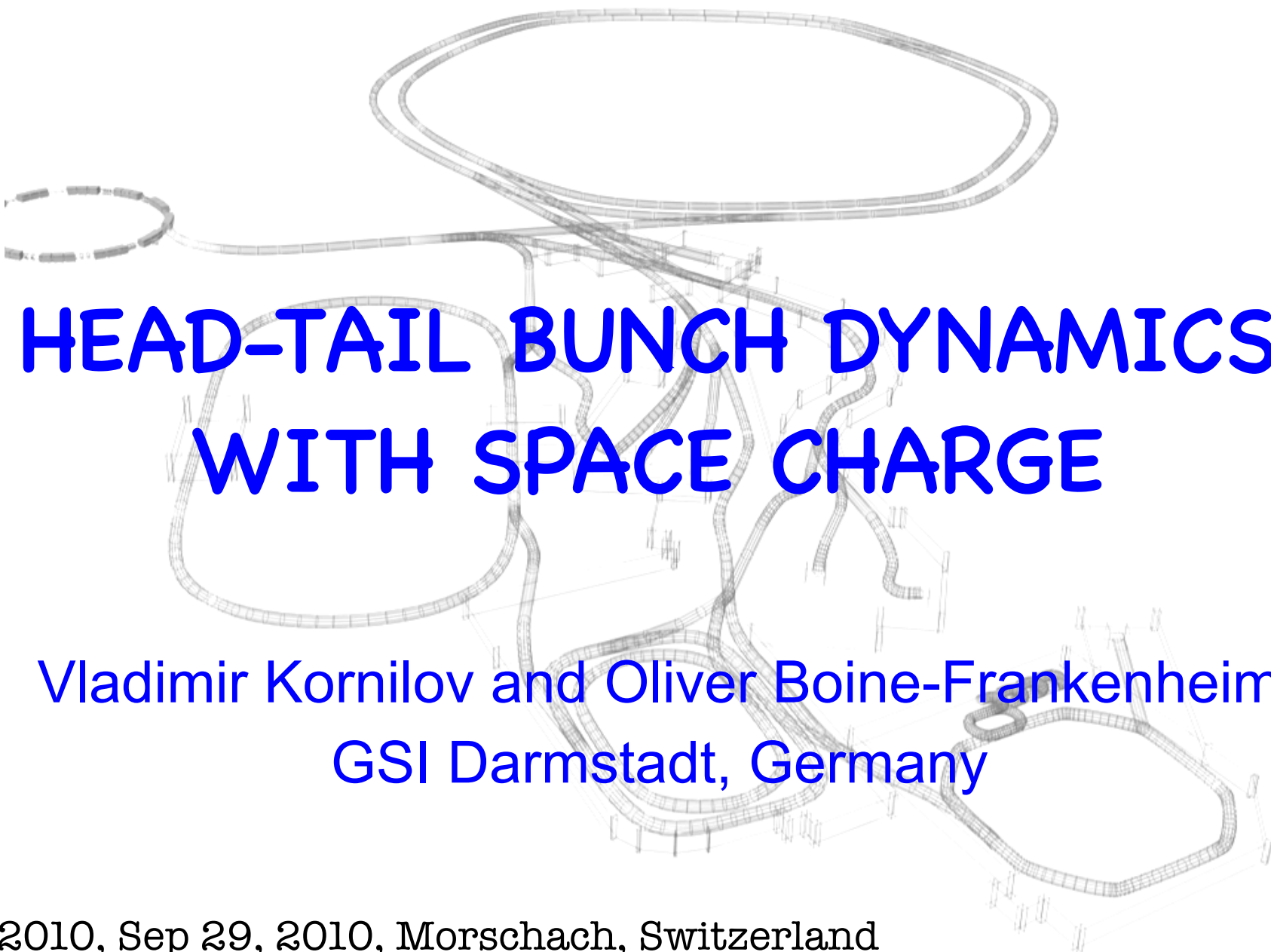
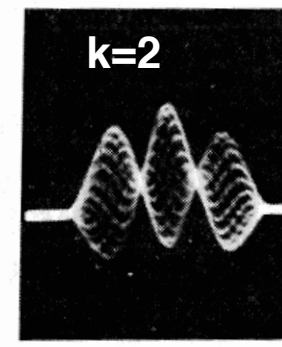
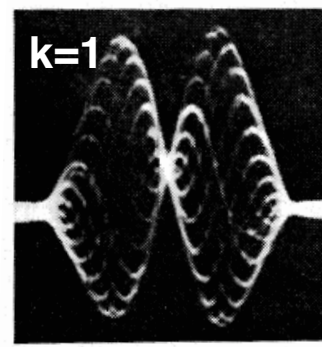
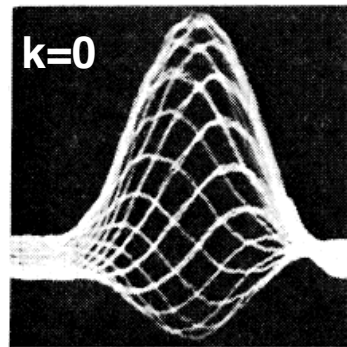




HEAD-TAIL BUNCH DYNAMICS WITH SPACE CHARGE

Vladimir Kornilov and Oliver Boine-Frankenheim
GSI Darmstadt, Germany



PS Booster 1974

Head-Tail Instability

SIS100 of FAIR: the major concern at the injection plato

modified by Space Charge (ΔQ_{SC} up to 0.25)

$$\Delta Q_{sc}(\tau) = \frac{\lambda(\tau) r_p R^2}{\gamma^3 \beta^2 Q_0 a^2}$$

not considered by classical (Sacherer) theories

here: single-bunch instability, below the coupling threshold

a non-realistic bunch, but a very useful model
M.Blaskiewicz, PRSTAB 1, 044201 (1998)

