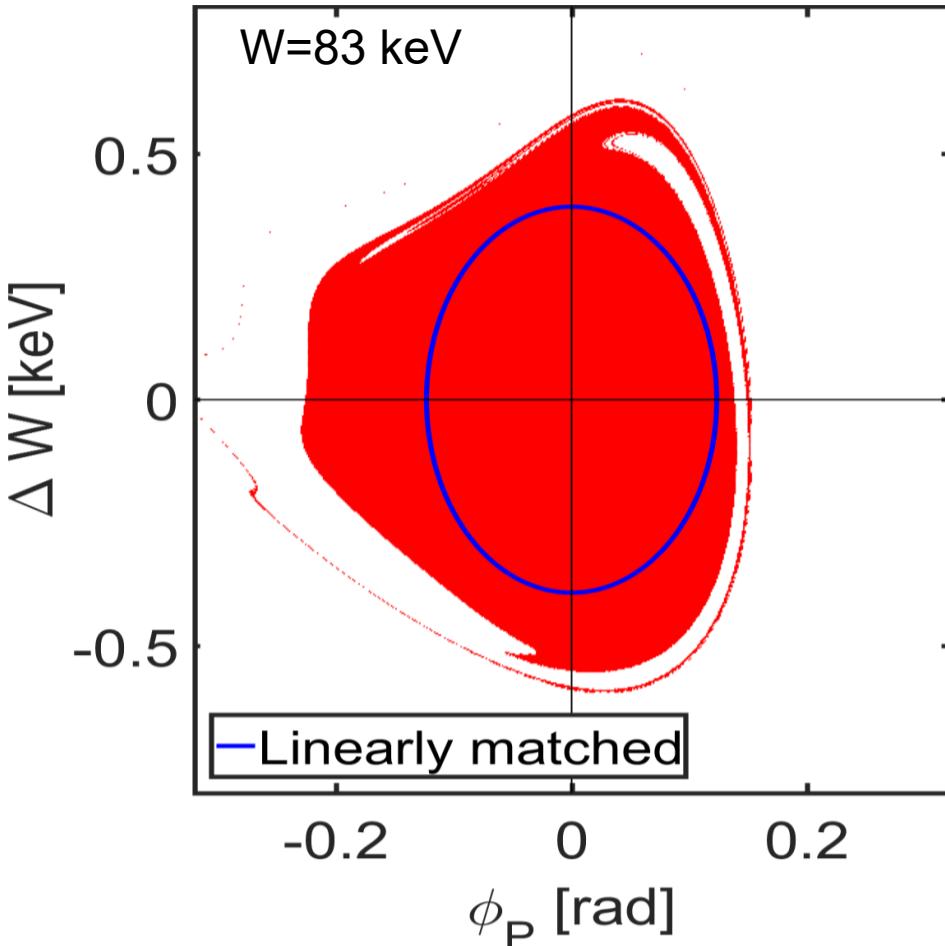
66% transmission at $\varepsilon = 0.025 \text{ nm}$

Initial mismatch:

- $y=10 \text{ nm}$
- $\Phi_P=0.1 \text{ rad}$

$$y'=1 \text{ mrad}$$
$$\Delta W=0$$

Initial phase space



Longitudinal matching condition:

$$\frac{\sigma_\varphi}{2\pi} = \frac{\hat{\beta}_L}{\beta^3 \gamma^3 \lambda_0} \frac{\sigma_{\Delta W}}{m_e c^2}$$

$$\hat{\alpha}_L = 0$$

Bunch length (numerically) :

