



Longitudinal beam diagnostics R&D

at GSI UNILAC

HIAT 2022, Darmstadt



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GSI Darmstadt

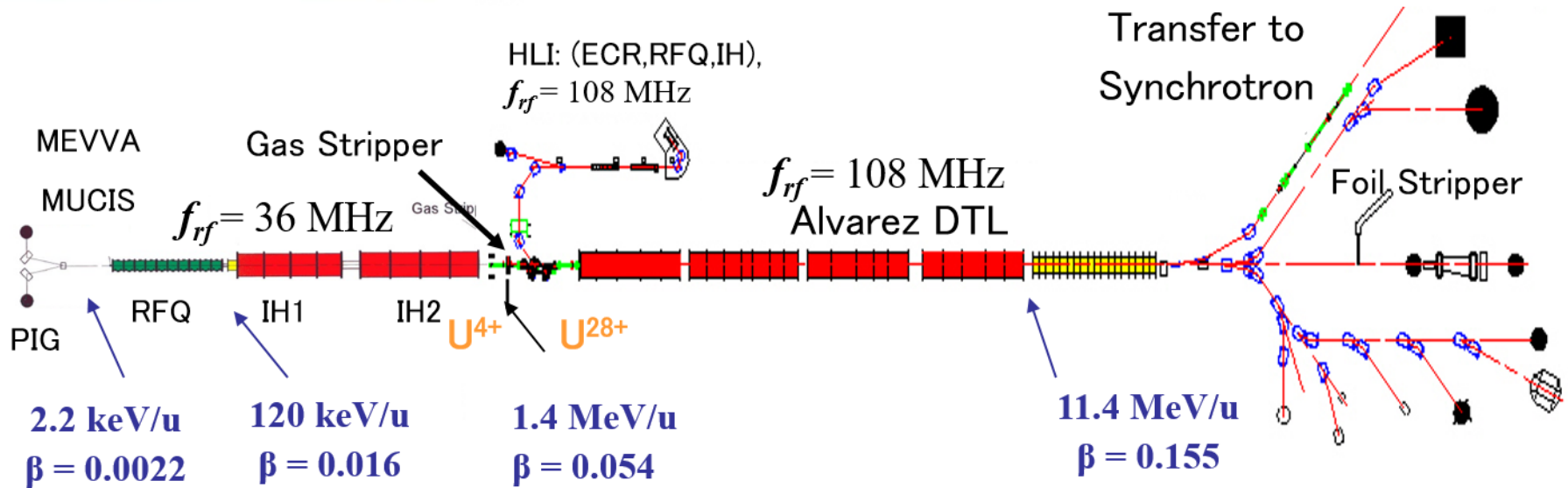
Contributors:

P. Forck, T. Reichert



- **Introduction to UNILAC longitudinal diagnostics**
 - Needs and challenges
 - Phase probes
- **Past/Existing devices:**
 - Particle detectors w.r.t. RF
 - Time-to-space conversion: Feshenko BSM , Gas ionization BSM
- **Recent developments: Direct time-domain**
 - Fast Faraday Cups (FFC)
 - GHz Transition Radiation (GTR)
- **Conclusion/Outlook**

UNILAC



- UNILAC: Complex set of resonators (RFQ, IH and Alvarez) with charge stripping sections. **Two injectors** HLI and HSI served by several ion sources
- Upto 50 Hz operation, several beam types and parameters available in time multiplexed or "parallel" operation
- Significant upgrades over the last years to achieve FAIR parameters → RFQ electrodes, H2 gas strippers, beam brilliance optimization

[1] W. Barth et al., High brilliance beam investigations at Universal linear accelerator, PRAB 25, 04101 (2022)

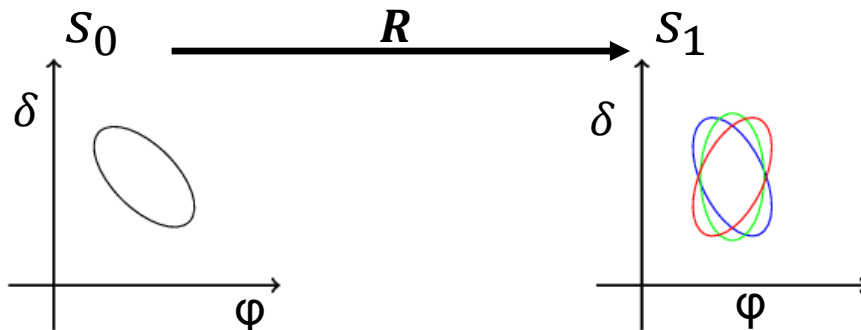
Longitudinal diagnostics and emittance



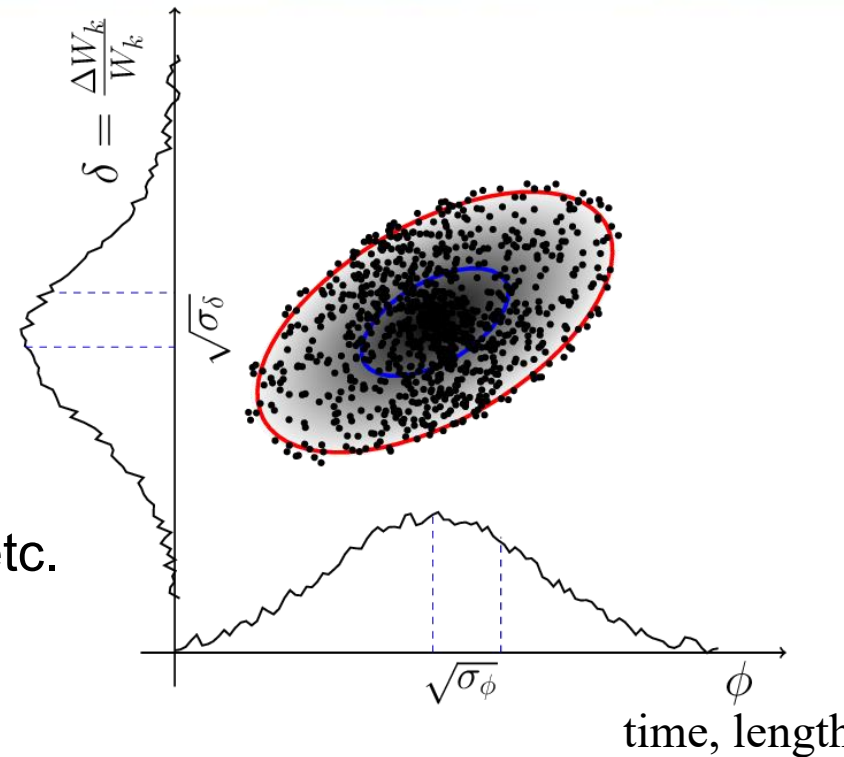
- Energy and time of arrival with respect to RF form correlated distributions
- Emittance represents the area of the phase space ellipse

$$\begin{pmatrix} \varphi \\ \delta \end{pmatrix}_{s_1} = \mathbf{R} \begin{pmatrix} \varphi \\ \delta \end{pmatrix}_{s_0}$$

Transfer matrix \mathbf{R} represents buncher, drift etc.

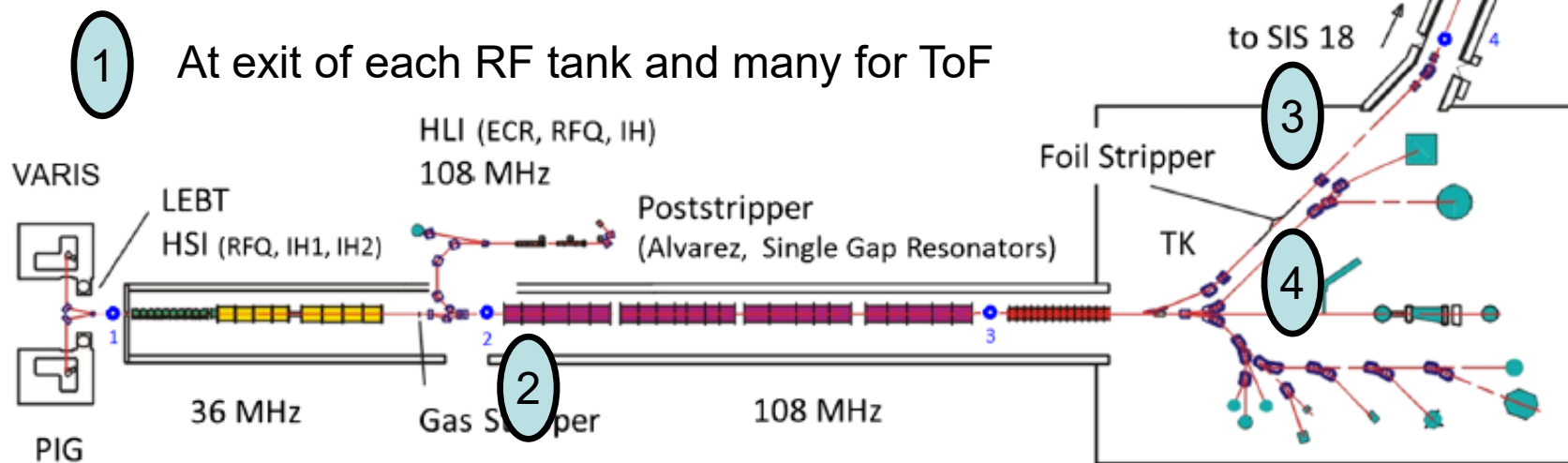


Rotation of the ellipse in phase space



- Multiple measurements of one projection enough to reconstruct both distributions → **Tomography**
- Typically only time/phase w.r.t to RF measured

Longitudinal diagnostics at UNILAC



Need for longitudinal diagnostics : Frequent changes

- Injection from HSI into the Alvarez
- Min. energy spread while injection into the SIS-18
- Max. energy spread at stripper for countering straggling effects

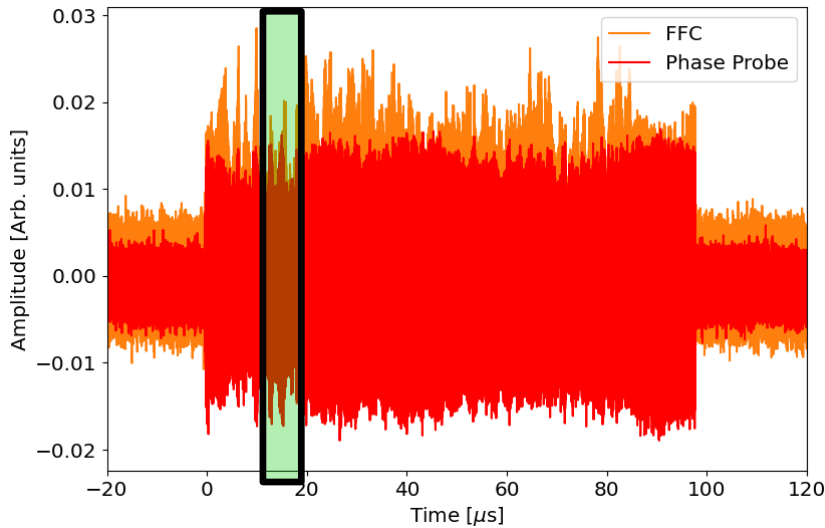
Devices for longitudinal diagnostics

- | | |
|---------------------------|---|
| 1 Phase probes / Pick-ups | 3 Residual gas ionization monitors |
| 2 Particle detectors | 4 RF deflector and dispersion with screen |

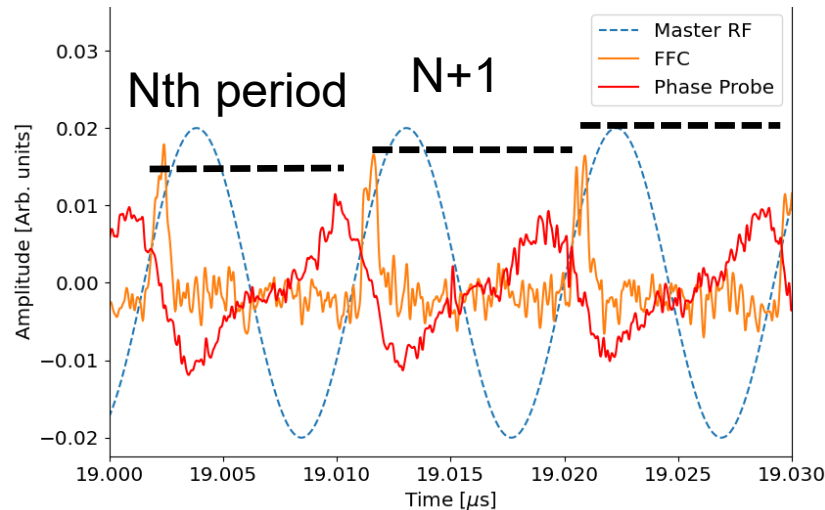
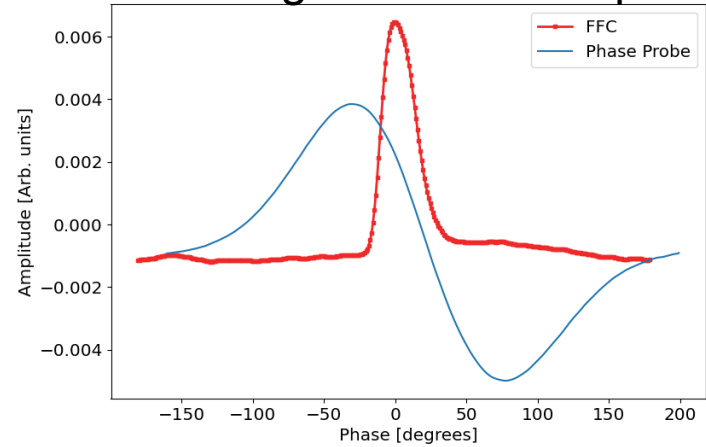
Macropulse and bunch phase w.r.t RF



