



INNOVATIVE SCHEMES OF PLASMA HEATING FOR FUTURE MULTIPLY-CHARGED IONS SOURCES: MODELING AND EXPERIMENTAL INVESTIGATION

*G. Castro*¹, *D. Mascali*¹, *G. Torrisi*¹, *L. Celona*¹, *O. Leonardi*¹, *M. Mazzaglia*^{1,2}, *D. Nicolosi*¹, *R. Reitano*^{1,2}, *G. Sorbello*^{1,3}, *C. Altana*¹, *F. P. Romano*^{1,4} and *S. Gammino*¹

¹LNS-INFN, Catania, Italy

²Università degli studi di Catania, dipartimento di Fisica, Catania, Italy

³Università degli studi di Catania, DIEEI, Catania, Italy

⁴CNR-IBAM, dipartimento di Fisica, Catania, Italy

Outline

- **Motivations:**
 - Overcoming the actual limitation of ECRIS

- **Introduction:**
 - The Electron Bernstein waves (generation and absorption)

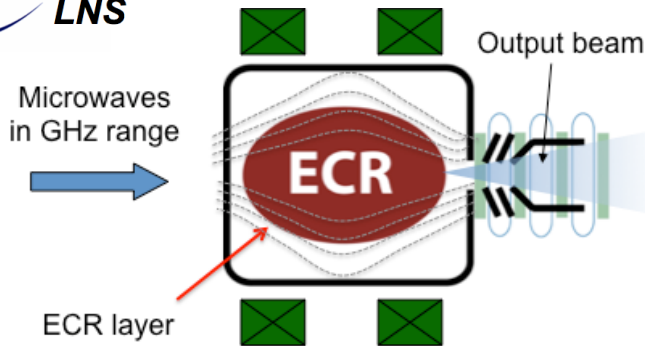
- **Experimental apparatus:**
 - The plasma reactor;
 - The diagnostics;

- **Experimental evidences:**
 - Overcome of the density cutoff (LP and interferometric techniques);
 - Ancillaries phenomena;

- **New perspectives at INFN-LNS:**
 - The FPT and launcher;

- **Conclusions**

Overcoming the actual limit of ECRIS



Beam characteristics

$$\langle q \rangle \propto n_e$$

$$I \propto n_e$$

$$\omega_p^2 < \omega^2 \rightarrow n_e < \frac{\epsilon_0 m_e}{e^2} \omega^2 = n_{cutoff}$$

ECRIS STD MODEL

$$n_e \propto \omega_{RF}^2 \rightarrow \text{INTRINSIC Density limitation}$$

1. **Higher Frequency Generators to increase the plasma density;**
2. **Higher Magnetic Fields to make longer the ions confining time;**

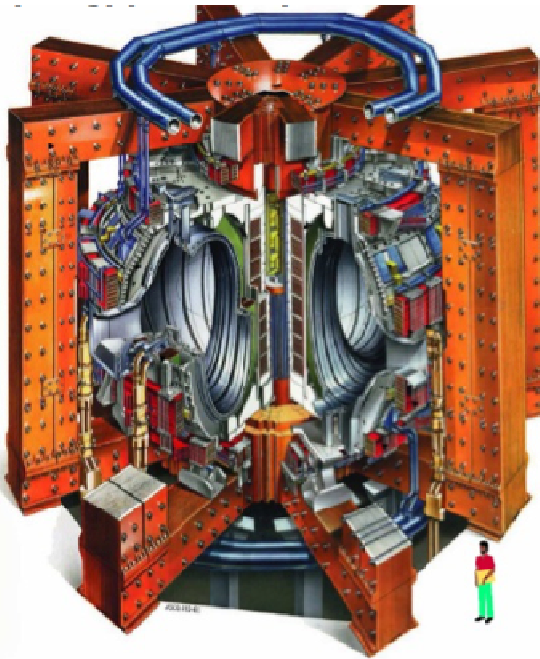
Brute force cannot be anymore used because of technological reasons (magnets, hot electrons generations, plasma overheating, cooling, ...)

Optimization of wave to plasma coupling

Alternative heating schemes

A look in the world of the fusion plasmas

Jet tokamak for nuclear fusion (energy production)



Also in the field of fusion, higher and higher density plasmas are required to satisfy the **Lawson criterion**;

Overdense plasmas have been generated by using Electron Static waves, and, in particular, Electron Bernstein Waves, instead of EM waves;

10 times overdense plasma have been generated at Greiswald (Ge) stellarator by using EBW heating

Y. Y. Podoba *et al.*, *Physical Review Letters* 98, 255003, 2007.

Outline

- **Motivations:**
 - Overcoming the actual limitation of ECRIS

- **Introduction:**
 - The Electron Bernstein waves (generation and absorption)

- **Experimental apparatus:**
 - The plasma reactor;
 - The diagnostics;

- **Experimental evidences:**
 - Overcome of the density cutoff (LP and interferometric techniques);
 - Ancillaries phenomena;

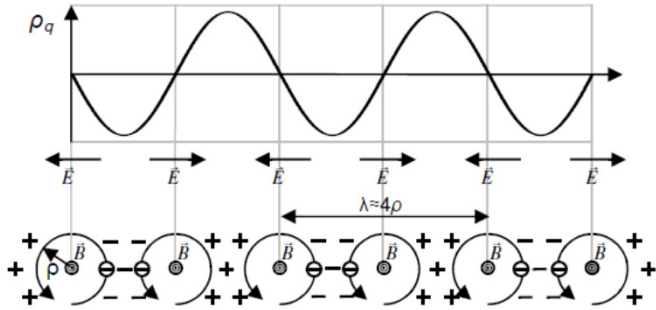
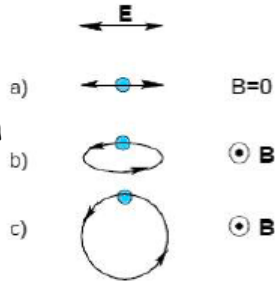
- **New perspectives at INFN-LNS:**
 - The FPT and launcher;

- **conclusions**

Bernstein waves

EBW consists in phase bunching of electrons in a magnetized plasma.

They generalized Langmuir waves

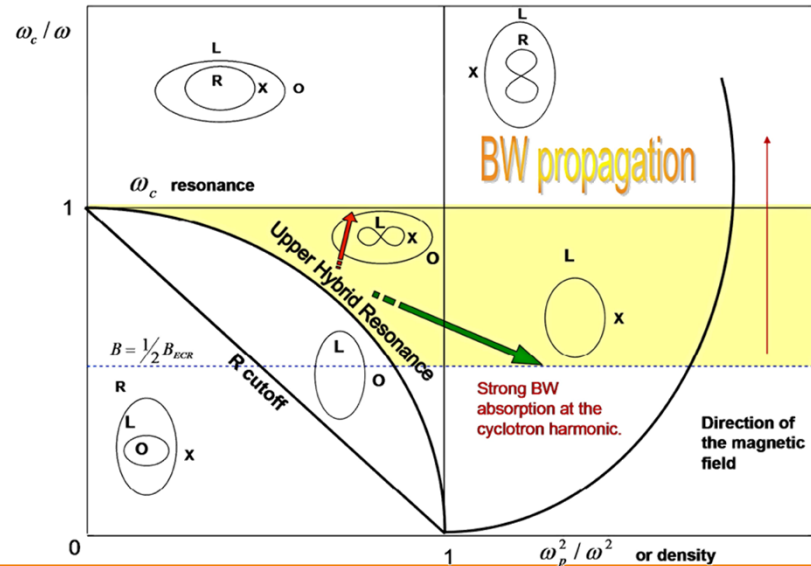


No Density cutoff

Resonance at cyclotron harmonics

Excited at UHR by X-waves

The BW propagation region



From EM waves to EBW

1. High field side launch

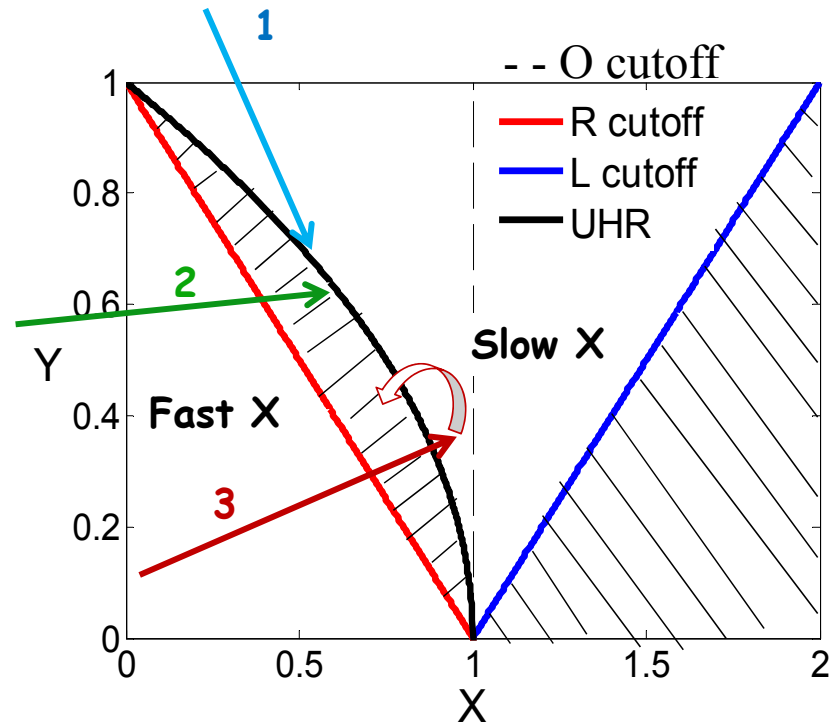
S-X wave is launched from high magnetic field to UHR

2. Direct FX-B conversion

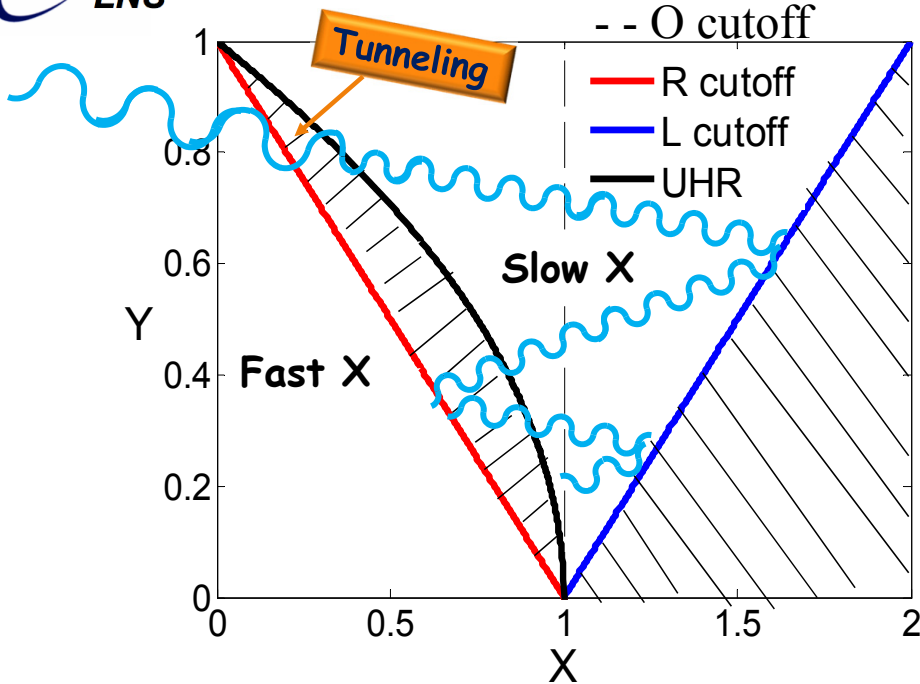
F-X tunnels trough R cutoff and reaches UHR

