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U.S. DEPARTMENT OF
ENERGY

Office of Science

Advanced Modeling of Beams & Accelerators

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NA-PAC 2013 – Pasadena, CA, U.S.A.

Importance of computer modeling is on the rise

computers capacity

with increasing:

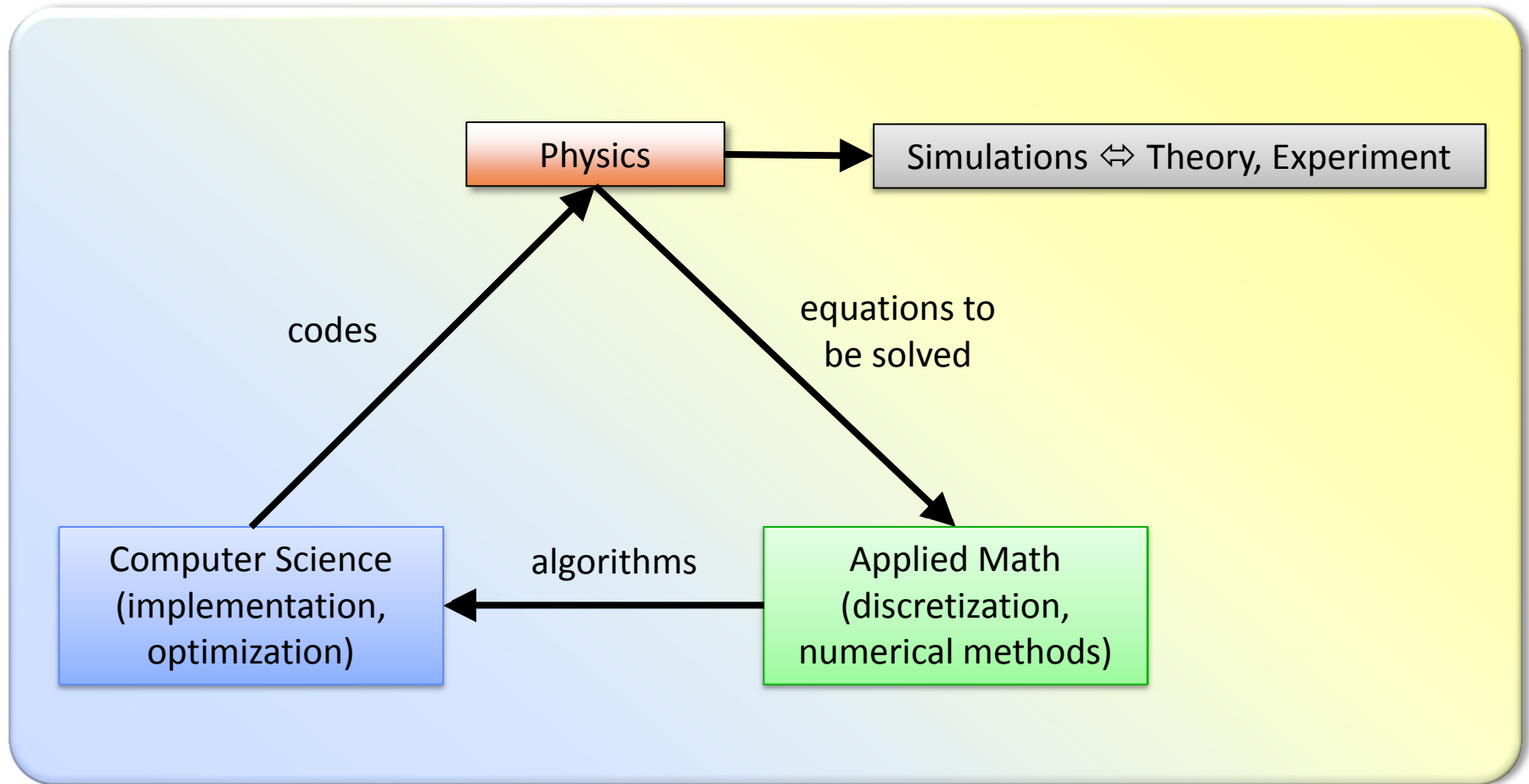
accuracy of codes with
better algorithms and
more physics



pressure for **reducing uncertainties** and
cost on development, construction &
commissioning of accelerator



Modeling of beams & accelerators



In practice, modeling activities can be significantly more complex.