100 μA He^{1+} after HLI @ 108 MHz rf

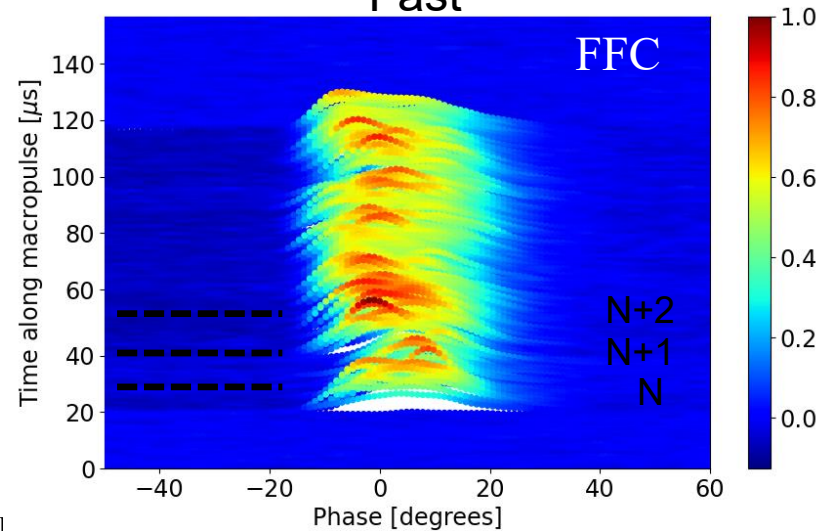


Slow averaged over macropulse

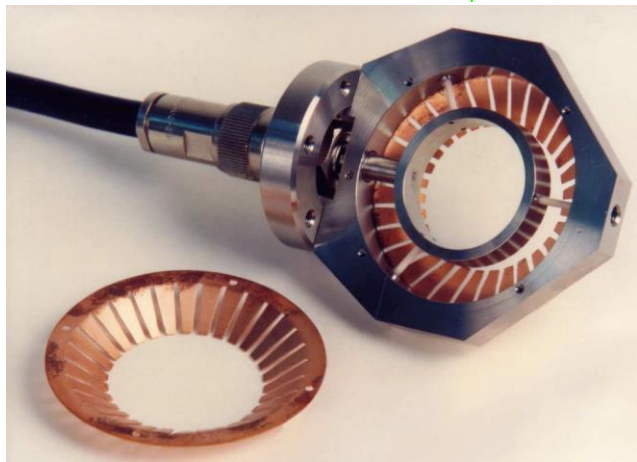
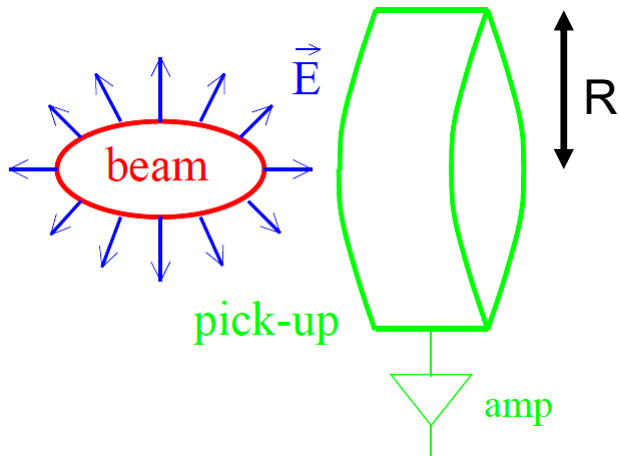


[2] R. Singh et al., Proc of IB...

Fast



Low velocity effect on phase probes

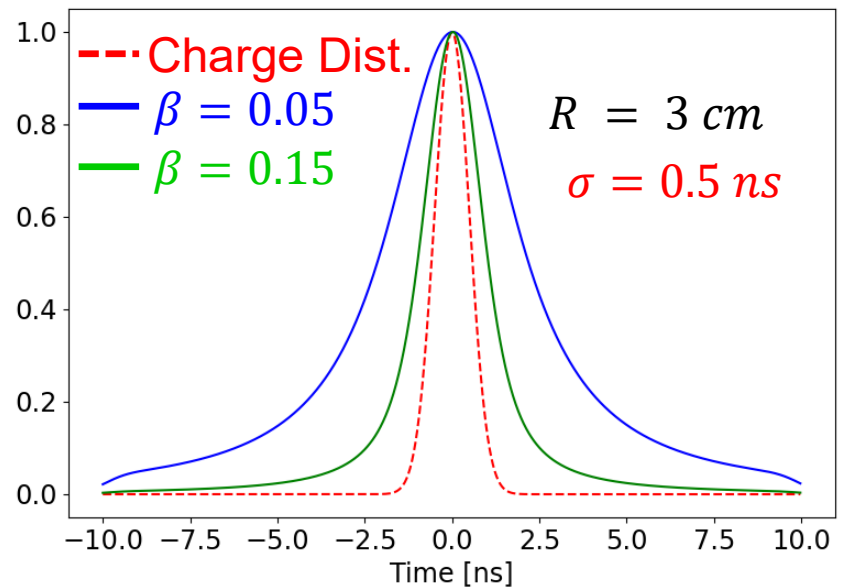


[4] P. Forck: Lecture notes on beam instrumentation

Lorentz boost **and** transformation of time for point charge:

Trans. \mathbf{E}_\perp lab.-frame of a point charge:

$$E_\perp(t) = \frac{e}{4\pi\epsilon_0} \cdot \frac{\gamma R}{\left[R^2 + (\gamma\beta ct)^2 \right]^{3/2}}$$

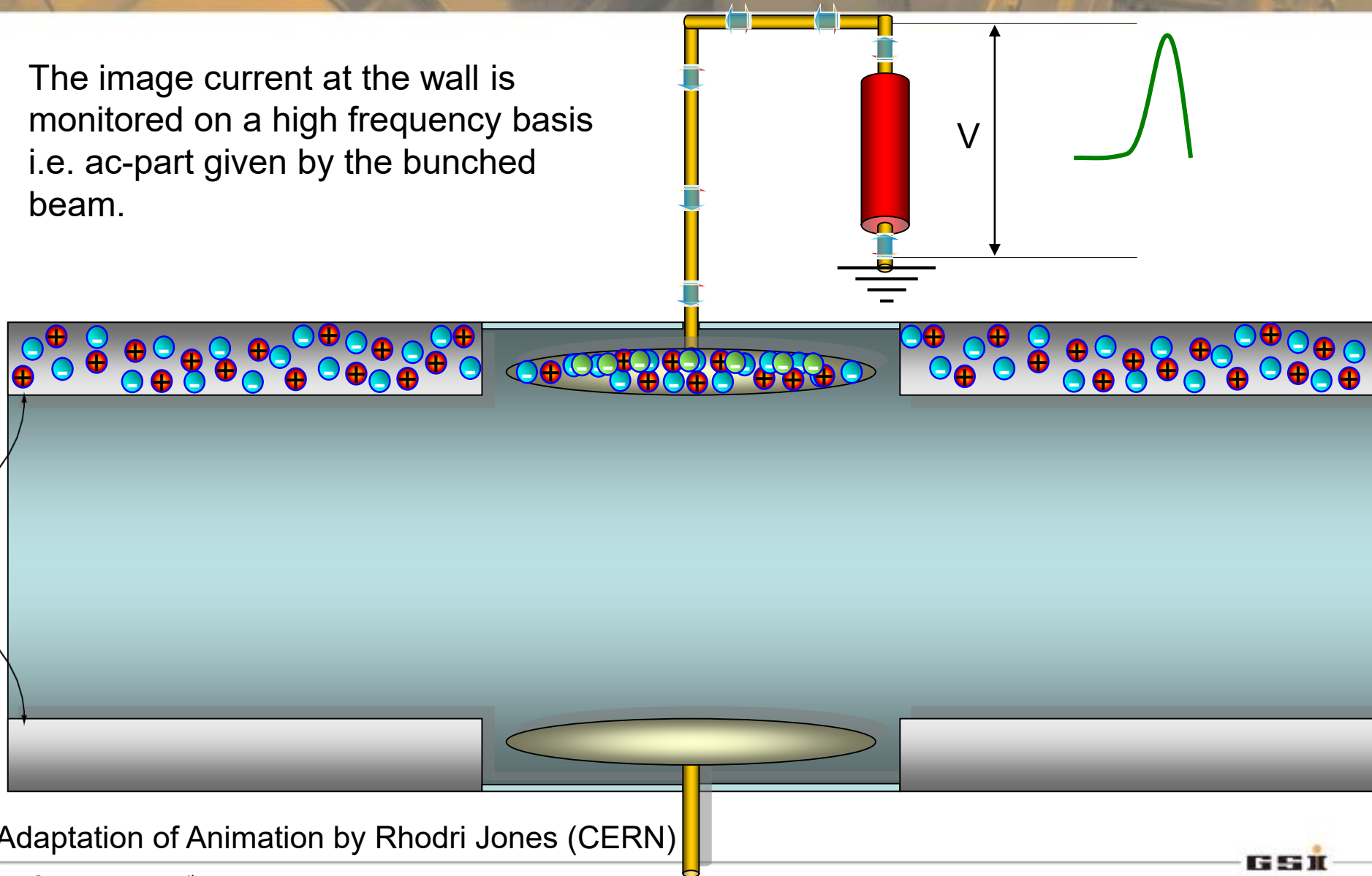


For $\beta < 1 \rightarrow$ Field distribution is not the same as charge distribution. Effect visible for shorter bunches $<$ few ns

Signal Generation in Phase Probes ($\beta < 1$)



The image current at the wall is monitored on a high frequency basis i.e. ac-part given by the bunched beam.

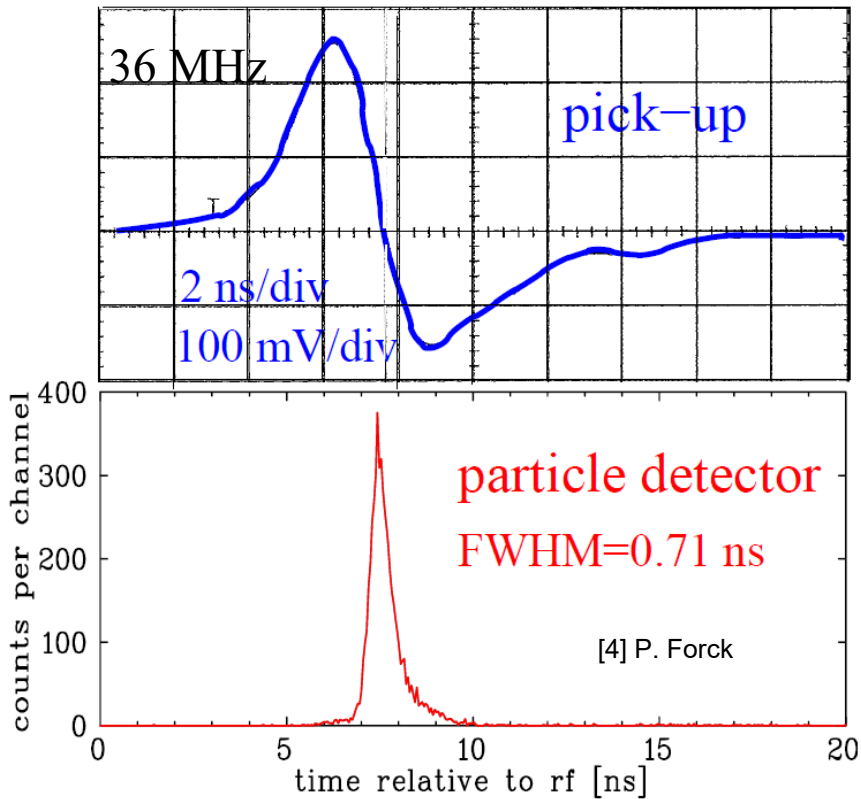


Adaptation of Animation by Rhodri Jones (CERN)

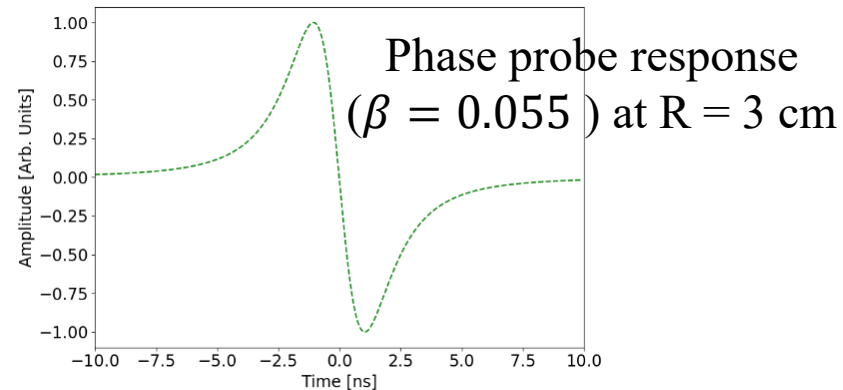
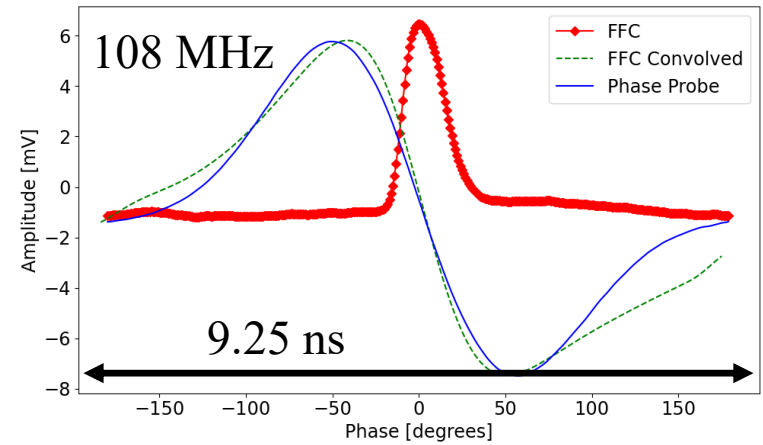
Bunch structure at low E_{kin}



pick-up – particle counter:
 Ar^{1+} with 1.4 MeV/u ($\beta = 5.5\%$)



pick-up – FFC:
 He^{1+} with 1.4 MeV/u ($\beta = 5.5\%$)