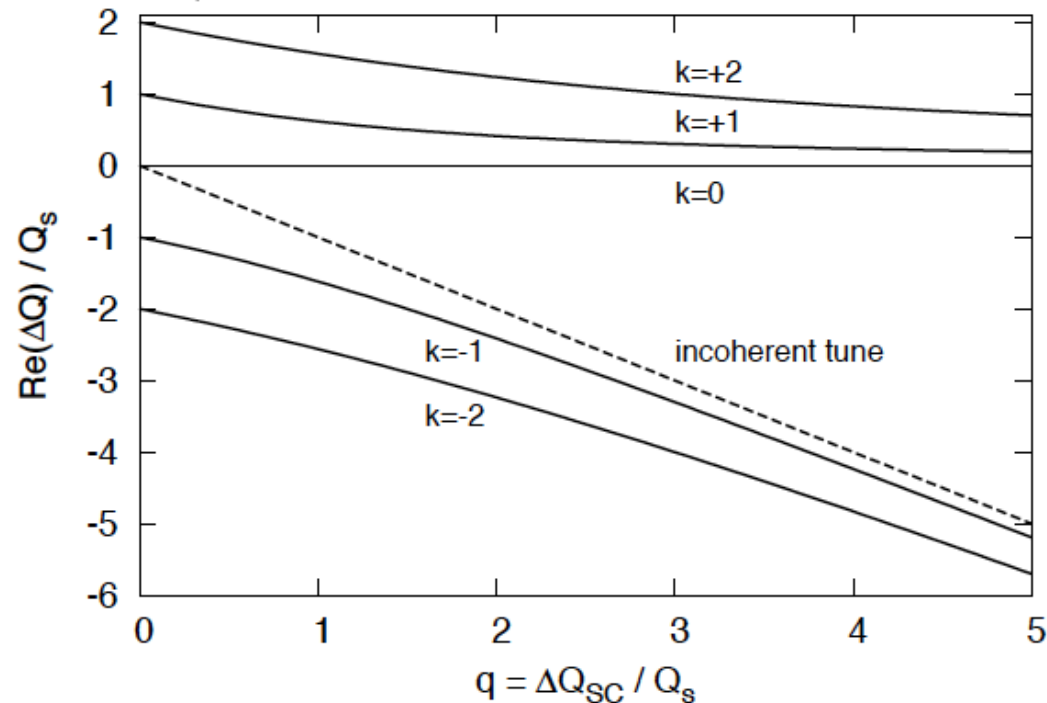
square-well (barrier) potential $\rightarrow$ constant line density $\rightarrow$ constant ΔQ_{sc}
longitudinal distribution $f(p) \propto \delta(p - p_0) + \delta(p + p_0)$

Assumptions: rigid flows, $Q_0 \gg |\Delta Q|$, arbitrary space charge

without wake: $\Delta Q = -\frac{\Delta Q_{sc}}{2} \pm \sqrt{\frac{\Delta Q_{sc}^2}{4} + k^2 Q_s^2}$ “+” for $k \geq 0$

space charge parameter

$$q = \frac{\Delta Q_{sc}}{Q_s}$$



Recent analytic works:

A.Burov, PRSTAB **12**, 044202 (2009); A.Burov, PRSTAB **12**, 109901(E) (2009)

V.Balbekov, PRSTAB **12**, 124402 (2009)

OUR APPROACH: PARTICLE TRACKING SIMULATIONS

PIC codes: PATRIC

O.Boine-Frankenheim, V.Kornilov, Proc. of ICAP2006 (2006)

HEADTAIL

G.Rumolo and F. Zimmermann, PRSTAB **5**, 121002 (2002)

this work: “semi-frozen” electric field, homogeneous transverse profile,

the space charge implementation verified using the airbag theory

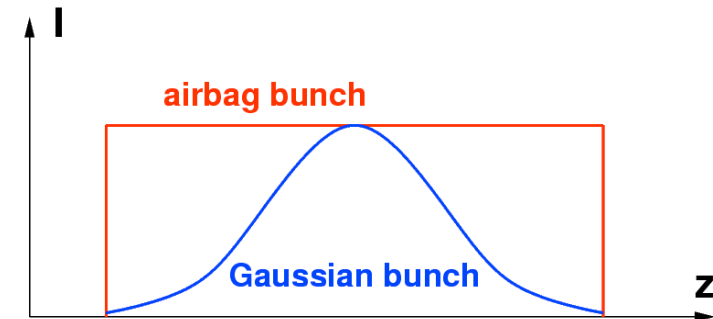
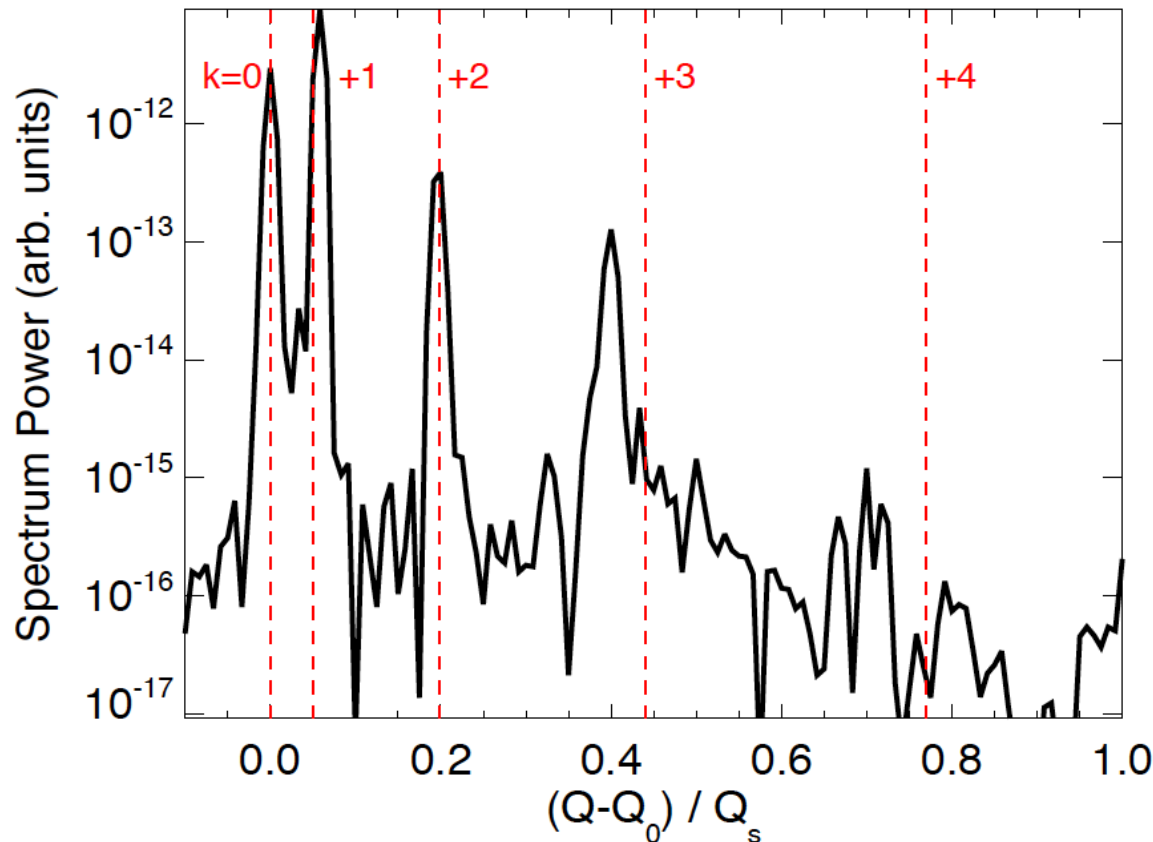
V. Kornilov and O. Boine-Frankenheim, Proc. of ICAP2009, San Francisco (2009)

Gaussian bunch: Gaussian line density, Gaussian momentum distr.