$$4\sigma_z = 40 \text{ nm}$$

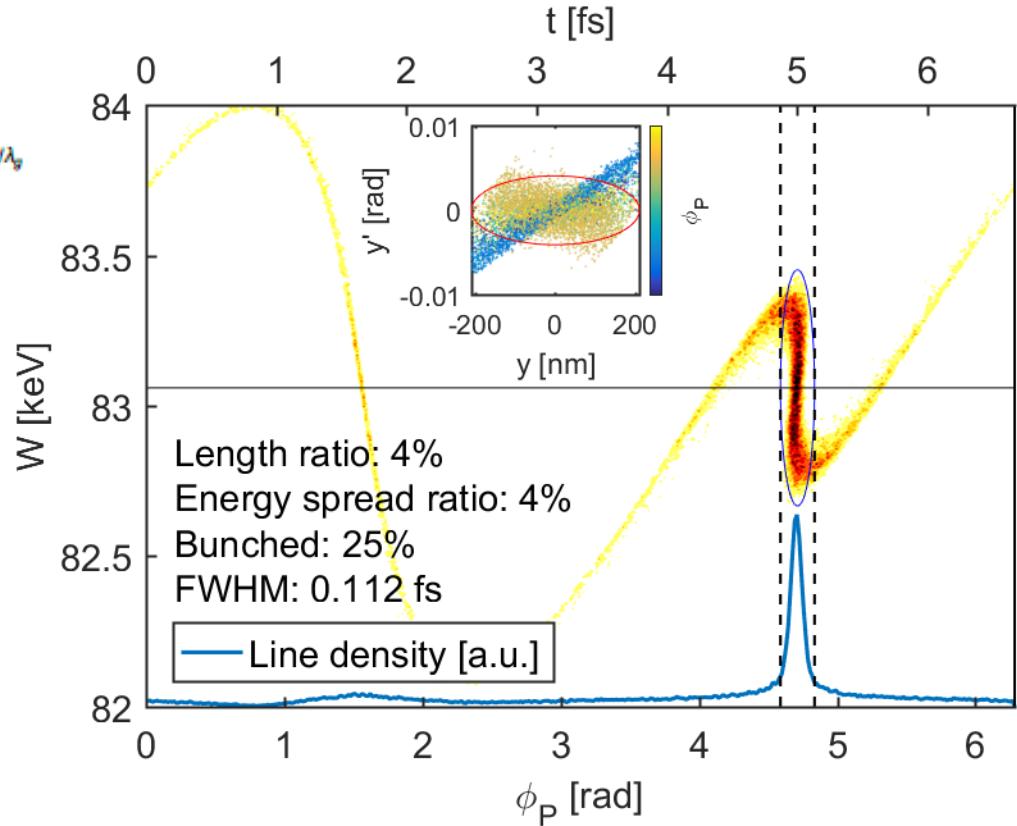
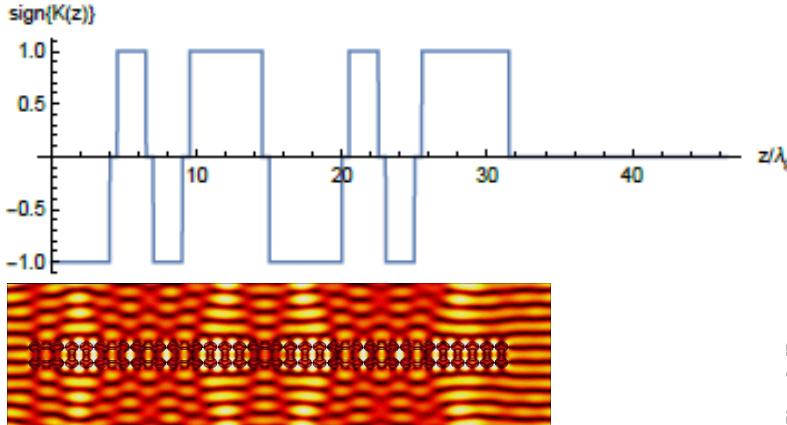
Transverse matching condition:

$$\sigma_y = \sqrt{\hat{\beta}\varepsilon}$$

$$\hat{\alpha} = 0 \Rightarrow \sigma_{y'} = \sqrt{\varepsilon/\hat{\beta}}$$

U. Niedermayer et al. accepted for publication in PRL. (see arXiv 1806.07287)

Bunching



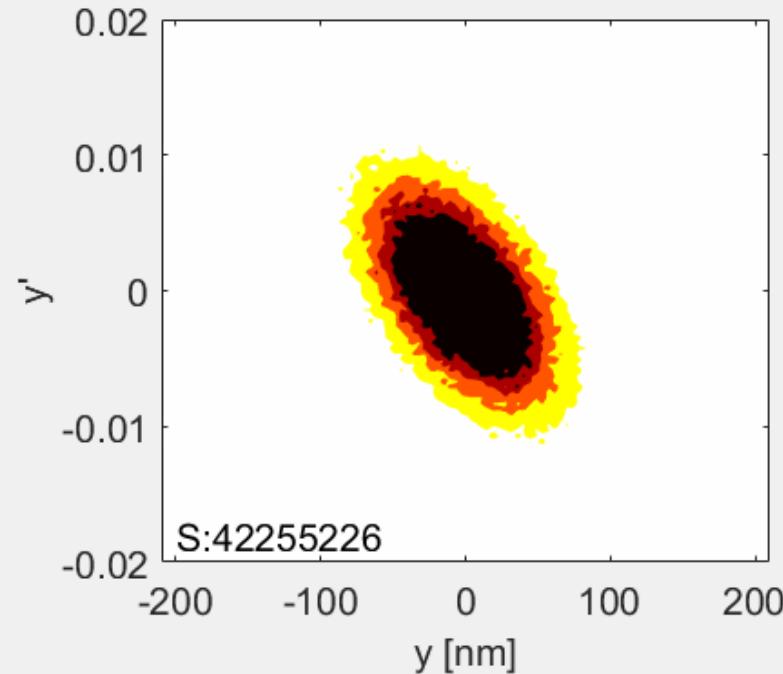
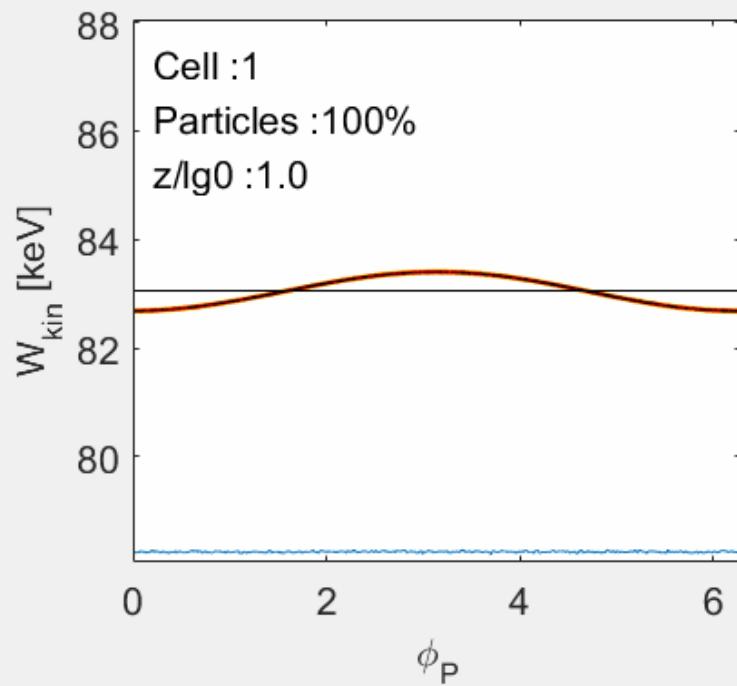
$$\mathbf{M}_b = \mathbf{M}_o^{30} \mathbf{M}_f^6 \mathbf{M}_o \mathbf{M}_d^2 \mathbf{M}_o \mathbf{M}_f^2 \mathbf{M}_o \mathbf{M}_d^5 \mathbf{M}_o \mathbf{M}_f^5 \mathbf{M}_o \mathbf{M}_d^2 \mathbf{M}_o \mathbf{M}_f^2 \mathbf{M}_o \mathbf{M}_d^4$$

U. Niedermayer et al. accepted for publication in PRL. (see arXiv 1806.07287)

Bunching



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U. Niedermayer et al. accepted for publication in PRL. (see arXiv 1806.07287)