3. O-SX-B conversion

An O wave is launched. O-SX conversion at O cutoff, S-X reach UHR



Direct FX-B conversion



Conversion factor from FX to B wave:

$$C_{FX-B} = 4e^{-\pi\eta}(1 - e^{-\pi\eta}) \cos\left(\frac{\phi}{2} + \theta\right)$$

$$\eta = 294 |BL_n|^{UHR} \quad L_n = \frac{n_e}{\partial n_e / \partial x}$$

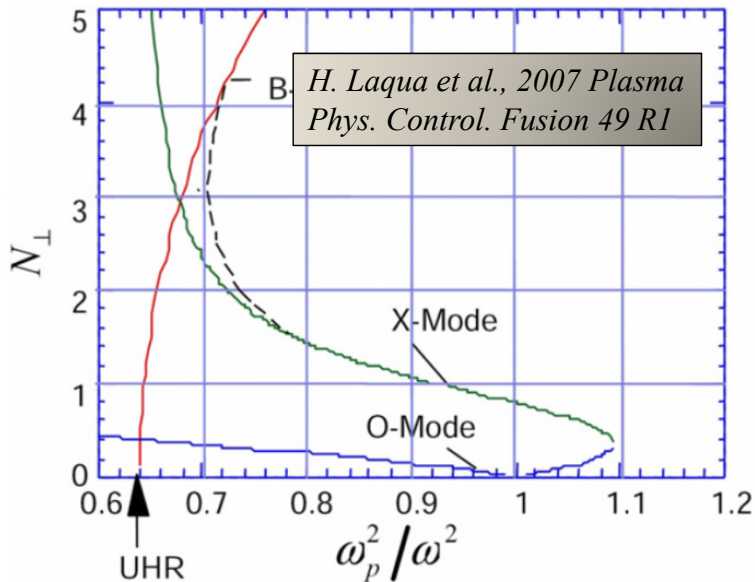
$$|C_{FX-B}| = 1 \text{ when } \eta \approx 0.22; \quad L_n \approx \frac{0.75}{B} [\text{mm}]$$

FX-B conversion optimized for
steep density gradient,
low magnetic field and **low frequency**:

Jones B et al 2003 *Phys. Rev. Lett* **90** 165001 [FX-B conversion in tokamak](#)

Chattopadhyay P K et al., 2002 *Phys. Plasma* **9** 752-5 [FX-B conversion in reversed pinch](#)

O-SX-B conversion



Conversion factor from O to SX wave:

$$T(N_{\parallel}, N_y) = \exp \left\{ -\pi k_0 L_n \sqrt{\frac{Y}{2}} \left[2(1+Y)(N_{\parallel, \text{opt}} - N_{\parallel})^2 + N_y^2 \right] \right\}$$

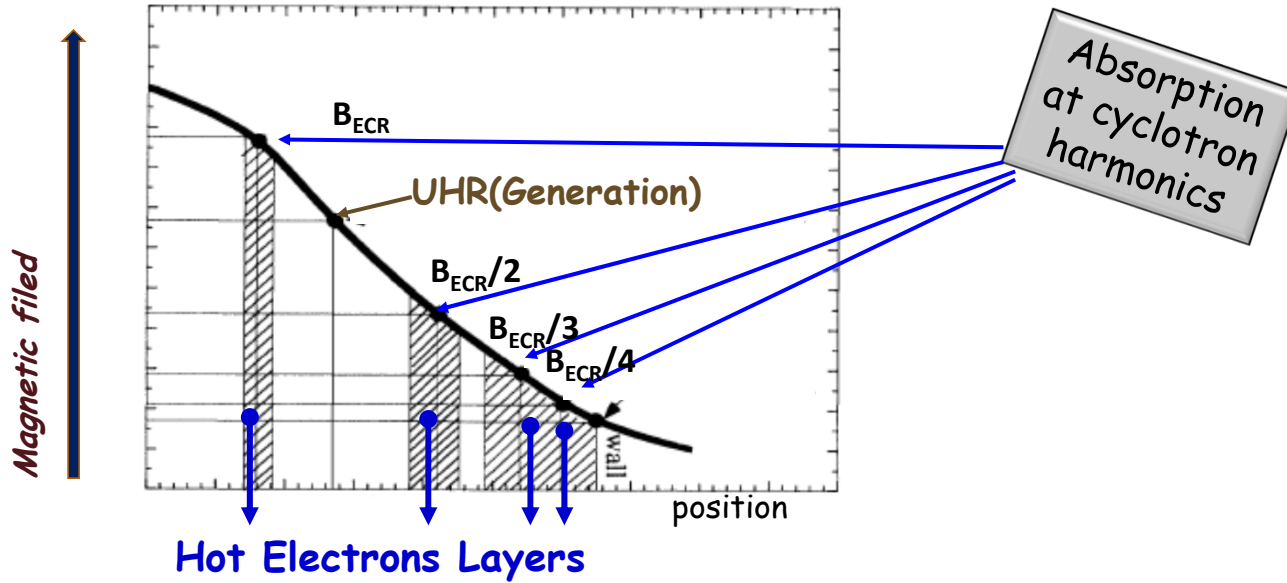
O-SX-B conversion optimized for
flat density profile and $\vartheta_{\text{launch.}} = \vartheta_{\text{opt.}}$

$$\sin^2 \theta_{\text{opt}} = N_{\parallel, \text{opt}}^2 = \frac{Y}{1+Y} \quad L_n \approx \frac{50-1000}{f [\text{GHz}]} [\text{mm}]$$

**Optimal launching angle and higher frequency
 required**

Y. Y. Podoba, et al., *Physical Review Letters* 98 (2007) 255003. [O-SX-B conversion in stellarator](#)

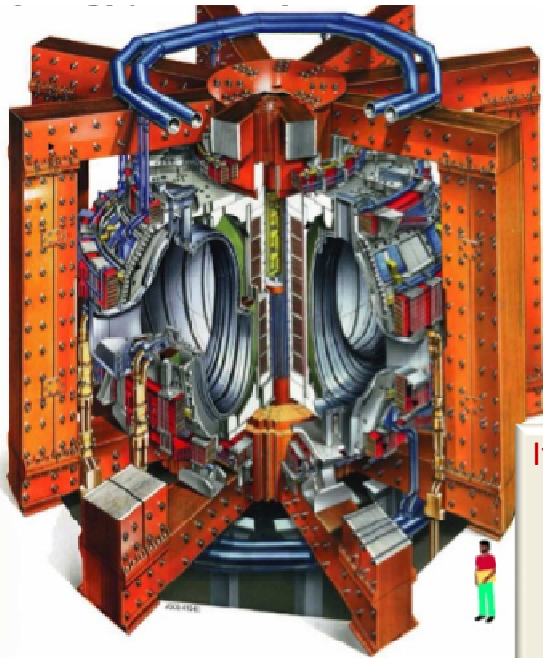
Absorption of Bernstein waves



Non linear behaviour

The main challenge: "downsizing"

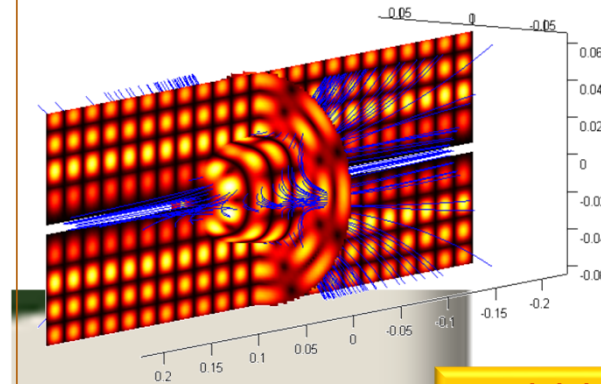
Jet tokamak for nuclear fusion (energy production)



$$\lambda \ll L$$

It is possible to launch E.M waves with optimized angles (optical approximation) for increase EM-to EBW conversion

ECR Ion Source: extremely compact plasma machine



$$\lambda \approx L$$

Modal dominated cavity.
Wave Launching is not possible («until now»)

Self-tuning needed to allows EBW heating

Outline

- **Motivations:**
 - Overcoming the actual limitation of ECRIS

- **Introduction:**
 - The Electron Bernstein waves (generation and absorption)

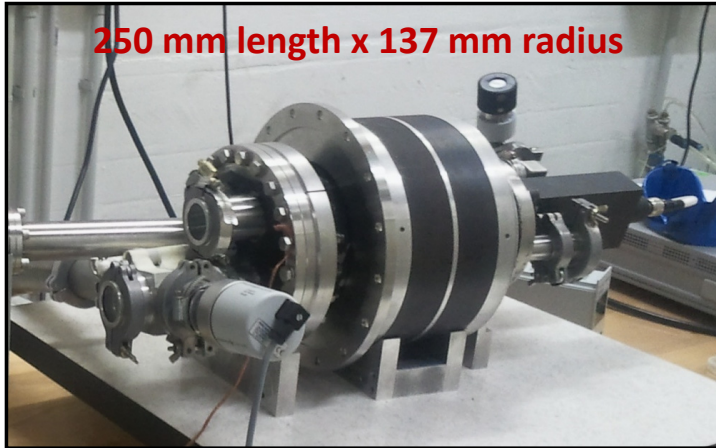
- **Experimental apparatus:**
 - The plasma reactor;
 - The diagnostics;

- **Experimental evidences:**
 - Overcome of the density cutoff (LP and interferometric techniques);
 - Ancillaries phenomena;

- **New perspectives at INFN-LNS:**
 - The FPT and launcher;