Simulations for the Gaussian bunch
Spectrum example for $q=20$,
red dashed line: the airbag theory

space charge parameter

$$q = \frac{\Delta Q_{sc}}{Q_s}$$



**Tune shifts are well predicted
by the air-bag theory,
especially for stronger SC**

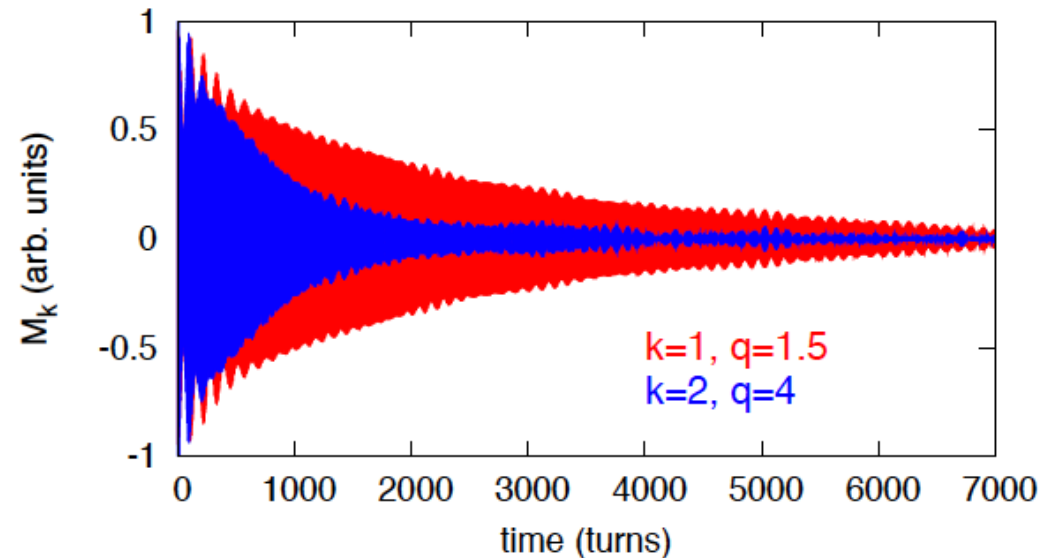
**Bunch form seems to be
not very important!**

Landau damping in bunches exclusively due to the effect of space charge discussed in [Burov 2009], [Balbekov 2009]

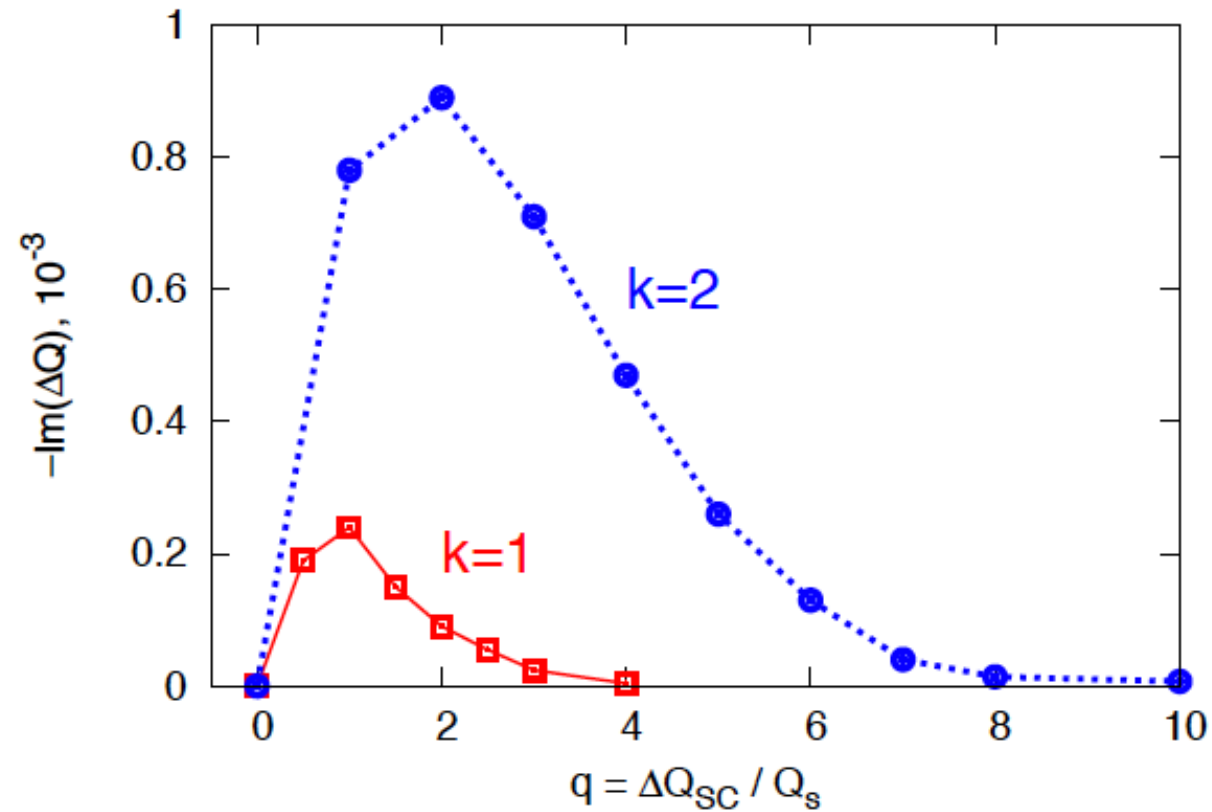
$$\Delta Q_{sc}(\tau) = \frac{\lambda(\tau) r_p R^2}{\gamma^3 \beta^2 Q_0 a^2}$$

note: in a coasting beam space charge CAN NOT produce Landau damping of its own!

