

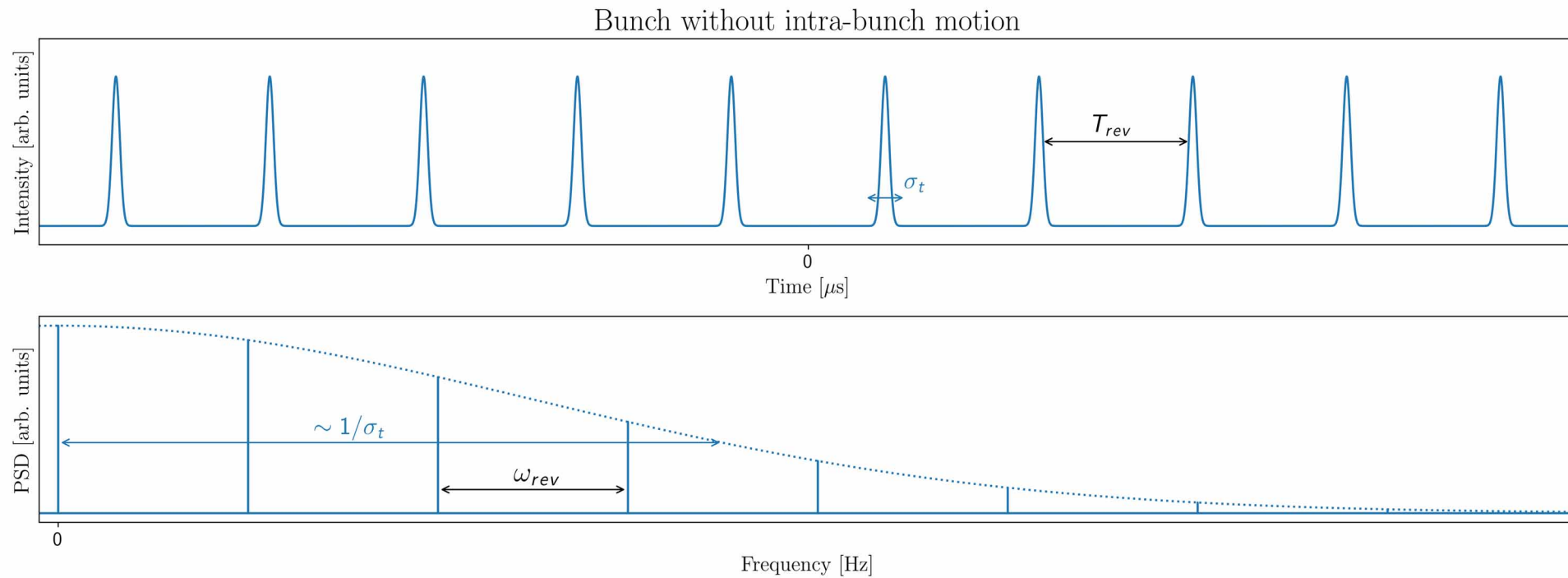
Analysis of the Transverse Schottky Signals in the LHC

Kacper Łasocha, Diogo Alves, CERN Beam Instrumentation Group

14.09.2023, Saskatoon, IBIC'23

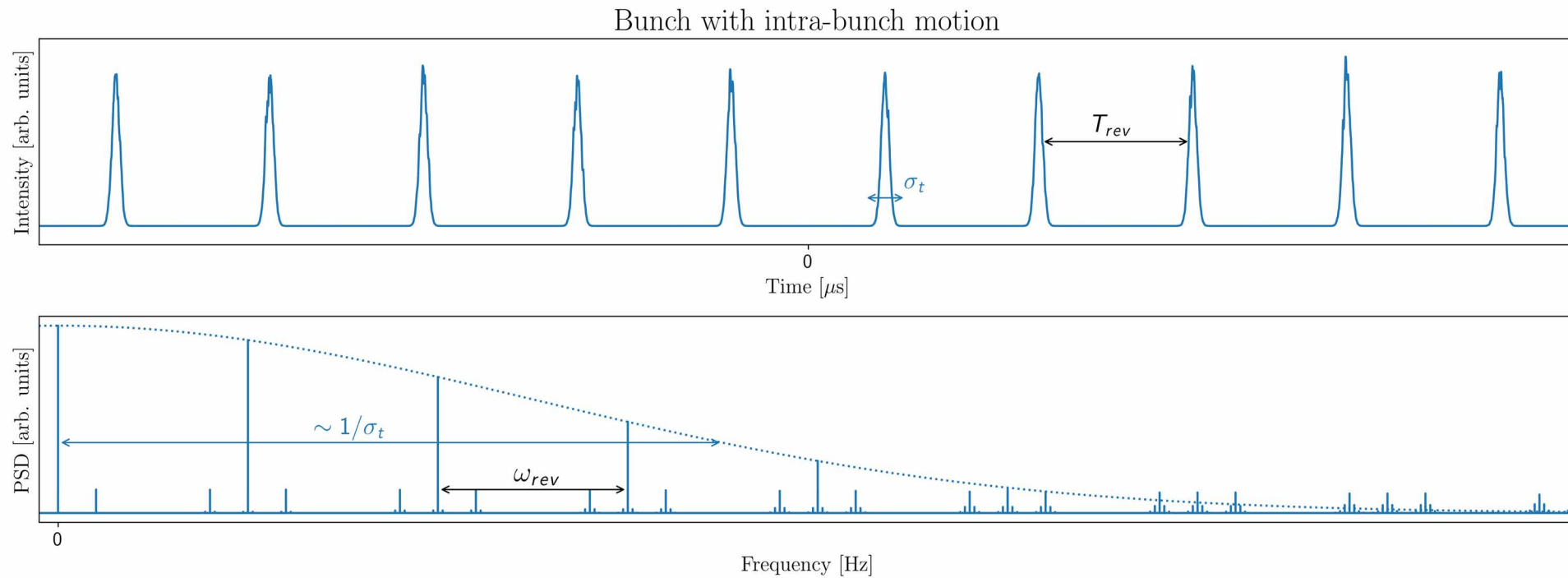
Schottky signals

- Fluctuations of the macroscopic beam signal due to **discrete motion (synchrotron or betatron)** of individual particles within the bunch