- **conclusions**

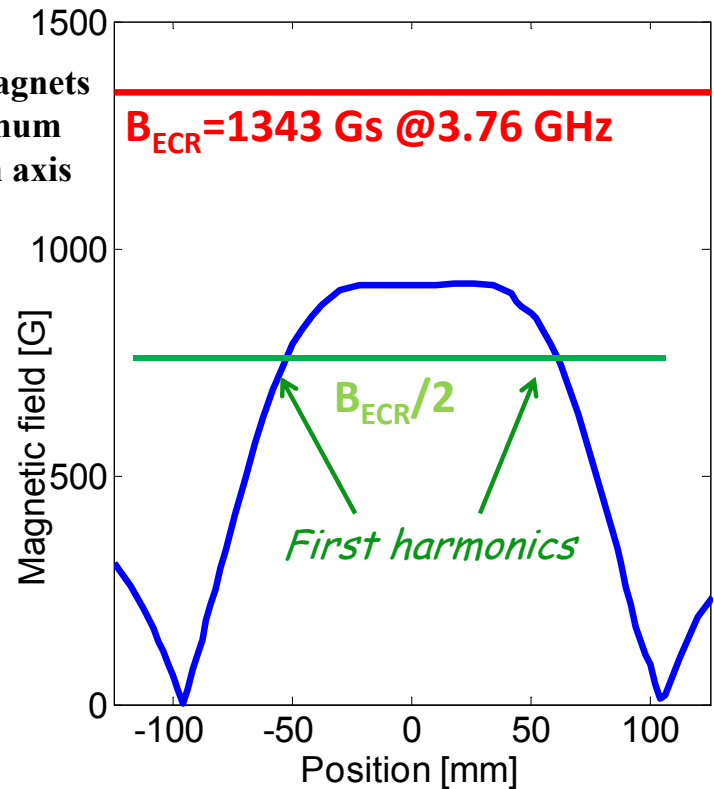
Set-up for EBW plasma generation



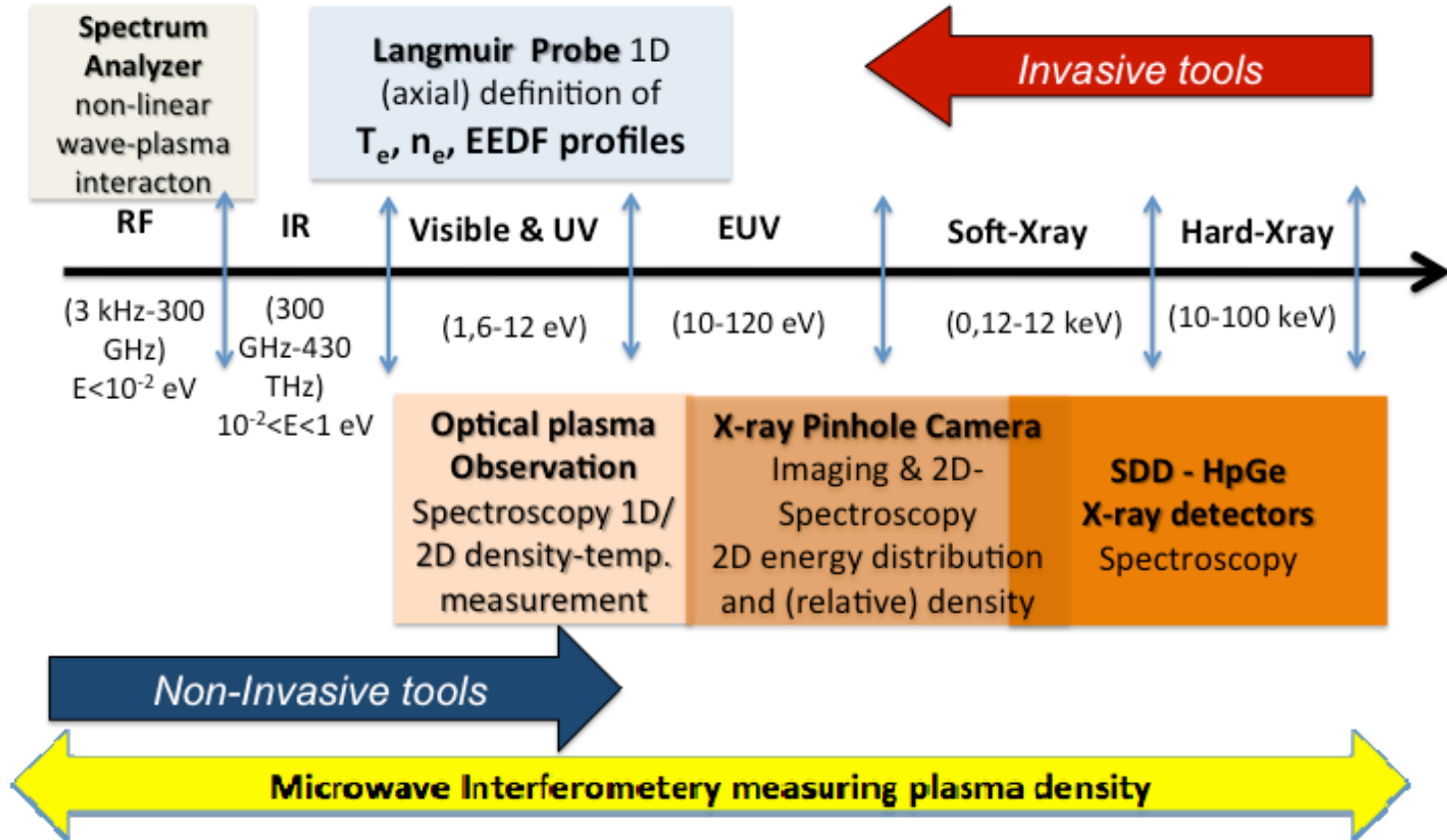
The permanent magnets produce a maximum field of ~ 0.1 T on axis

These devices typically work at **2.45 GHz** for proton generation (**ECR @ 875 Gs**), Plasma reactor has been operated at **3.76 GHz** to focus on EBWs

No ECR possible



Plasma diagnostics methods



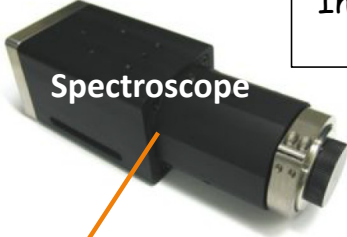
Plasma diagnostics tools



X ray detectors



X ray imaging
CCD camera



Information about ion waves and **ion heating**

Information about **total electron density** and electron density distribution

Information about expected Current, emittance, and brilliance

Information about **Electron Energy Distribution Function** and 2D X-ray distribution.

Outline

- **Motivations:**
 - Overcoming the actual limitation of ECRIS

- **Introduction:**
 - The Electron Bernstein waves (generation and absorption)

- **Experimental apparatus:**
 - The plasma reactor;
 - The diagnostics;

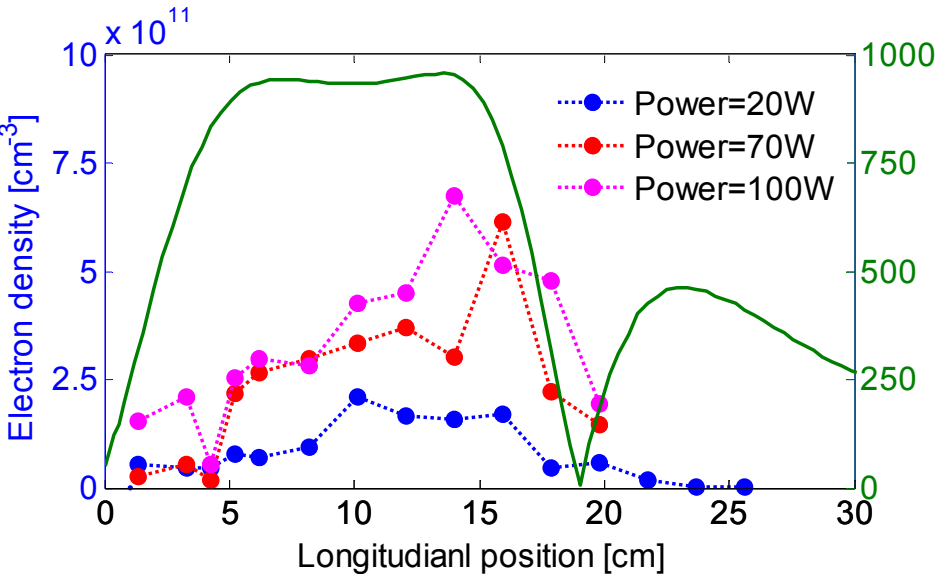
- **Experimental evidences:**
 - Overcome of the density cutoff (LP and interferometric techniques);
 - Ancillaries phenomena;

- **New perspectives at INFN-LNS:**
 - The FPT and launcher;

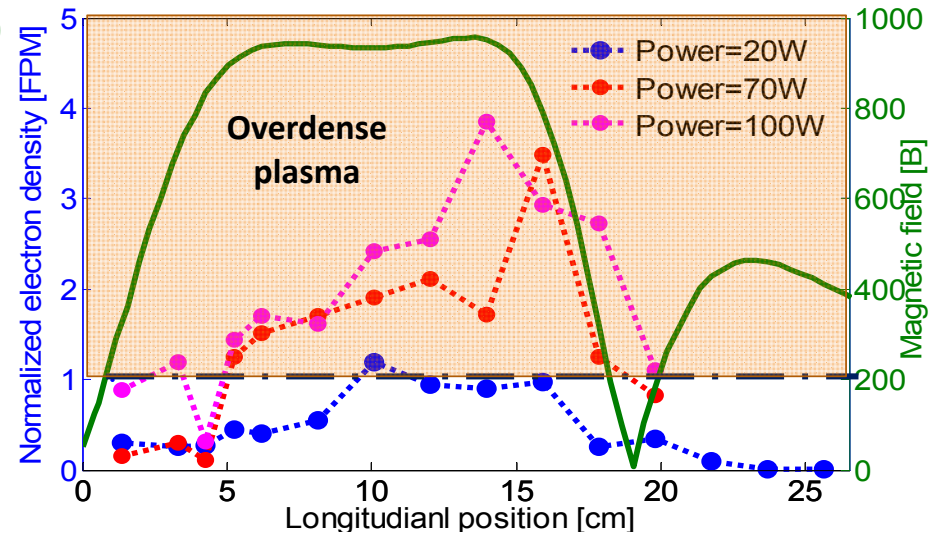
- **conclusions**

Results from LP measurements-1

Nitrogen gas @ $1.5 \cdot 10^{-4}$ mbar
3.76 GHz MicroWave frequency

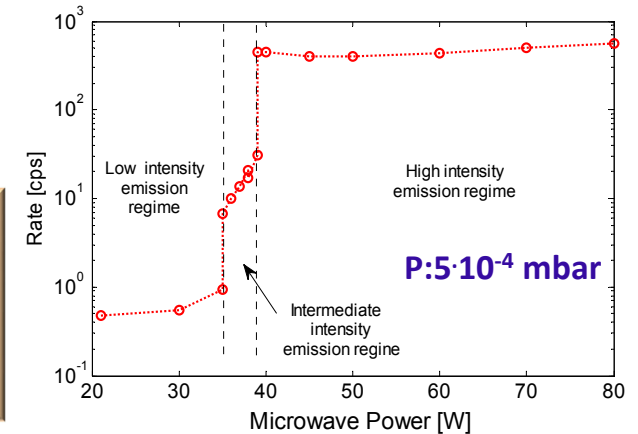
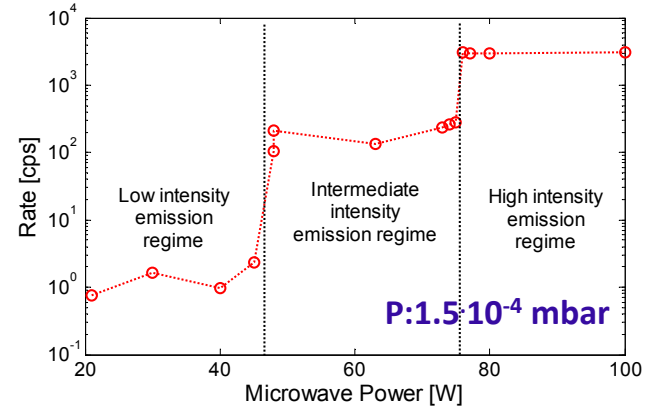
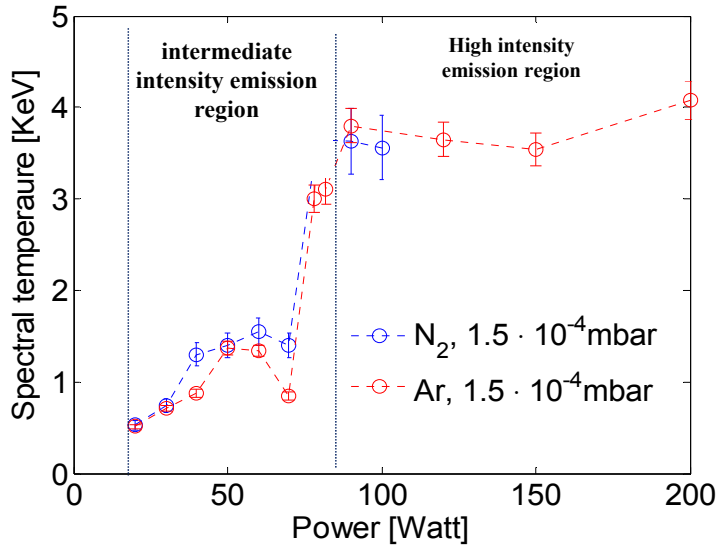


Density cut-off @3.76 GHz = $1.75 \cdot 10^{11} \text{ cm}^{-3}$



different models validated in high density (10^{10} - 10^{12} cm^{-3}) magnetized plasmas (RF)

X ray measurements from SDD

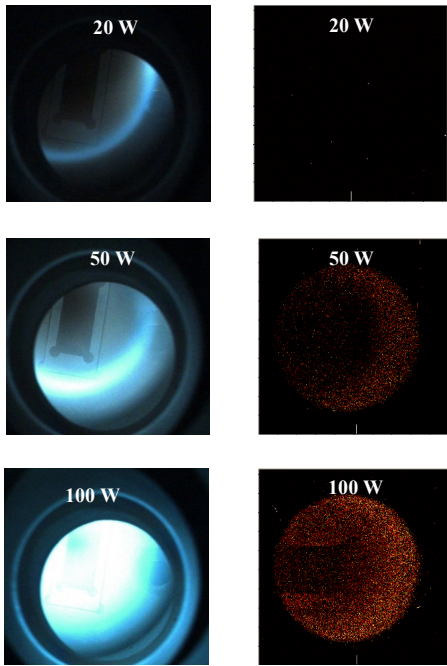


- Spectral temperature increases from ≈ 1 to 4 keV with MW power;
- Three emission regimes identified, depending on power and pressure;

From optical to X ray inspection

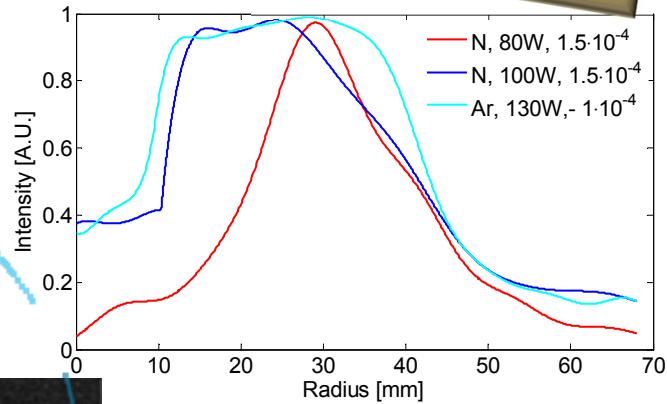
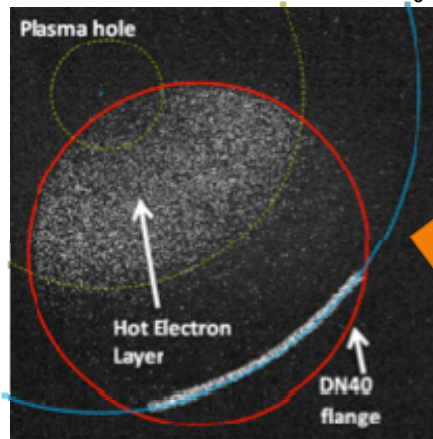
R. Racz talk: MODO01
for more information

Imaging in optical and X-ray domain



Transversal reconstruction of the plasma structure in X-ray domain (1-30 keV).