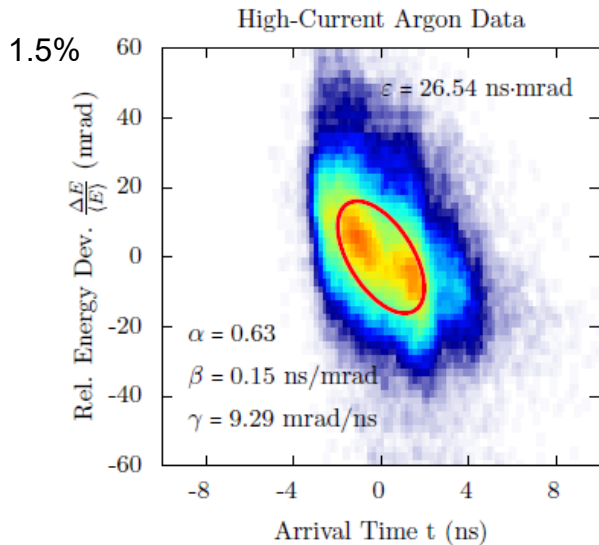
Pick-ups are used for:

- Bunch-center relative to rf for the whole macropulse
- Does not show the detailed bunch structure although **directly related it**

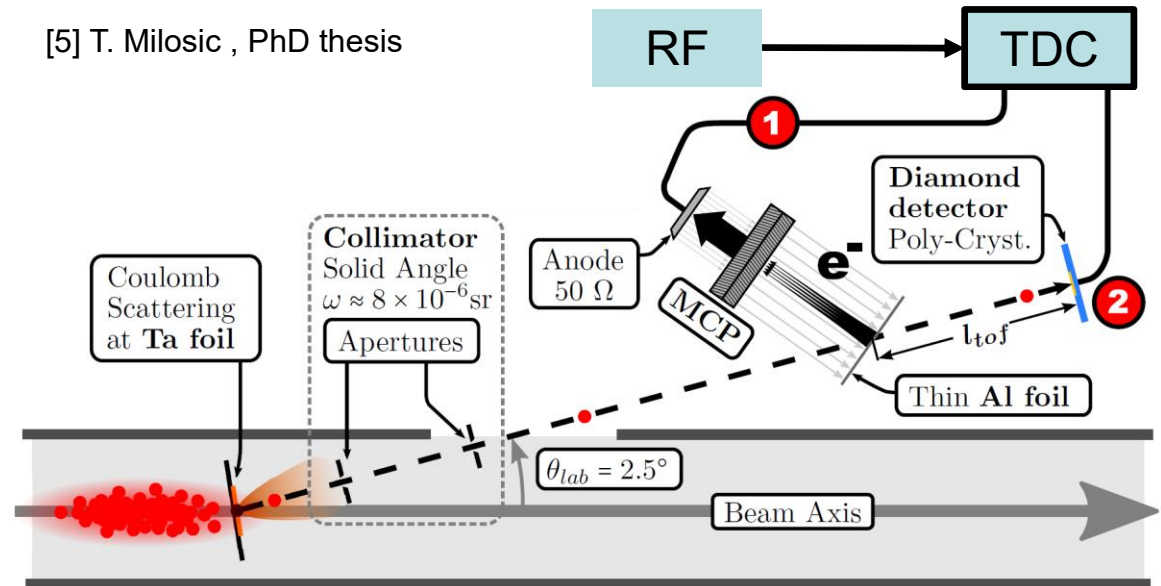
Particle detectors as bunch shape monitor



The time of arrival of the particle is determined relative to the accelerating rf:



[5] T. Milosic , PhD thesis



Foil (130 nm): attenuation $\approx 10^{-9}$ by Rutherford scat.

Start-detector: Thin Al foil (50 nm) for secondary e^- acc. toward an MCP +50 Ω anode

Stop-detector: Diamond detector with 1 ns pulse width

TDC: Time relative to rf, resolution less than 25 ps

Result: Correct determination of phase spread *but measured* energy spread much larger than expected values \rightarrow Foil non-uniformities, straggling on apertures.

Bunch structure using secondary electrons



Bunch Shape Monitor (BSM): Secondary e^- liberated from a wire or gas ionization carrying the time information.

[4] P. Forck, [6] B. Zwicker, PhD thesis

Working principle:

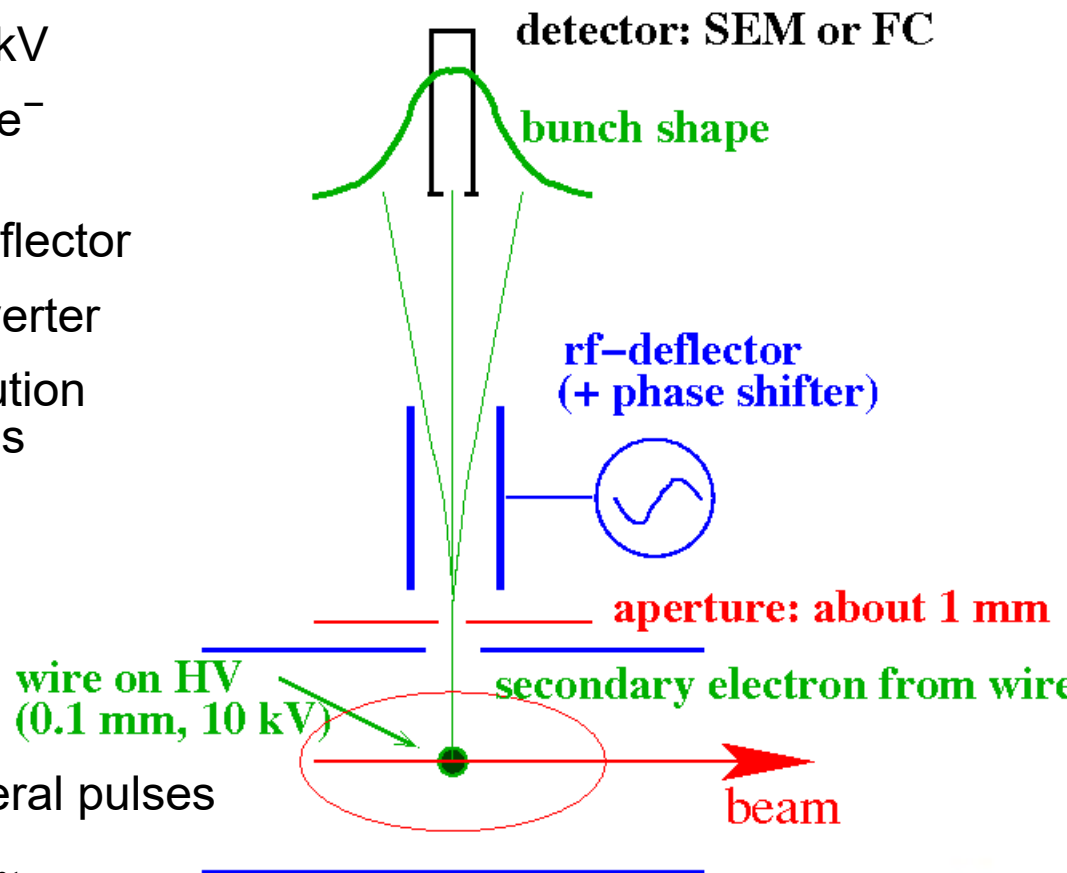
- insertion of a 0.1 mm wire at ≈ 10 kV
- immediate emission of secondary e^- (within < 10 ps)
- e^- are accelerated toward an rf-deflector
- rf-deflector as 'time-to-space' converter
- either slow shift of the phase resolution $\approx 1^\circ < 10$ ps or "streak" the electrons

Challenges:

- Stray fields and beam fields
- Complex installation/maintenance
- Current designs: Average over several pulses

SEM: secondary electron multiplier

detector: SEM or FC



[2] Comparison with FFC, R. Singh et al., Proc of IBIC 2021

Fast Faraday Cups (FFC)



Faraday cup designed to measure fast bunch structures

Challenges:

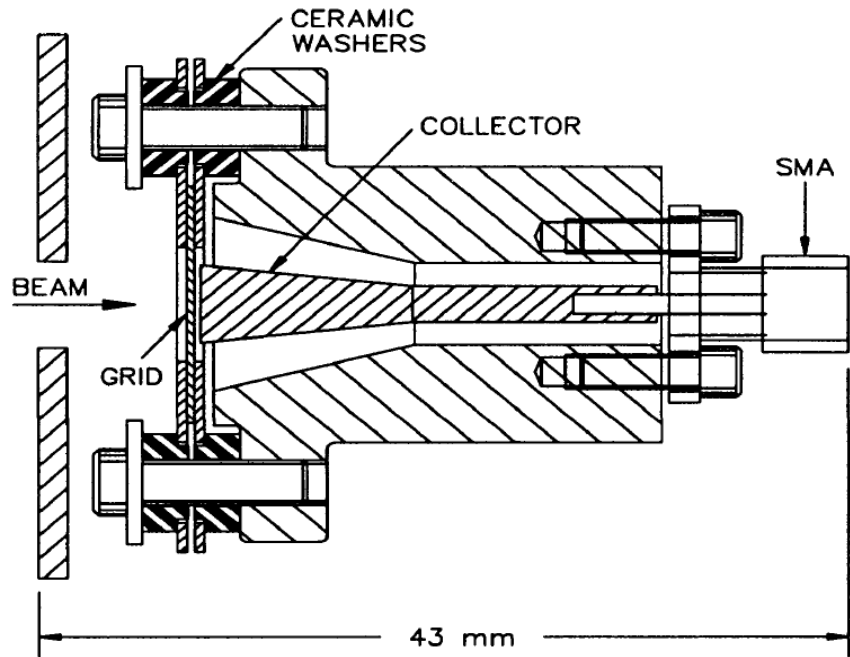
- Signal out-coupling should be very well matched until high frequencies , fast digitizers

$$\text{i.e. } BW > 5\sigma_f = \frac{5}{2\pi\sigma_t}$$

- Avoid measuring the self-field of the bunch
- Suppress distortion due to e⁻ secondaries
- Heating/Melting of cups

[9] J. M. Bogaty et al. (1990): A very wide bandwidth Faraday cup suitable for measuring GHz structure on ion beams with velocities to beta < 0.01

Axially coupled co-axial (AC-Co FFC)

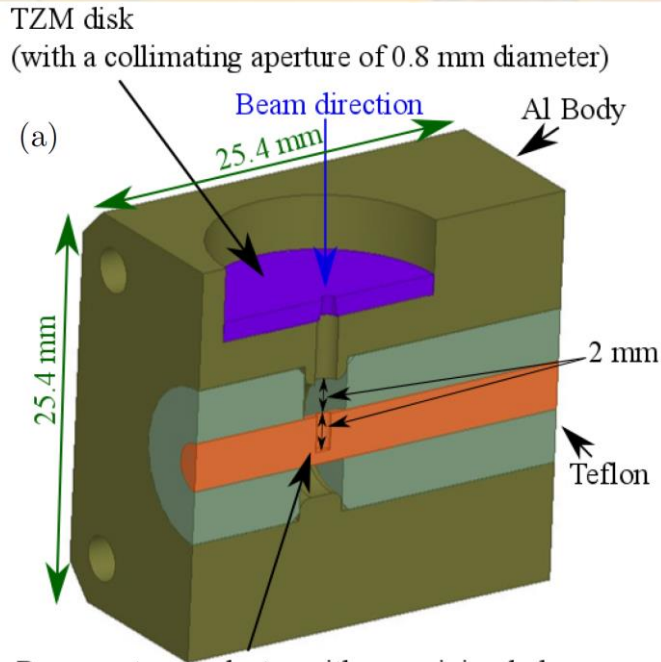


[7] P.Strehl, Beam instrumentation and diagnostics

[8] Rawnsley et al. <https://doi.org/10.1063/1.1342629>

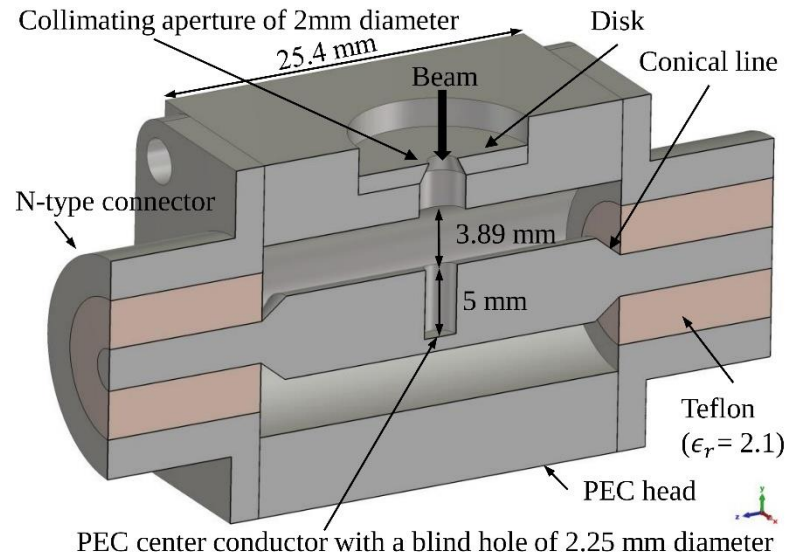
About 18 US patent applications!

Radially coupled co-axial (RC-Co FFC)



- Hole from the side in a co-axial cable
- 2mm distance from the ground to the central conductor → avoid pre-field
- Large depth to width ratio to avoid emission of secondaries

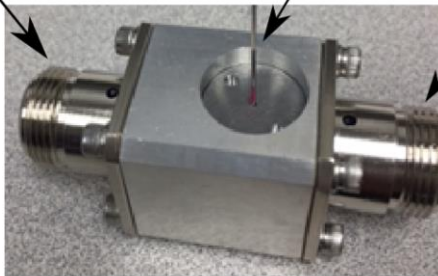
Minor modification of previous design



IUAC+GSI

[3] K. Mal et al.