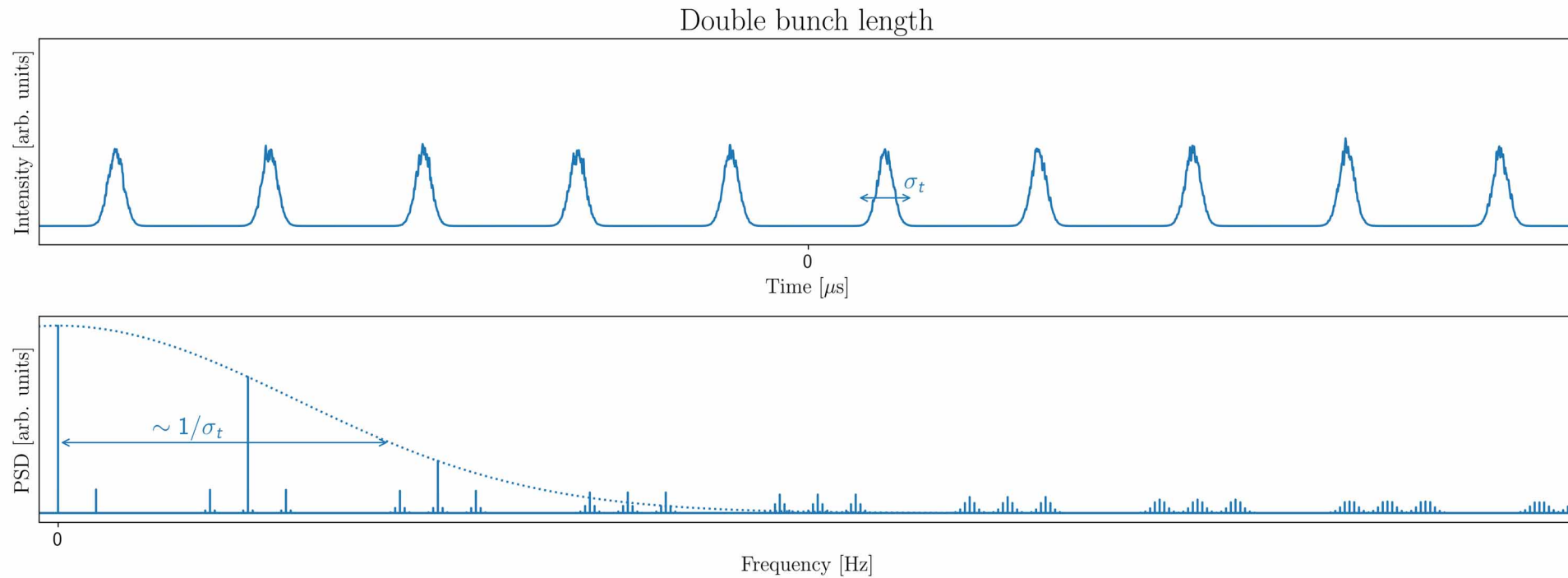
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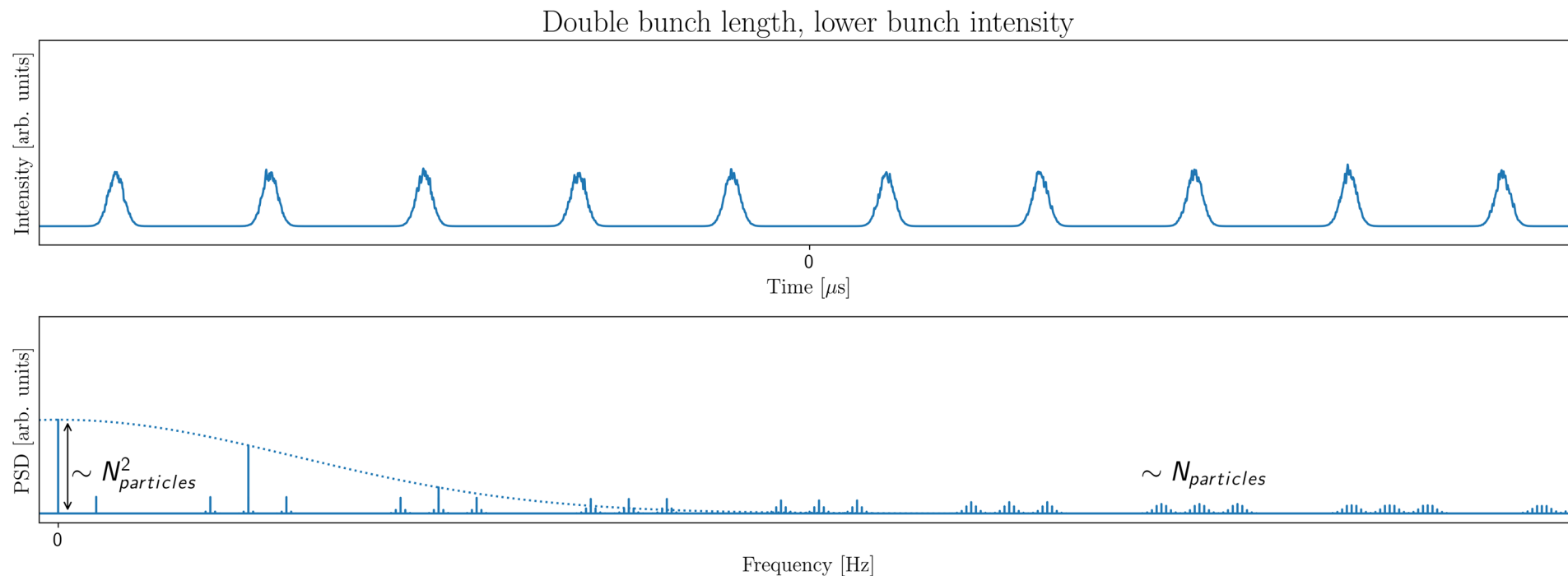
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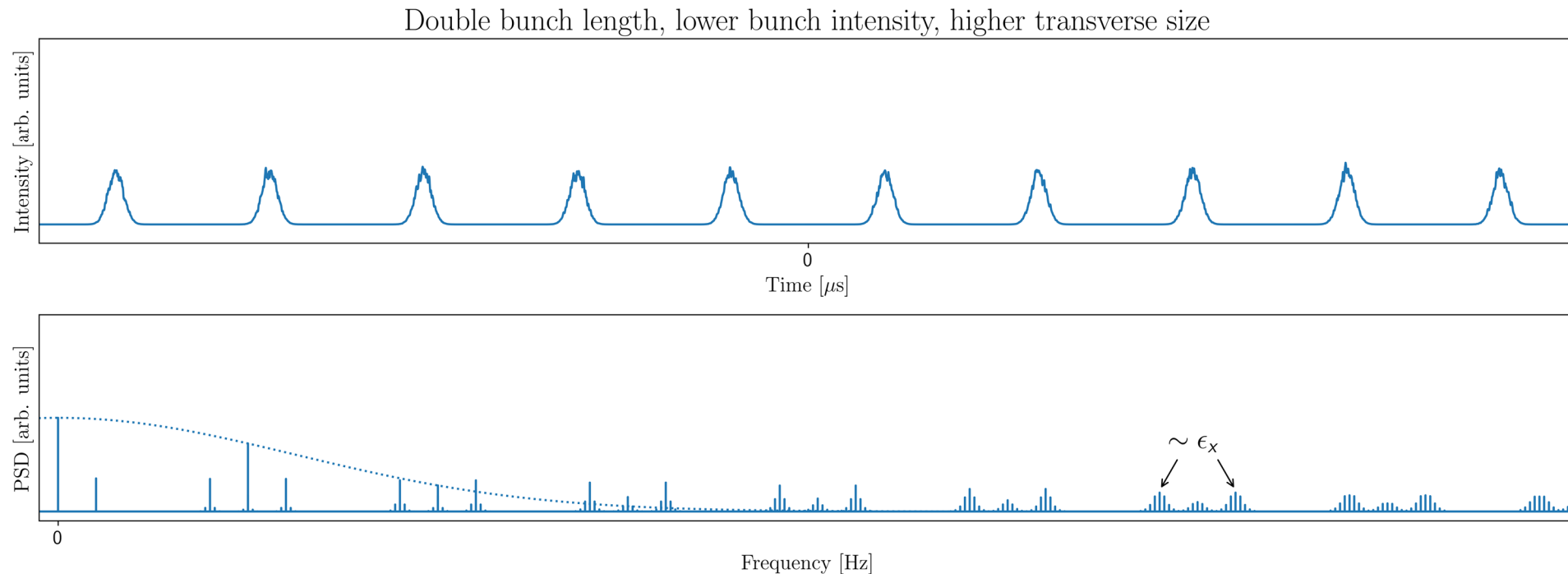
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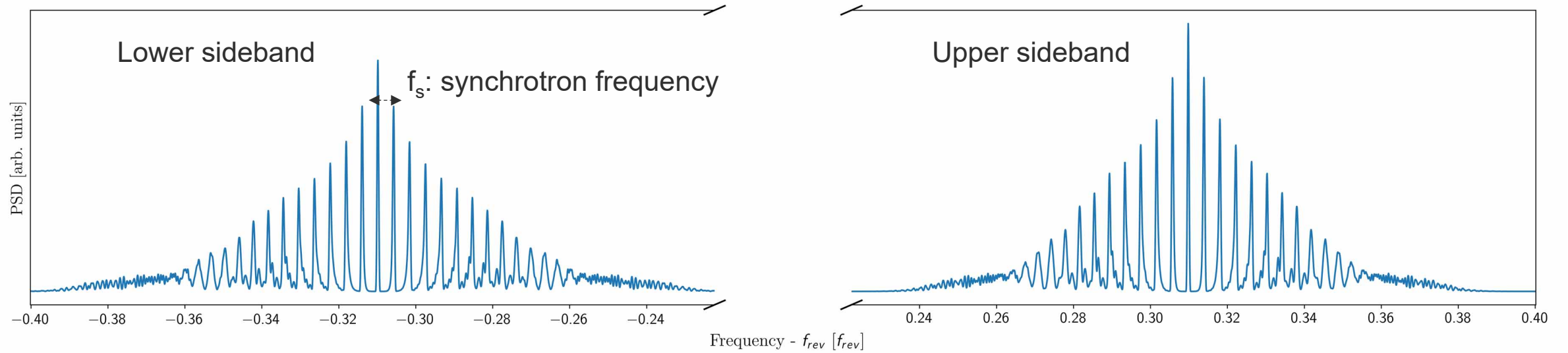


Schottky signals

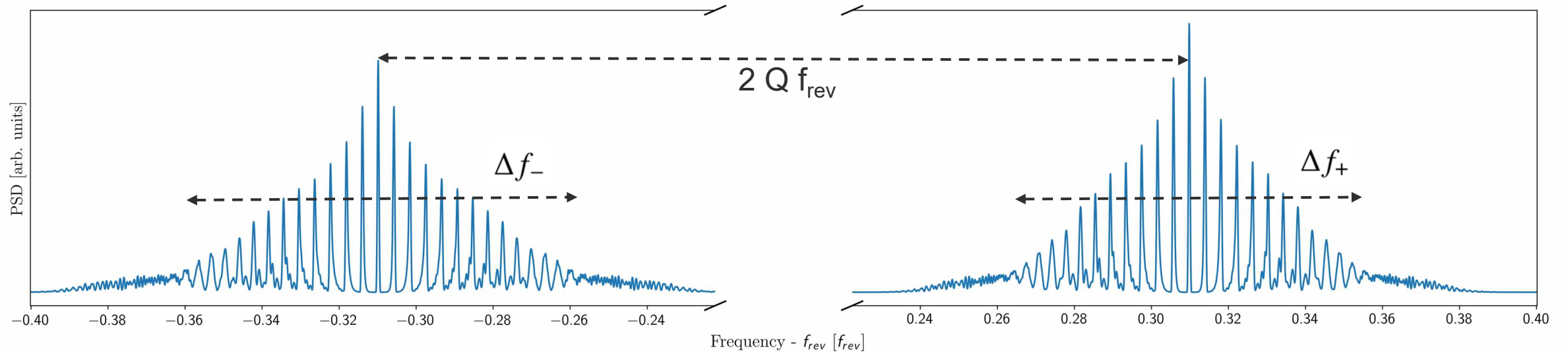
- Fluctuations of the macroscopic beam signal due to **discrete motion (synchrotron or betatron)** of individual particles within the bunch
- Most pronounce for **long, low intensity, transversely large** bunches



Transverse Schottky signals



Transverse Schottky signals



Betatron tune

Mirrored difference method,
minimize the cost function:

$$C_{MD}(k) = \sum_{i=1}^{i=N} |P_T^{\pm}(\omega_{k-i}) - P_T^{\pm}(\omega_{k+i})|$$

Chromaticity

$$Q\xi = -\eta \left(n \frac{\Delta f_- - \Delta f_+}{\Delta f_- + \Delta f_+} - Q_I \right)$$

Δf_{\pm} : RMS width of upper/lower sideband

In certain conditions +/- signs flip, see
K. Lasocha and D. Alves, Phys. Rev. Accel. Beams 25, 062801

Transverse Schottky signals: assumptions

Synchrotron and betatron motion

Synchrotron motion – harmonic, with amplitude dependent frequency:

$$\tau_i(t) = \hat{\tau}_i \sin(\Omega_{s_i} t + \varphi_{s_i})$$

$$\Omega_{s_i} = \frac{\pi}{2\mathcal{K} \left[\sin\left(\frac{h\omega_0 \hat{\tau}_i}{2}\right) \right]} \Omega_{s_0}$$

Betatron motion – harmonic, with frequency modulated wrt to the momentum:

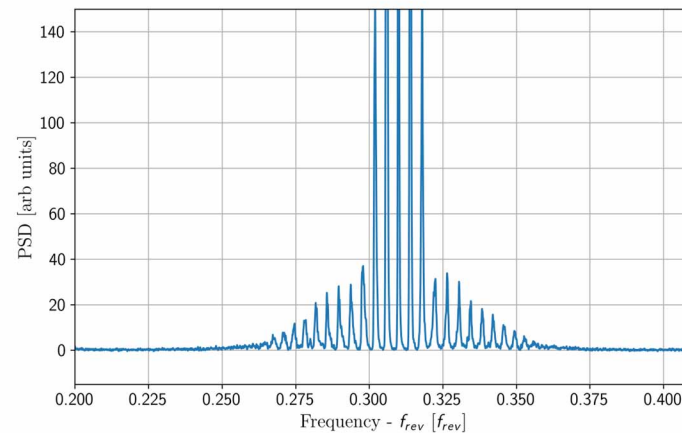
$$x_i(t) = \hat{x}_i \cos \left[Q\omega_0 t + \frac{\hat{Q}_i \omega_0}{\Omega_{s_i}} \sin(\Omega_{s_i} t + \varphi_{s_i}) + \varphi_{\beta_i} \right]$$

$$\hat{Q}_i = Q_5^\xi \frac{\hat{p}_i}{p_0}$$

Uniform distribution of phases;
no "coherent" components

Uniform distribution of φ_{s_i} and φ_{β_i} implies
PSD proportional to the number of particles N.

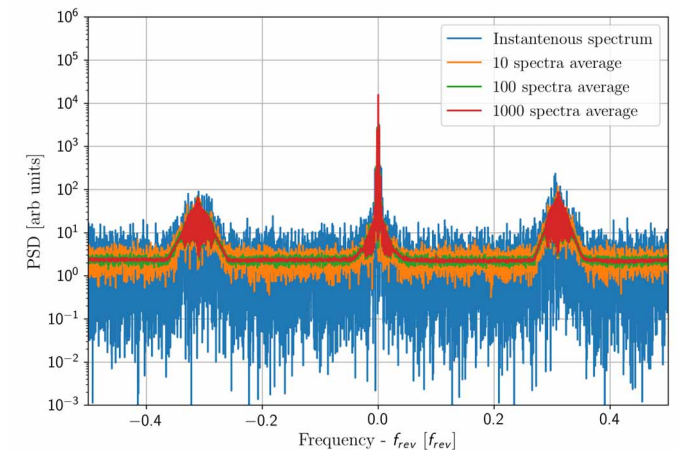
Otherwise the power can be proportional to N²:



Sufficiently long time averaging

The theory predicts only the expected,
ensemble averaged spectrum. Time averaging
required to have a correspondence.