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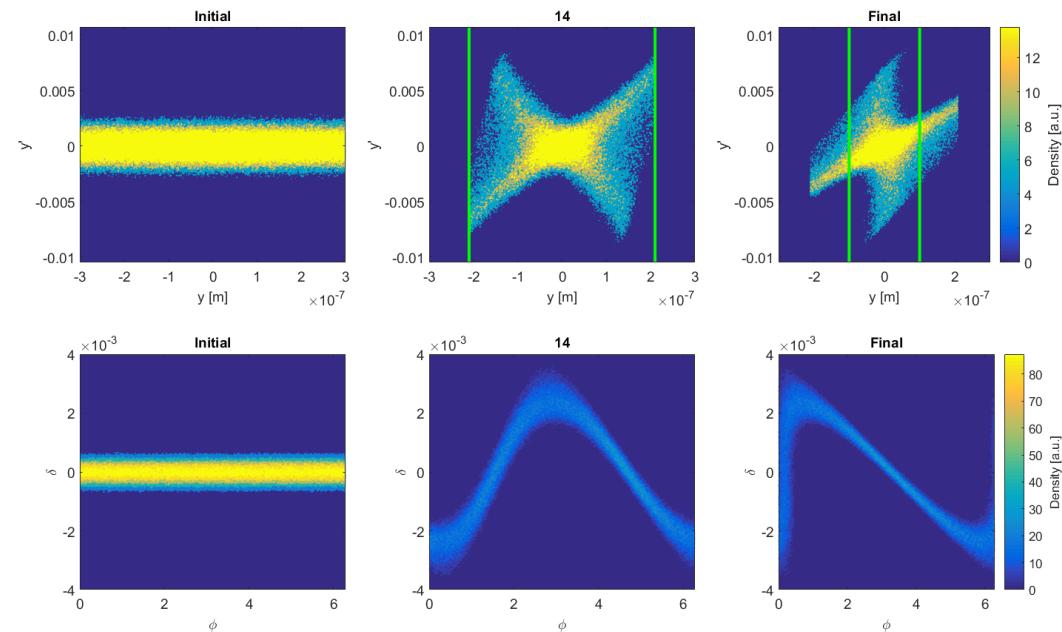
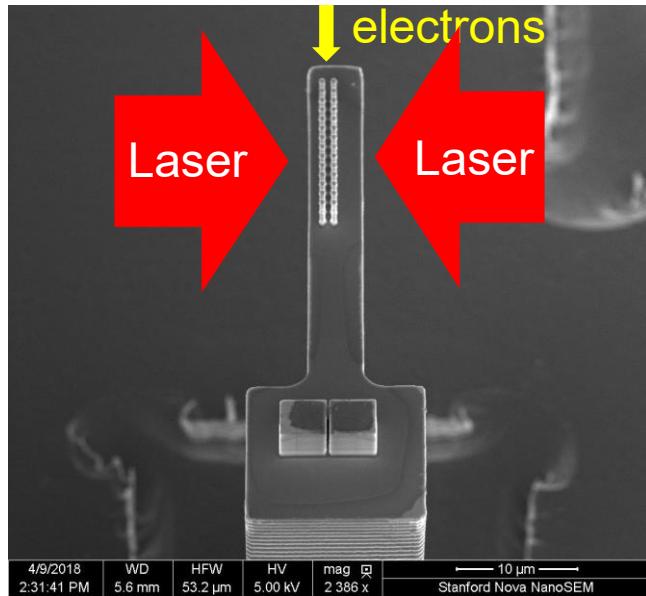
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Single stage focusing experiment with aperture



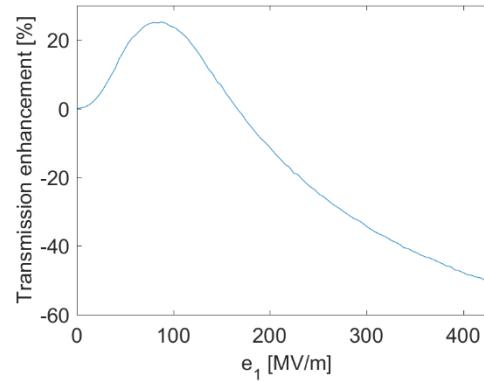
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Data has been taken, publication almost ready



Dylan Black,
Stanford



U. N.
DLAtrack6D
simulation

preliminary

Buncher / streaker experiment at Stanford



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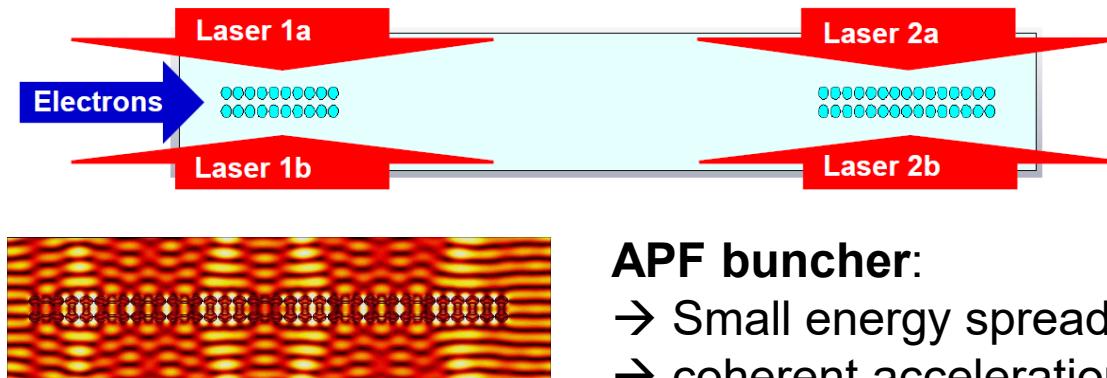
Symmetric DLA in out of phase mode:

$$\underline{e}_1(x, y) = \underline{e}_1(0, 0) \sinh(i k_y y) e^{i k_x x}$$

$$\Delta y' = \frac{-ik_y \lambda_0^2 q \beta}{2\pi p_{z0} c} \cosh(i k_y y) \Im \{ \underline{e}_1 e^{i\varphi} \}$$

$$\Delta \delta = \frac{q \lambda_{gz}}{\gamma m_e c^2} \sinh(i k_y y) \Re \{ \underline{e}_1 e^{i\varphi} \}$$