X ray imaging evidences that the pumping power is deposited in the annulus where the energetic electrons are generated

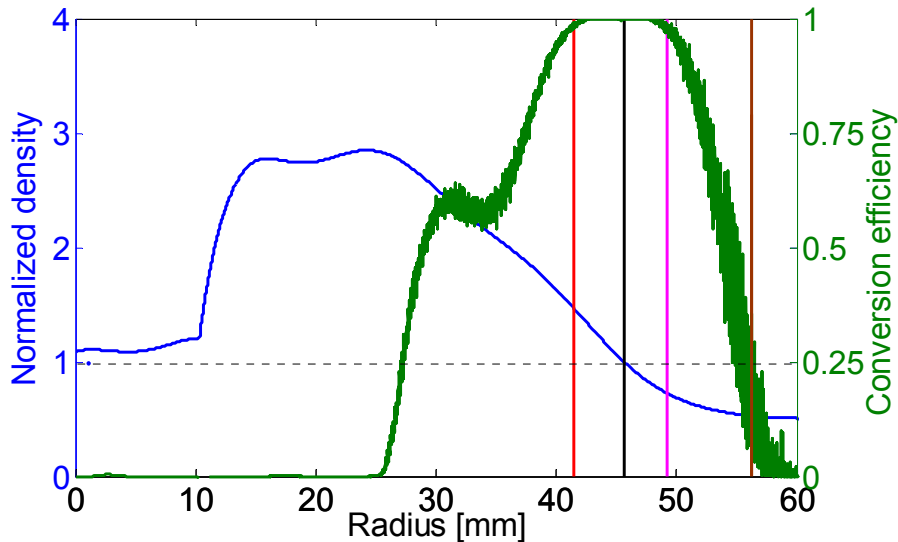


Radial density reconstruction

reconstruction

Radial profile and FX-B conversion

The knowledge of the radial profile allows us to calculate the absolute value of the conversion factor FX-B



Density gradient assumes the value which maximizes FX-B conversion

$$\eta = 294 |BL_n|^{UHR}$$

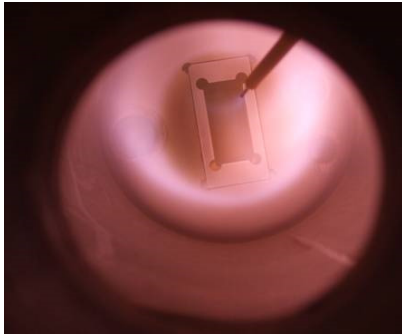
$$|C_{FX-B}| = 4e^{-\pi\eta}(1 - e^{-\pi\eta})$$

Radial gradient is the optimal one to allow FX-B conversion

FX-B conversion is the best candidate to explain the generation of **overdense plasma** in the plasma reactor

From LP diagnostics to interferometry

LP diagnostics



- i) Local plasma parameters measurement ✓
- ii) the density measurement is "model-dependent";
- iii) LP undergoes to plasma damage when density overcomes 10^{12} cm^{-3} ;
- iv) LP perturbs the plasma chamber, in particular by an electromagnetic point of view.
- v) Only cold electron population is characterized



Interferometry

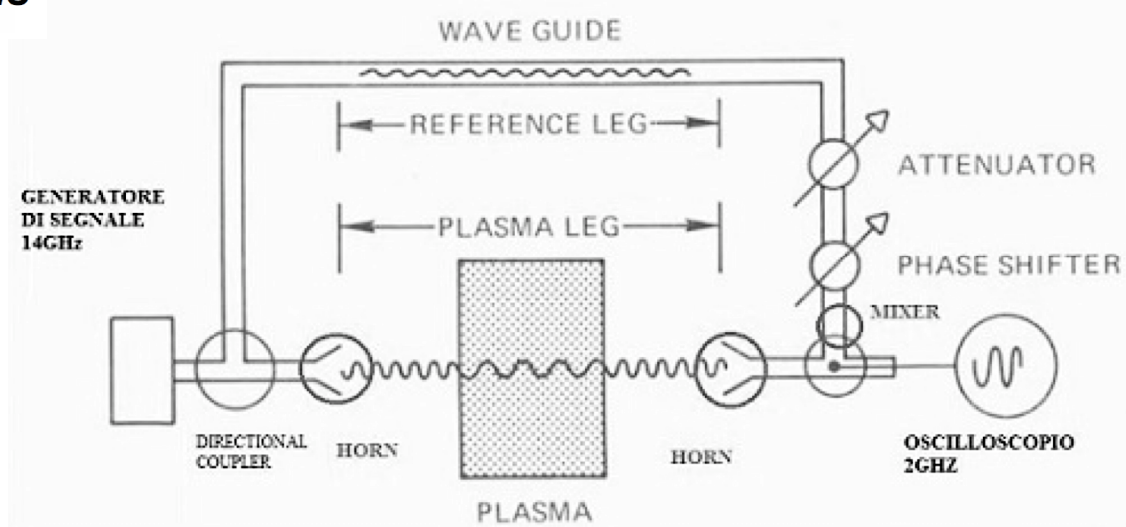


- i) Allows whole plasma density probing;
- ii) Non-invasive diagnostics; ✓
- iii) Well-known technique;
- iv) Probing wavelength $\lambda \sim L$ (plasma chamber dimension)

Modal behaviour of the probing wave



Classical interferometry for plasmas



Classical Scheme of Interferometer

$$\Delta\varphi = \frac{\omega}{c} \left[1 - \left(1 - \frac{\omega_p^2}{\omega^2} \right)^{\frac{1}{2}} \right] L$$

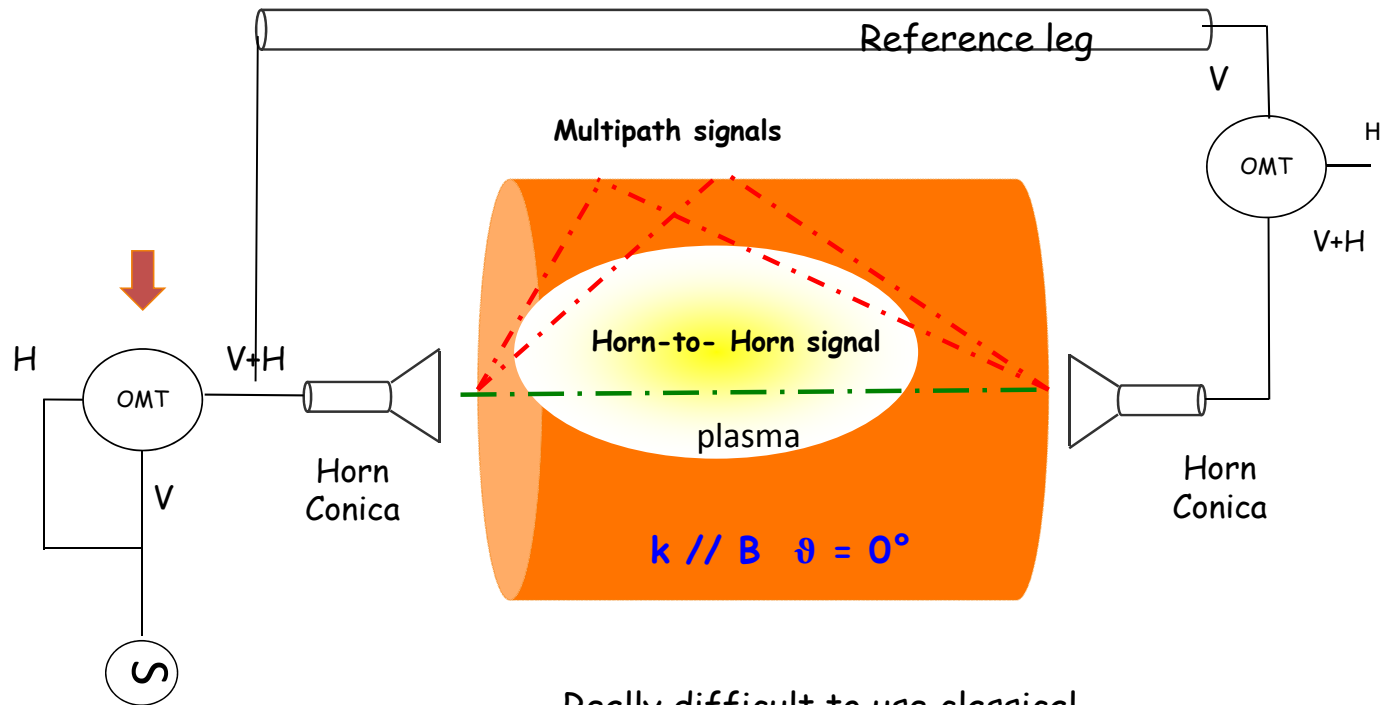
In plasmas the phase variation depends on the "natural plasma frequency"

$$\omega_p^2 = \frac{4\pi n e^2}{m \epsilon_0}$$

The plasma frequency depends on the density

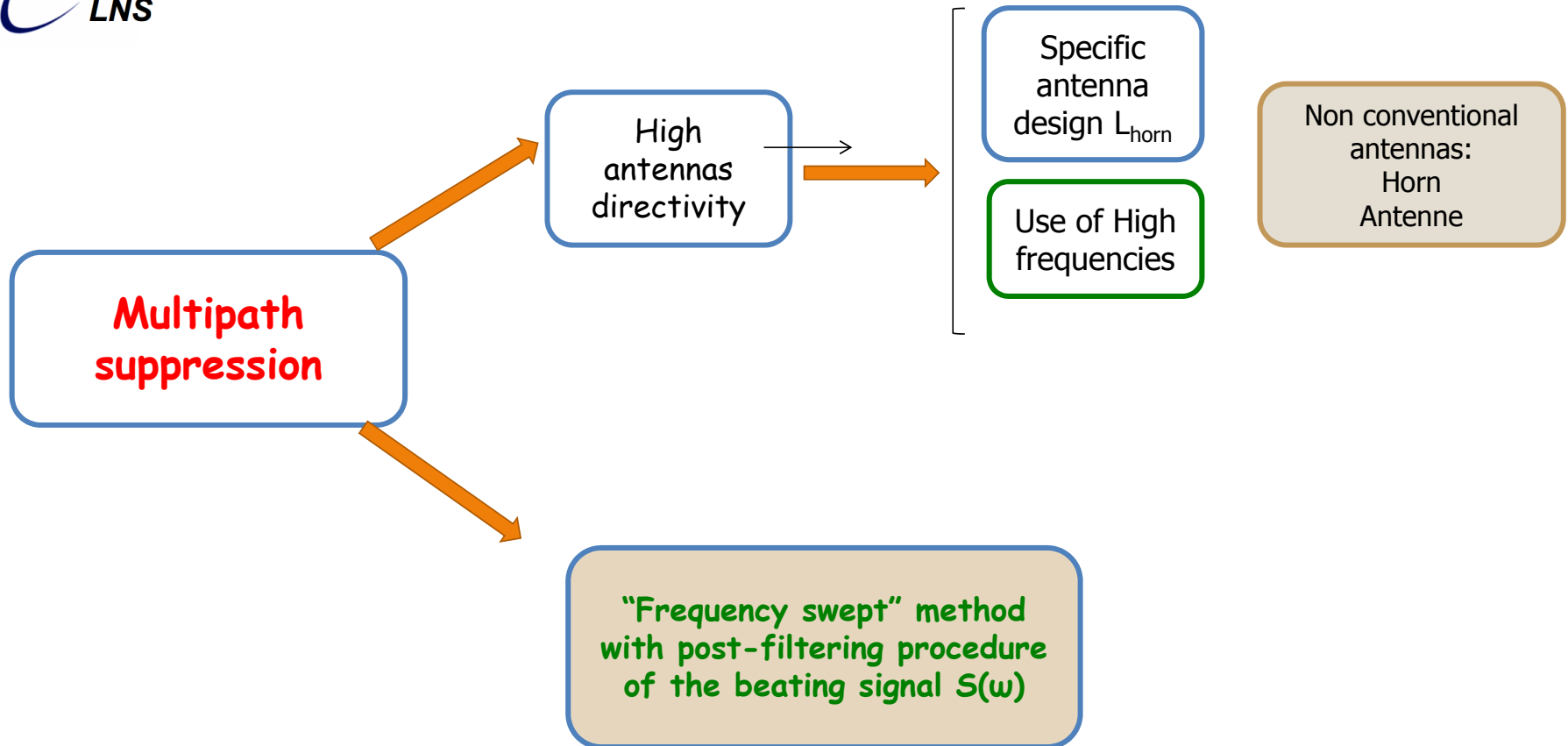
Microwave interferometry measures plasma density through a measurement of phase shift.

