# COMPARISON OF THE METHODS FOR BEAM ENERGY SPREAD MEASUREMENT AT THE VEPP-4M

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#### Abstract

The VEPP-4M electron-positron collider is now operating with the KEDR detector for the experiment of precise measurement of tau-lepton mass. In this experiment, monitoring of beam energy spread is important to know the energy spread contribution into the total systematic error. Information about the energy spread gives an opportunity to reduce the error of the taulepton mass measuring. Several techniques for measuring the energy spread are described in the paper. Width of the  $\psi'$  resonance measured with the KEDR detector is used as a reference.

# **INTRODUCTION**

The basic physical program of the VEPP-4M collider consists in precise measurements of mass of  $J/\psi$ ,  $\psi'$ ,  $\psi''$ narrow resonances and  $c \cdot \tau$  lepton mass on the threshold production. Value of beam energy spread  $\sigma_E$  is directly included into accuracy of the mass measurement. Knowledge of exact value of the beam energy spread enables us to reduce significantly a systematical error in the experiment of  $c \cdot \tau$  lepton mass measurement. It is also an essential supplement to the precise measurement of average beam energy.

Clear understanding of the reasons influencing the beam spread and the ability to control this value are important tasks for our experiments.

Energy spread of the beam can be increased with the 3pole snakes and depends on strength of Robinson wigglers. Snakes have a length about 1 m and 1.8 T maximal strength of magnetic field. They were applied only during the beam energy spread measurements described below.

# **MODES OF THE VEPP-4M OPERATION**

Table 1: Operation modes of the VEPP-4M used for energy spread measurements

Name	E, MeV	I <sub>WG</sub> , A	I <sub>SN</sub> , A	Comments	
PNT4	1843	1055	0	$c - \tau$ lepton production. KEDR magnetic field is on.	
PSIS	1843	1055	0	KEDR magnetic field is off.	
ZMEJ	1843	1055	2000	KEDR magnetic field is off.	
JPSI	1548	620	0	$J/\psi$ meson peak. KEDR magnetic field is off	

The target of our experiments was not only the definition of the beam energy spread for basic modes of the collider operation, but also the comparison of several procedures for measurement of relative energy spread  $\delta_{\rm E} = \sigma_{\rm E} / E$ . The experiments were carried out at the four modes of the collider operation (Tab. 1). The modes differ both in the energy *E* and in the energy spread  $\sigma_{\rm E}$  value. Application of several methods to determine the energy spread for different modes of the VEPP-4M operation enables us to realize a cross-validation of the measurements and to compare the diagnostics considering convenience and efficiency.

# **METHODS**

# Spectrum of chromatic sideband peak of beam betatron oscillation (I)

Optical system [3] was applied to measure the beam dimensions  $\sigma_{x,y,z}$  and spectrum of vertical betatron oscillations.

Chromaticity of a storage ring causes appearing of synchrotron sideband peaks in a spectrum of beam oscillation. The amplitude of the central betatron frequency and the synchrotron satellites is [4]:

$$R_m(y) = \frac{1}{y^2} \int_0^\infty J_m^2(x) e^{-\frac{x^2}{2y^2}} x dx,$$
$$y = \left(\frac{\omega_{\beta}\alpha}{\omega_s} + \frac{\omega_0 C_y}{\omega_s}\right) \delta_E \quad ,$$

m is the number of harmonic,  $\delta_E$  is the relative energy spread.

Determination of energy spread is based on the measurement of the ratio of synchrotron satellites to the main peak height.

# Current dependence of energy spread (II)

The experiments with method (*I*) were carried out at the small beam current  $I_0 = 10-50$  mkA, when collective effects are negligible. In the course of experiments with mesons mass measurements the beam currents were close to beam-beam effect threshold restriction. This value was from 1.5 to 3.5 MA depending on the beam energy spread. Radial and longitudinal beam dimensions  $\sigma_{xz}$  were taken to determine current dependence of the beam energy spread.