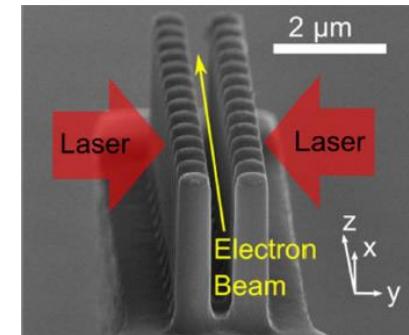
4 lasers allow bunching and coherent streaking:



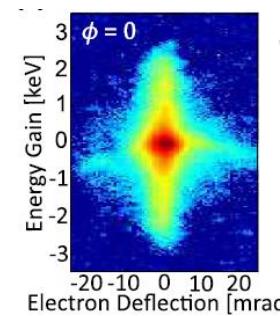
APF buncher:

- Small energy spread
- coherent acceleration
- **All 4 quadrants coherent!**

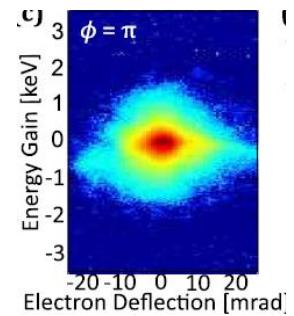
2 lasers: “incoherent”
Acceleration/streaking



Acceleration



Streaking

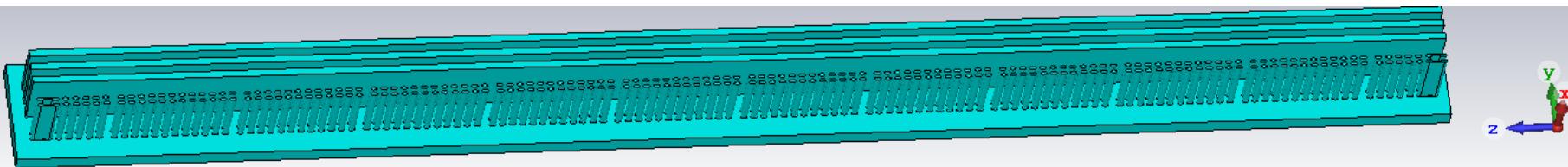


K. Leedle et al.
Optics Letters, 2018

Transport Experiment at FAU Erlangen

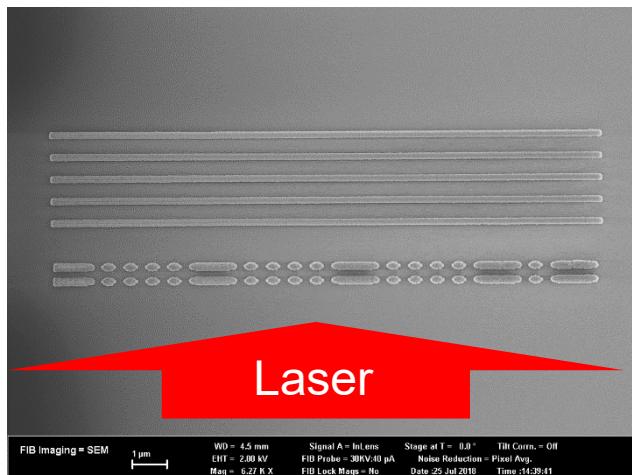


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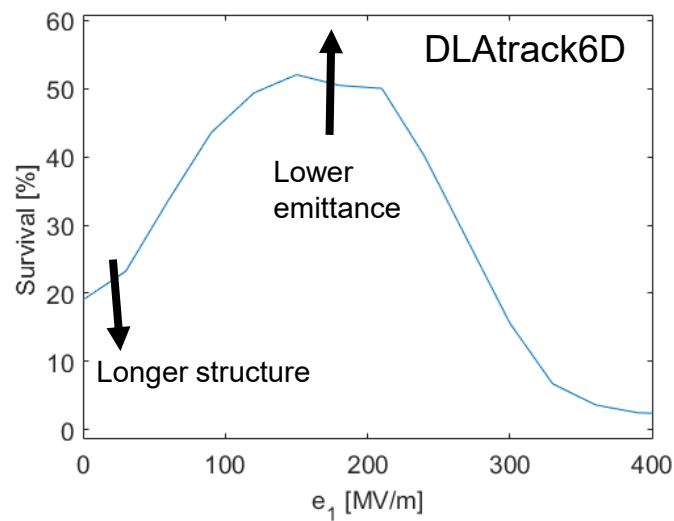
- Fully scalable (length independent transmission)
- Limited by x-Rayleigh length $800\mu\text{m}$
- Transversely matched (min betafct. $\sim 20\mu\text{m}$ @ $E_L=0.8\text{GV/m}$)