Analyzed LHC spectra are averaged for 100 s.



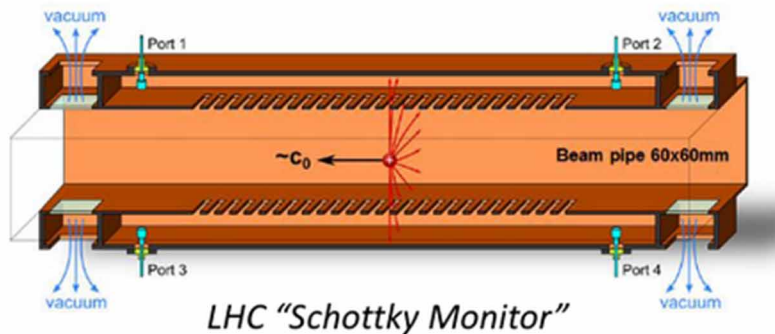
See C. Lannoy et al., WEP035, this conference.

Schottky signals in LHC

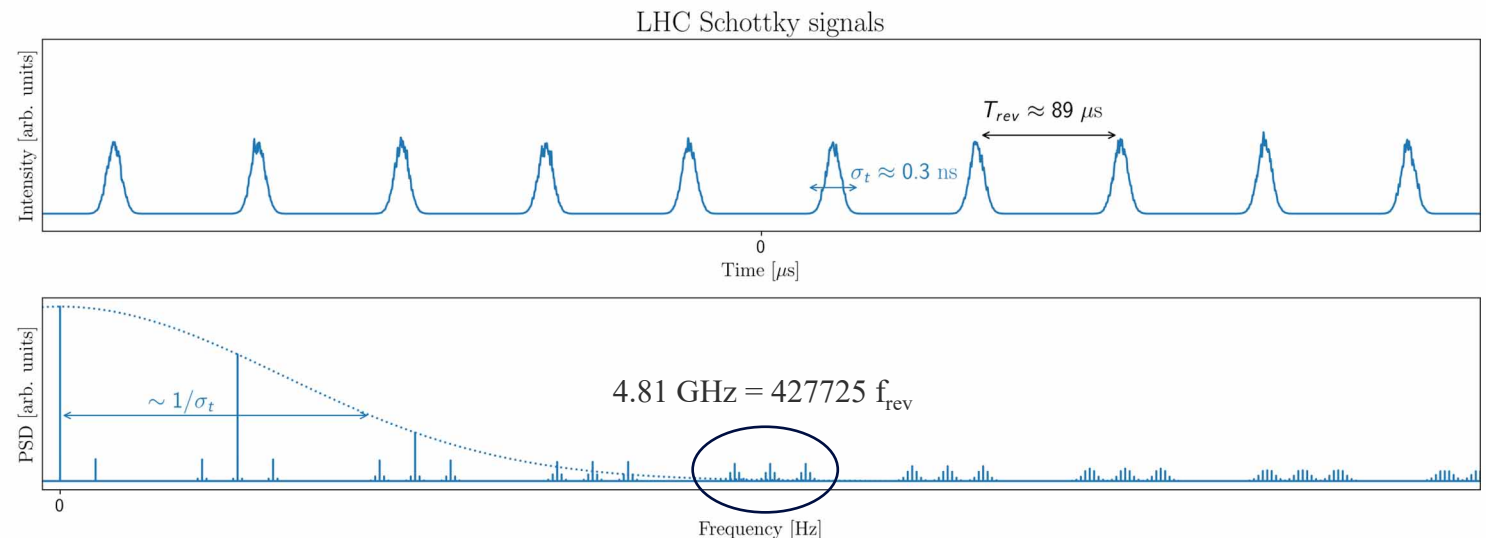
- One system for two particle species: protons and Pb^{82+} ions, one device per beam and plane
- Pair of slotted waveguides, probing beam field at 4.81 GHz, filtering and downmixing signal to 11.2 kHz
- Gating system enables observation of single bunches
- **The only instrument measuring the LHC chromaticity in the non-invasive way**

Typical LHC beam parameters:

	p^+	Pb^{82+}
$N_{\text{particles}}$ (per bunch)	10^{11}	10^8
Bunch length (4σ)	1-1.4 ns	
Normalized transverse emittance	1.5-2.5 μm	
Energy Inj/Flattop (per nucleon)	0.45 - 6.8 TeV	0.18 - 2.6 TeV



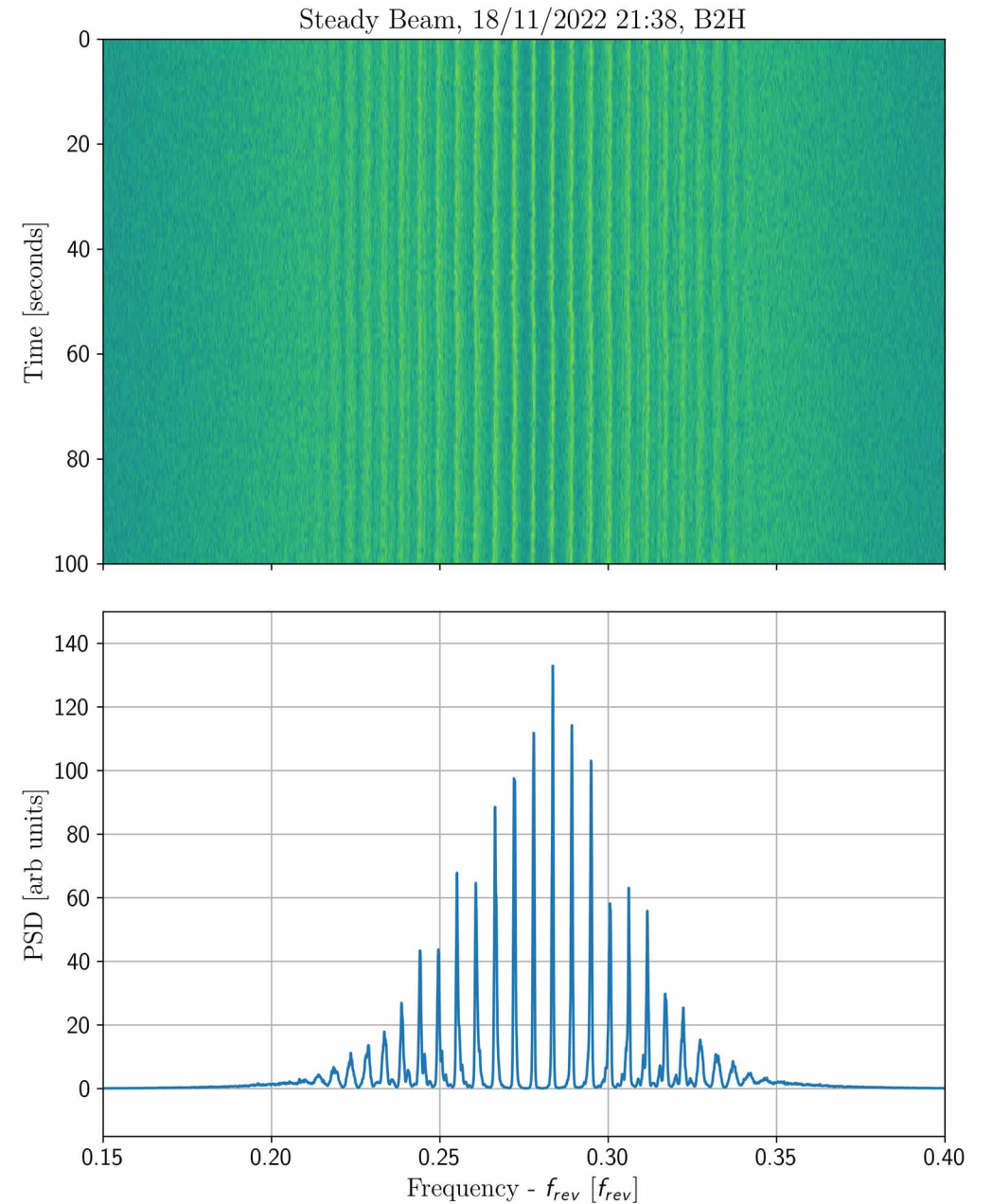
Details on the LHC Schottky system in
M. Betz et al., NIM, vol. 874, pp 113-126, 2017



Schottky spectra examples

I. Steady beam conditions:

- Mostly at flattop energy of ion fills, shorter periods at flatbottom
- Easy to analyse: just use the theory



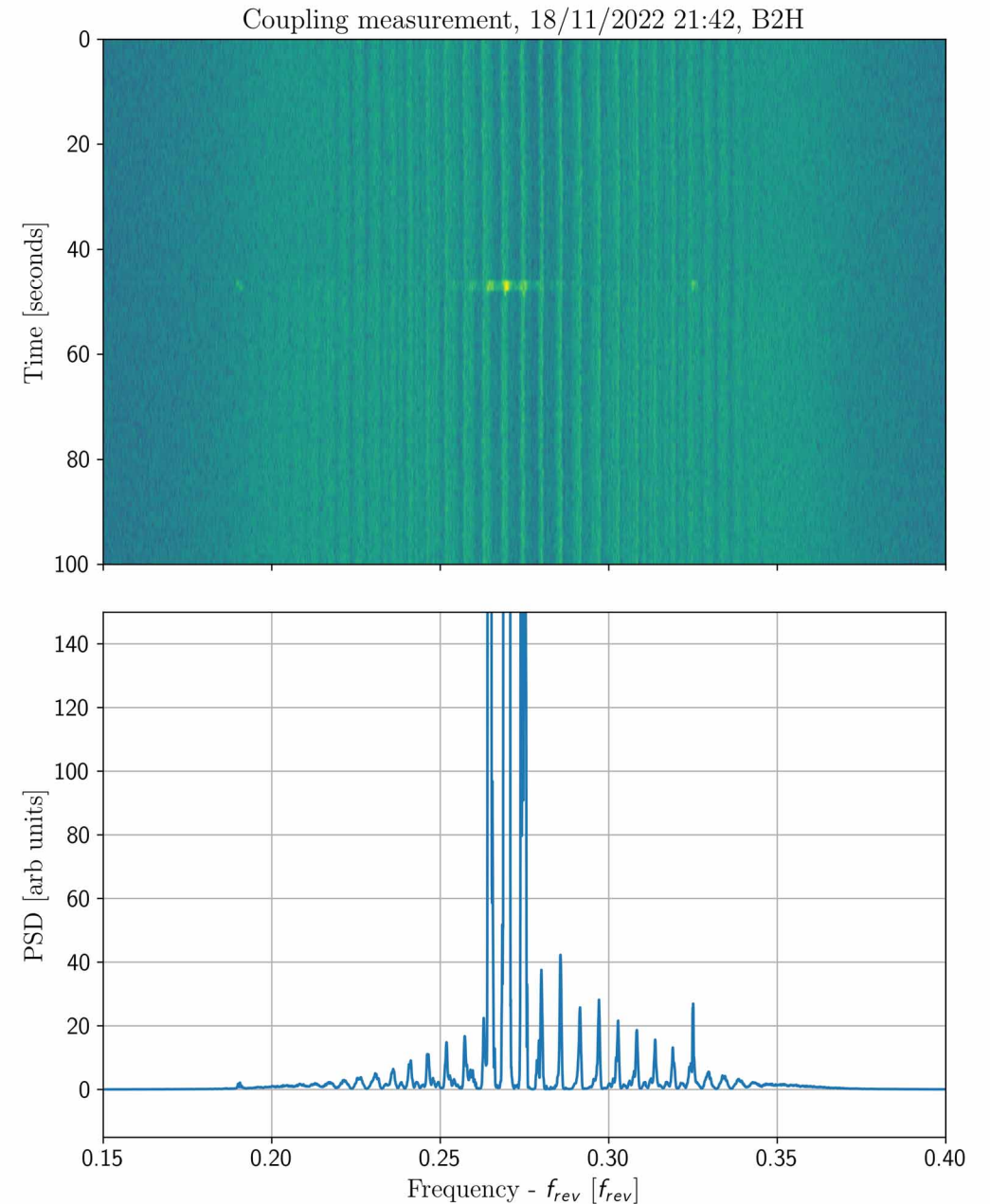
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II. Local distortions:

- Caused by residual coherence, beam parameter measurements, beam interaction with the surrounding
- Theory cannot be directly used



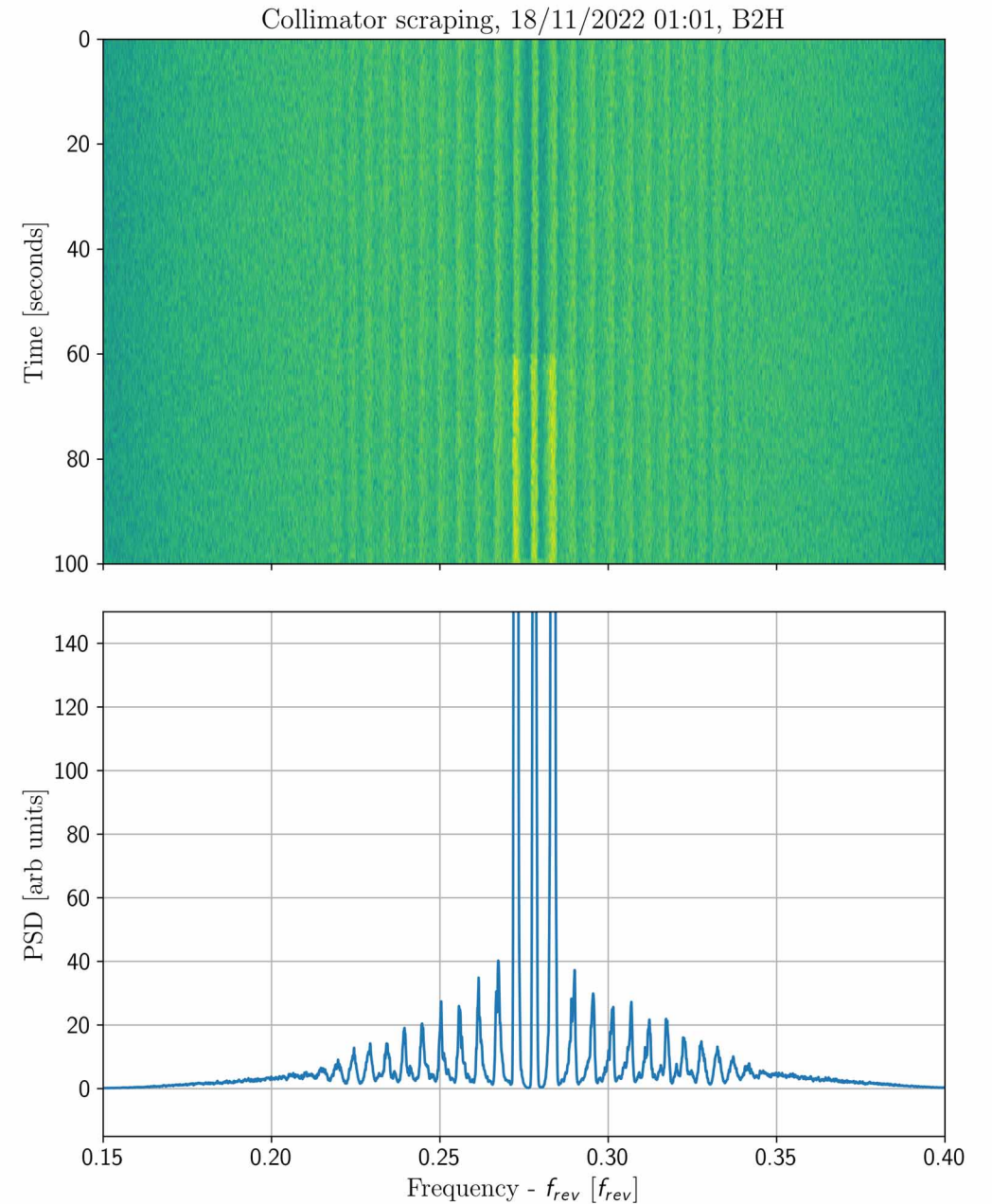
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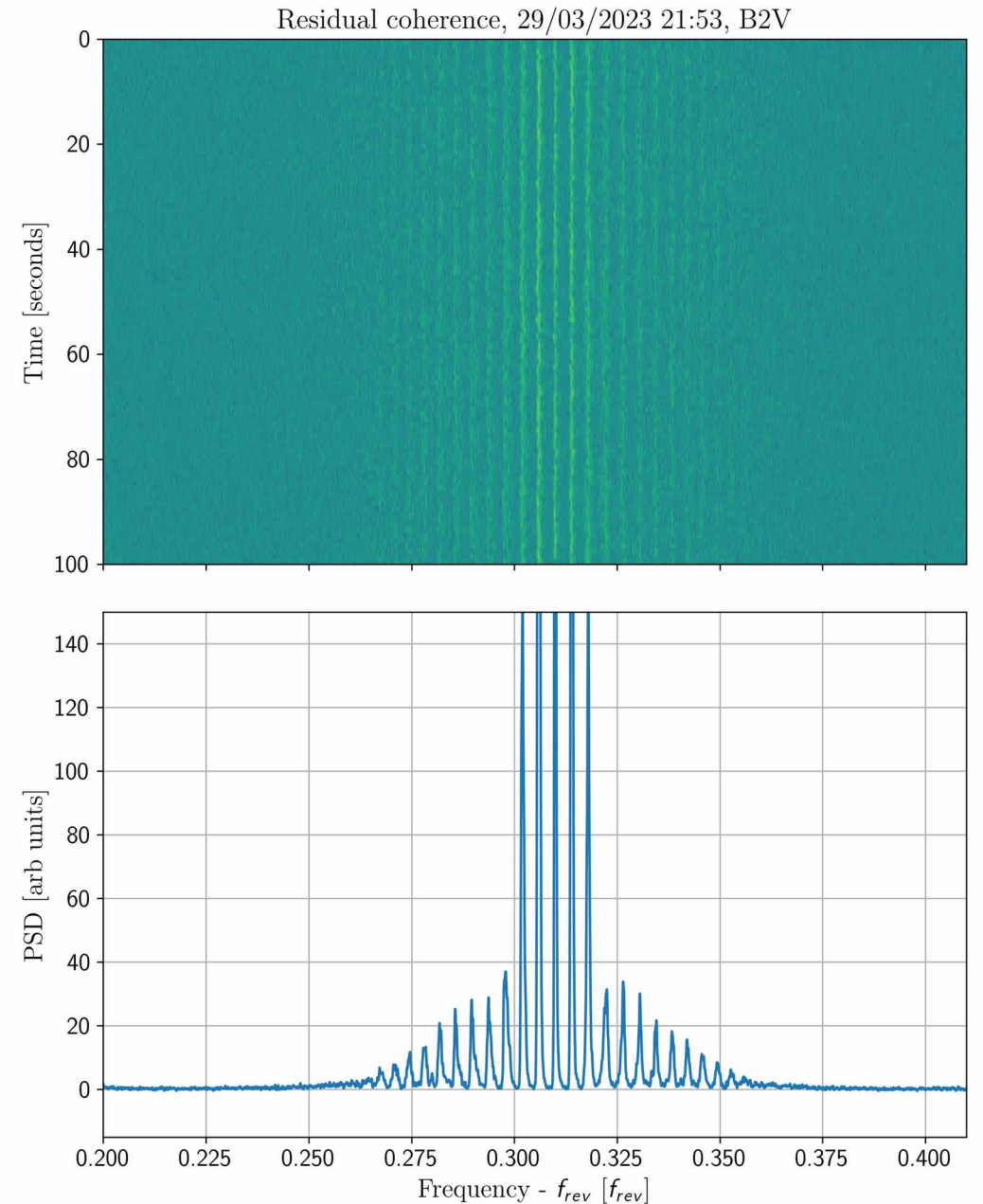
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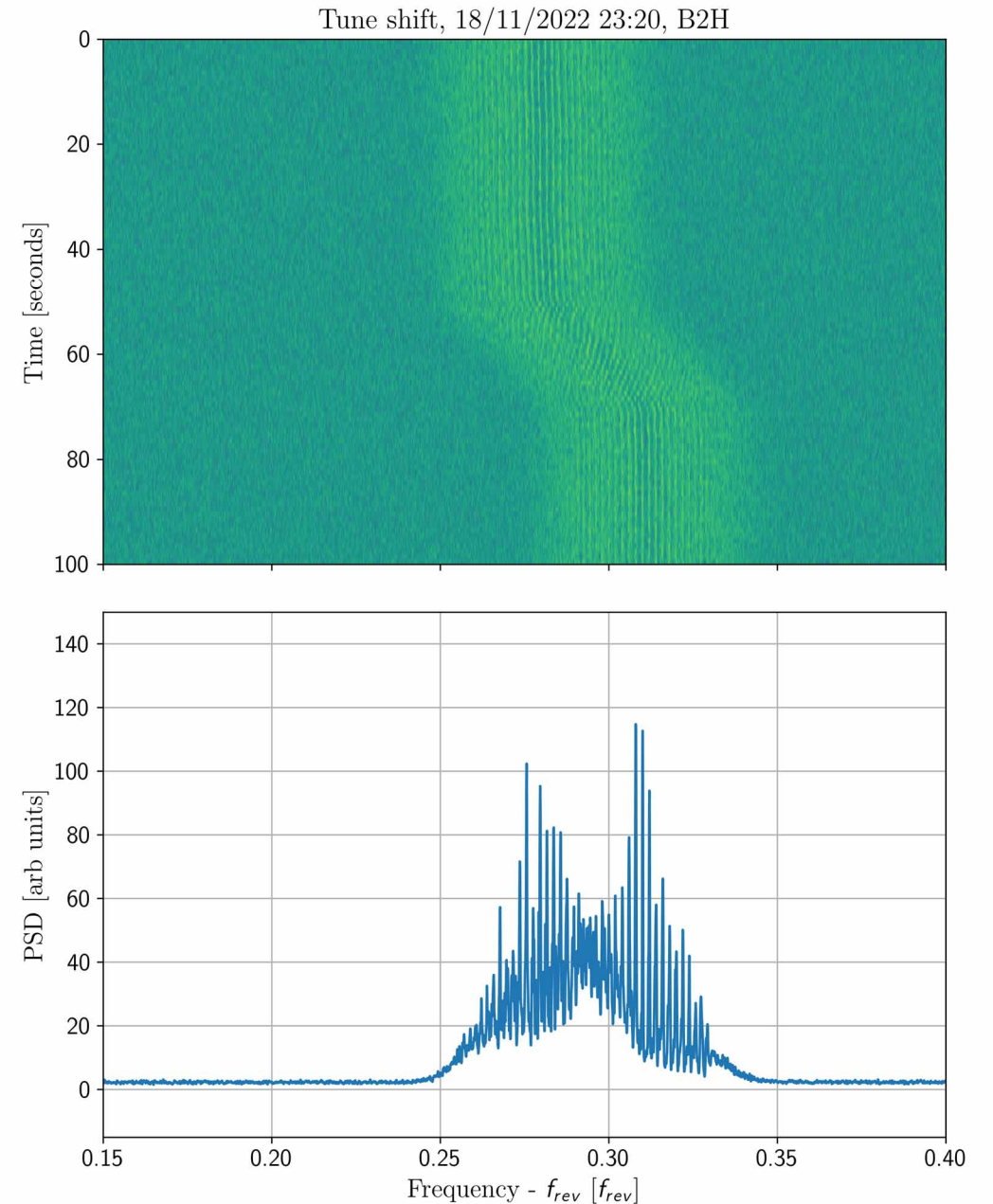
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- Tune shifts, RF modulation, energy ramp
- Theory cannot be directly used to averaged spectra, but spectrograms are easy to read