Brass center conductor with a receiving hole
N-connector 0.8 mm hole



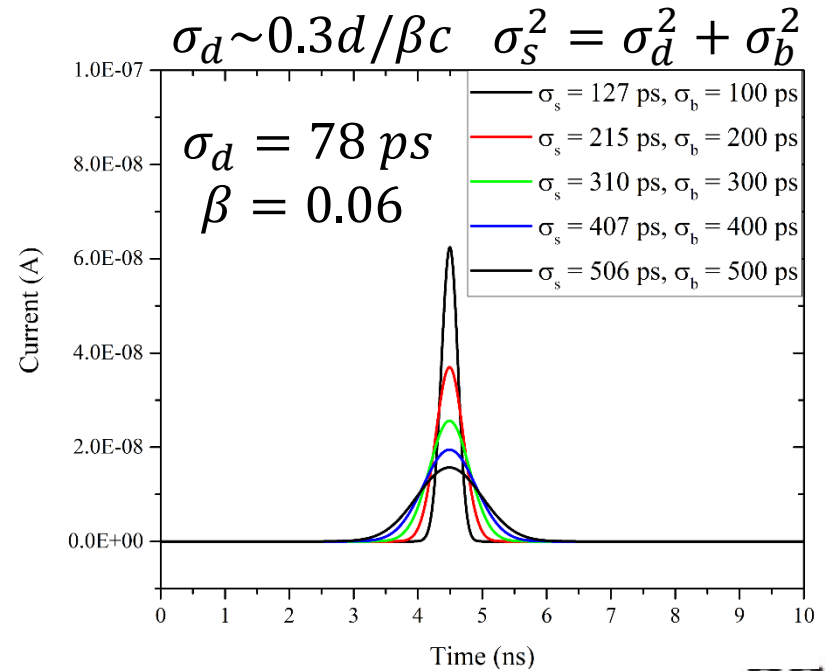
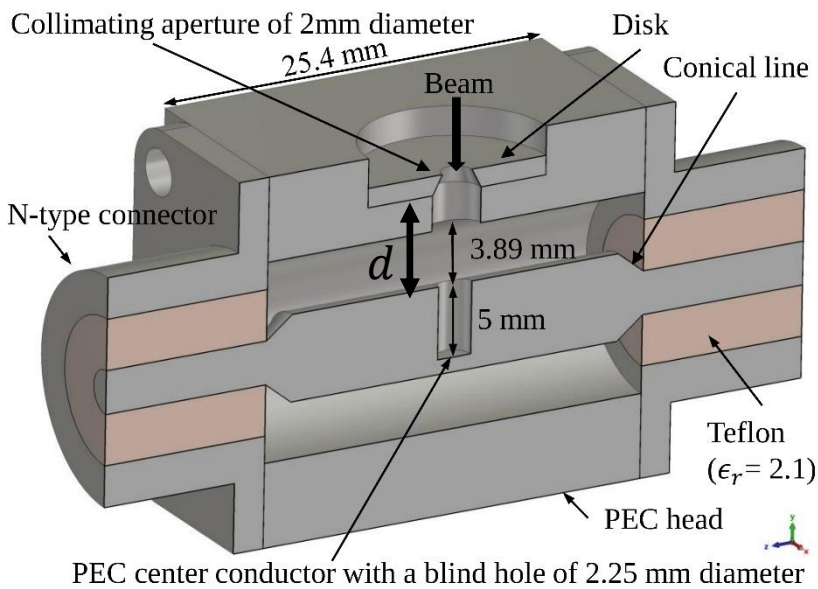
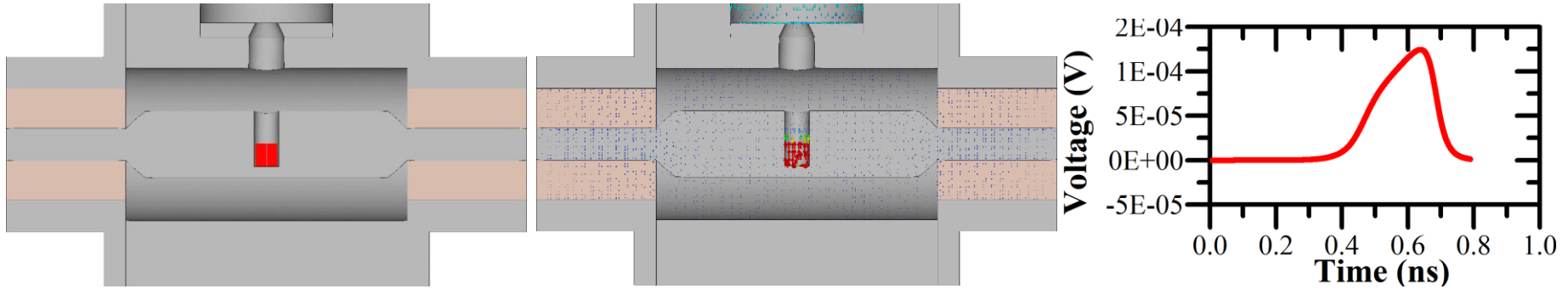
[11] J. -P Carniero et al., (2019)

D. Sun and A. Shemyakin

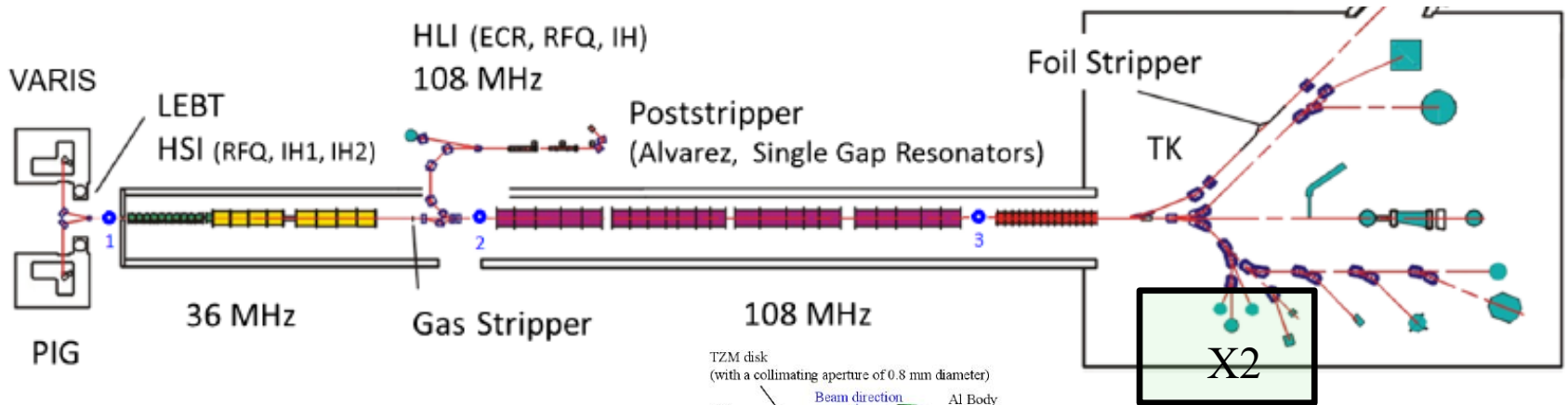
RC Co-ax FFC simulations



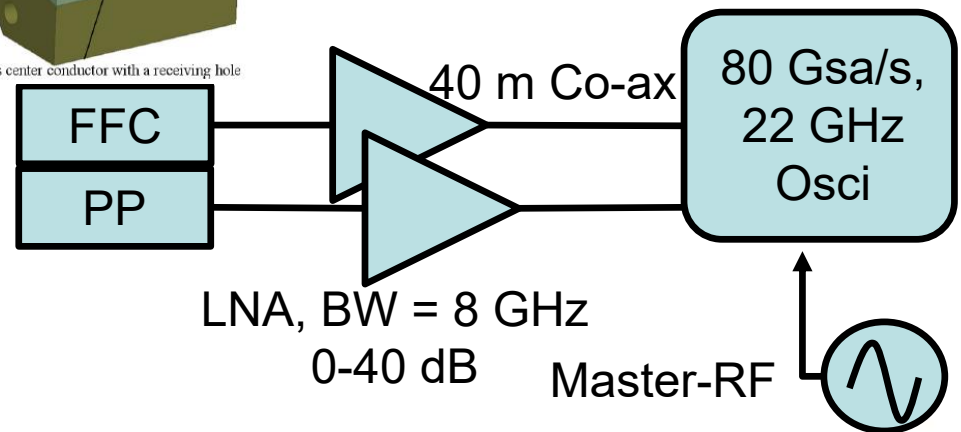
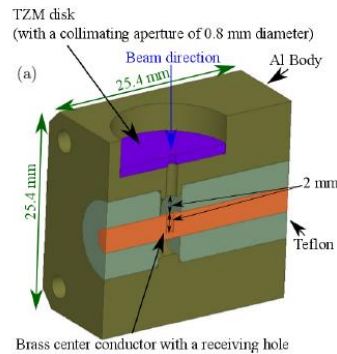
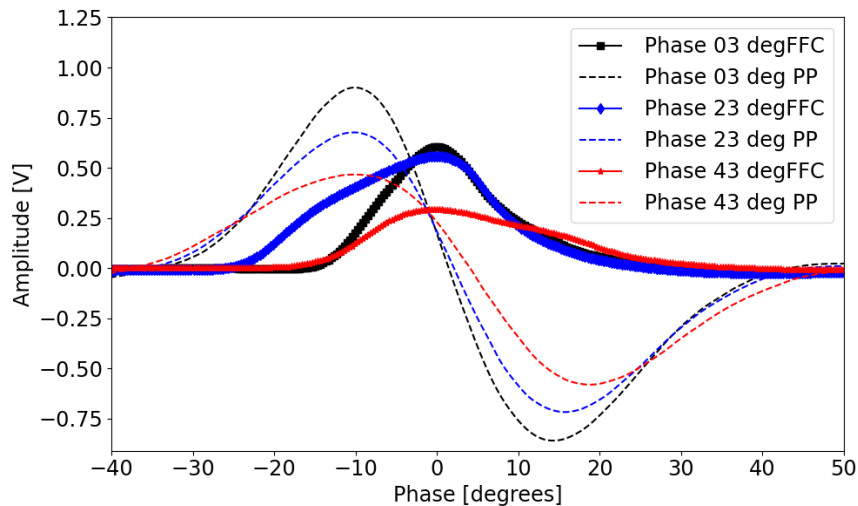
Simulations by [13] K. Mal et al. (IUAC)