Criticality: multi-paths introduce spurious signals in small devices



Really difficult to use classical microwave interferometry

Design challenges



Design of a “frequency swept” interferometer

The method is based on the frequency shift of the beating signal given by the superposition of the reference leg plus the plasma leg waves.

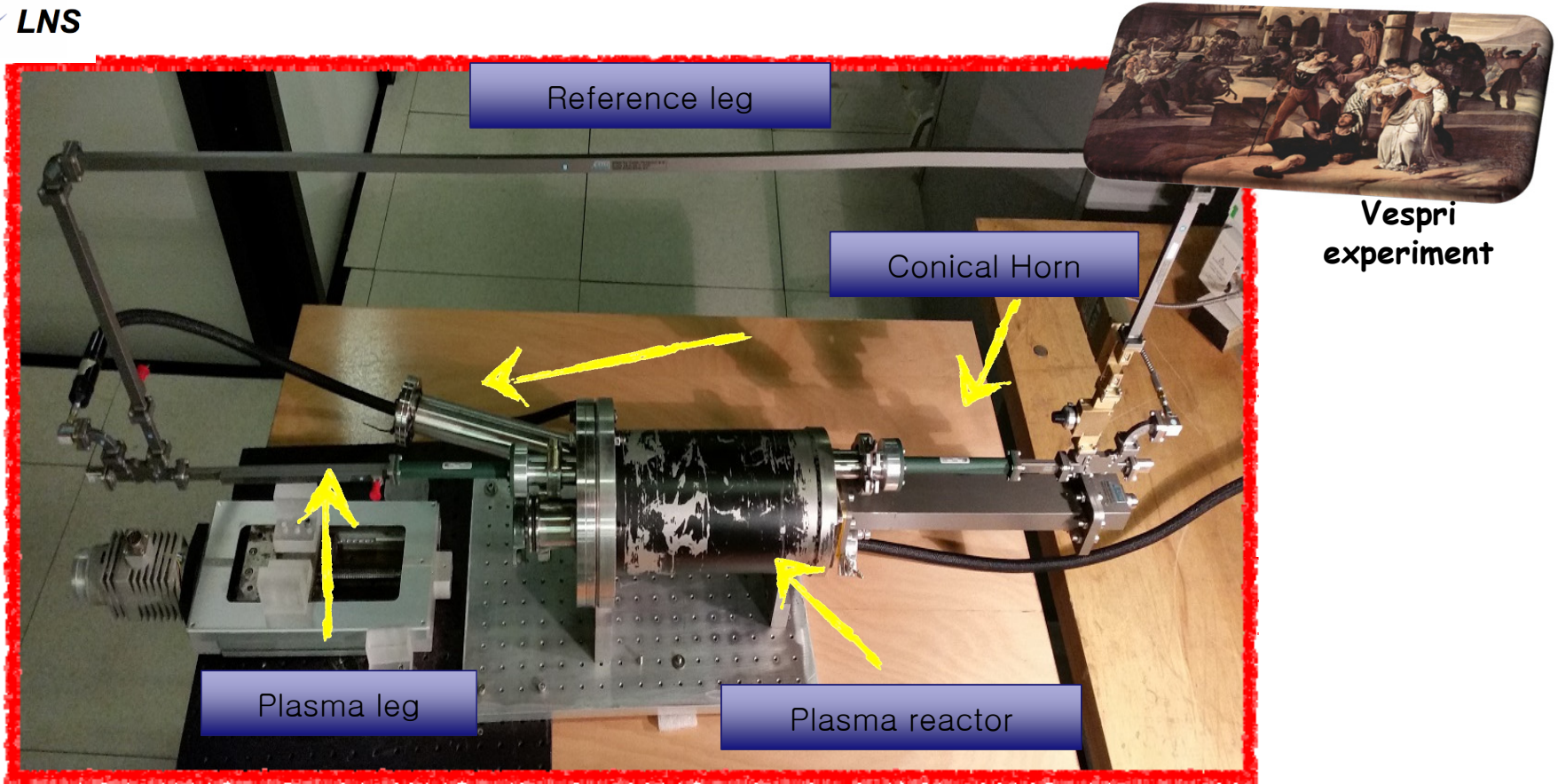
The beating frequency can be fixed as long as the ramp relation “freq. vs. time” is chosen in the following way:

$$\omega_{B_0} = \frac{\partial \omega}{\partial t} \left(\Delta L \frac{\partial k_g}{\partial \omega} + \frac{2a}{c} \right) = \text{constant}$$

$$S(\omega) \propto 2A^2 \cos^2 \left\{ \left[\Delta L \sqrt{\omega^2 - \omega_c^2} + \int_{-a}^a \sqrt{\omega^2 - \omega_p^2(l)} dl \right] / 2c \right\}$$

The presence of plasma (accounted by the plasma frequency ω_p) only shifts the beating frequency, while multipath introduce spurious components in the spectrum

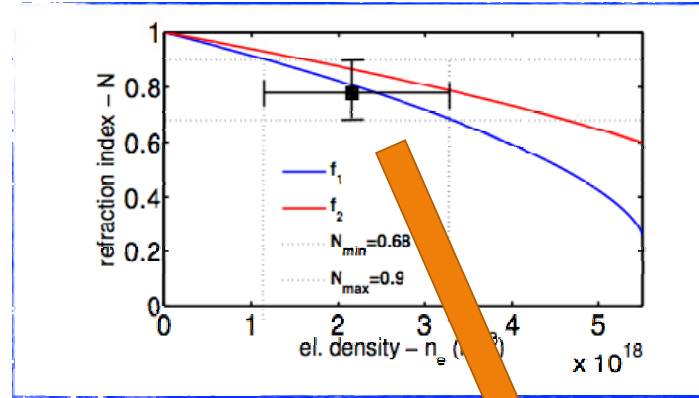
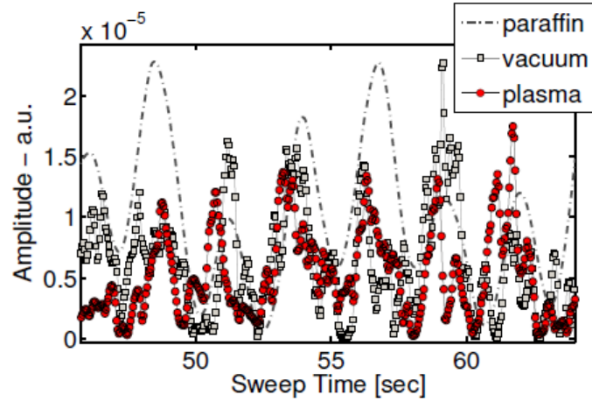
Experimental setup @ INFN-LNS



Interferometry in ECRIS:

the measurement by Freq. Sweep method

Courtesy of G. Torrissi



Medium	P_{RF} [W]	ω_{beat} [rad/s]	Refr. Index N
Empty cavity	0	$2\pi * 0.397$	1
Bulk paraffin	0	$2\pi * 0.380$	1.43
Plasma	150	$2\pi * 0.407$	0.792 ± 0.11

Interferometric method:
 $n_e = (2 \pm 1) \cdot 10^{12} \text{ cm}^{-3}$

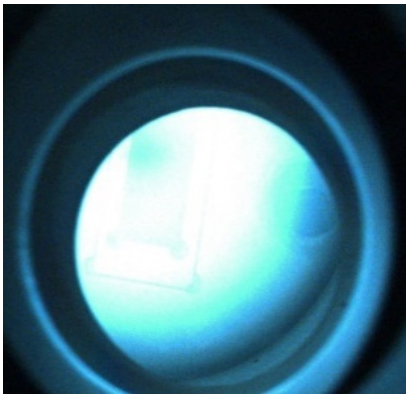
By literature $N_{paraffin} = 1.45$

Remarks of density measurements on plasma reactor

Overdense plasma generated by means of FX-B conversion:

- **4 times overdense** according to LP diagnostics
- **10 times overdense** according to interferometry

	Electron density (cm^{-3})
Interferometry	$(2.1 \pm 1) \cdot 10^{12}$
LP diagnostics	$(5.5 \pm 1.5) \cdot 10^{11}$
$n_{cut-off}$ (at 3.76 GHz)	$1.75 \cdot 10^{11}$



This result has been obtained in a small plasma reactor:

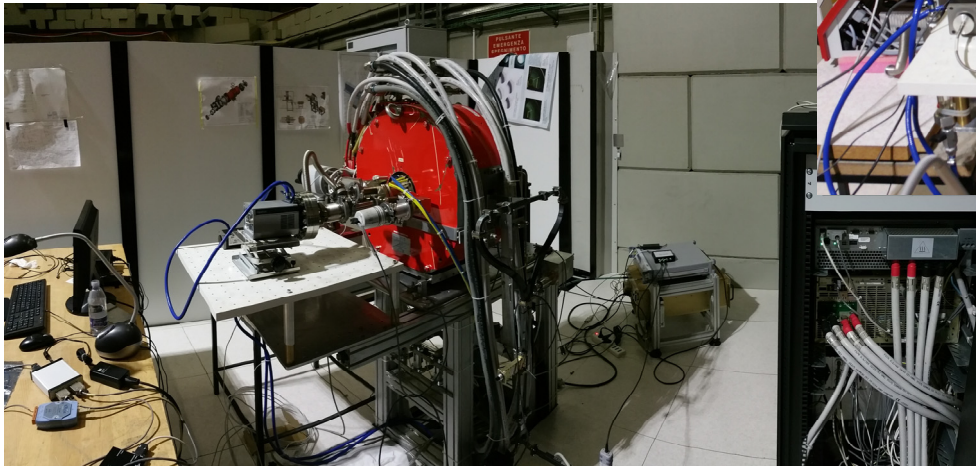
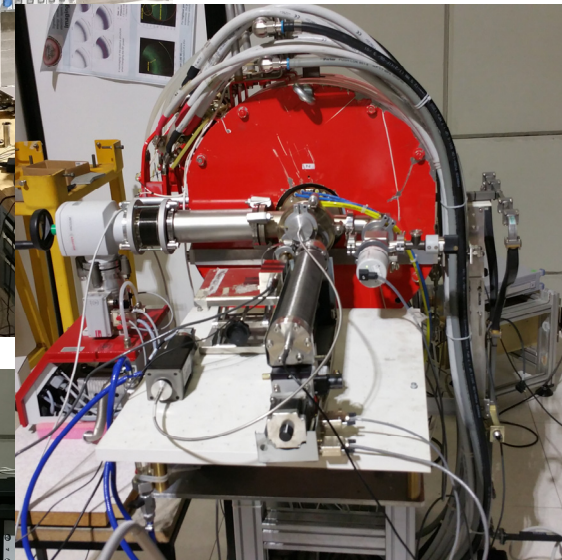
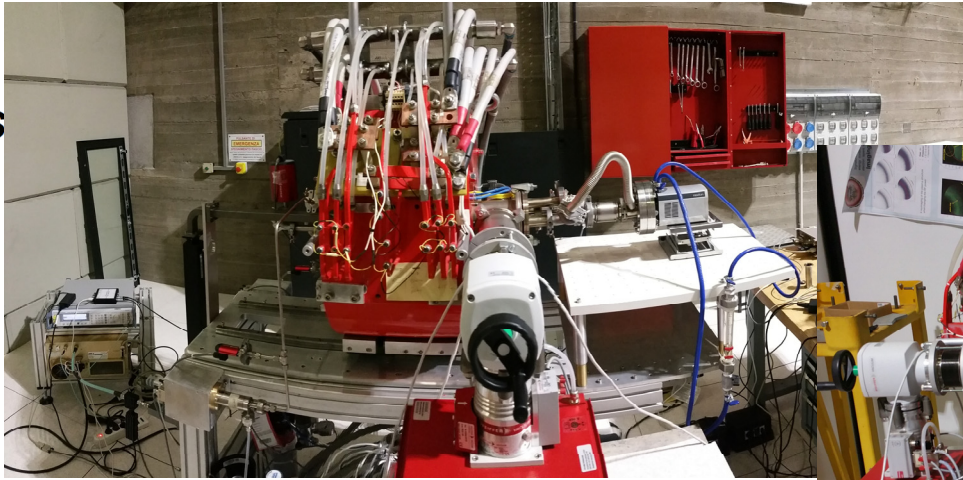
- **No chance to modify magnetic field;**
- **No chance to launch the wave with wanted angles (cavity mode dominated system);**
- **Modal dependent conversion: Fine frequency tuning needed to couple X-mode to EBW;**