we start a simulation with an initial perturbation of a head-tail mode



$$M_k = \int \bar{x}_{code} \cos(k\pi\tau/\tau_b) d\tau$$



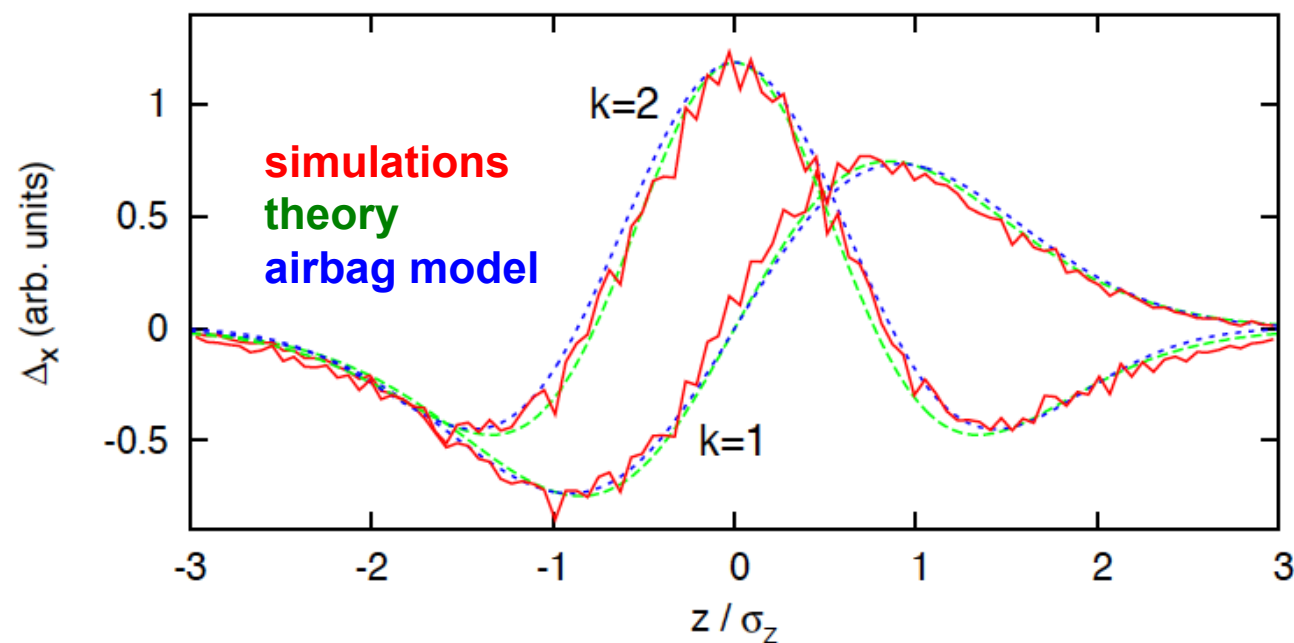
Summary of Landau damping simulations:
damping decrement for a Gaussian bunch

Eigenfunctions for a Gaussian bunch:
extracted in simulations using Landau damping (at $q=6$)

theory [Burov 2009]: $\bar{y}'' + \nu \exp(-\tau^2/2)\bar{y} = 0$

airbag head-tail eigenfunctions
[Blaskiewicz 1998]

$$\bar{x}_k(\tau) = A_0 \exp(-i\xi Q_0 \tau / \eta) \cos(k\pi \tau / \tau_b)$$



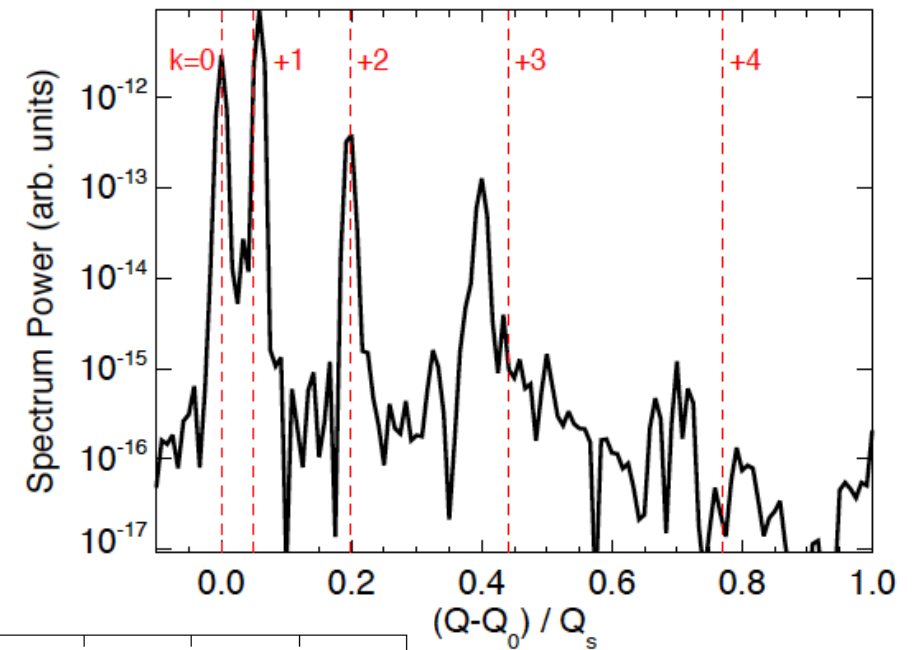
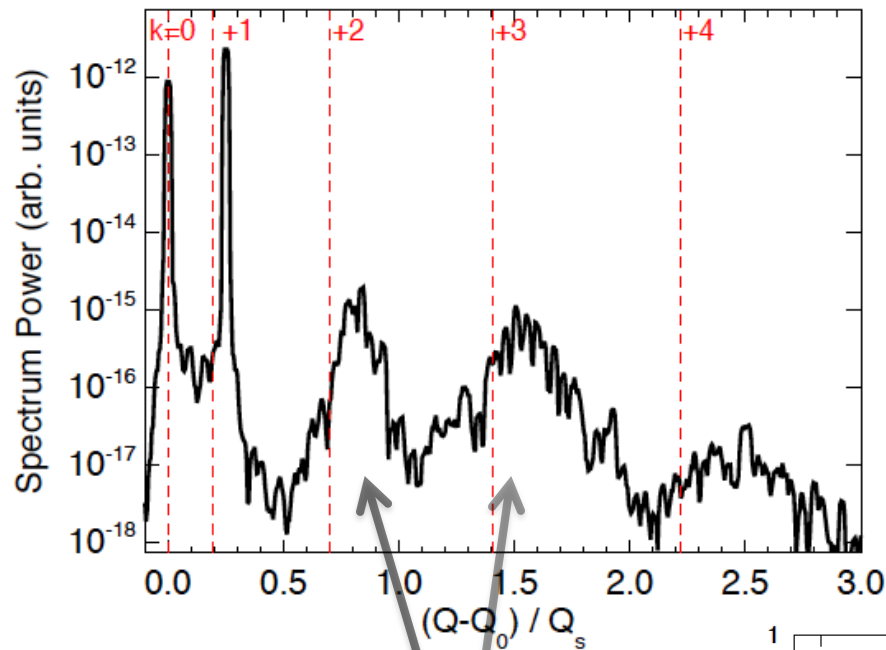
$$\tau_b = \frac{4\sigma_z}{R}$$

eigenfunctions of a Gaussian bunch are very close to the airbag modes!