RC- Coax FFC validation measurements@ X2



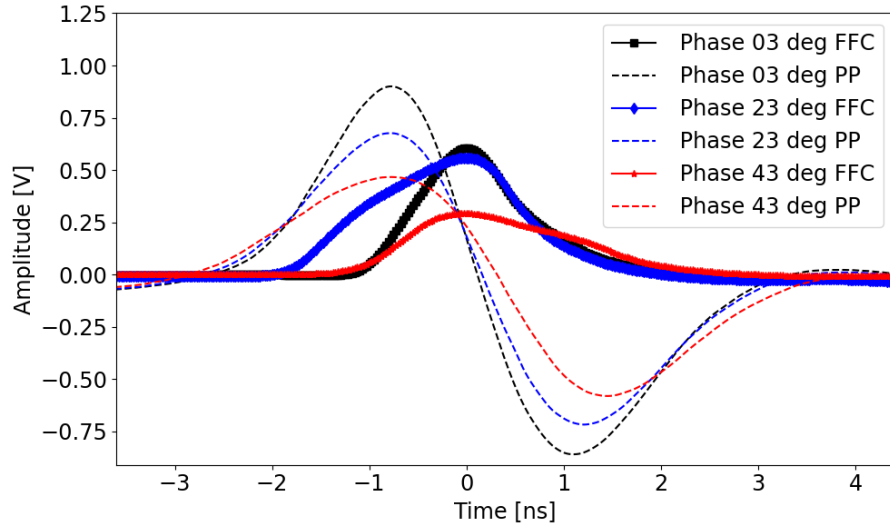
Ar¹⁰⁺, 0.6 mA, 8.6 MeV/u



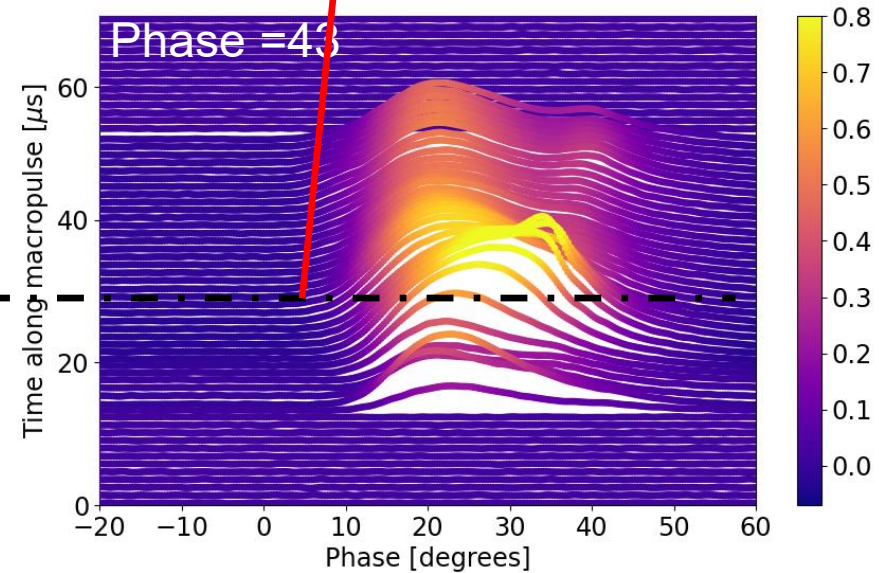
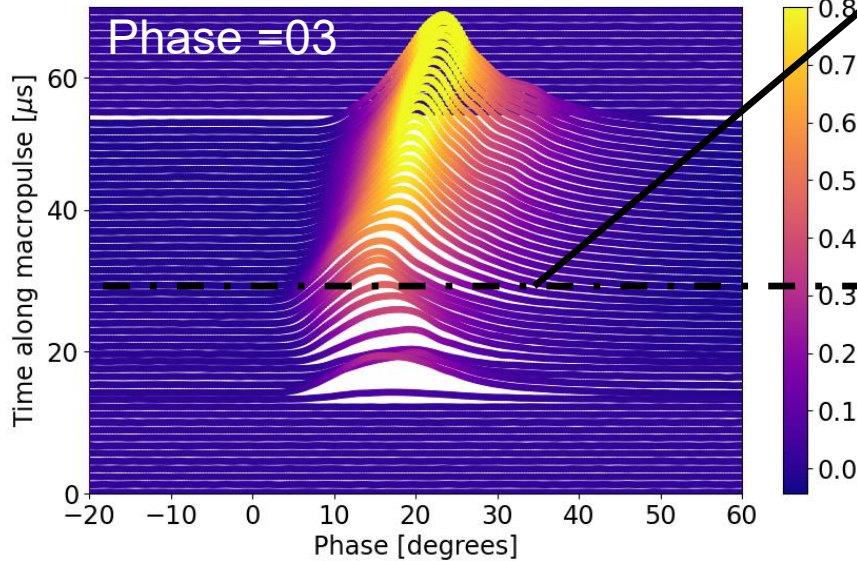
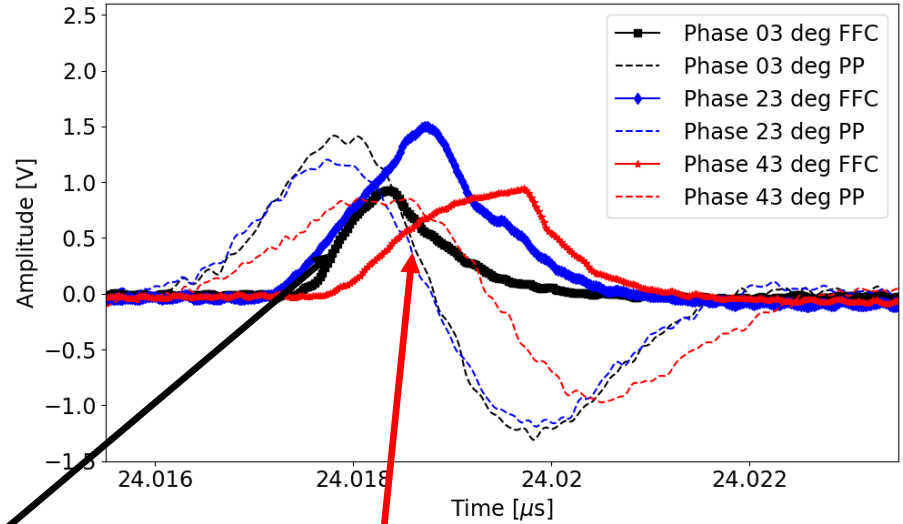
RC Coax FFC validation measurements



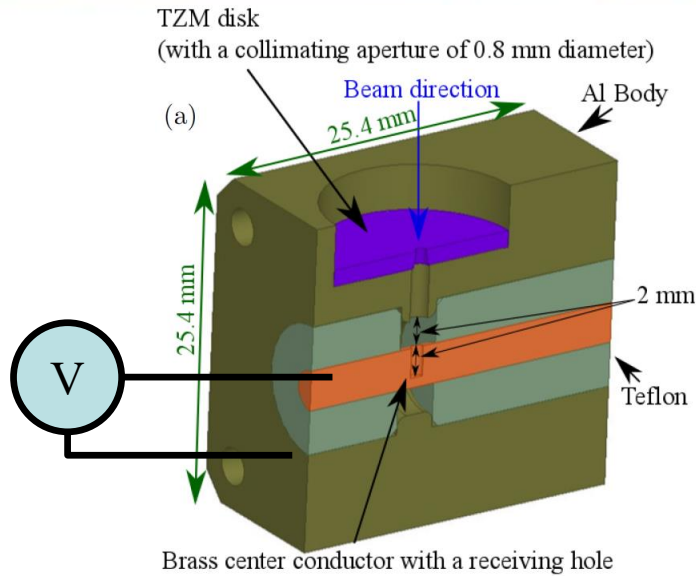
Averaged along macropulse



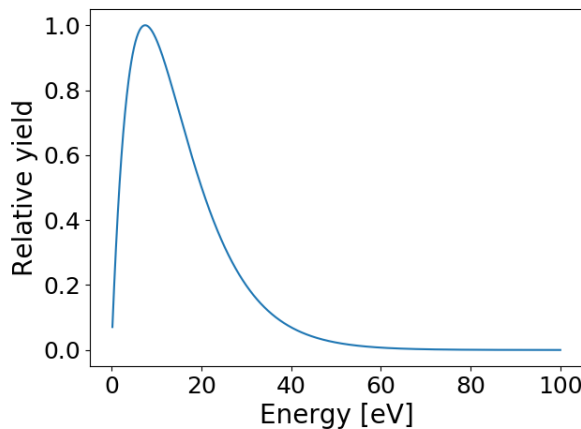
Single shot



RC- Coax FFC SEE measurements



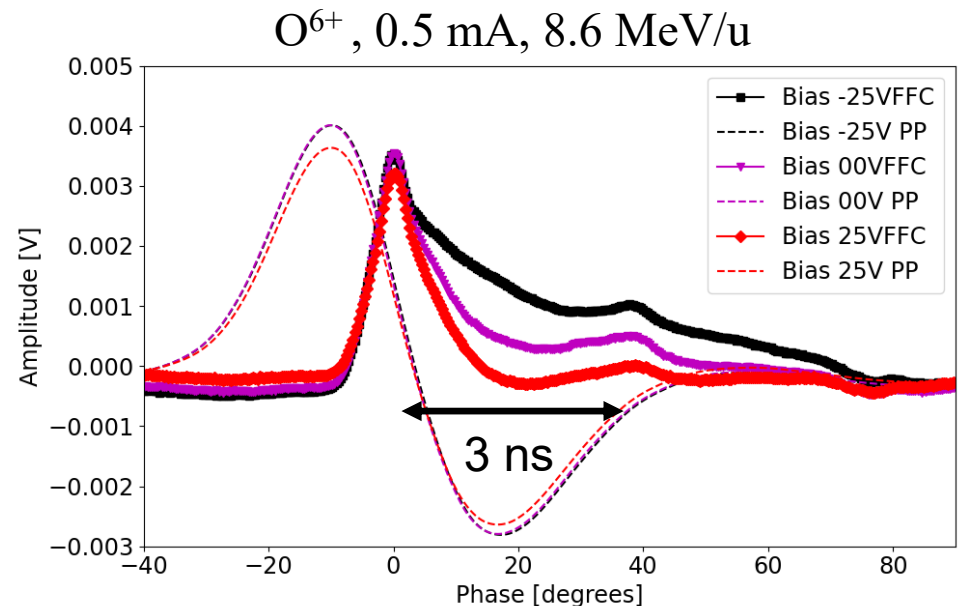
- Empirical estimates : 95% SEEs < 50 eV, Integrated SEE yield for 1-10 MeV/u p+ beam ~0.6-0.1 on metals. **Scales with ion charge state**
- Measured bunch shape is a strong function of DC bias voltage



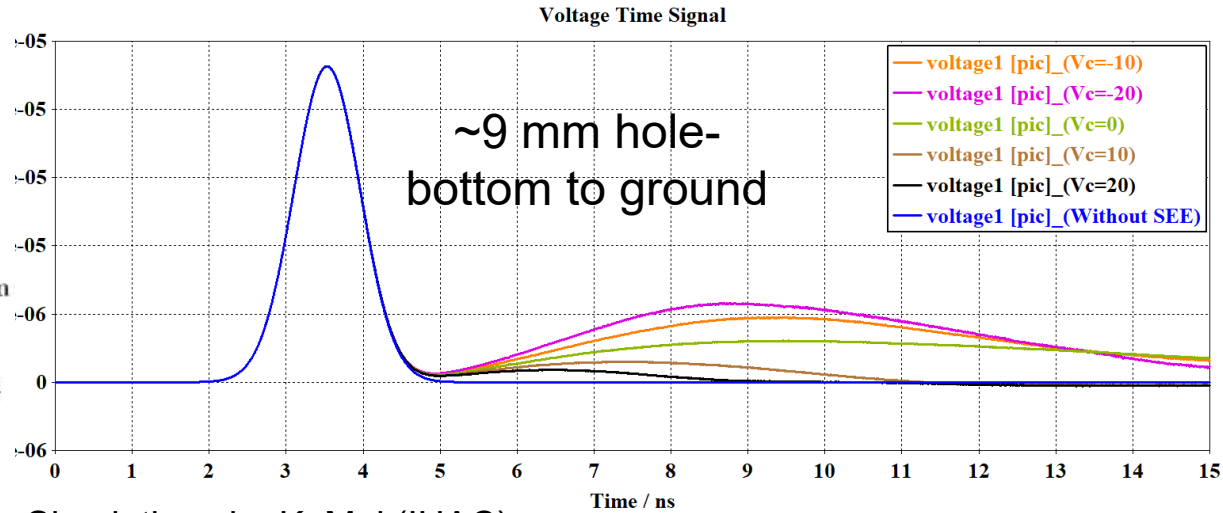
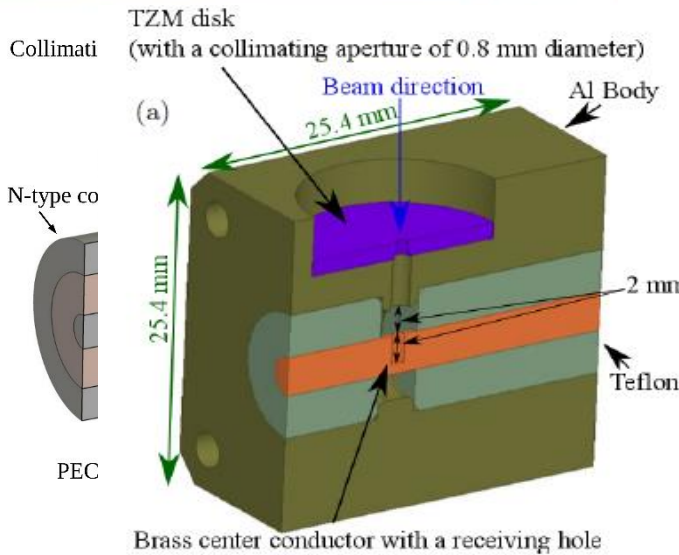
[14] J. R. Vaughan. Secondary Emission Formulas

[8] W. R. Rawnsley et al. AIP Proceedings (2000)

[15] A.Reiter et al., Investigation of Cross Talk in Secondary Electron Emission Grids. Technical Note LOBI-TN-SEM-2012-001

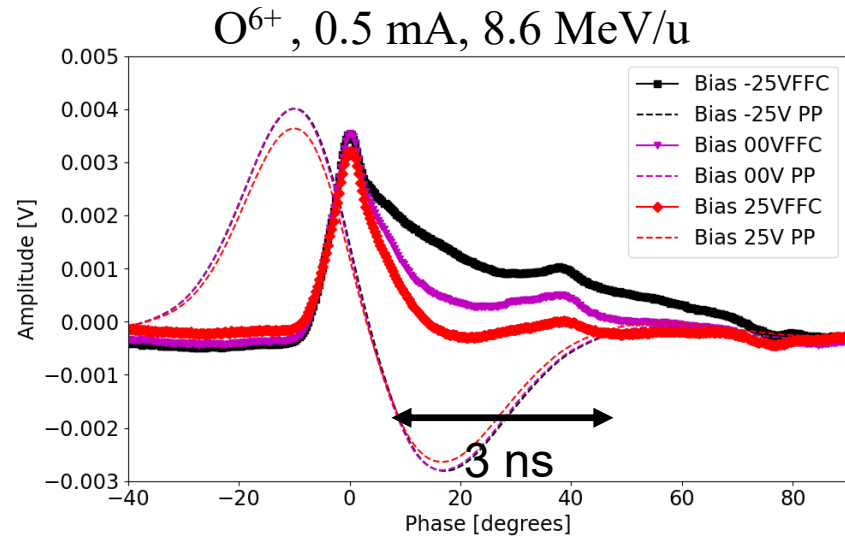


RC- Coax FFC SEE measurements

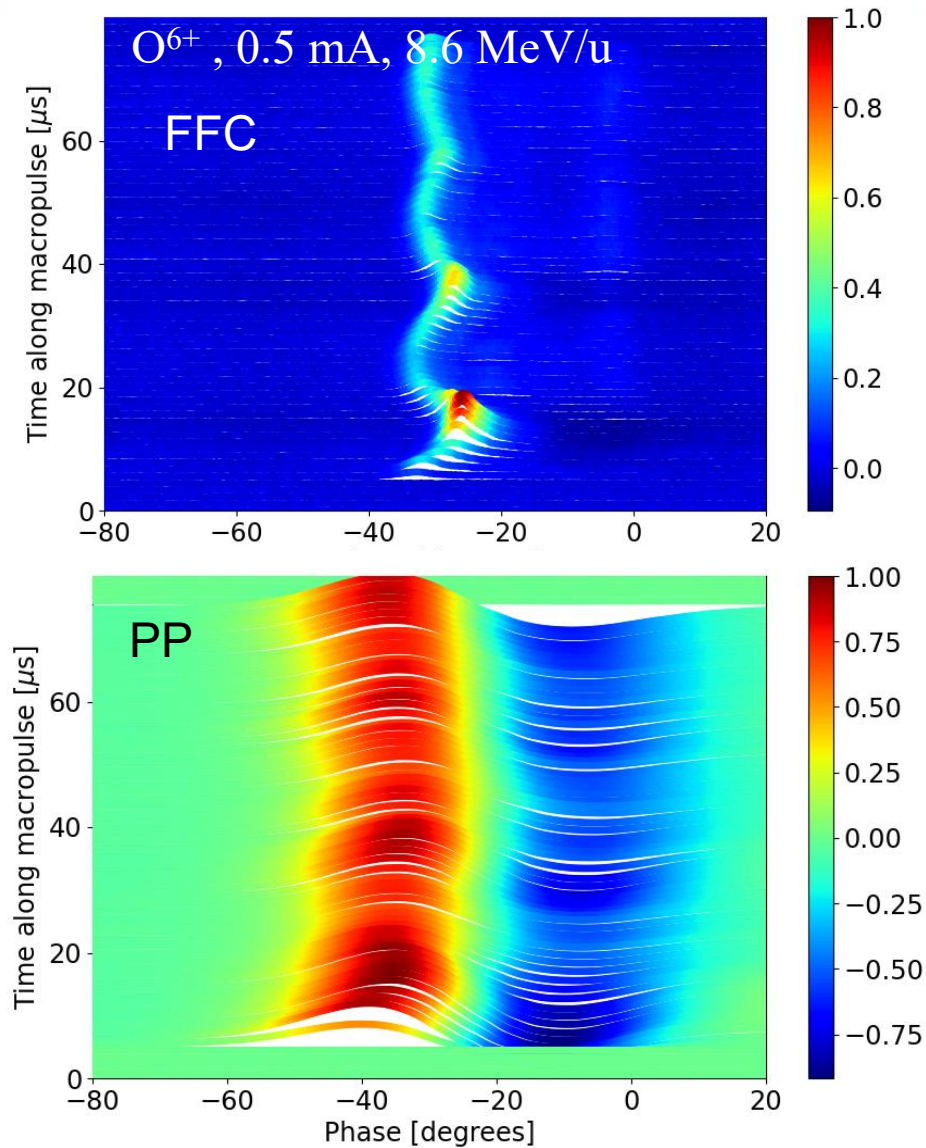


Simulations by K. Mal (IUAC)