Multistage focusing:
(Laser on/off contrast)



$e_1/E_L \sim 0.25$
(incl. DBR)

Geo. Emittance 0.3nm
 $E_{\text{kin}}=26.47\text{keV}$



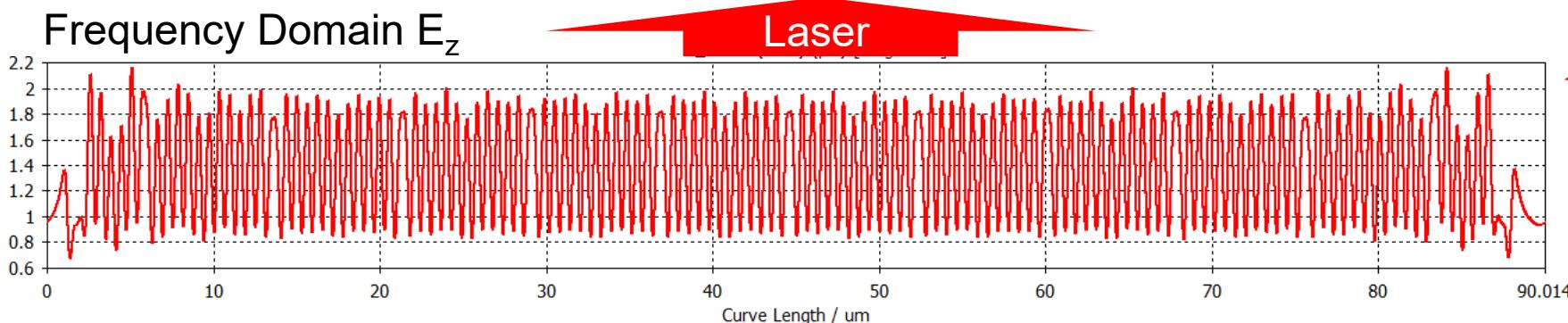
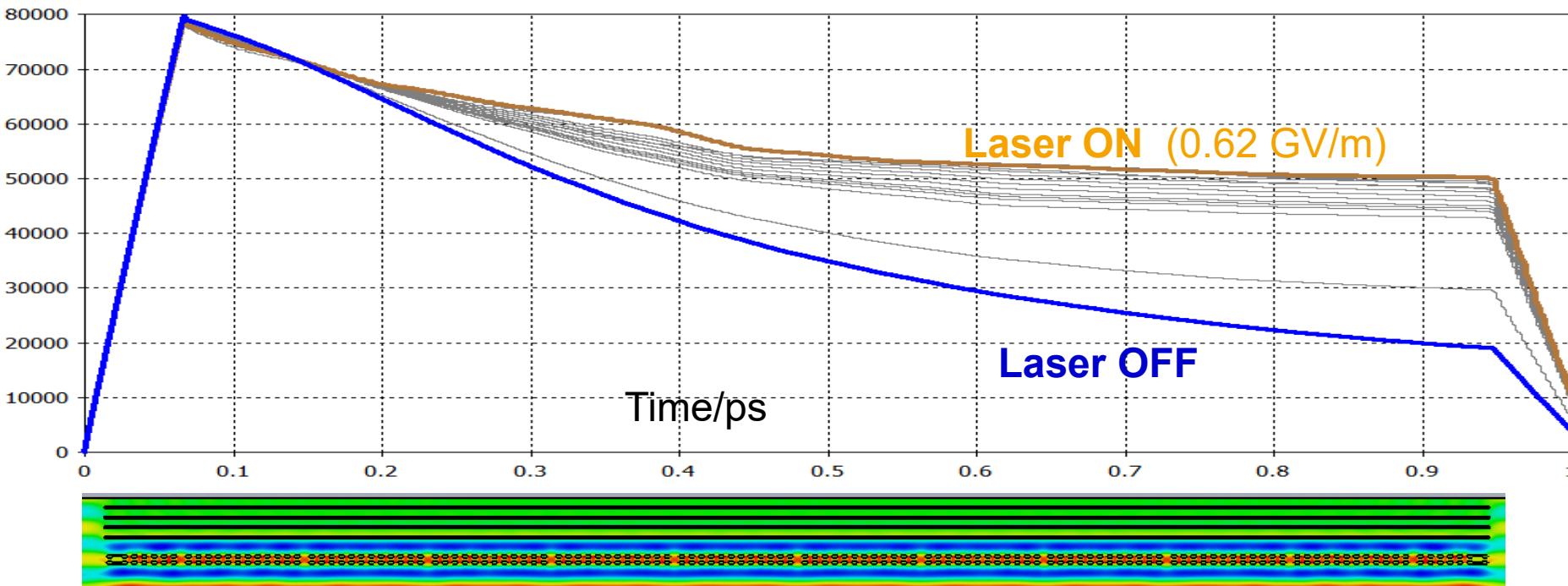
P. Yousefi, J. Illmer:
Fabrication test. Final fab ongoing.