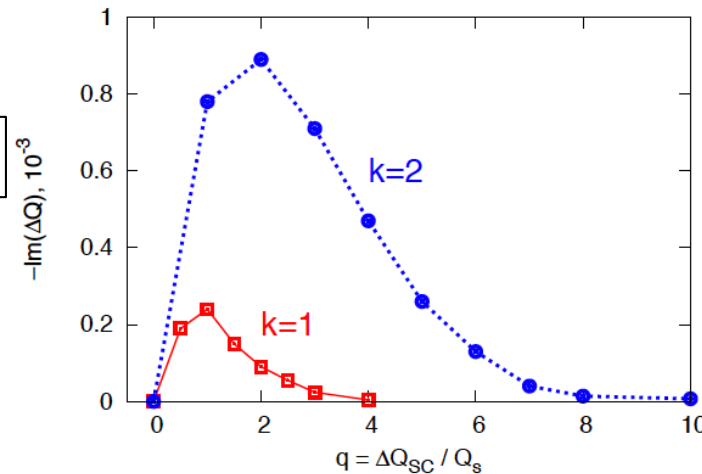
Landau damping can be seen in the bunch spectrum

$q=5$

$q=20$



the modes $k=2, k=3$ suppressed

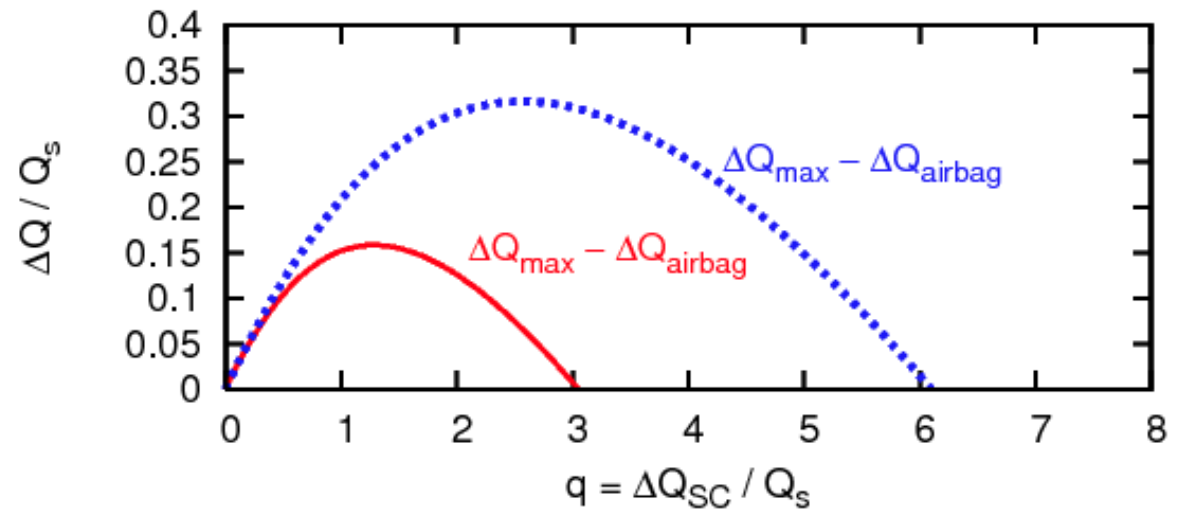
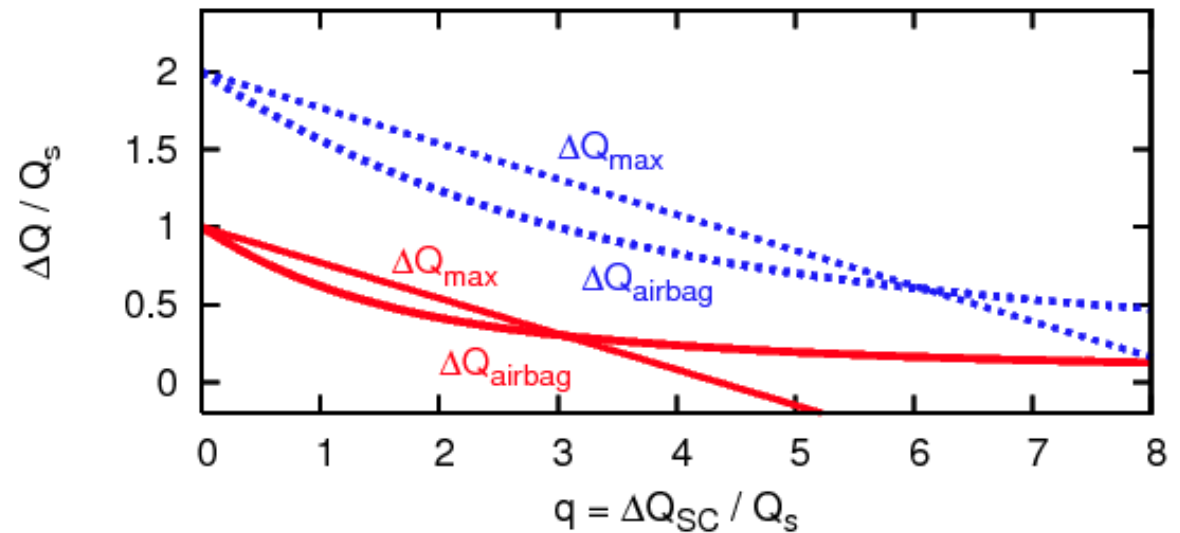


here not

a physical interpretation

upper boundary of the
incoherent effective spectrum:
small ΔQ_{SC} ,
large synchrotron amplitudes