We want to move from a cavity modes dominated system towards a microwave orientend system:



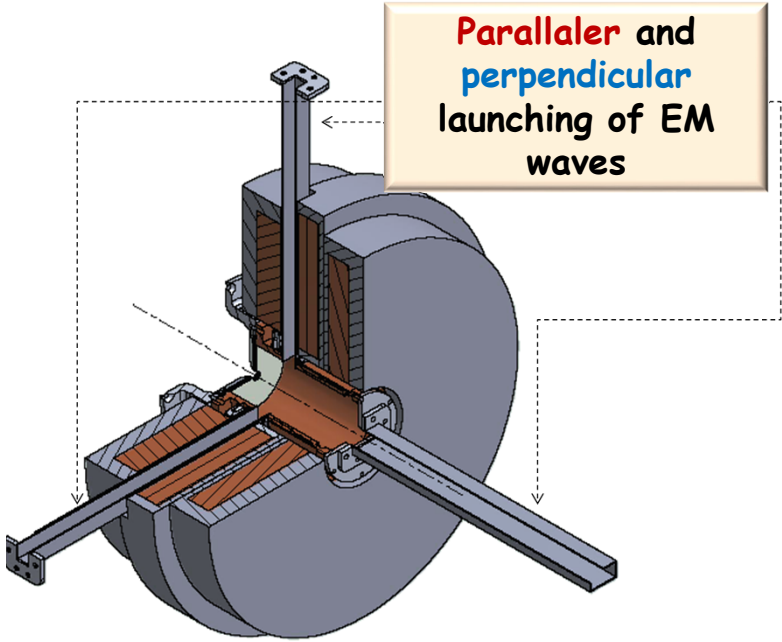
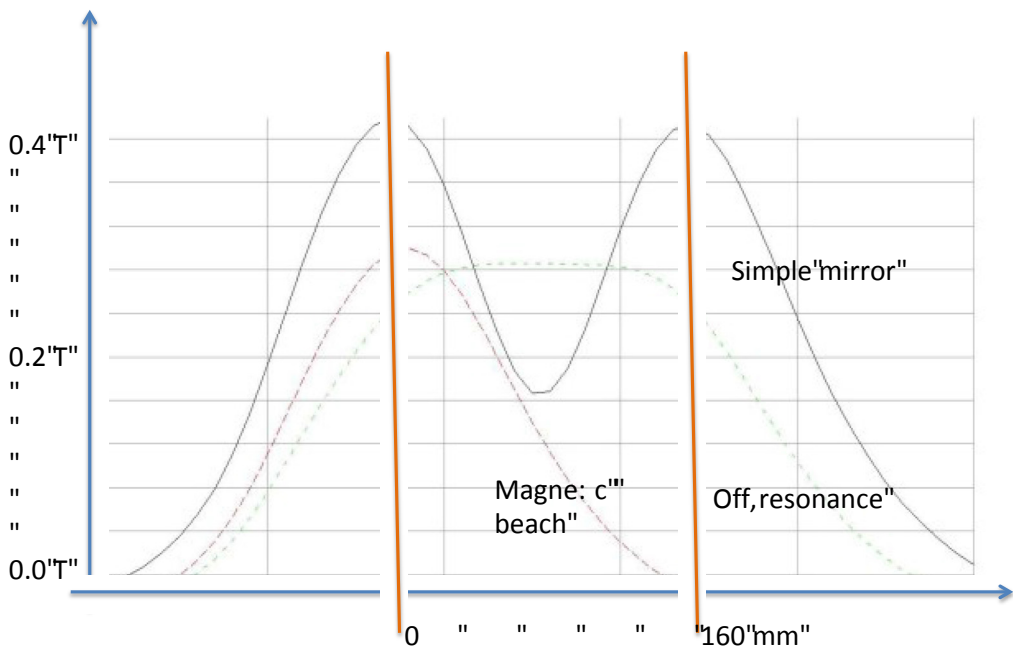
The New Flexible Plasma Trap and launcher at INFN-LNS



Flexible
Plasma
Trap

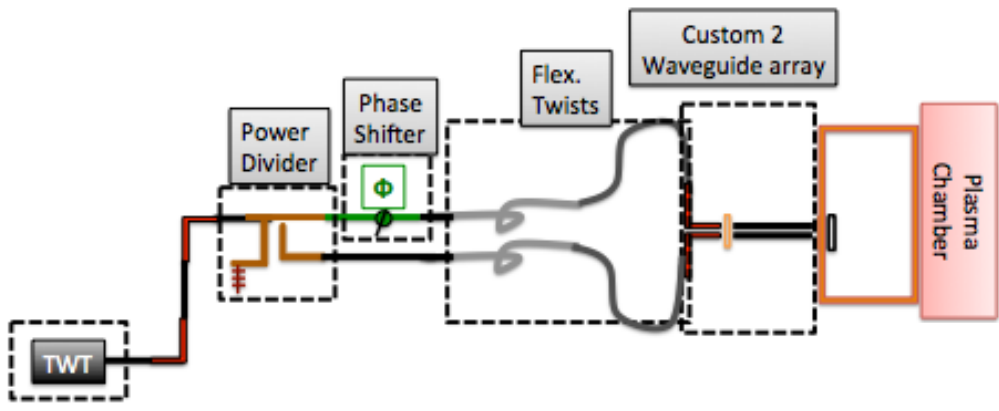
Now operating
at LNS

The FPT magnetic field and the microwave injections



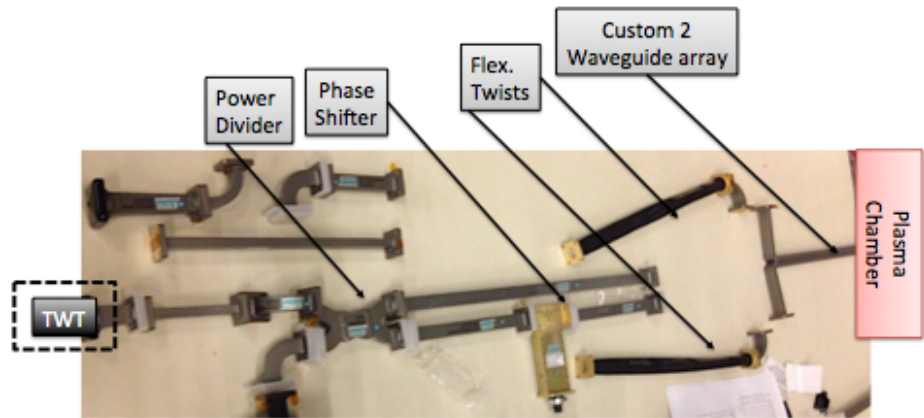
The launcher: mechanical Implementation at INFN-LNS

Phased waveguide array of two elements



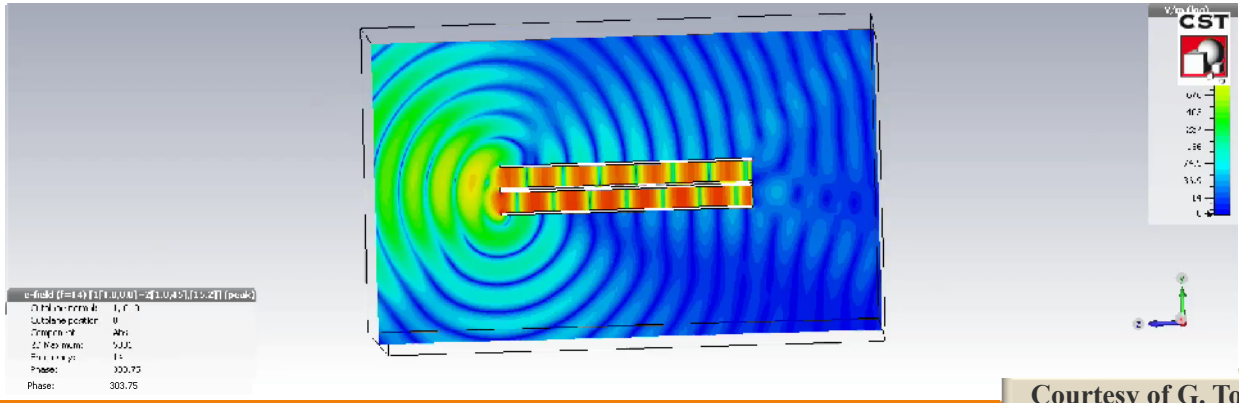
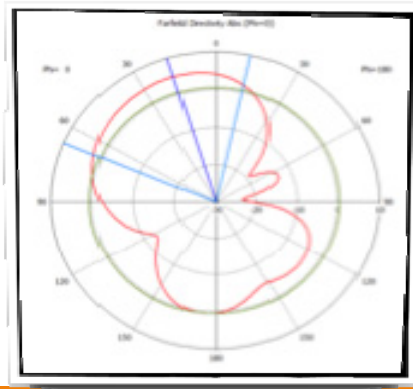
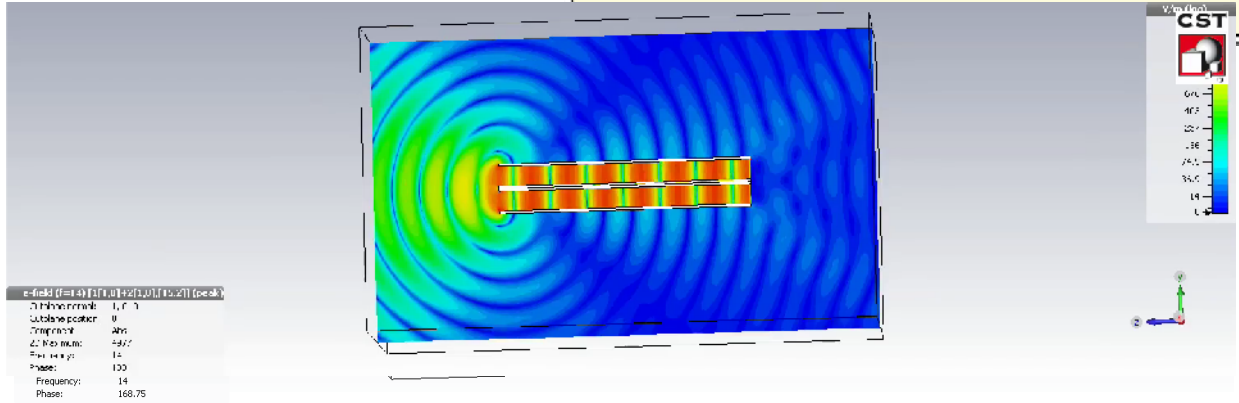
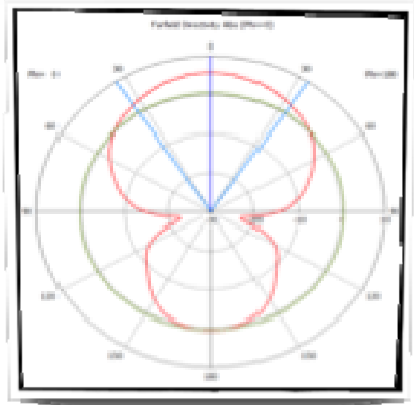
microwave launcher:
proposed setup

- TWT ampl. 13.75-14.5 GHz
- Power divider
- Phase shifters
- flexible waveguides