PIC simulation (CST Particle Studio)

100,000 electrons with random phase (uniform dist.)



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Experiments: Shoebox at Stanford

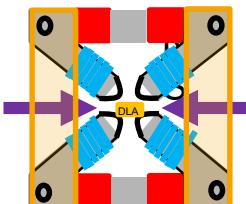
(a slightly more conservative design)



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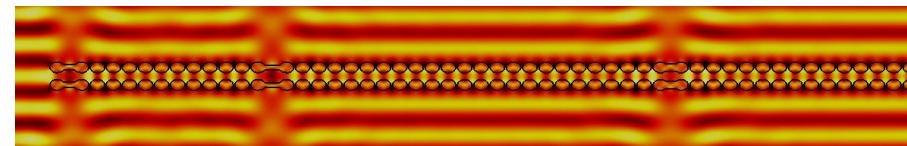
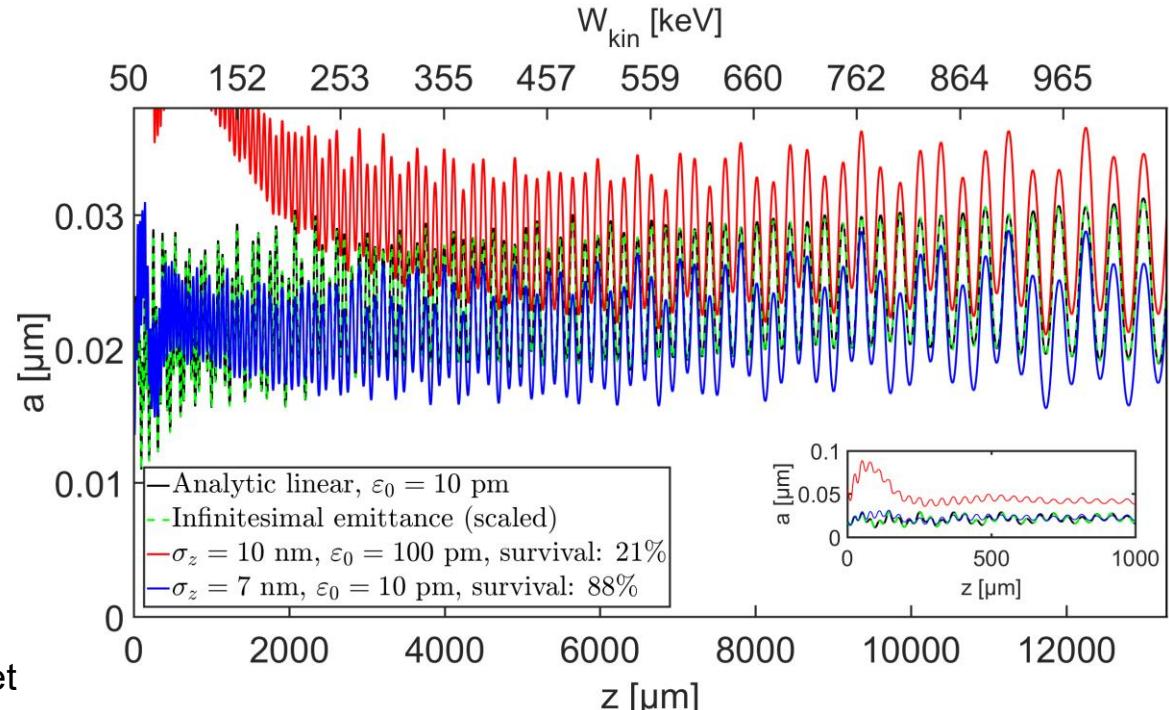
Chirped APF dual pillar
grating for acceleration
50keV → 1MeV