$$\Delta Q_{\max} \approx -0.23Q_s q + kQ_s$$

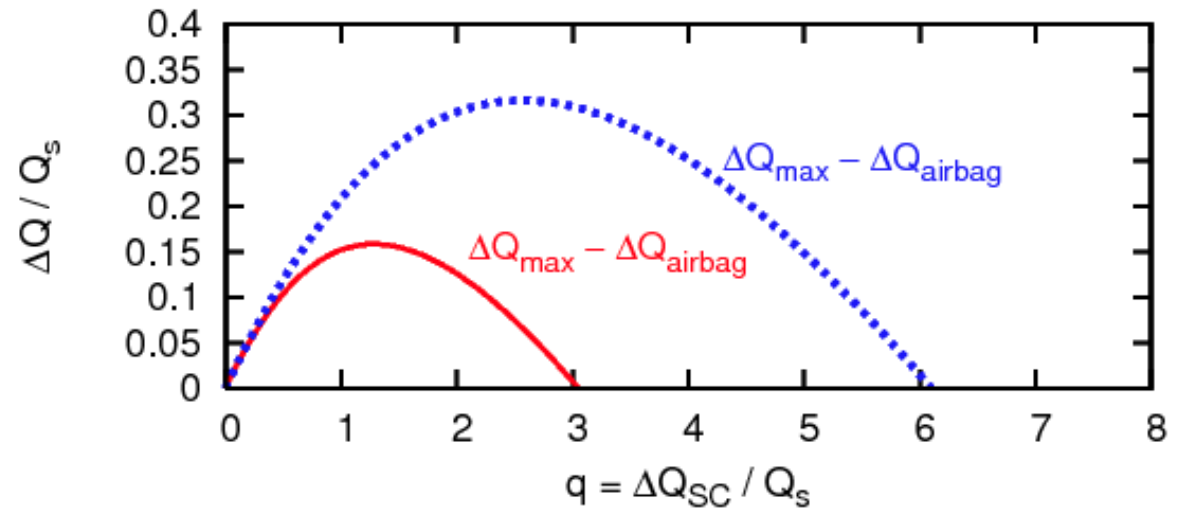
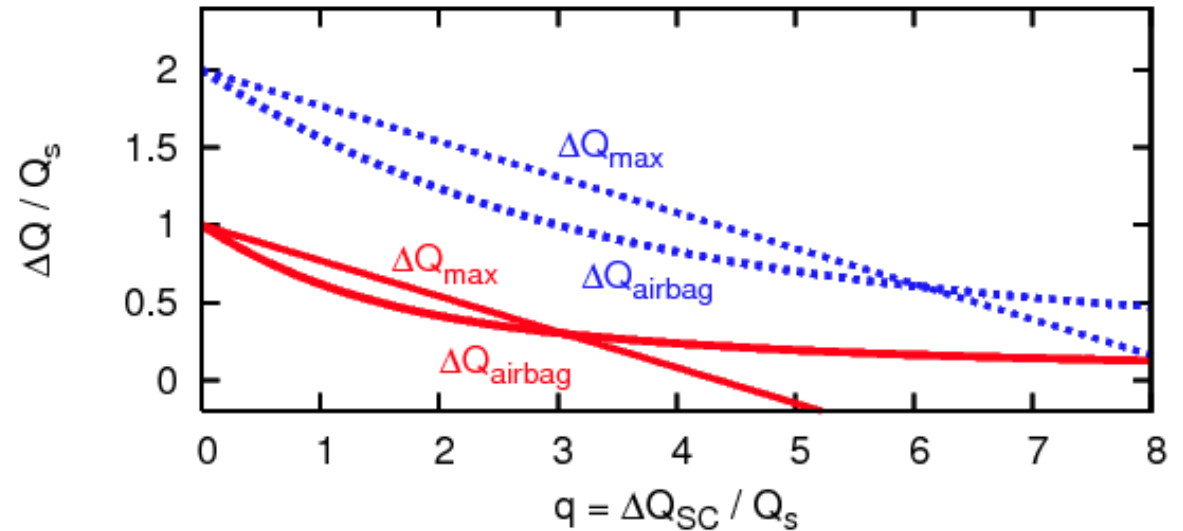
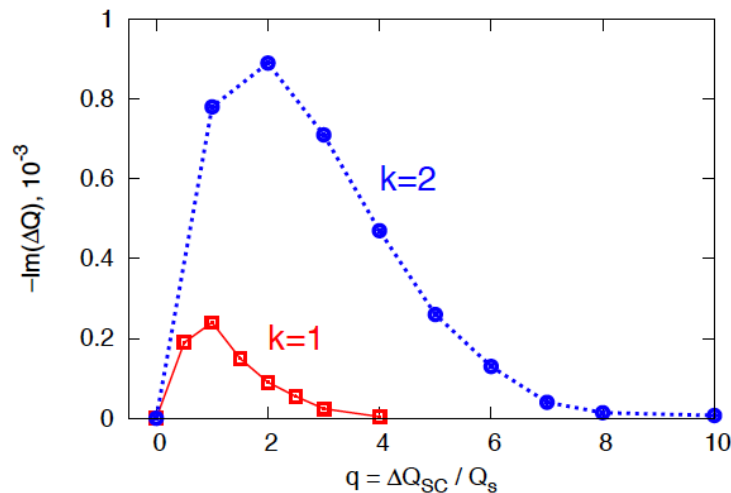


red: $k=1$; blue: $k=2$

a physical interpretation

upper boundary of the
incoherent spectrum:
small ΔQ_{SC} ,
large synchrotron amplitudes

$$\Delta Q_{\max} \approx -0.23Q_s q + kQ_s$$

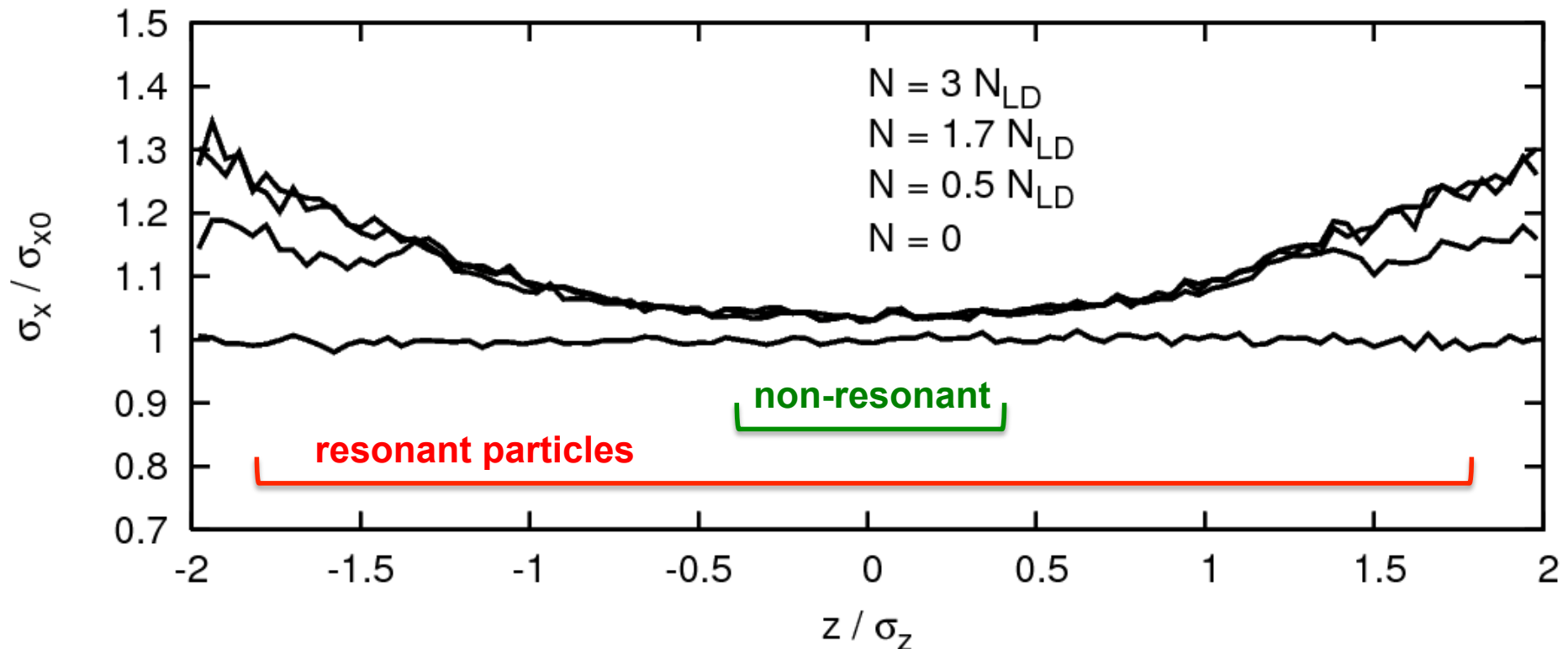


red: $k=1$; blue: $k=2$

Landau damping:

resonant particles with small ΔQ_{sc} have large synchrotron amplitudes, energy transfer happens in the bunch tails

this leads to the local emittance increase:



horizontal rms beam size along the bunch,
a simulation for $k=2$, $q=3$

airbag bunch:
analytical results
for unstable head-tail modes
[Blaskiewicz 1998]

$$W(\tau) = W_0 \exp(-\alpha\tau)$$

$$\alpha\tau_b \gg 1$$

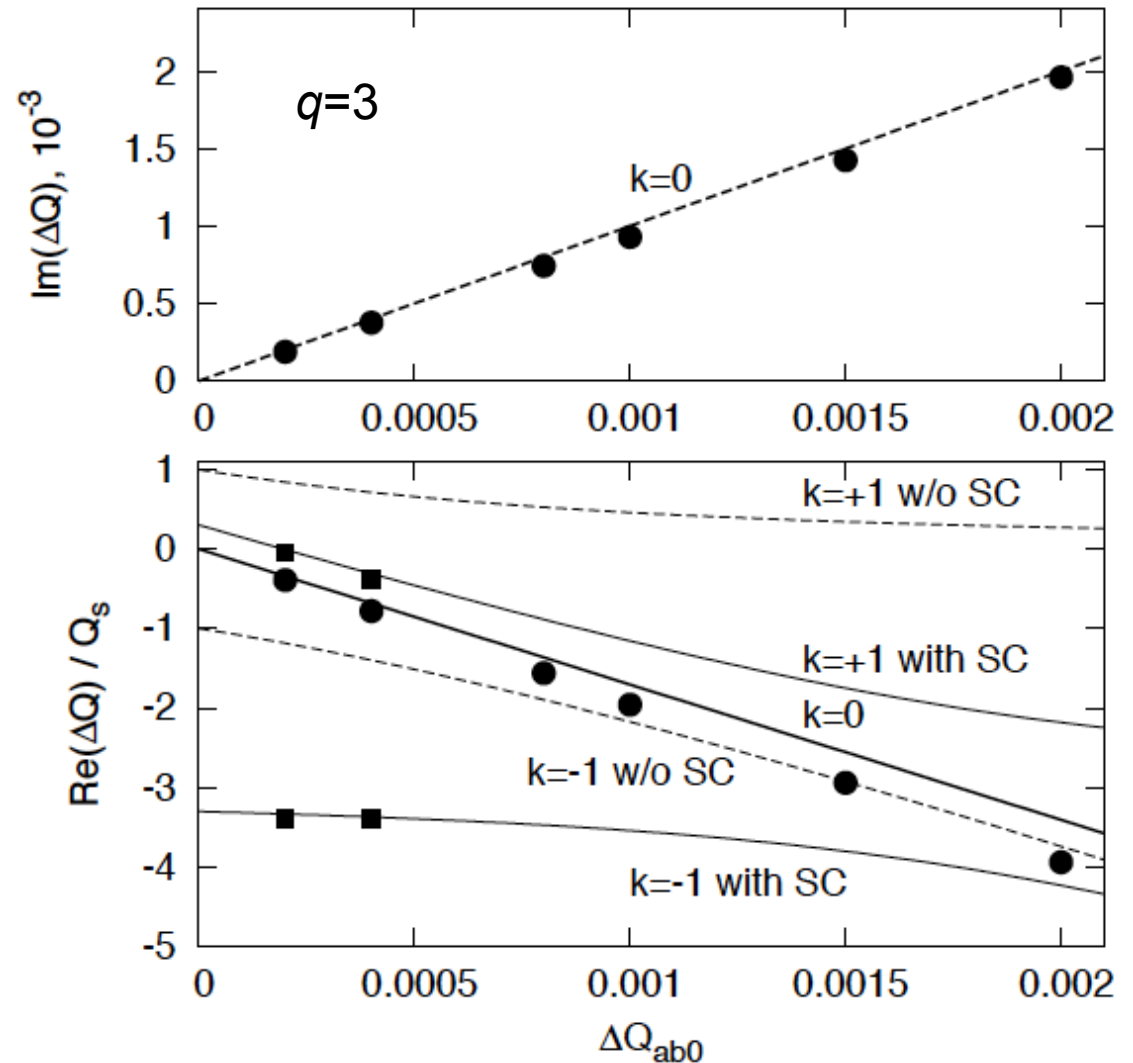
for $k=0$
(not affected by space charge)

$$\Delta Q = \Delta Q_0(\alpha/\zeta + i)$$

$$\Delta Q_0 = -\frac{\kappa\zeta}{\alpha^2}W_0$$

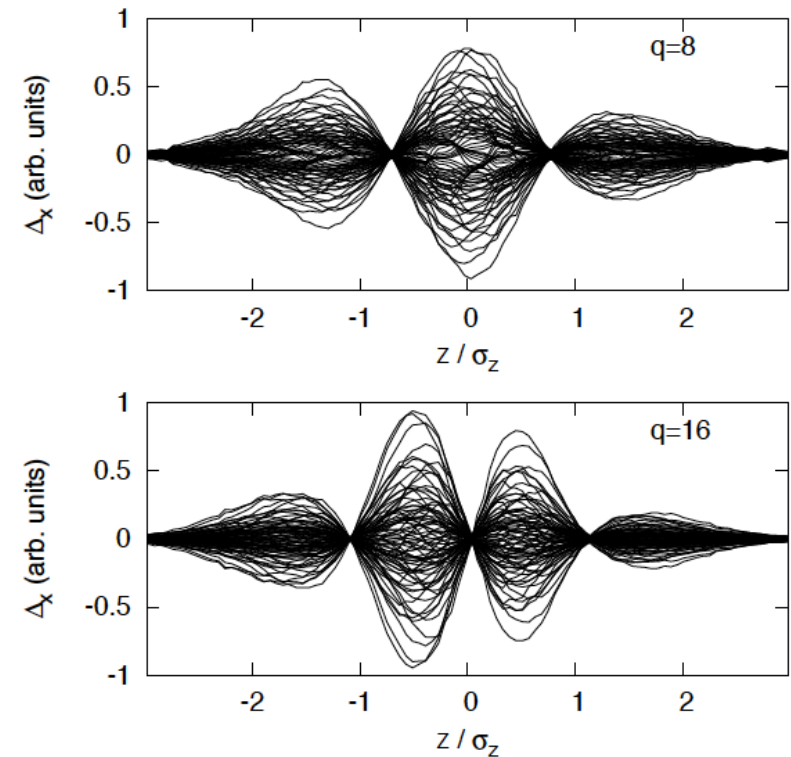
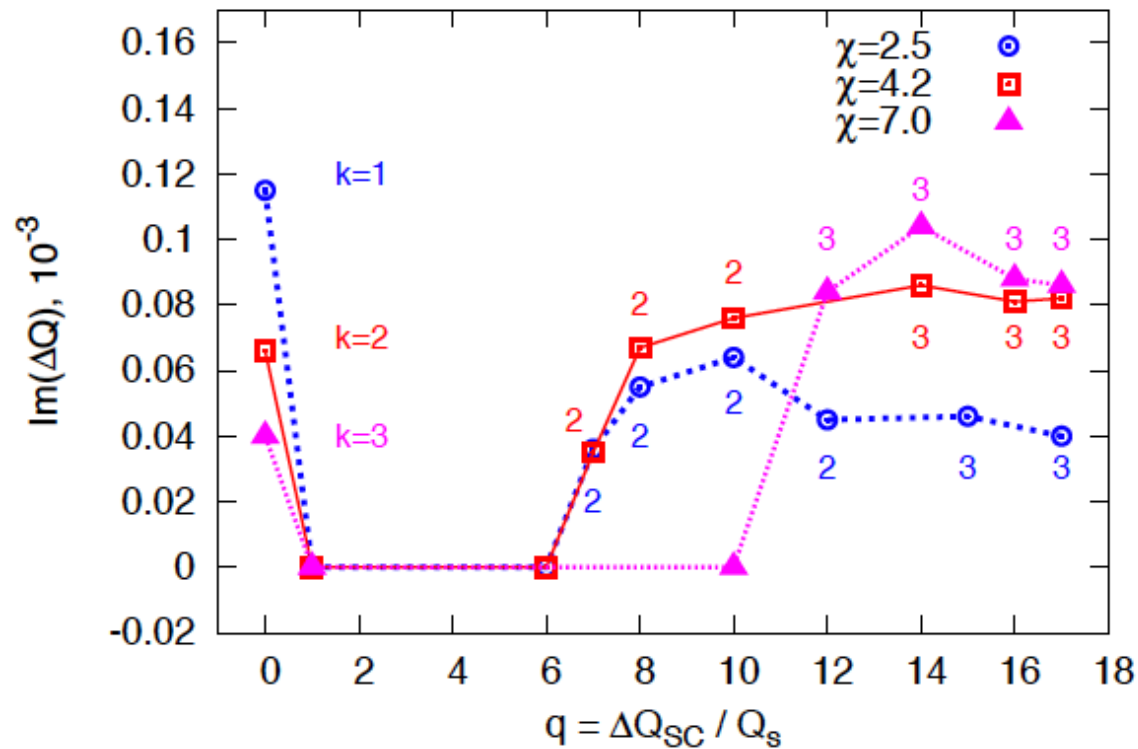
for $k>0$