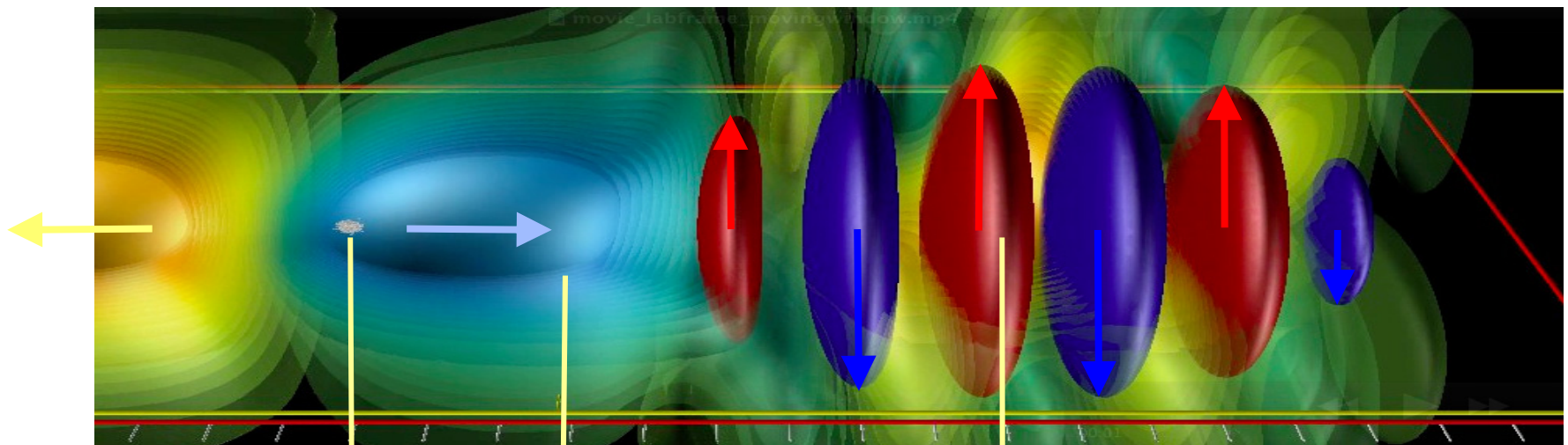
Laser plasma acceleration (LPA)



surfer

wake

boat



e- beam

wake

laser

Modeling from first principle is very challenging



For a 10 GeV scale stage:

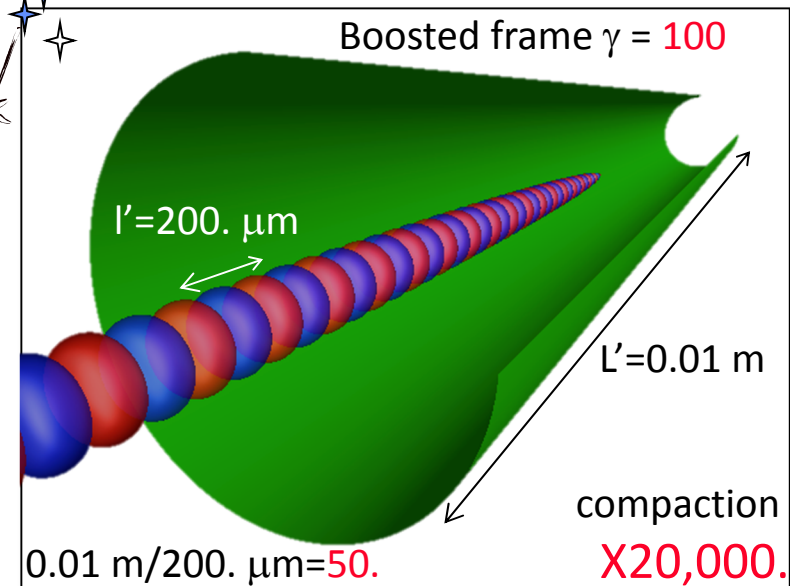
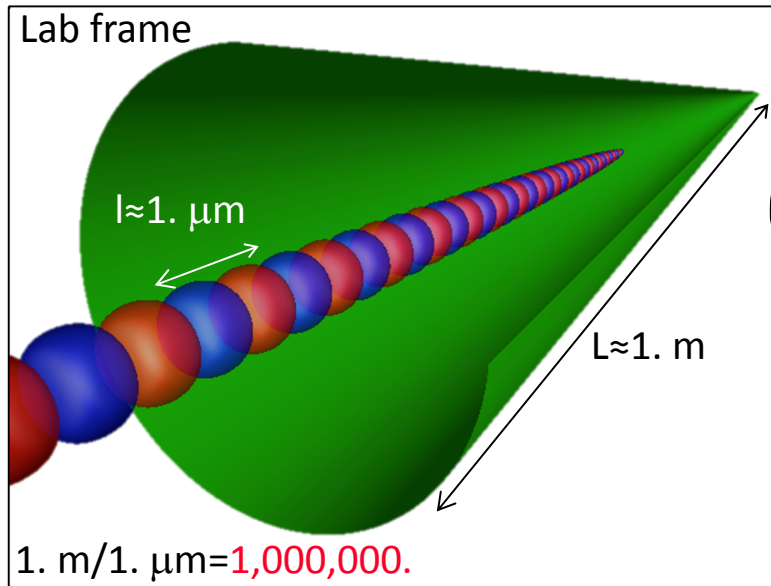
~1 μ m wavelength laser propagates into ~1m plasma