“Microwave-absorption-oriented” design

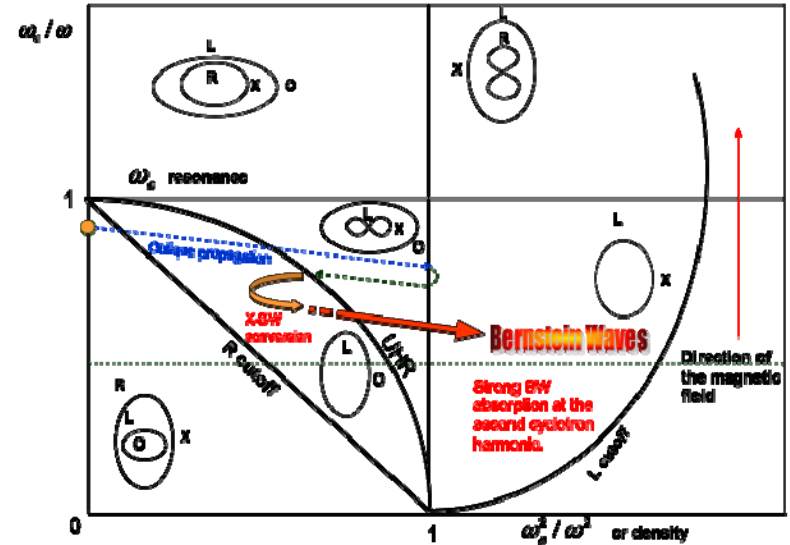
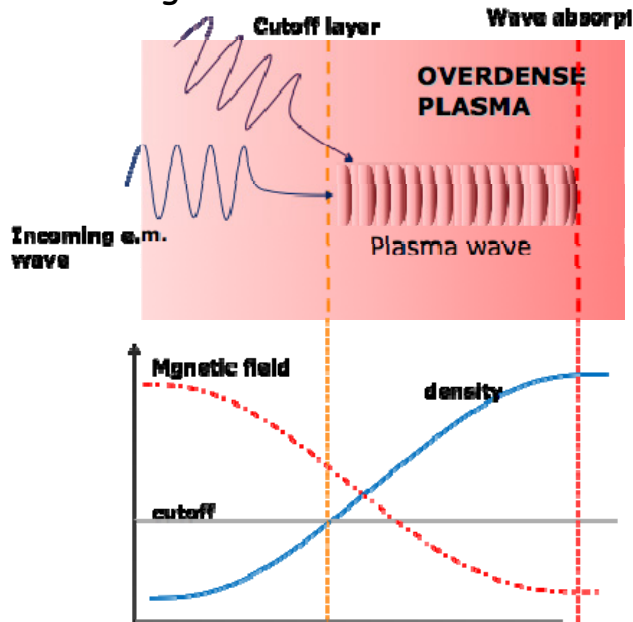
Launcher by “two-waveguides-array”: lobe tilt by phase shift for optimizing oblique coupling of O-modes -->



Courtesy of G. Torrisi

“Microwave-absorption-oriented” design

- 1) Primary plasma generated at 7 GHz by ECR.
- 2) A secondary plasma generated by perpendicular launching at 14 GHz



Oblique launching allow a more efficient matching of the incoming microwave radiation with the self-oscillations of the magnetized plasma

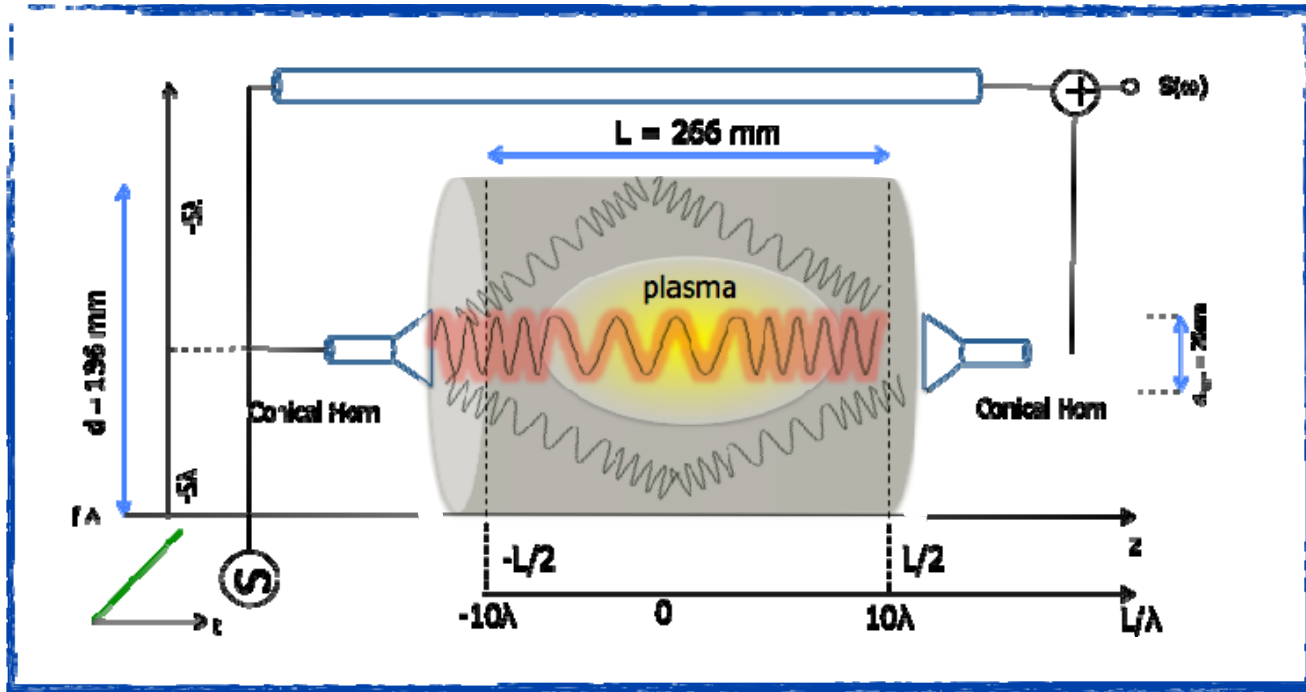
Conclusions

- **New schemes of plasma heating can allow to overcome ECRIS density limitation:**
- At infn-lns, an overdense plasma has been generated in a small plasma reactor. Density was **~5 times overdense** according to LP, **~10 times overdense** according to interferometry;
- **Characteristic signatures of Bernstein waves generation and absorption was found:**
- Non linear plasma behaviour (different plasma regimes identified), overdense generation and ion noise; FX-B conversion is the candidate to explain Bernstein generation;
- **The development of new diagnostics tools represents a fundamental task for the comprehension of the plasma behaviour. A set of diagnostic tools has been installed INFN-LNS** (spectroscopy, interferometer, X-ray detector and X-ray pinhole camera,...). They form a unique diagnostics set for the plasma characterization in different temperature and density ranges; **Information about expected current - emittance and brilliance**
- **A Microwave-absorption-oriented design** is being developed at INFN. The new FPT consists of three coils, three microwave injections (1 parallel, 2 orthogonal). A new launcher will launch microwave with desired angles within the plasma chamber;

Thank you
for your attention



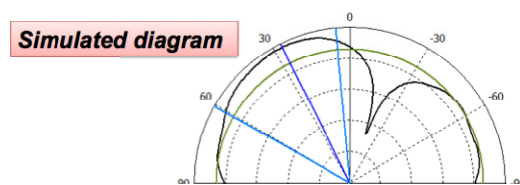
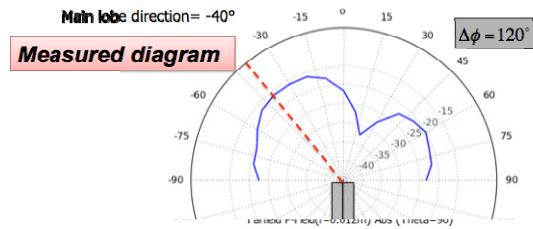
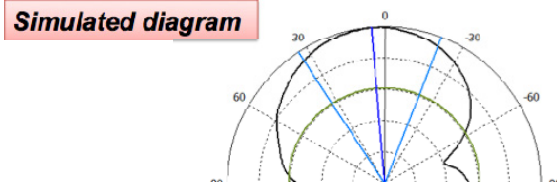
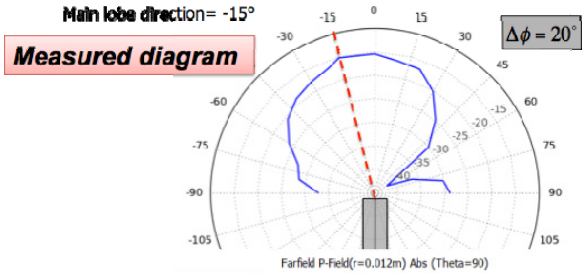
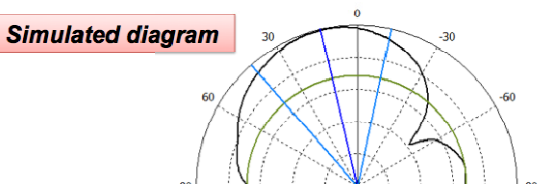
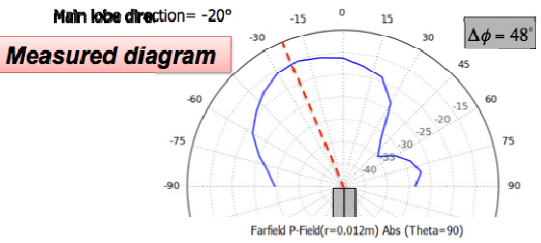
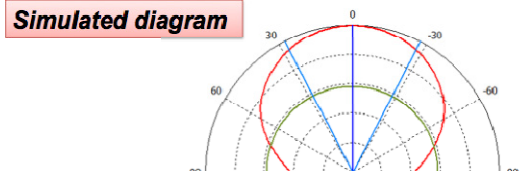
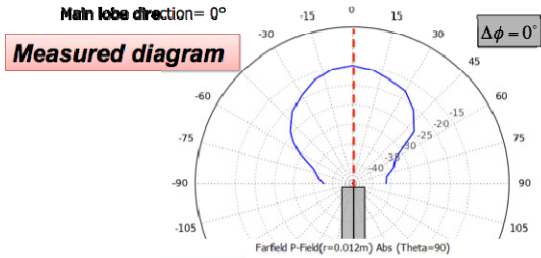
Probing the plasma density in all the energy domains



Drawbacks



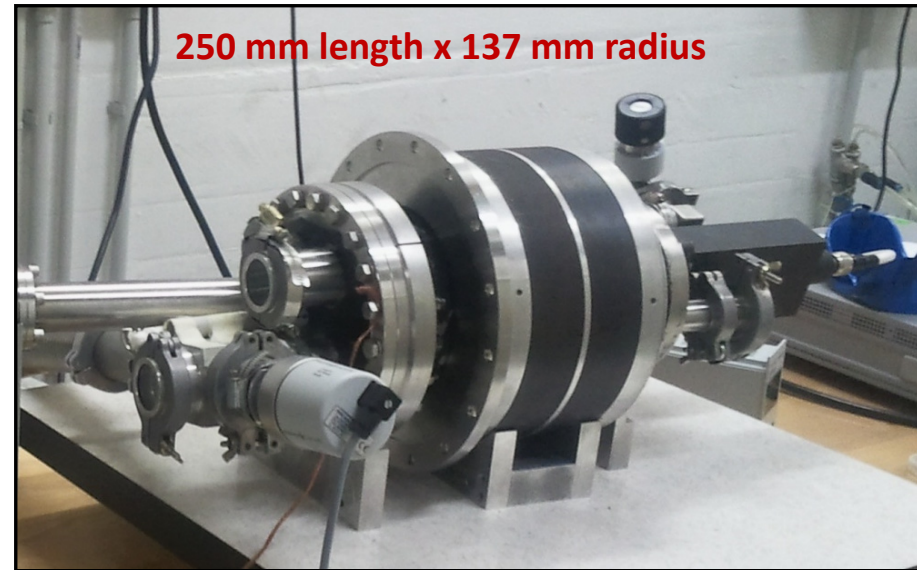
- Limited ECRIS access probing port
- Multi-paths introduce spurious signals