$$\Delta Q = -\Delta Q_{sc} + \frac{\Delta Q_0(\alpha/\zeta + i) + \Delta Q_{sc} \pm \sqrt{[\Delta Q_0(\alpha/\zeta + i) + \Delta Q_{sc}]^2 + 4k^2Q_s^2[1 - (\Delta Q_0\pi/2\zeta Q_s\tau_b)^2]}}{2[1 - (\Delta Q_0\pi/2\zeta Q_s\tau_b)^2]}$$



Simulations for a Gaussian bunch:
growth rates of the most unstable head-tail mode

$$\chi = \xi Q_0 \tau_b / \eta = 4.2$$



Landau Damping!

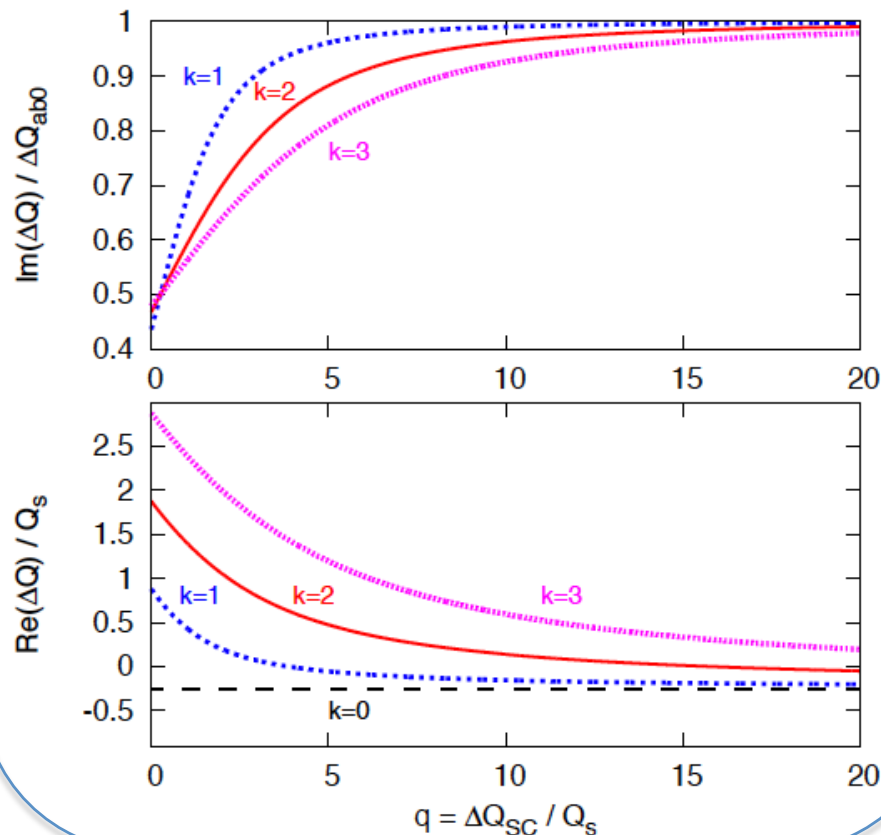
wake function

of the thick resistive wall:

$$W_{rw}(z) = -\frac{cL_{rw}}{b^3} \left(\frac{\beta}{\pi}\right)^{3/2} \sqrt{\frac{Z_0}{z \sigma_{rw}}}$$

simulations: the growth rates saturate at strong space charge

analytic theory
for the airbag bunch
[Blaskiewicz 1998]



calculations in [Burov 2009]:
treating a wake as a perturbation
→ diagonal element of the wake operator

$$\Delta Q = \frac{\kappa}{N_{\text{ion}} \lambda_0 R} \int_0^{z_b} dz \int_z^{z_b} ds W(s-z) d_k(s) d_k^*(z)$$

$$d_k(s) = \lambda(s) \bar{x}_k(s)$$

provides the growth rate
at strong space charge

comparisons with the simulation results
show a good agreement!

Landau damping in bunches exclusively due to space charge observed in particle tracking simulations.

The damping strength depends on $q = \Delta Q_{sc}/Q_s$ and on the mode index k . The energy transfer happens in the bunch tails.

Gaussian bunch: the transverse functions and frequencies are very similar to that predicted by the airbag theory [Blaskiewicz 1998]

Simulations of the head-tail instability with space charge:

Landau damping can stabilize the bunch.

The instability growth rates saturate at strong space charge (compared to [Blaskiewicz 1998] and [Burov 2009]).

Examples:

- ♦ CERN PS (E.Métral et al, PAC07)
 - $q \sim 150$: far above Landau damping, growth rates saturated
- ♦ FAIR SIS100
 - $q \sim 20$: Landau damping might contribute to the stability