➔ millions of time steps needed

Physics insight

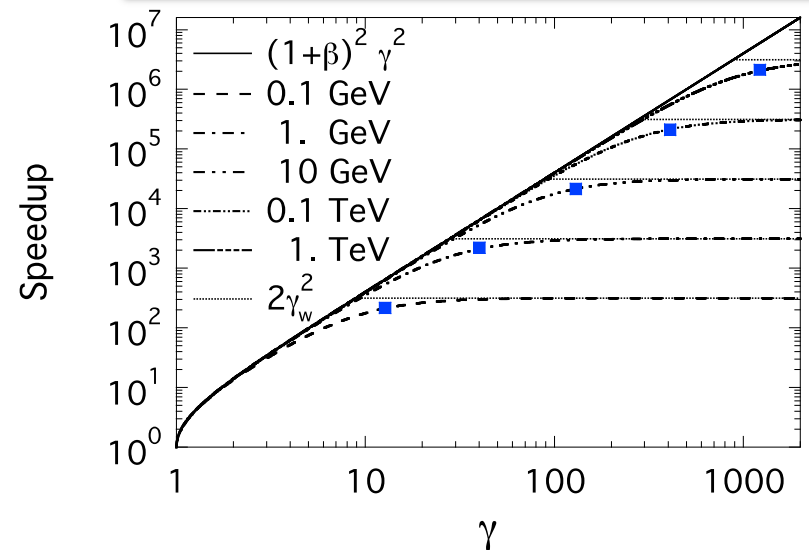
Lorentz boosted frame → lower # time steps

Lorentz boosted frame reduces scale range by orders of magnitude*



Predicted speedup*:

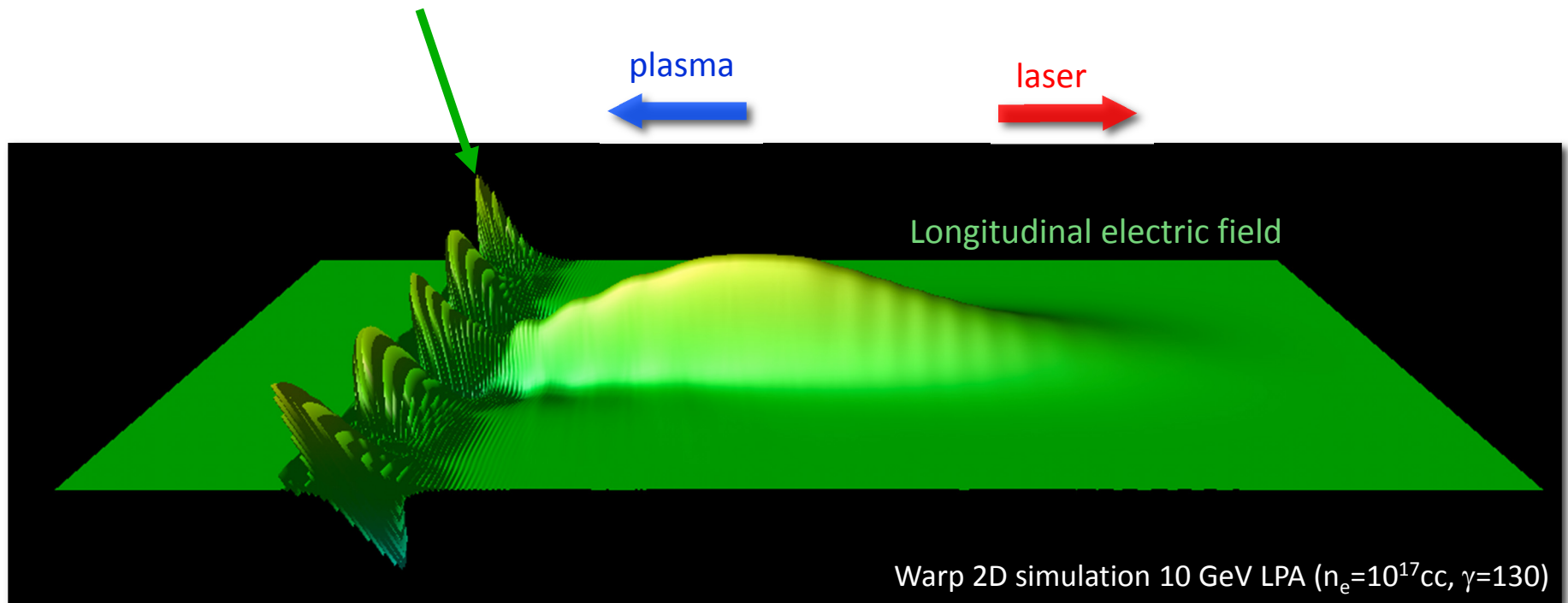
- > 10,000 for single 10 GeV (Bella) stage,
- > 1,000,000 for single 1 TeV stage.



*J.-L. Vay, *Phys. Rev. Lett.* **98**, 130405 (2007)