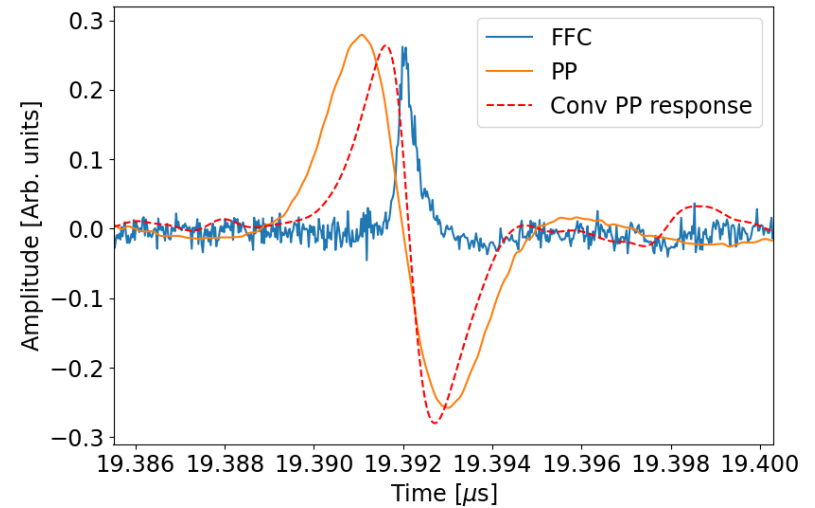
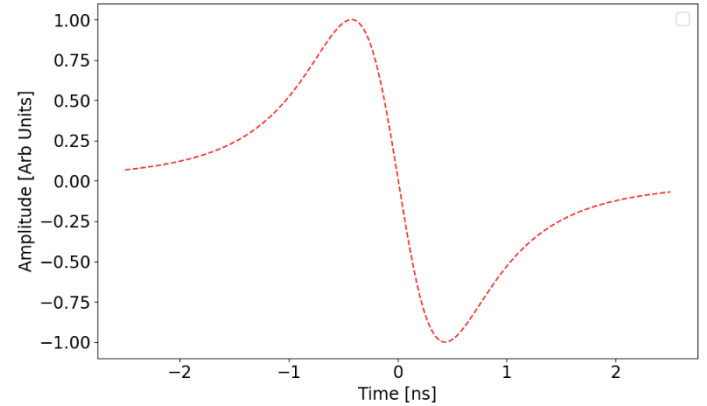
- Speed of 1-10 eV electrons is 0.6- 2 mm/ ns.
- Bottom of the hole to ground distance in the installed device is 4mm, e.g. 1 to 2 ns delay in peak sec. e- emission
- 30 V DC bias on central conductor enough to suppress the secondaries, however risk of pulling secondaries generated in the blind hole



FFC at dispersive location

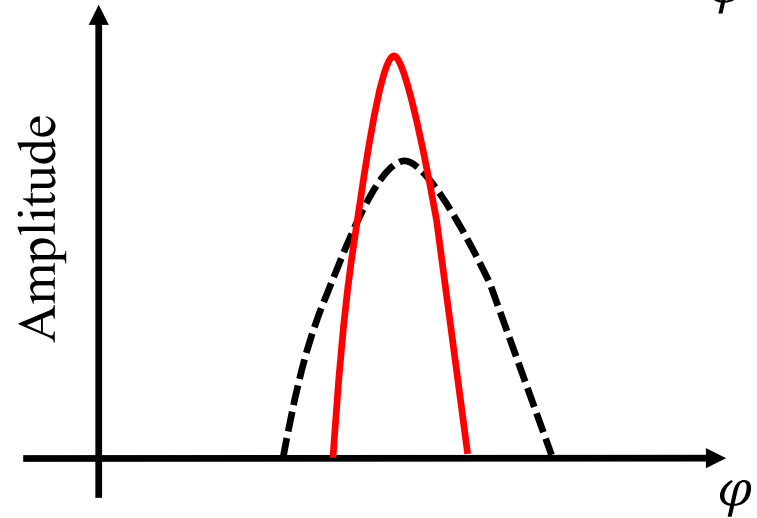
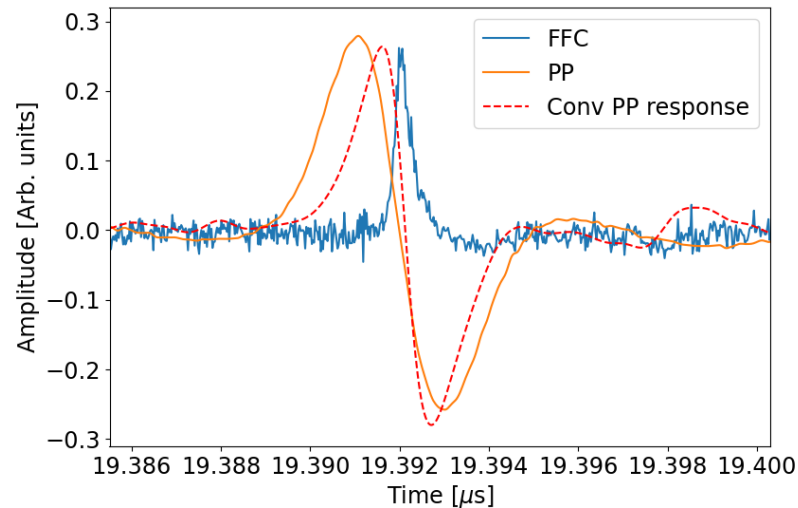
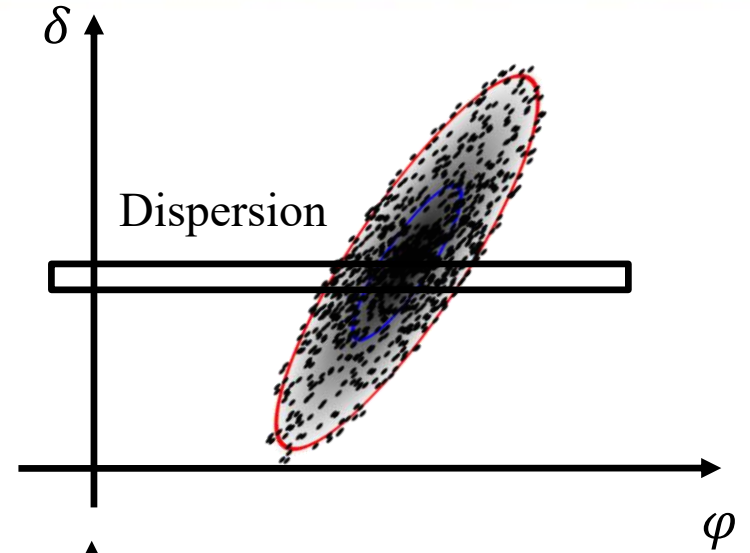
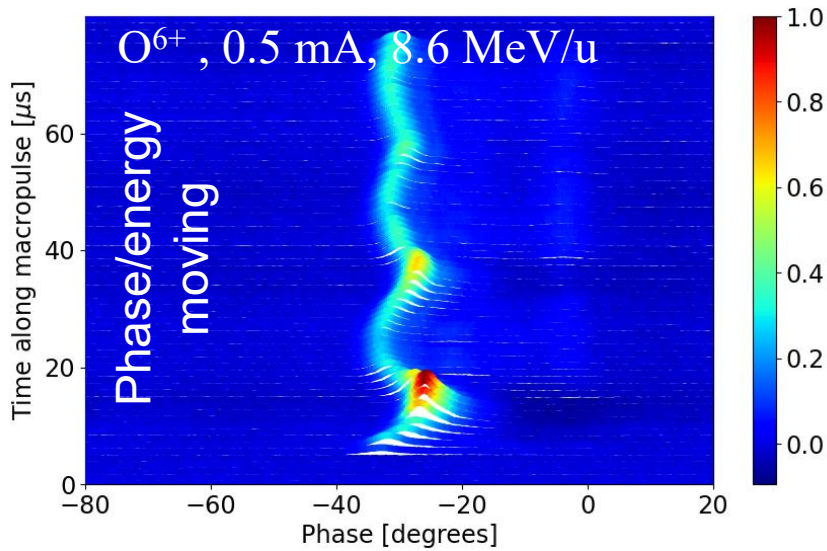


Phase probe low velocity beam response ($\beta = 0.134$) at $R = 3$ cm



Phase probe remains a validation device!

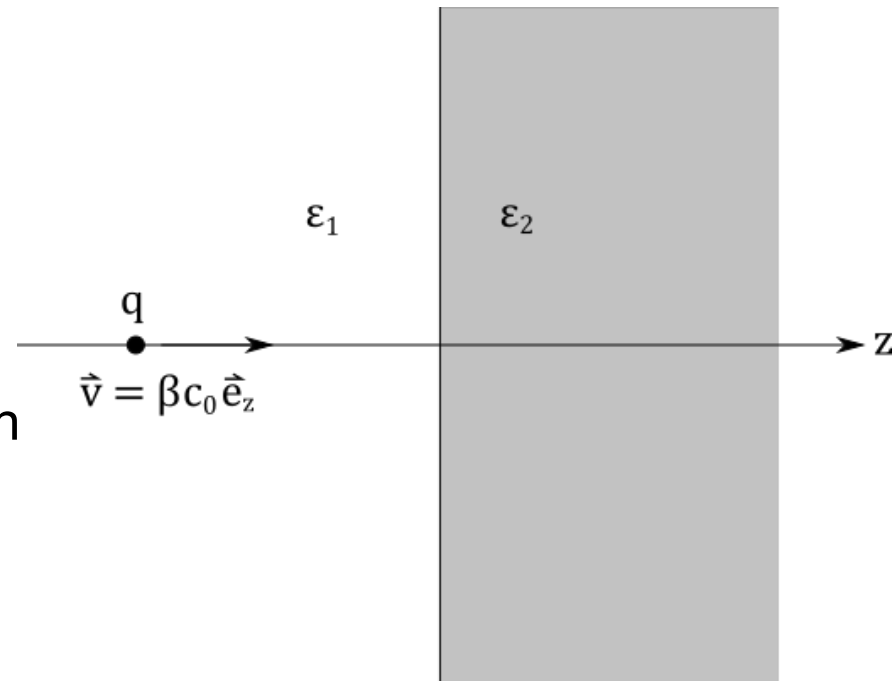
FFC at dispersive location



Careful placement of the FFC or simulation supported interpretation!

Transition Radiation: A charge with velocity $v = \text{const.}$ crossing an interface between two media radiates.