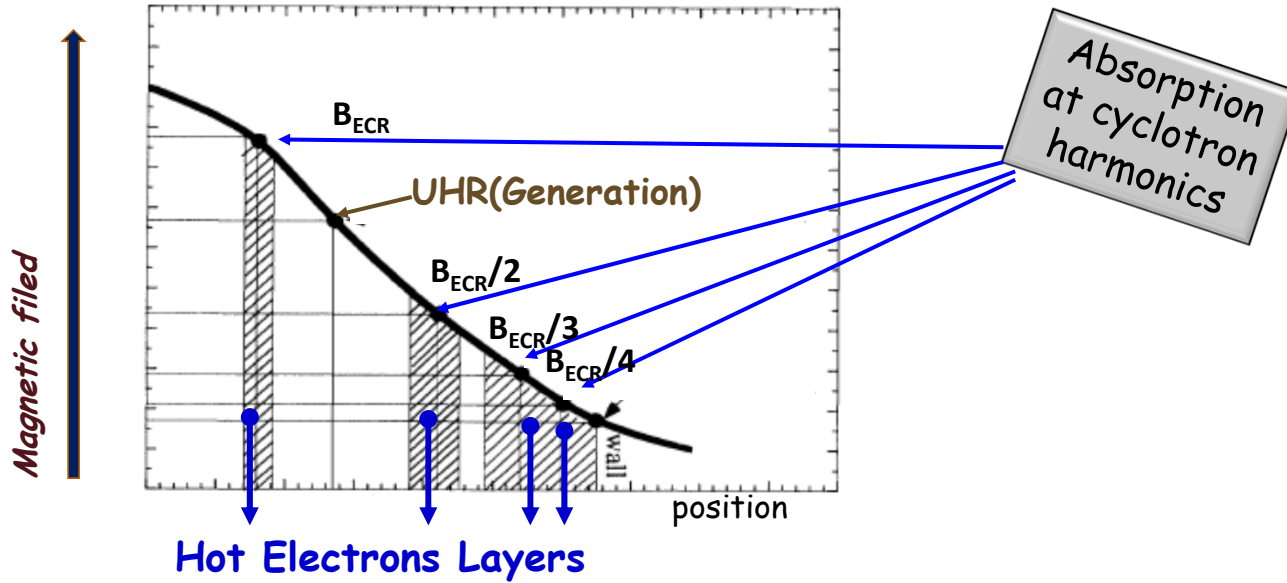
Set-up for EBW plasma generation

For the first time EBWs have
been studied in a compact
plasma based ion source:
The Plasma reactor

Limitations:
EM field is cavity mode dominated;
EM field and plasma must self-organize;
No direct control from external;

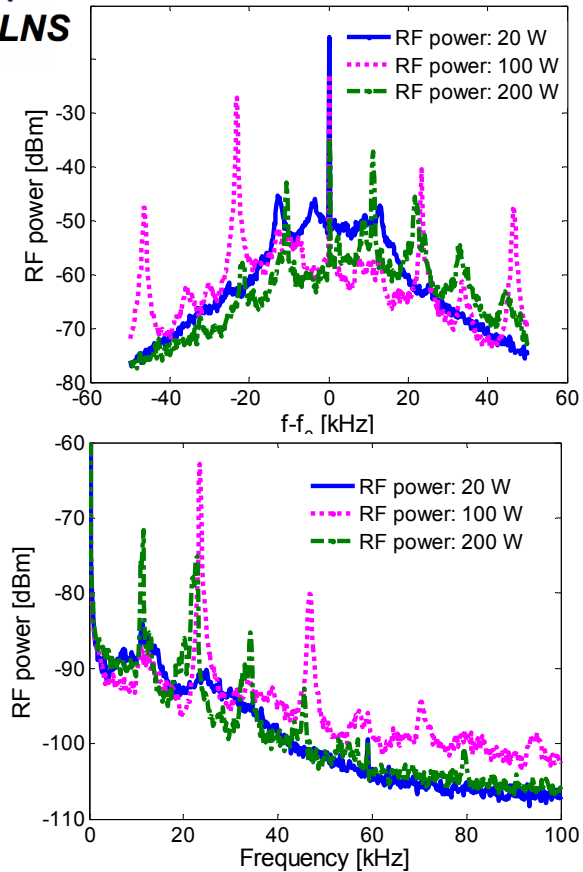


Absorption of Bernstein waves

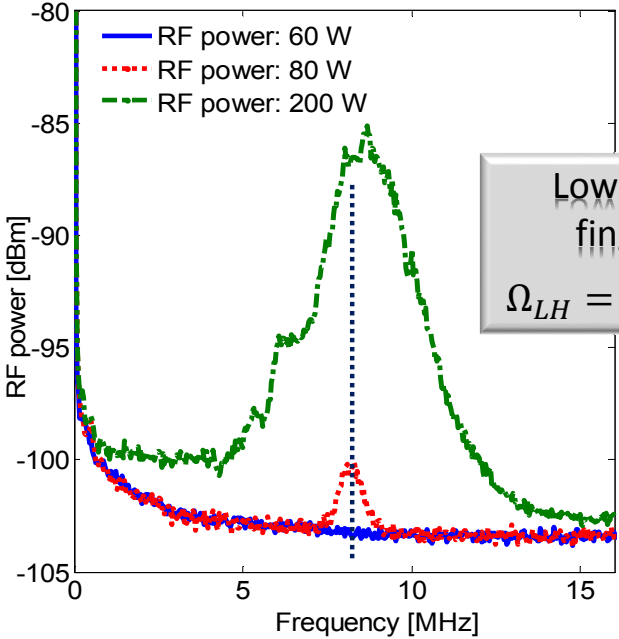


- | | | | |
|---|-----------------------|---|--|
| 1 | No cutoff density | → | Overdense plasma |
| 2 | Non-linear effects | → | <ul style="list-style-type: none"> • Sidebands around pumping wave • Ion waves (MHz range) |
| 3 | High energy electrons | → | X ray emission |

EM spectrum from Spectrum analyzer



Generation of sidebands and ion noise in E.M spectrum over the threshold.



Lower hybrid oscillations:
fingerprint of ES decay
$$\Omega_{LH} = [(\Omega_i \Omega_e)^{-1} + \Omega_p^{-2}]^{-1/2}$$

Models for LP diagnostics

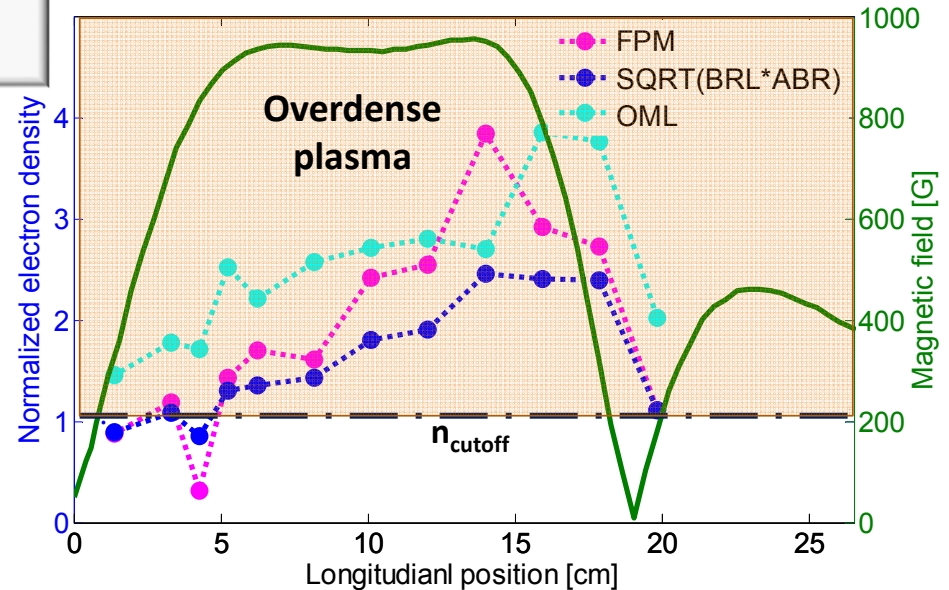
different models validated in **high density** (10^{10} - 10^{12} cm^{-3}) **magnetized plasmas (RF)**

FPM: Floating Potential Method (F. Chen - 2002)

BRL: Bernstein, Reynolds and Lafambroises model

ABR: Allen, Boyd and Reynolds model;

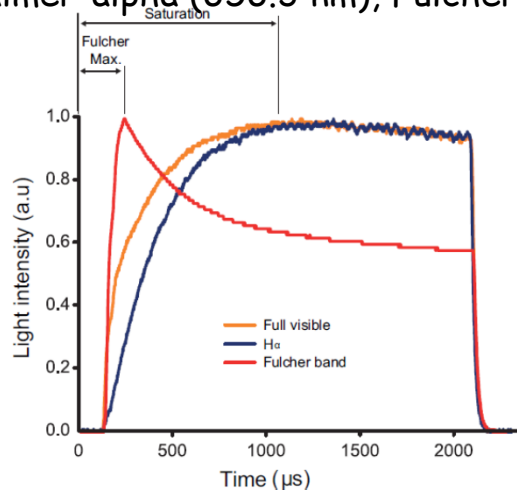
OML: Orbited motion limited



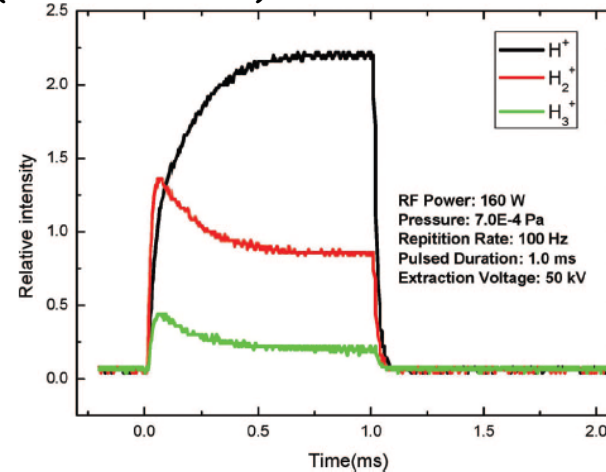
State of art

Correspondence between balmer alpha and fulcher band trend and H and H₂⁺ trend was found by literature

Balmer-alpha (656.3 nm), Fulcher band (around 600 nm)



Cortazar et al Nucl. Instrum. and Methods A 781 (2015) 50–56

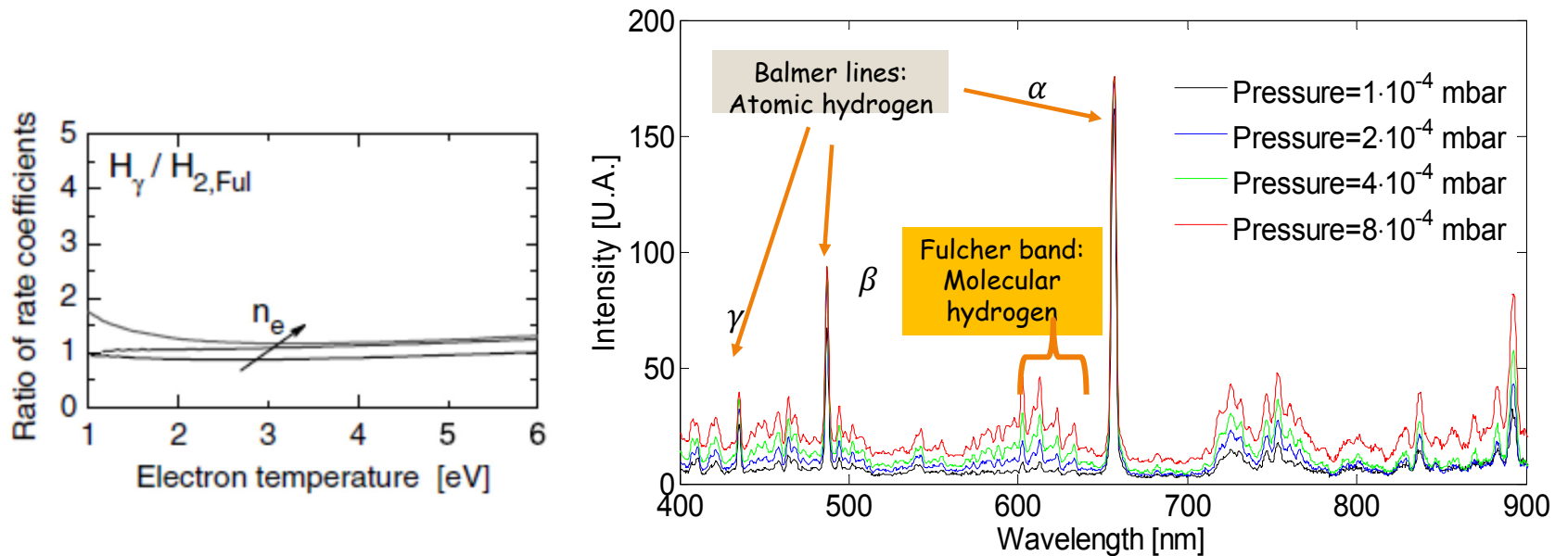


Y. Xu et al. Rev. of Scie. Instrum. **85**, 02A943 (2014);

Room for improvements

- 39° ECPM 23-26 September 2015 – Château de Limelette

Spectroscopic studies suggest that the ratio between H_γ and $H_{2\text{ ful}}$ allows to obtain the ratio between the H and H_2 population, in turn, connected to the proton fraction.



Courtesy of U. Fantz

Preliminary results from the FPT testbench