Complication

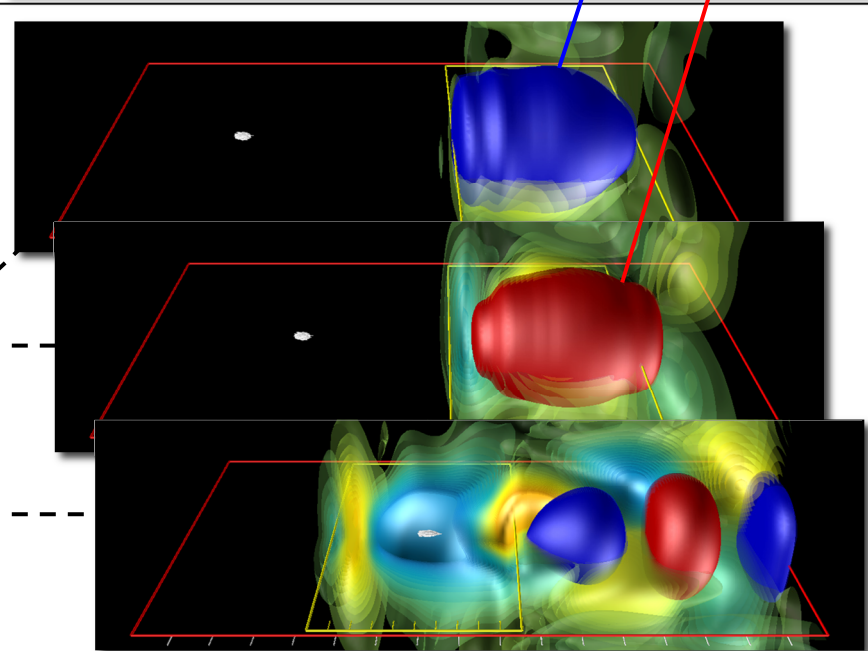
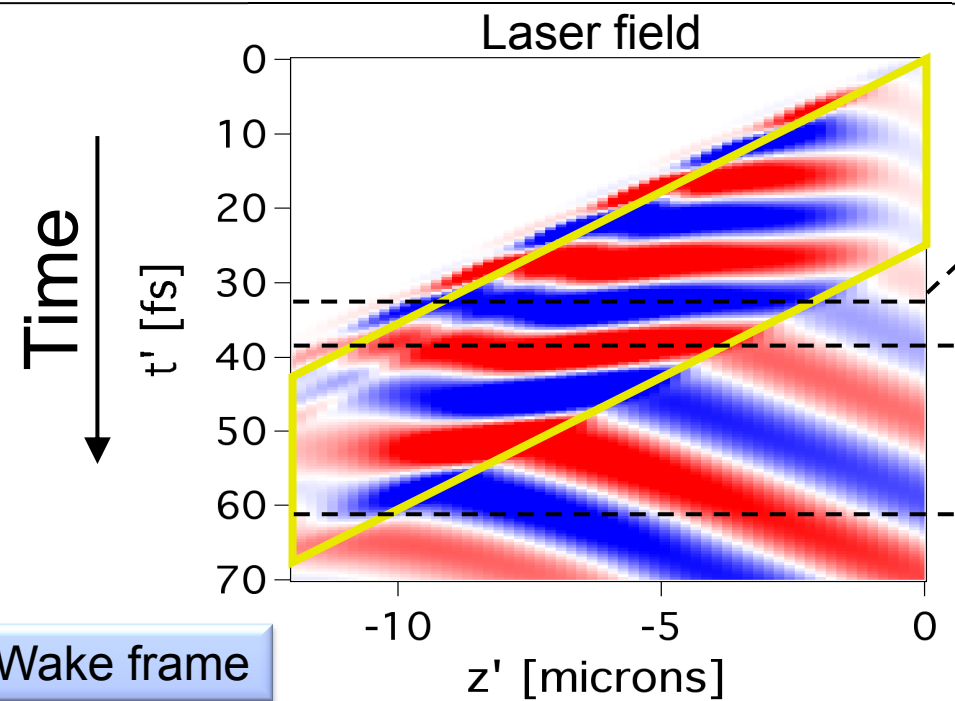
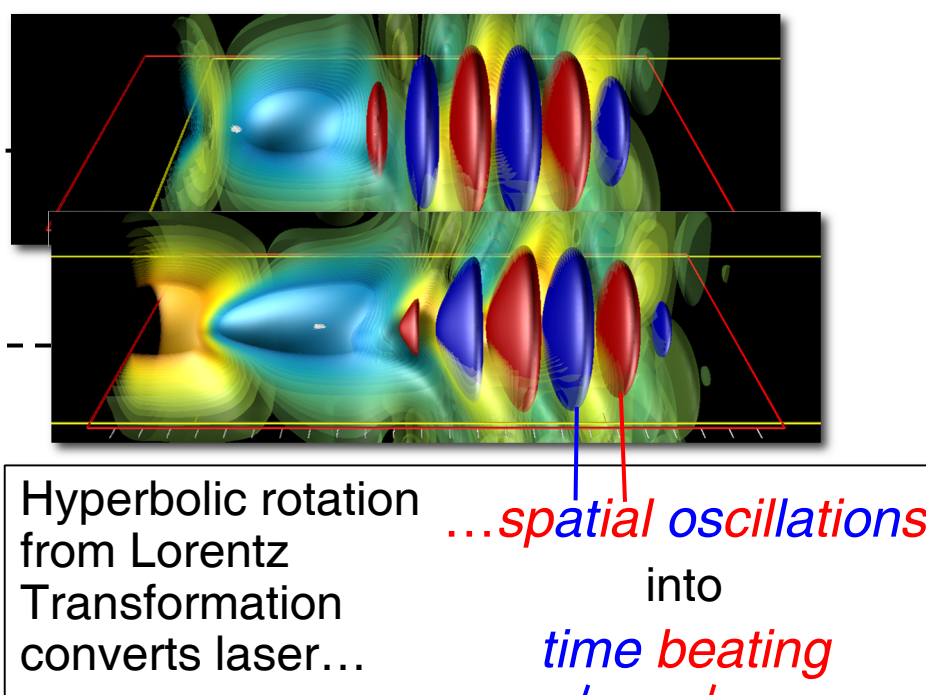
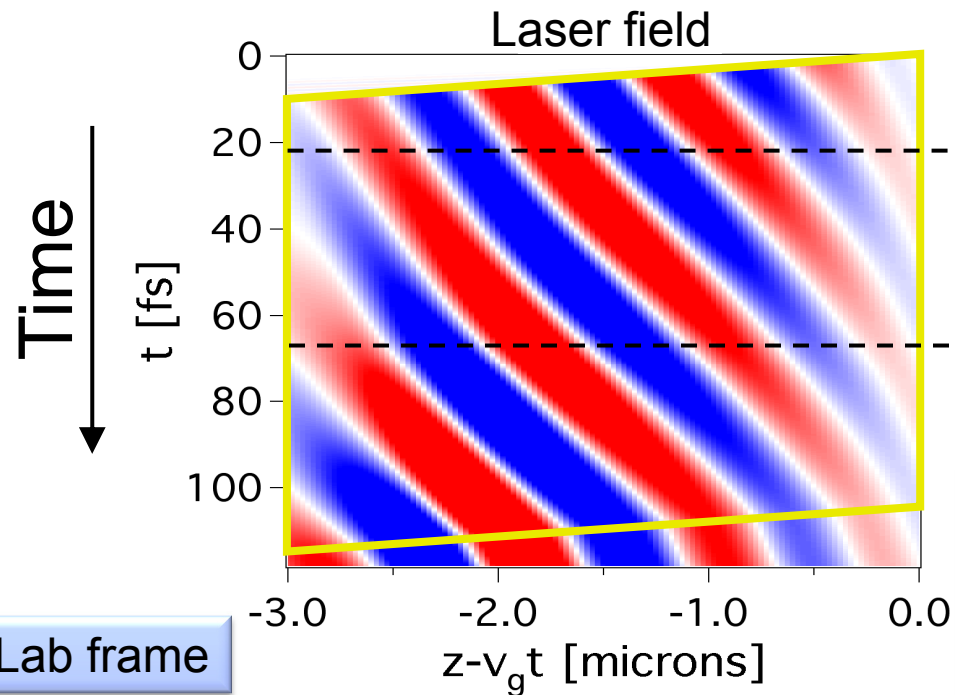
Short wavelength instability observed at entrance of plasma for large γ (≈ 100)