Energy spread of the beam was derived from the measured radial size  $\sigma_x$  and known amplitude functions

 $\beta_x$ =620 cm,  $\eta_x$ = 94 cm at the observation point. It was supposed that the main reason, which caused the growth of size  $\sigma_x$  and energy spread  $\sigma_E$ , was Touschek effect.

The experiments performed at [2] allowed us to obtain a semi-empirical formula to estimate the contribution of Touschek effect into beam energy spread. This formula (3) was used to adjust the energy spread dependence with the beam current:

$$\sigma_{x} = \left[\beta_{x}\varepsilon_{x} + (\eta_{x}\delta_{ET})^{2}\right]^{1/2}$$
(3),  
where  $\delta_{ET} \approx \frac{5.1 \cdot 10^{-4}}{E[Gev]} \left(\frac{I_{0}[mA] \cdot v_{s}}{K}\right)^{\frac{1}{6}}, K = \sqrt{\frac{\varepsilon_{y}}{\varepsilon_{x}}} -$ 

coupling coefficient,  $\varepsilon_v, \varepsilon_x$  – beam emittance,

$$\delta_E = \sqrt{\delta_o^2 + \delta_{ET}^2}$$
,  $\delta_0$  - energy spread at  $I_0 = 0$ .

We must mention that all the data of method (II) were obtained under collider modes very similar to the modes listed in Table 1, but with the KEDR field switched on.

#### Compton Back Scattering (III)

VEPP-4M collider has a system of Compton Back Scattering for permanent measurement of average beam energy and energy spread [5].

The scattered photons with maximal energy  $W_{\text{max}}$  form a narrow edge in the spectrum. The value of  $W_{\text{max}}$  is strictly coupled with the average energy of the beam electrons by a simple equation:

$$W_{\rm max} \approx 4w_0 \cdot \left(\frac{E}{m_e c^2}\right)^2$$
.

The width of the edge is mostly determined by the resolution of the photons detector and the energy spread in the electron beam. Thus, direct measurement of the energy spectrum of scattered photons allows us to measure both the electron beam average energy and the energy spread.

At the VEPP-4M collider we use the infrared laser with  $w_0=0.117$  eV. In this case,  $W_{\text{max}}$  is from 4 MeV to 7 MeV for the beam energy in the range from 1.5 GeV to 2 GeV.

The energy spectrum of backscattered photons is measured by High Purity Germanium (HPGe) coaxial detector Canberra GC2518 with the energy resolution around  $W_{\text{max}}$  about  $\delta w/w = 4 \cdot 10^{-4}$ . The spectrum in Fig. 1 was gathered with the average electron beam current about Ie=1 mA and the electron beam energy  $\varepsilon = 1842$ MeV at the ZMEJ mode of the VEPP-4M. The resulting value of the measured energy spread is  $\delta E = 1.13 \pm 0.04$ MeV.



Figure 1: The edge of the energy spectrum of backscattered photons with fitting function.

# **EXPERIMENT**

# Method (I)

Beam oscillation was excited by a short kick with amplitude  $b \ge \sigma_y$ . Spectrum of betatron motion was derived with FFT. Frequency and amplitude of the peaks were defined more exactly with the approach proposed in [6]. Blackman-Harris window was also applied. All the measurements were made during 1024 beam turns following the kick. An example of measured spectrum is presented in Fig.2.



Three sideband satellites are clearly seen. The same measurements were made for various values of vertical chromaticity  $Cy=5\div20$ . Chromaticity was changed with sextupole magnets and was measured from the dependence of betatron tune on the rf frequency shift.

The best fit of experimental points corresponds to the energy spread  $\delta_E$ =3.2·10<sup>-4</sup> for JPSI mode,  $\delta_E$  = 4.6·10<sup>-4</sup> for PSIS mode and  $\delta_E$  = 6.6·10<sup>-4</sup> for ZMEJ mode.