- Laser 300 MV/m (2 sides)
- Av. gradient 73 MeV/m
- Length 1.3 cm
- Emittance requirement:
<0.1nm (geo. @ 50keV)
- Bunch length: < 40nm



Quadrupole magnet
for focusing in the
invariant direction

$$B'_{\text{quad}} = 1 \text{ kT/m} \ll B'^{\text{equiv.}}_{\text{APF}} = 5 \text{ MT/m}$$



Summary

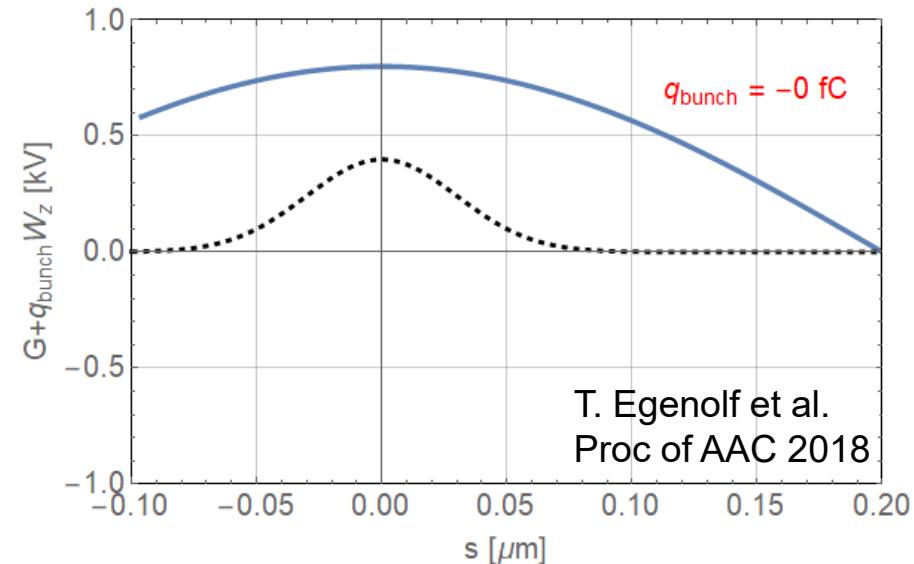
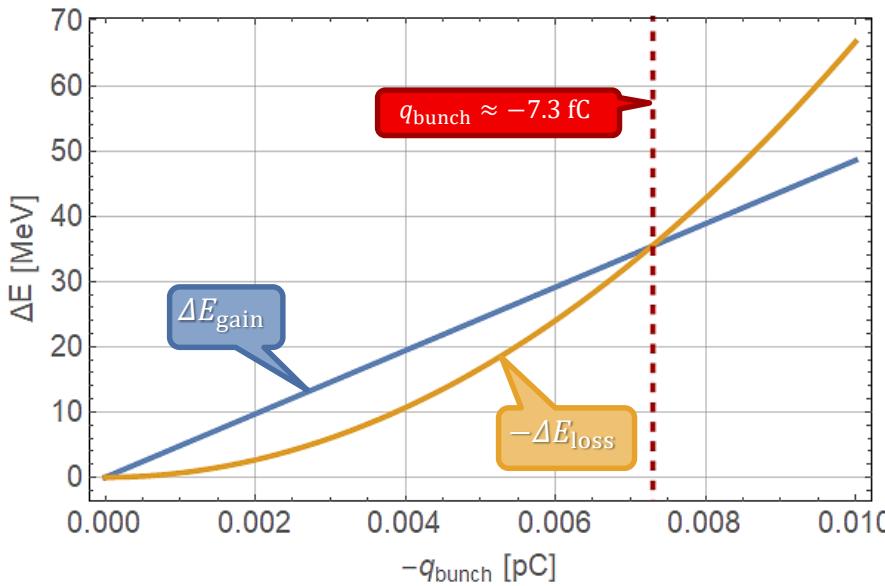


- Simplified simulations with DLAttrack6D in place
- APF beam dynamics scheme works in theory
 - Overcomes resonant acceleration defocusing problem
 - Can also used for bunching in the attosecond range
- Field flatness needs to be tuned (optimized)
- Quadrupole magnet is being developed

Outlook and Wish List



- Plans are to include wake field kicks in DLAttrack6D



A moving window 3D track/PIC code would be nice!
(incl. dielectrics, pulse front tilted laser, open BC)

Thank you for your attention!

The End



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Questions?