- an interface ($z = 0$) separating two half-spaces of different media
- solving MW-equations subject to interface conditions exhibit radiation field
- Surface electromagnetic phenomenon
→ prompt radiation
- In GHz regime, **coherent** transition radiation for $\sim\text{ns}$ bunches



A potential method un-affected by pre-field and secondary emission

Properties of (GHz) Transition Radiation



GTR electric field for single charge:

$$\vec{E} = \frac{q\beta}{2\pi\epsilon_0 cR} \frac{\sin\theta \delta\left(\frac{R}{c} - t\right)}{1 - \beta^2 \cos^2\theta} (\hat{e}_x \cos\theta + \hat{e}_z \sin\theta)$$

- Linear q and β dependence
- Parallel polarization for normal incidence
- Good signal: 10pC charges in 100 ps (σ) with $\beta=0.15 \rightarrow 10$ mV peak

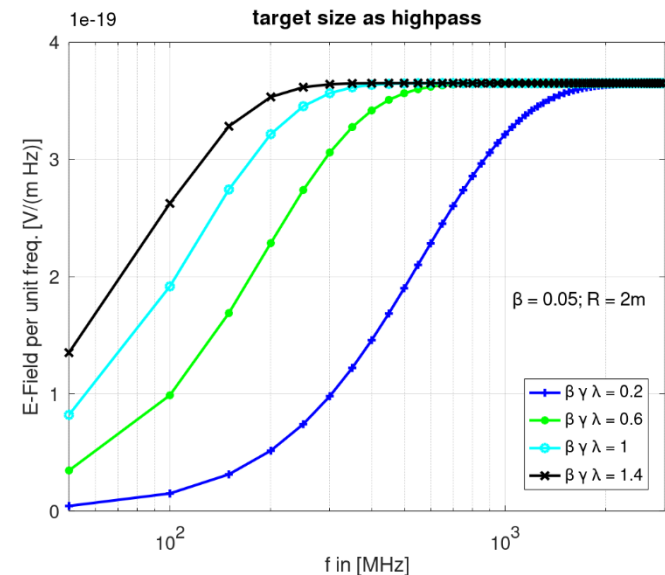
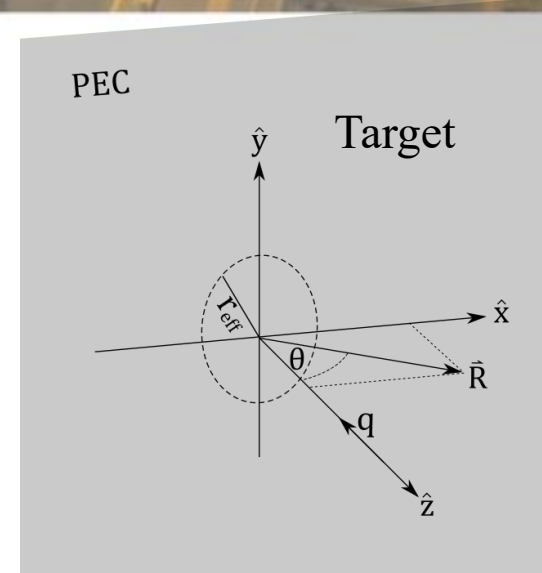
Target Size: Effective trans. extent of incident field: $r_{\text{eff}} \sim \beta \gamma \lambda$

For targets $< r_{\text{eff}}$: strong deviation from Far-Field

In practice: finite targets yield f -dependence

Formation Length: Distance from interface to reach Far-Field distribution (spherical wave)

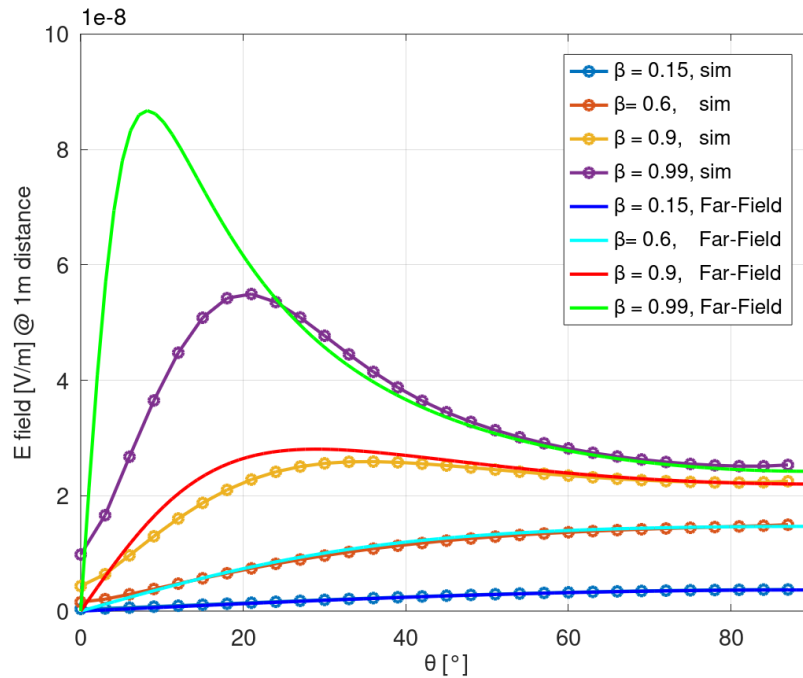
conservative estimate: $R \gg \gamma^2 \lambda$, also depends on θ



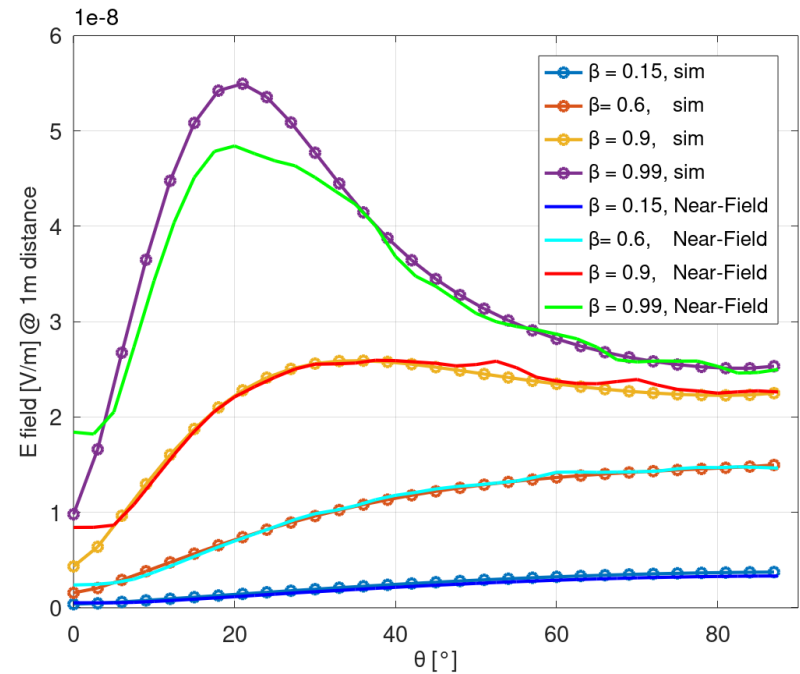
Angular distribution: CST simulation and analytical



Far-field, $R > \text{Formation length}$



Near-field, $R < \text{Formation length}$



- Simulation: $\sigma_t = 100\text{ps}$, $L_z = L_{tr} = 1\text{m}$
- Far-Field Distribution recovered for all β @ $\theta > 55^\circ$
- Up to $\beta \sim 0.8$ Far-Field Distribution for all θ

CST Simulations match with near and far field analytical models for all distances.

Near field models:

[16] A. G. Shkvarunets and R. B. Fiorito, Phys. Rev. ST Accel. Beams 11, 012801

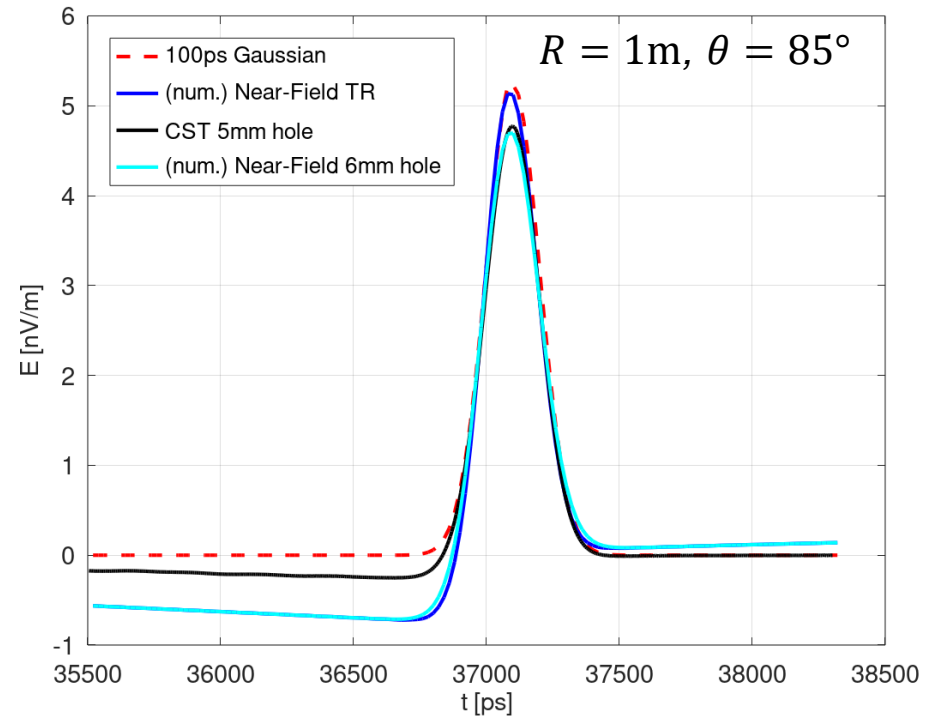
[17] R. Singh and T. Reichert, Phys. Rev. Accel. Beams 25, 032801



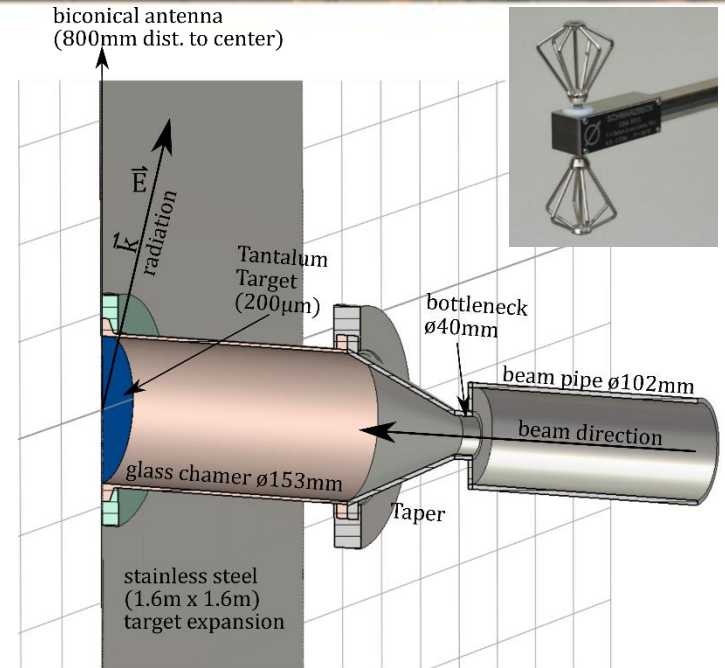
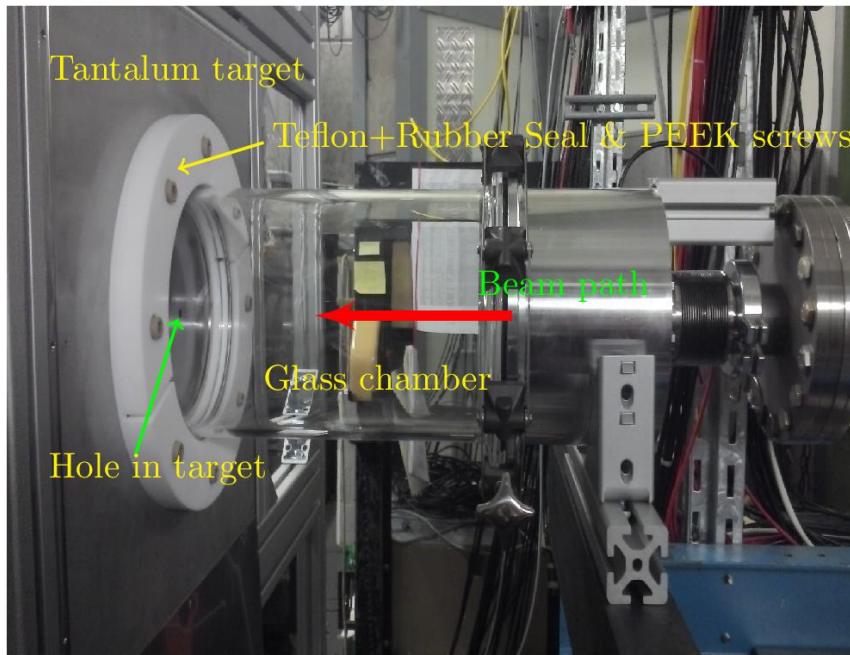
Diffraction Radiation (DR) makes it non-invasive



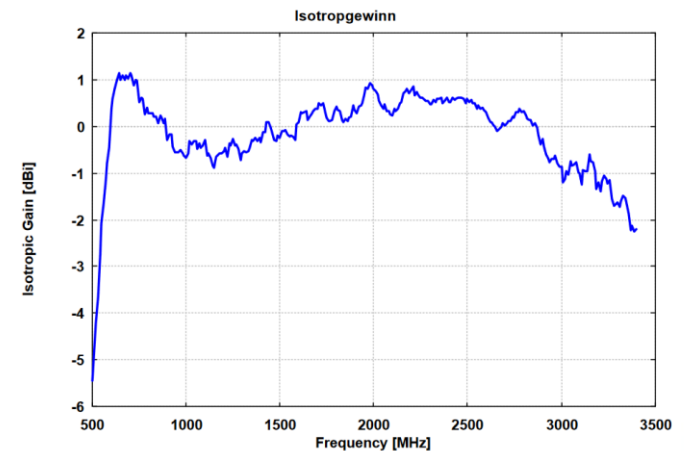
- DR is very similar to TR but charge traverses close to the media interface
- Here: Instead of impacting on the target bunch can go through hole
- Allowable hole size: $\varnothing \downarrow$ for $\beta \downarrow$
- For $\beta \sim 0.15$, $\varnothing \leq 6\text{mm}$
- **Non-destructive measurements possible!**