# Method (II)

Fit of radial size  $\sigma_x$  vs beam current  $I_0$  with (3) is shown in Fig.3. The fit declines from the experimental points at the value of  $I_0 = 4$  MA. Further dependence  $\sigma_y(I_0)$ has a threshold behavior and requires additional studying. Declination might be caused by microwave instability with the threshold depending on accelerating voltage  $V_{\rm rf}$ . Method (*I*) was applied with reduced value of  $V_{\rm rf}$  = 150 ÷ 250 kV to decrease a synchrotron frequency  $v_{\rm s}$ , which improves a resolution of the measurements described above. The collider runs in 2004-2006 were performed at  $V_{\rm rf} \ge 400$  kV and instability threshold was significantly higher than the currents of the operated beams restricted by the beam-beam effects.



Figure 3: Dependence of beam radial size  $\sigma_x$  on beam current  $I_0$  at PNT4 mode; O - CCD matrix measurement; — - derived with energy spread variation (3).

Measurement of the longitudinal beam size  $\sigma z$  enables us to derive the energy spread at  $I_0 = 0$ . Further beam lengthening  $\sigma z \alpha I_0^{1/3}$  is caused by the ring longitudinal impedance about 6 Om of an inductive type [2].

#### DISCUSSION

One can note that measurements by method (II) are in a good agreement with data of  $\psi'$  resonance scan at  $E_0=1843\pm10$  MeV performed in 2006 (Tab. 2). This method has the same good agreement for  $J/\psi$  meson scan (Tab.2) with reduced wiggler current  $I_{WG}=620$ A.

	Е,	dW,	σ <sub>E</sub> ·,	$I_{\rm WG}$	I <sub>0</sub>	Year
	MeV	MeV	10-4	[A]	[mA]	
ψ	1843	1.24	4.77	1135	2.0	2004
width		1.15	4.42	1055	2.5	2005
		1.09	4.19	1055	2.5	2006
$J/\psi$	1548	0.858	3.93	952	1.7	2002
width		0.664	3.04	652	1	2002

Table 2: Data of  $\psi'$  and  $J/\psi$  resonance scanning

Methods (*I*) and (*II*) demonstrate a good agreement for PSIS, ZMEJ and JPSI mode of the collider as well.

Fig. 4 represents the collected data of energy spread measurements of the beam for collider modes listed in Table 1. Unfortunately, we did not have enough time for machine study and did not perform the measurements with method (*II*) with the KEDR magnetic field switched off. Nevertheless, the data of all the three methods are in a satisfactory agreement.



Figure 4: Data comparison of all the methods. Lines – data of (*II*);  $\blacksquare$ - data of (*I*) for ZMEJ mode;  $\blacktriangle$ - data of (*I*) for PSIS mode;  $\blacklozenge$ - data of (*I*) for JPSI mode.

One can see the distinctions of  $\psi$ ' resonance width obtained with the equal wiggler current (Tab.2). It indicates the necessity of radial orbit stability and the control of dispersion function in the wigglers because the value of energy spread depends on these parameters as well as on the current of the wigglers.

We suppose that variation of beam orbit could be the reason of distinction of energy spread for PSIS and PNT4 modes. It should be mentioned that reduction of  $\delta_E$  by CBS data took place only during the run in 2006 and requires additional investigation. The data of (*II*) and CBS are in a good agreement with the runs in 2004/2005.

In conclusion we should admit that beam energy spread measurement is a nontrivial task and requires the comparison of data of several diagnostics.

#### **SUMMARY**

The measurements of the beam energy spread of the VEPP-4M collider have been done. The collider modes correspond to energy area of 1.5 - 1.8 GeV. Several experimental methods have been applied for each mode. The methods data are in a satisfactory agreement. The following experiments of precise measurement of taulepton mass should be performed with an accurate control of beam radial position as well as with the control of the dispersion function in the wigglers.

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