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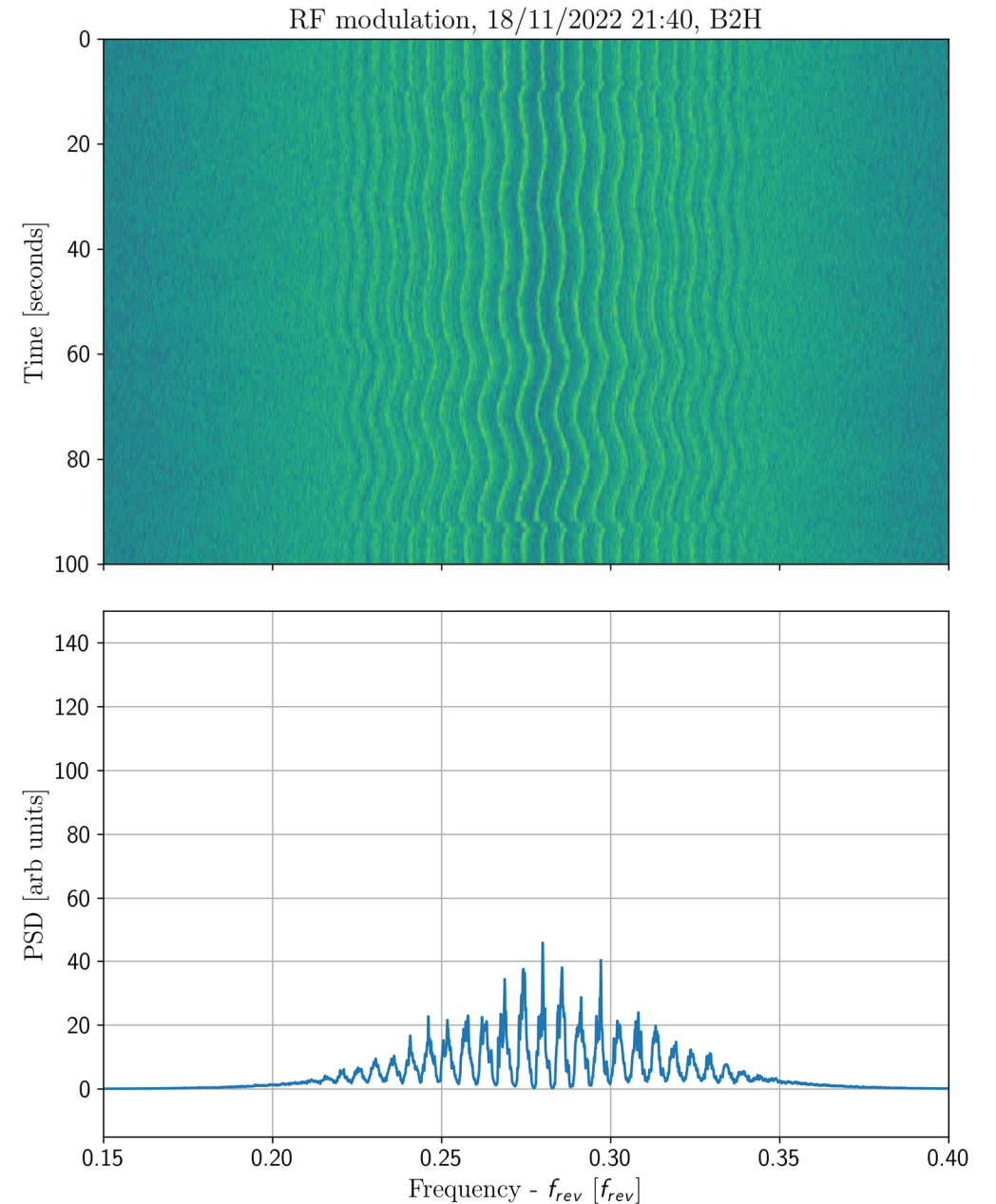
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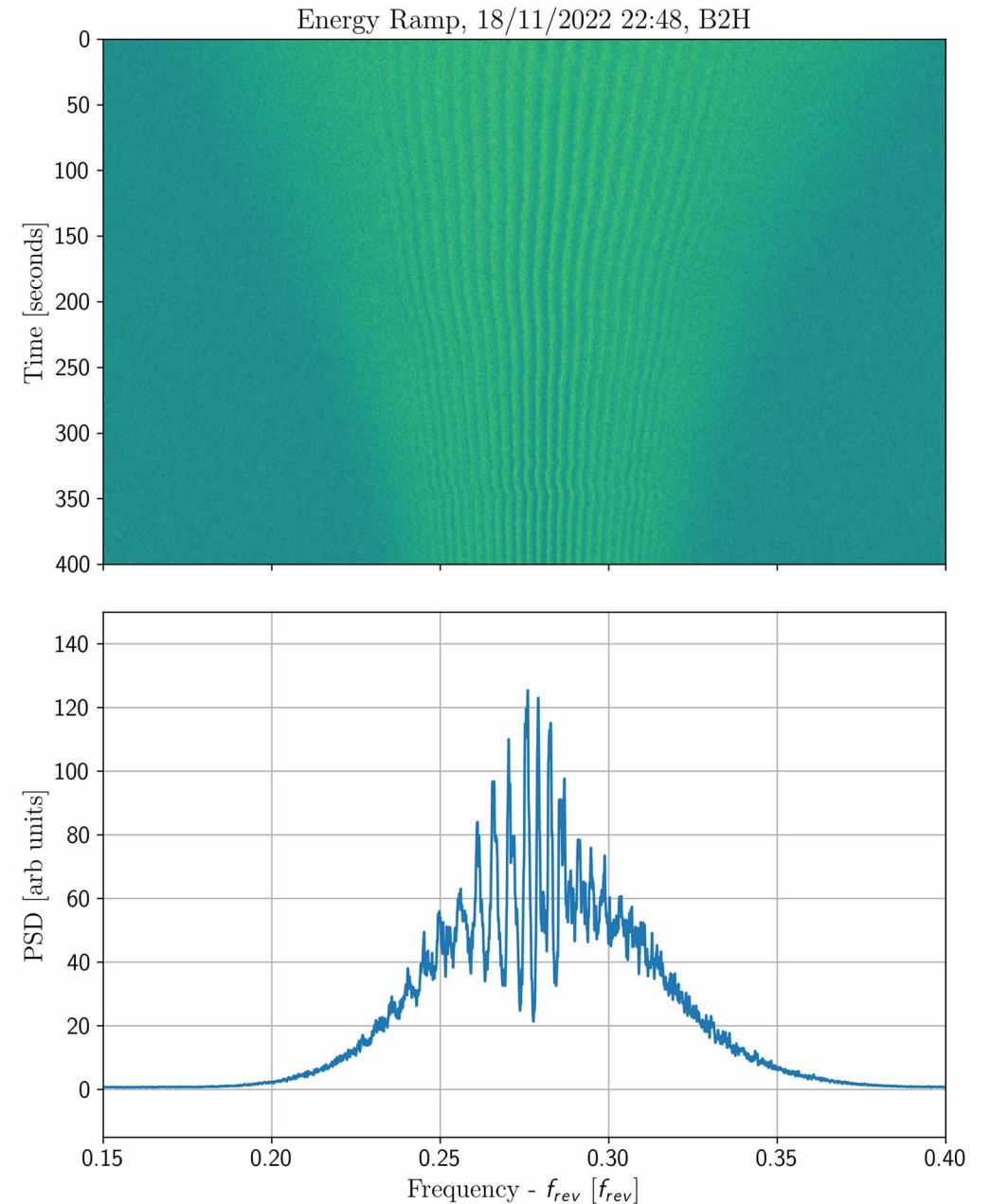
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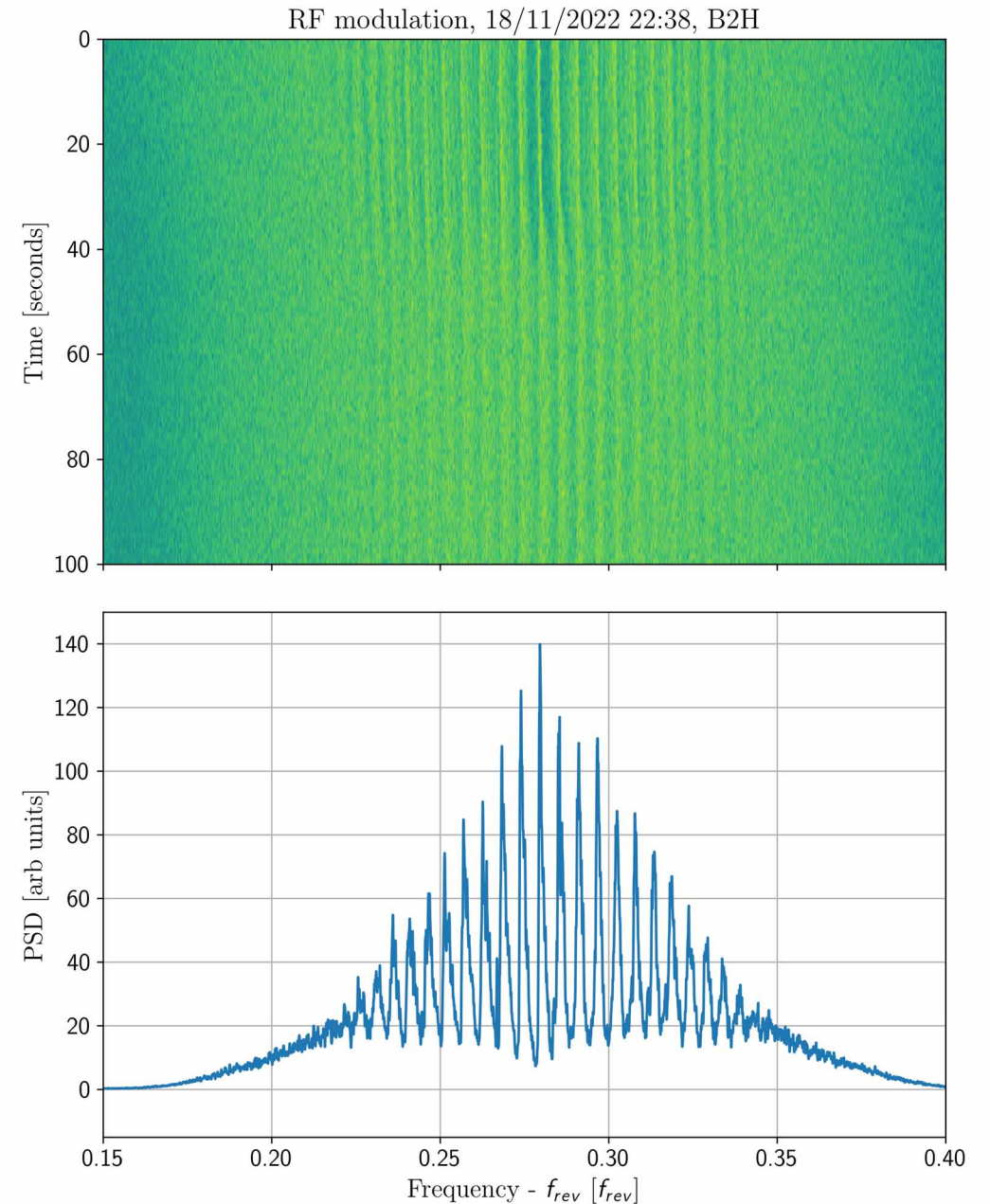
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IV. Beyond current theory effects:

- Octupole magnets, betatron coupling, impedance, all instrument problems
- Theory to analyse such spectra is still to be developed



Schottky spectra examples

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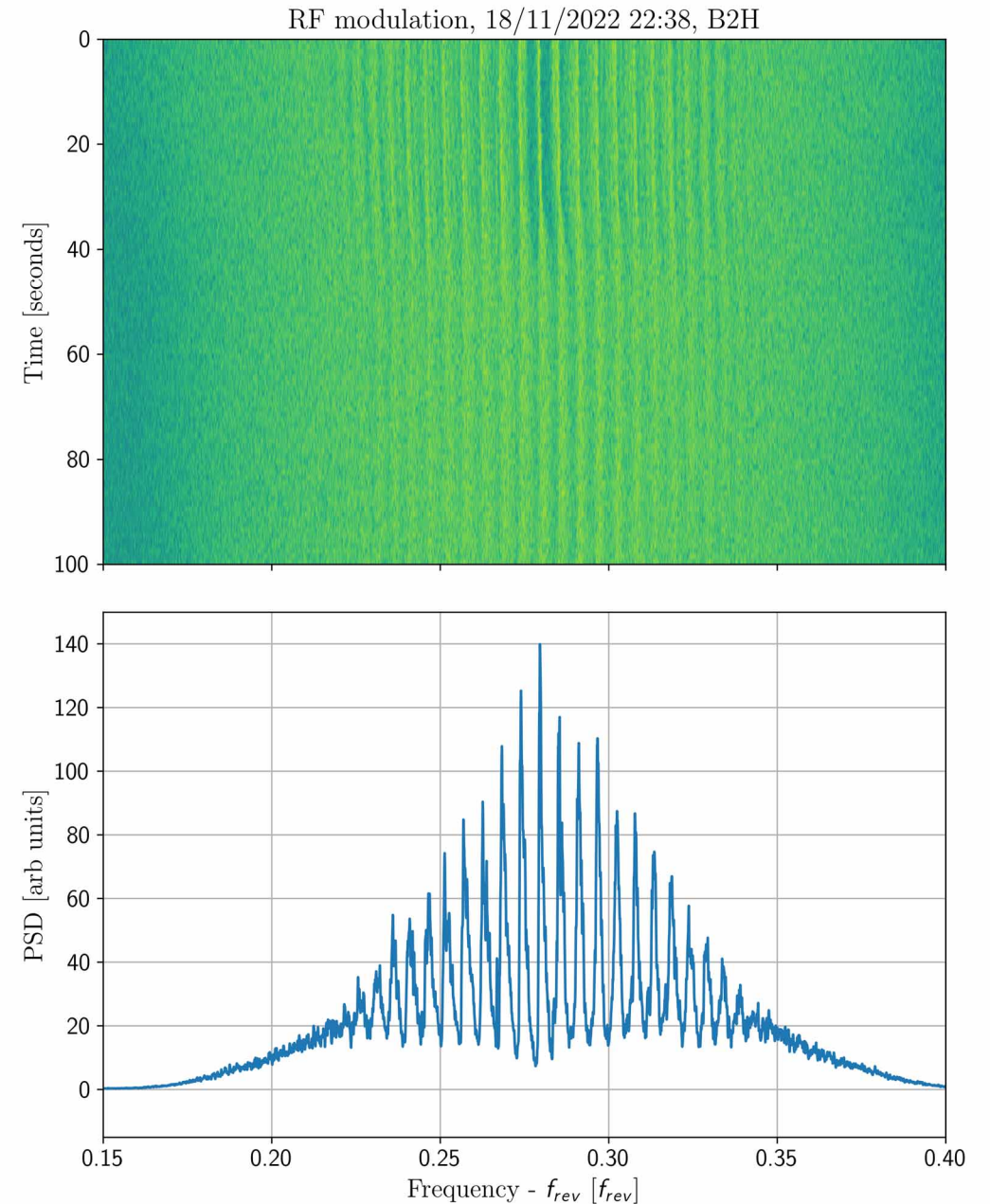
- Caused by residual coherence, beam parameter measurements, beam interaction with the surrounding
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III. Transient effects: Under development

- Tune shifts, RF modulation, energy ramp
- Theory cannot be directly used to averaged spectra, but spectrograms are easy to read

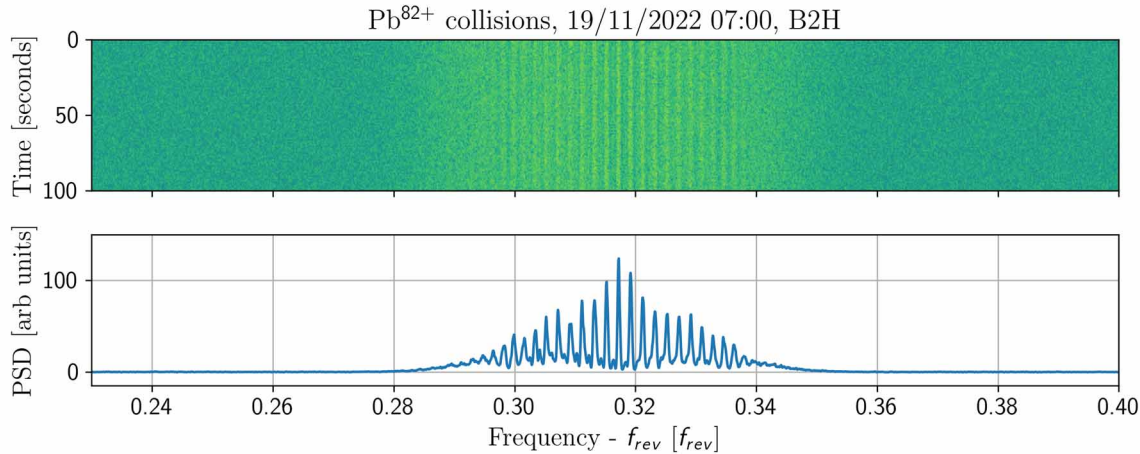
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Signal analysis in stable conditions

Example: 8 hour long ion collisions in Nov 2022

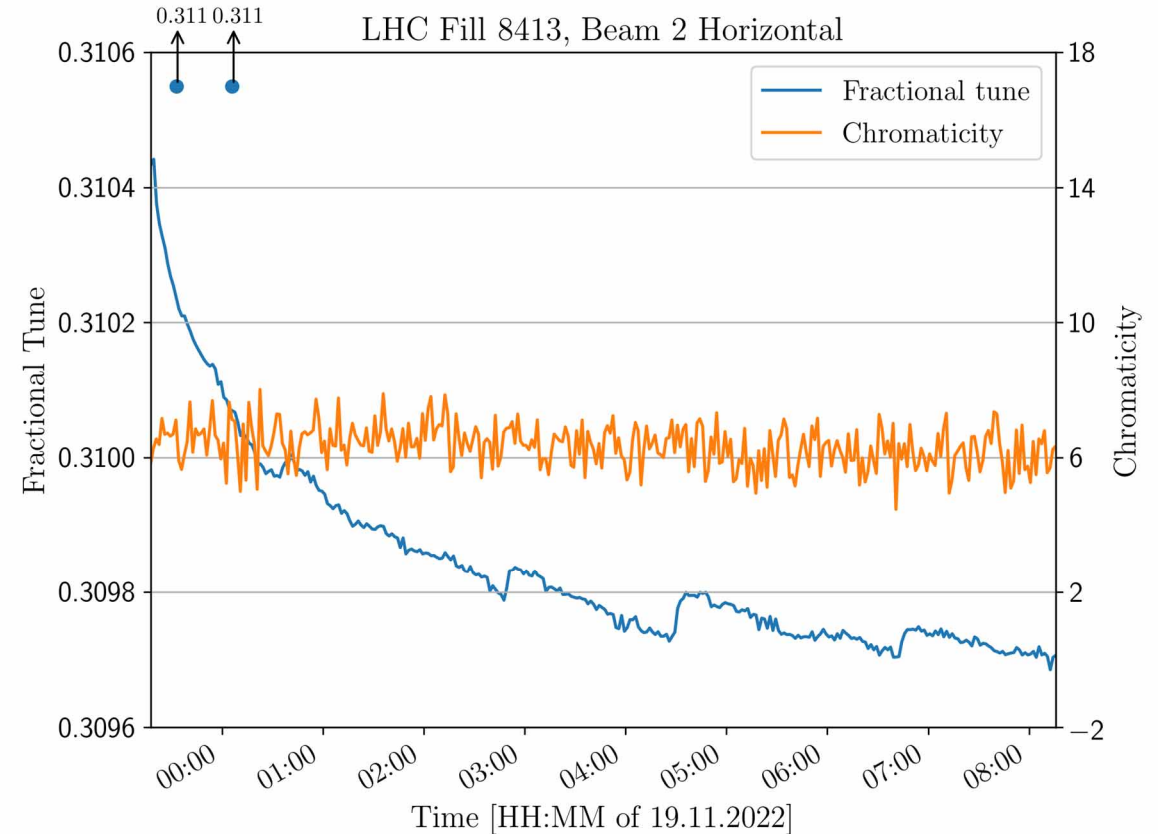


Betatron tune

$$C_{MD}(k) = \sum_{i=1}^{i=N} |P_T^{\pm}(\omega_{k-i}) - P_T^{\pm}(\omega_{k+i})|$$