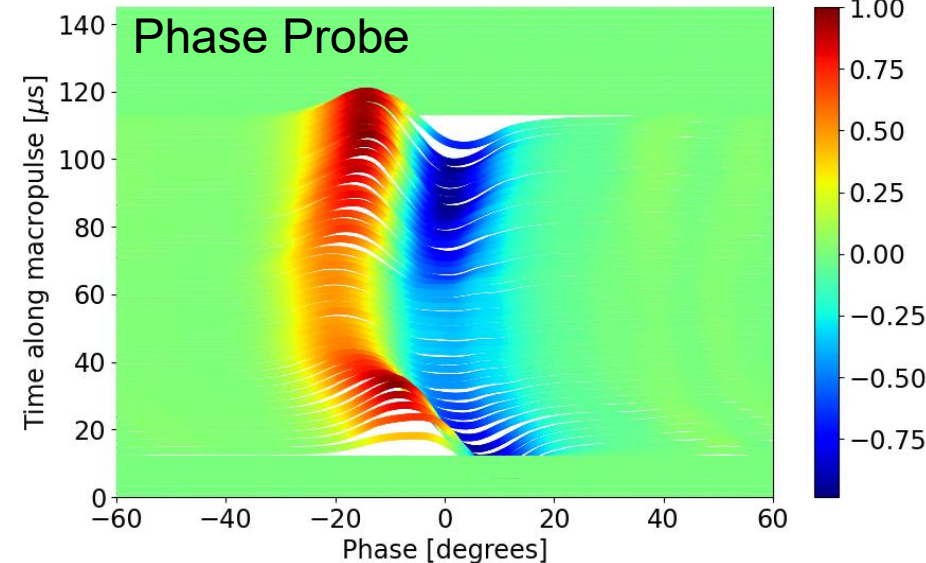
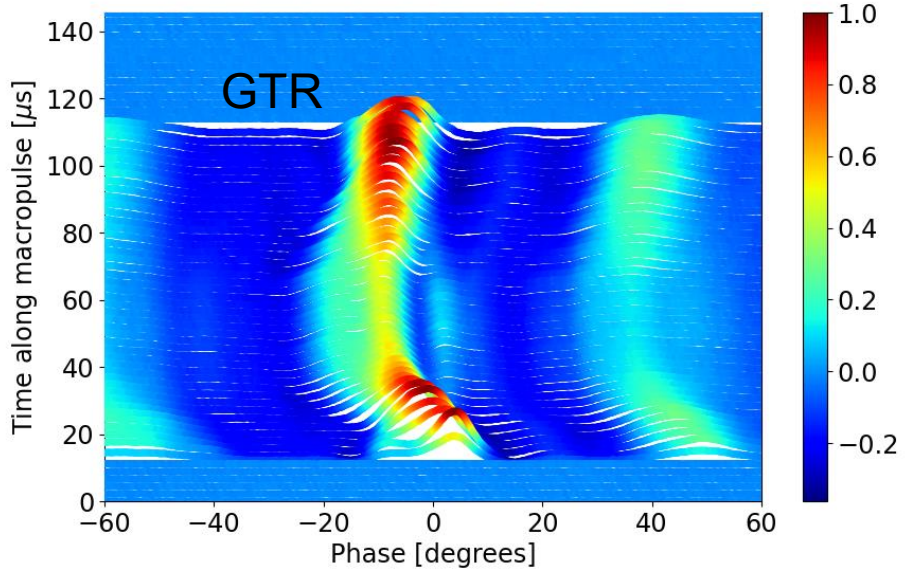
GTR setup in X2



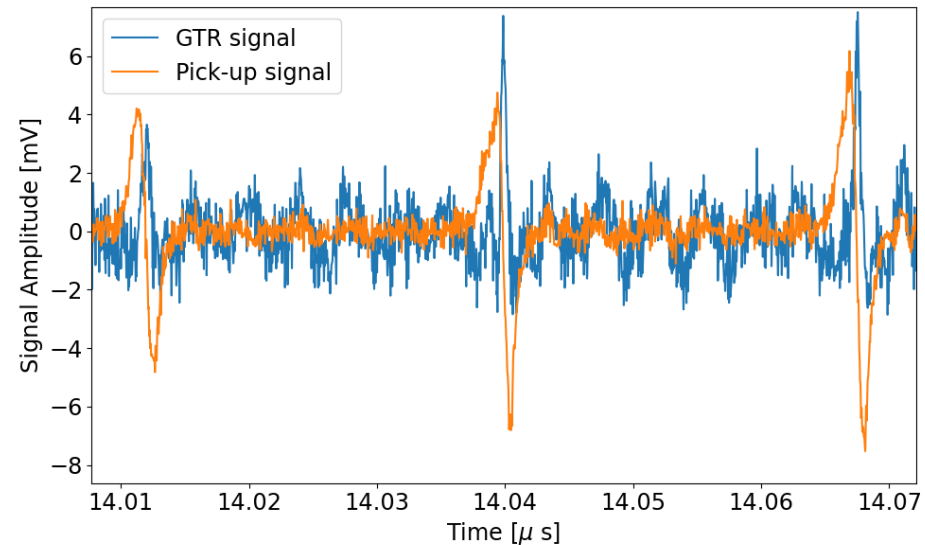
- An EM window to couple out the TR signal
→ Vacuum tolerance → critical
- Absorbers to avoid reflections
- Linear phase antenna designs



First results



Bi^{26+} 11.4 MeV/u, $\sim 400 \mu\text{A}$,
100 μs pulse length, 36 MHz RF
Antenna angle (θ) = 40 deg,
Antenna distance to target (R) = 1.0 m

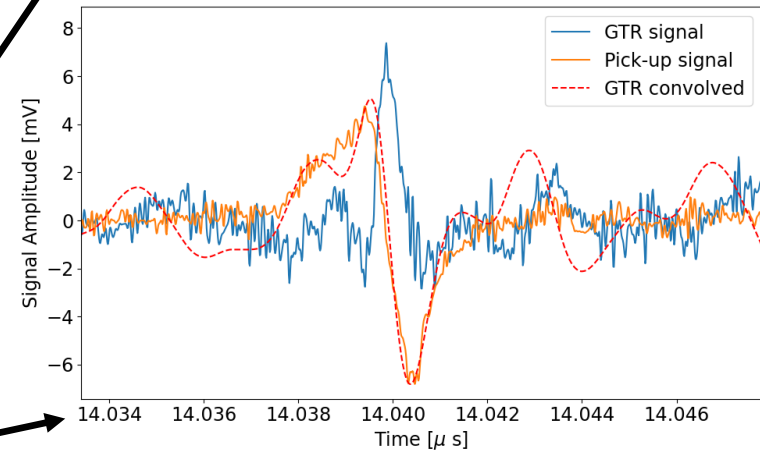
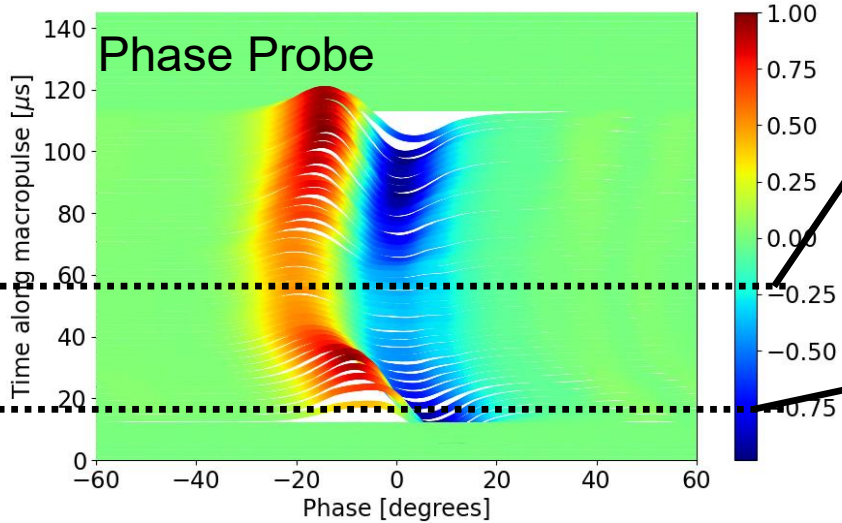
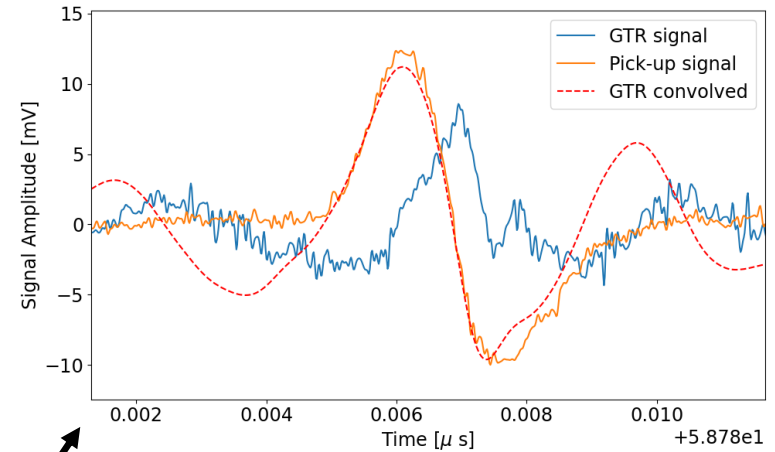
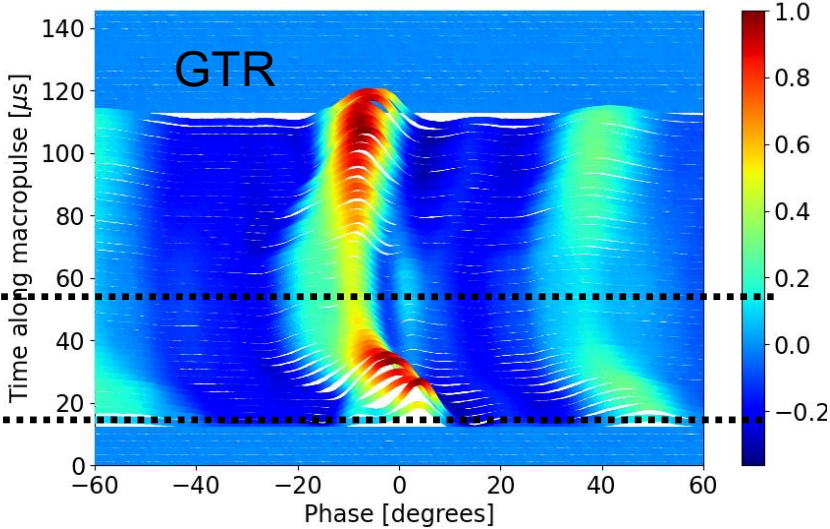


- Good correlation with the pick-up data
- Mean beam energy matches with ToF between pick-up and GTR

Comparison to phase probe data



Bi²⁶⁺ 11.4MeV/u, ~400μA, 100μs pulse length, 36MHz RF, theta = 40 deg,(R) = 1.0 m



Convolved GTR has **precise** agreement with phase probe signal!

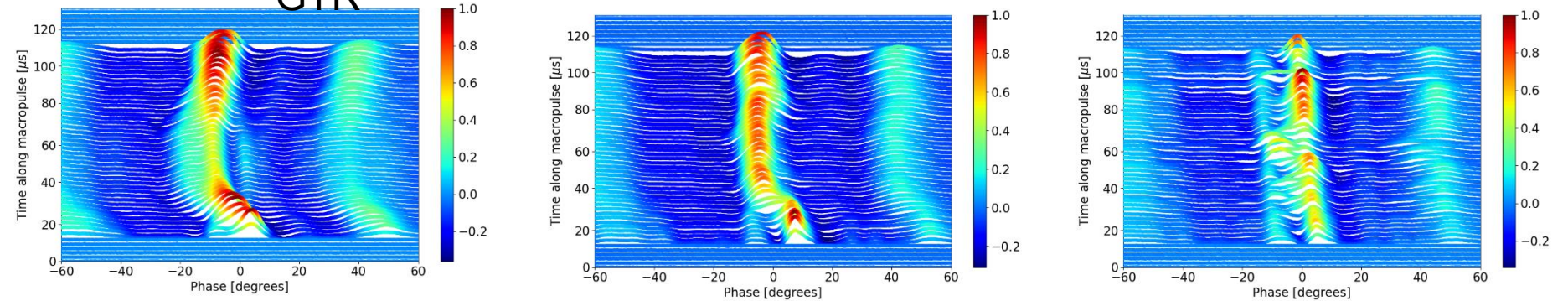


Pulse-to-pulse variation in bunch shapes

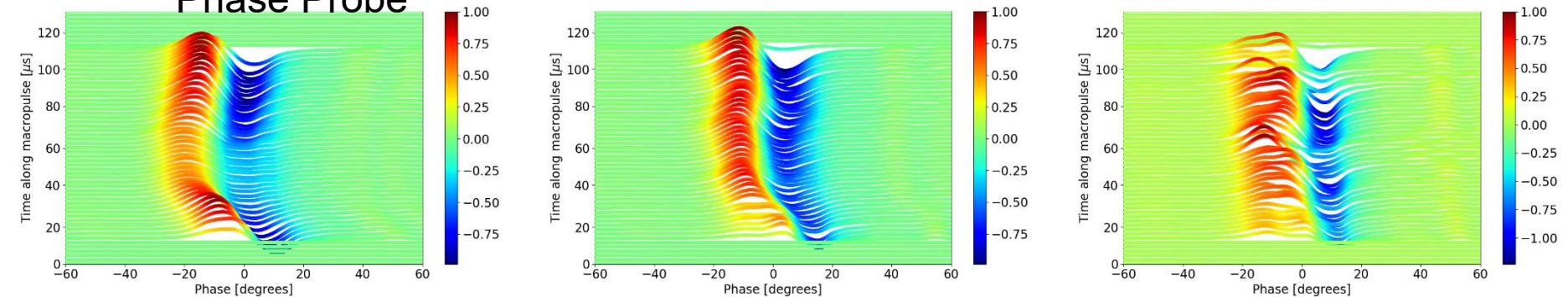


Bi²⁶⁺ 11.4MeV/u, ~400μA, 100μs pulse length, 36MHz RF, theta = 40 deg,(R) = 1.0 m

GTR



Phase Probe



- Three consecutive macropulses show different charge distributions
- Longitudinal diagnostics need to be prepared for such fast changes

[17] R. Singh and T. Reichert, Phys. Rev. Accel. Beams 25, 032801



Conclusion and Outlook



- Fast and robust longitudinal diagnostics is important for various alignments in UNILAC.
- Lot of interesting efforts in the past → limitations for intra-macropulse and high intensity beam monitoring
- Advent of fast acquisition electronics allow devices with time domain monitoring → FFC and GTR with minimal user parameters
- FFC is a promising **compact option** *but* requires careful placement and biasing is essential in UNILAC energy regimes. New designs being tested, comparison with calculated phase space needed
- GTR a promising **non-invasive option for high currents** *but* not a compact installation. Further investigation under BMBF project ongoing

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