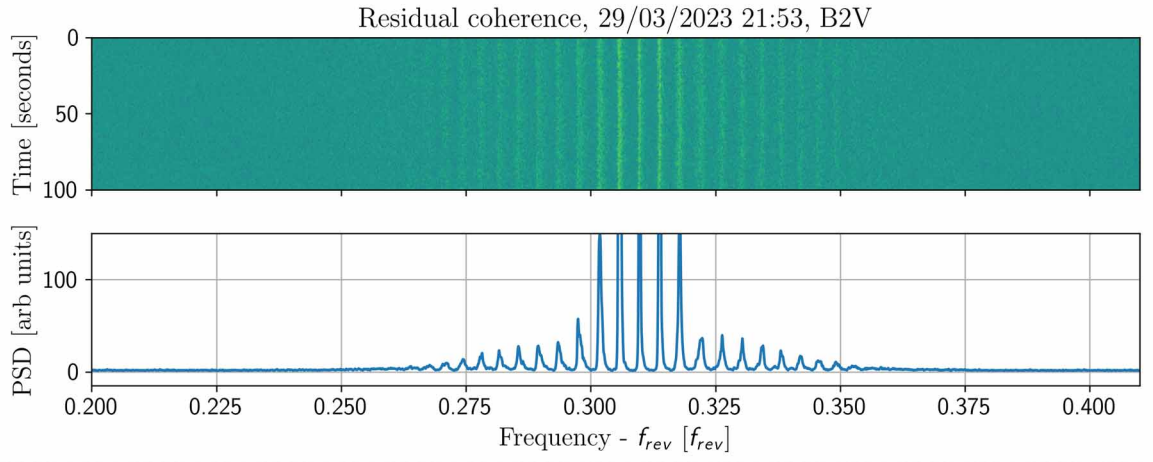
Chromaticity

$$Q_{\xi} = -\eta \left(n \frac{\Delta f_{-} - \Delta f_{+}}{\Delta f_{-} + \Delta f_{+}} - Q_I \right)$$



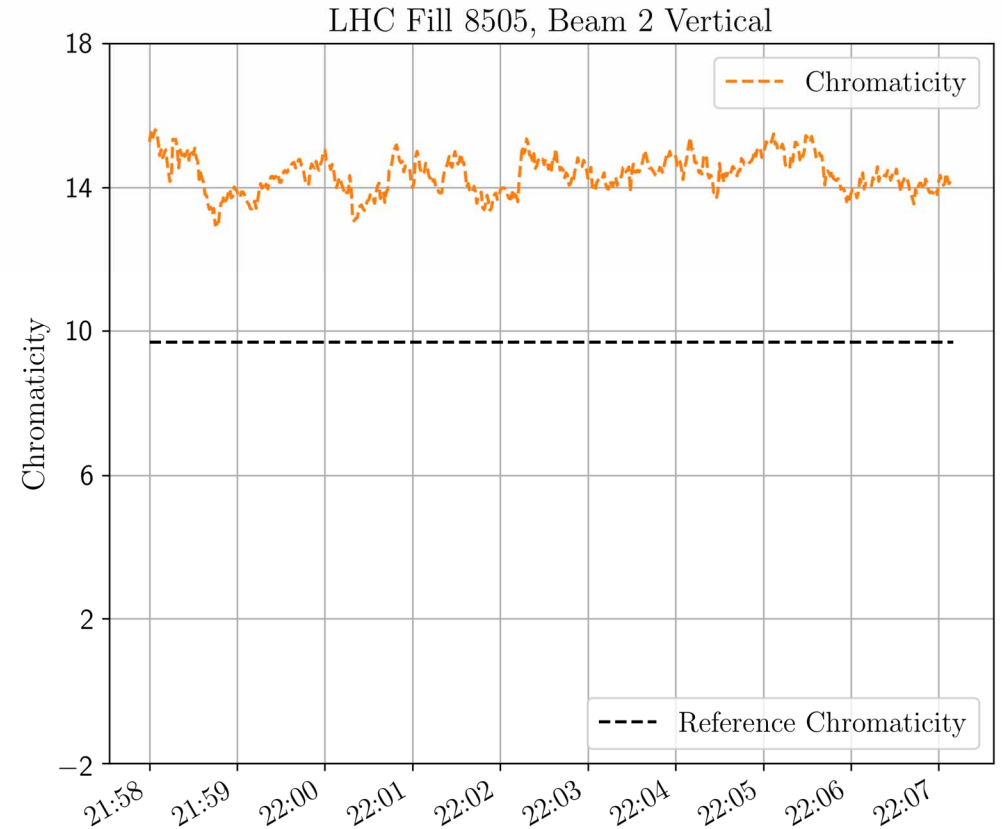
Signal analysis in presence of local distortions

Example: Early proton fill in March 2023



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Offset of over 4 units...

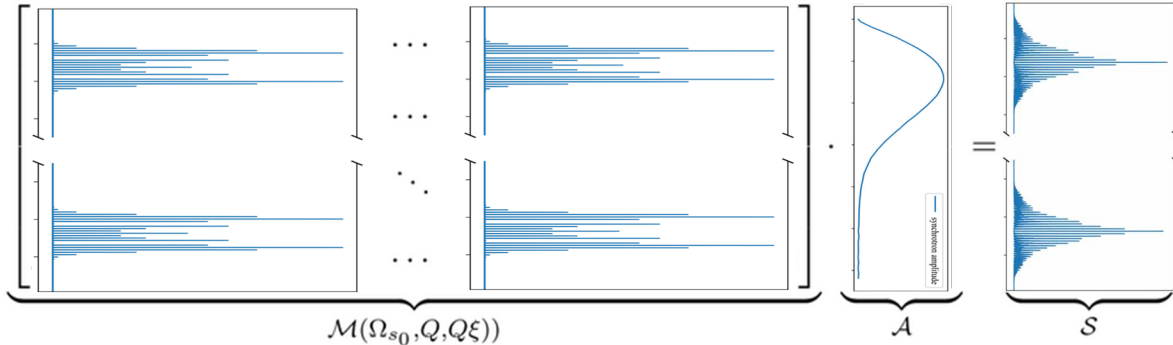
Matrix formalism for Schottky spectra

Mathematically, Schottky spectra are given as a function of:

- Synchrotron amplitude distribution - 2 parameters
- Nominal synchrotron frequency - 1 parameter
- Betatron tune - 1 parameter
- Chromaticity - 1 parameter

For given parameters the spectrum can be calculated with a simple matrix transform.

$$\underbrace{\begin{bmatrix} P_T^\pm(\omega_1, \hat{\tau}_1, \Omega_{s0}, Q, Q\xi) & \cdots & P_T^\pm(\omega_1, \hat{\tau}_n, \Omega_{s0}, Q, Q\xi) \\ P_T^\pm(\omega_2, \hat{\tau}_1, \Omega_{s0}, Q, Q\xi) & \cdots & P_T^\pm(\omega_2, \hat{\tau}_n, \Omega_{s0}, Q, Q\xi) \\ \vdots & \ddots & \vdots \\ P_T^\pm(\omega_m, \hat{\tau}_1, \Omega_{s0}, Q, Q\xi) & \cdots & P_T^\pm(\omega_m, \hat{\tau}_n, \Omega_{s0}, Q, Q\xi) \end{bmatrix}}_{\mathcal{M}(\Omega_{s0}, Q, Q\xi)} \cdot \underbrace{\begin{bmatrix} \tilde{g}(\hat{\tau}_1) \\ \tilde{g}(\hat{\tau}_2) \\ \vdots \\ \tilde{g}(\hat{\tau}_n) \end{bmatrix}}_{\mathcal{A}} = \underbrace{\begin{bmatrix} P_T^\pm(\omega_1) \\ P_T^\pm(\omega_2) \\ \vdots \\ P_T^\pm(\omega_m) \end{bmatrix}}_{\mathcal{S}}$$



Details in: K. Lasocha and D. Alves, Phys. Rev. Accel. Beams 23, 062803
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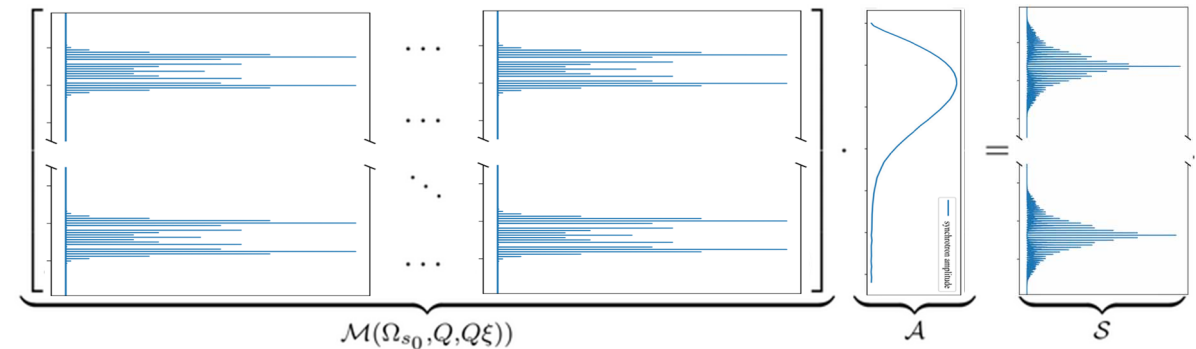
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Use case 1: fast Schottky spectra simulation

Use case 2 (spectra fitting): given an experimentally measured spectrum, true parameters would minimize the cost function:

$$C(\Omega_{s0}, Q, Q\xi, \mathcal{A}) = |\mathcal{M}(\Omega_{s0}, Q, Q\xi) \cdot \mathcal{A} - [\mathcal{S}_{exp}]|^2$$

Minimizing routines iteratively simulate Schottky spectra and compare them with the measurement.

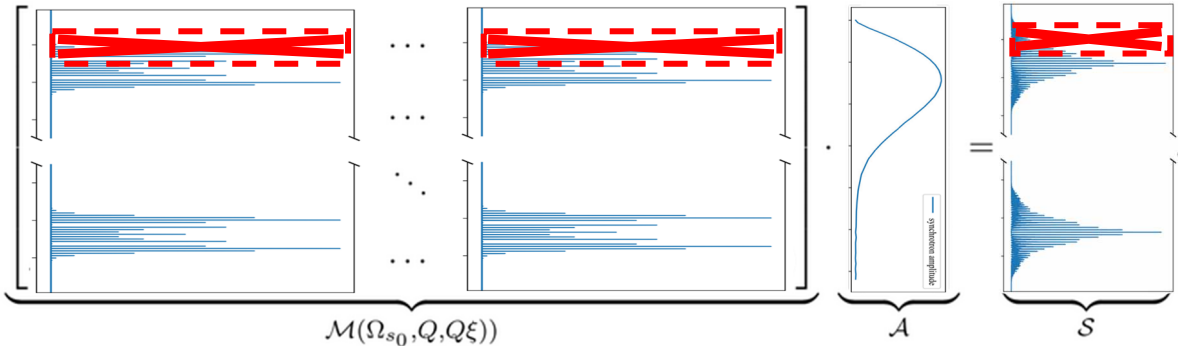


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Matrix formalism: excluding frequency bins

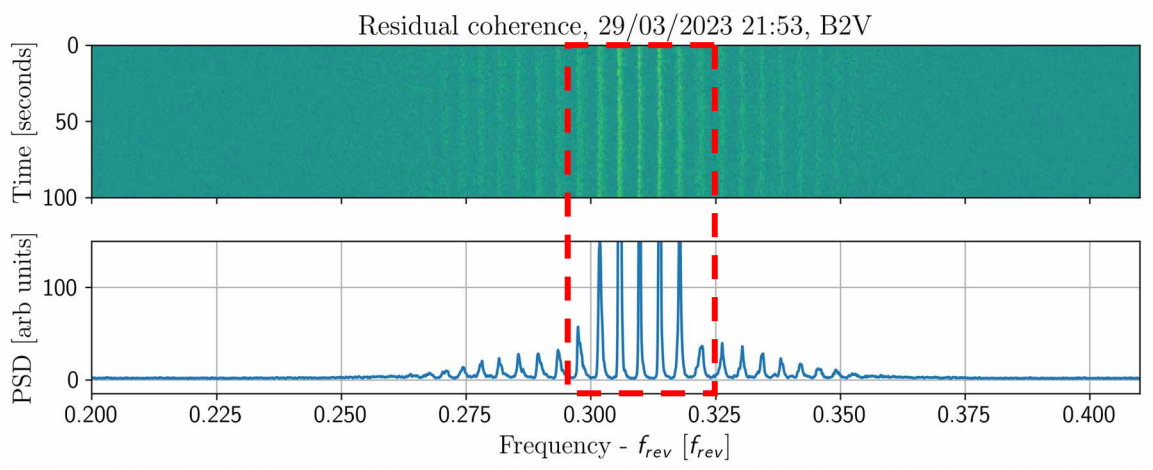
Contrary to previous methods, fitting procedure allow to exclude the spectral regions with undesired components.

$$\underbrace{\begin{bmatrix} P_T^\pm(\omega_1, \hat{\tau}_1, \Omega_{s0}, Q, Q\xi) & \cdots & P_T^\pm(\omega_1, \hat{\tau}_n, \Omega_{s0}, Q, Q\xi) \\ \text{---} & & \text{---} \\ \vdots & \ddots & \vdots \\ P_T^\pm(\omega_m, \hat{\tau}_1, \Omega_{s0}, Q, Q\xi) & \cdots & P_T^\pm(\omega_m, \hat{\tau}_n, \Omega_{s0}, Q, Q\xi) \end{bmatrix}}_{\mathcal{M}(\Omega_{s0}, Q, Q\xi)} \cdot \underbrace{\begin{bmatrix} \tilde{g}(\hat{\tau}_1) \\ \tilde{g}(\hat{\tau}_2) \\ \vdots \\ \tilde{g}(\hat{\tau}_n) \end{bmatrix}}_A = \underbrace{\begin{bmatrix} P_T^\pm(\omega_1) \\ \text{---} \\ \vdots \\ P_T^\pm(\omega_m) \end{bmatrix}}_S$$



Signal analysis in presence of local distortions

Example: Early proton fill in March 2023



Betatron tune

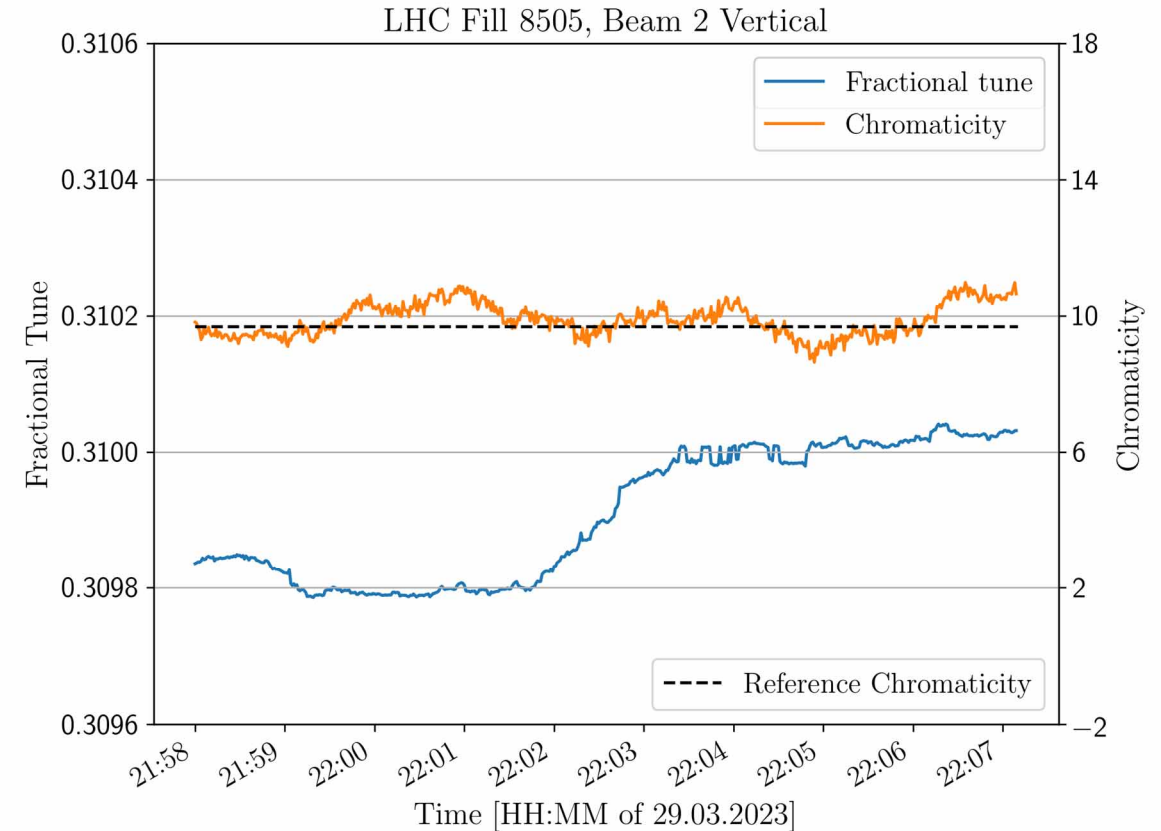
$$C_{MD}(k) = \sum_{i=1}^{i=N} |P_T^{\pm}(\omega_{k-i}) - P_T^{\pm}(\omega_{k+i})|$$

Only "valid" frequencies taken into sum

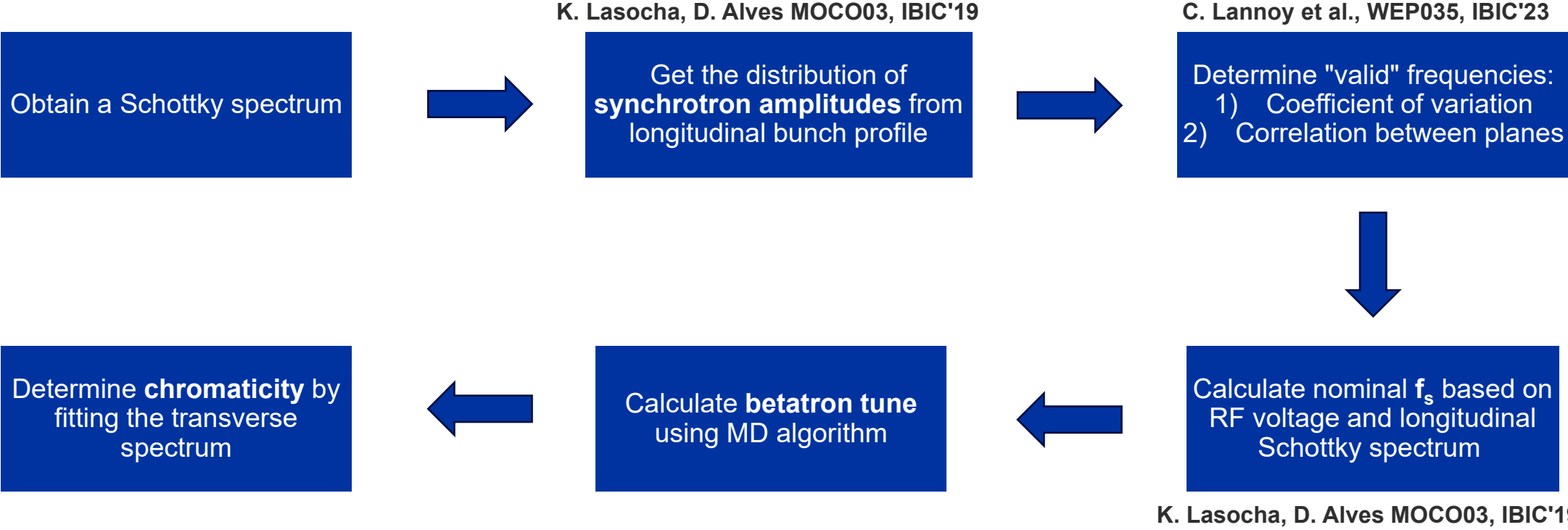
Chromaticity

$$C(\mathcal{A}, Q\xi) = |\mathcal{M}(Q\xi) \cdot \mathcal{A} - \mathcal{S}_{exp}|^2$$

Nominal synchrotron tune calculated independently,
Cost function minimization using Differential Evolution algorithm.



LHC Schottky online signal analysis pipeline



Implementation in the final stage of development, planned be in use in the end of 2023

Conclusions

Summary of the talk:

- Theory of transverse Schottky spectra reviewed and successfully applied to **stable beams**,
- **Local spectral distortions** mitigation technique proposed and demonstrated on proton spectra,
- Automation of the analysis will be tested in the coming days.

Possible next steps:

- Handling of **transient effects**: 1D ---> 2D analysis; image recognition techniques?
- Expanding the theory of Schottky spectra: **impedance**, **octupoles**, ... See C. Lannoy et al., WEP034, IBIC'23

Thanks for your attention!