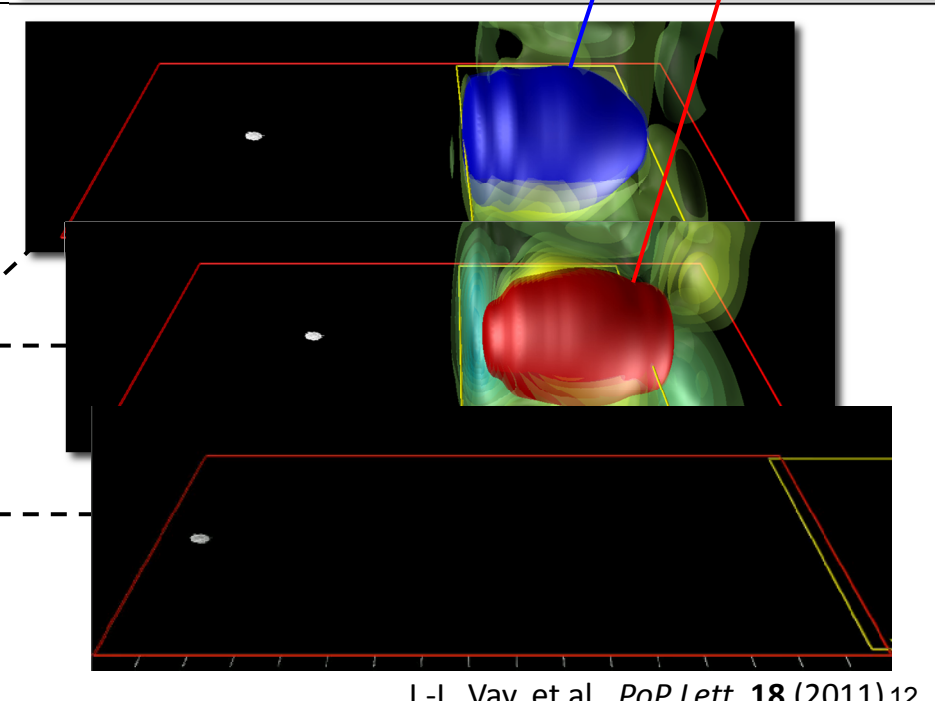
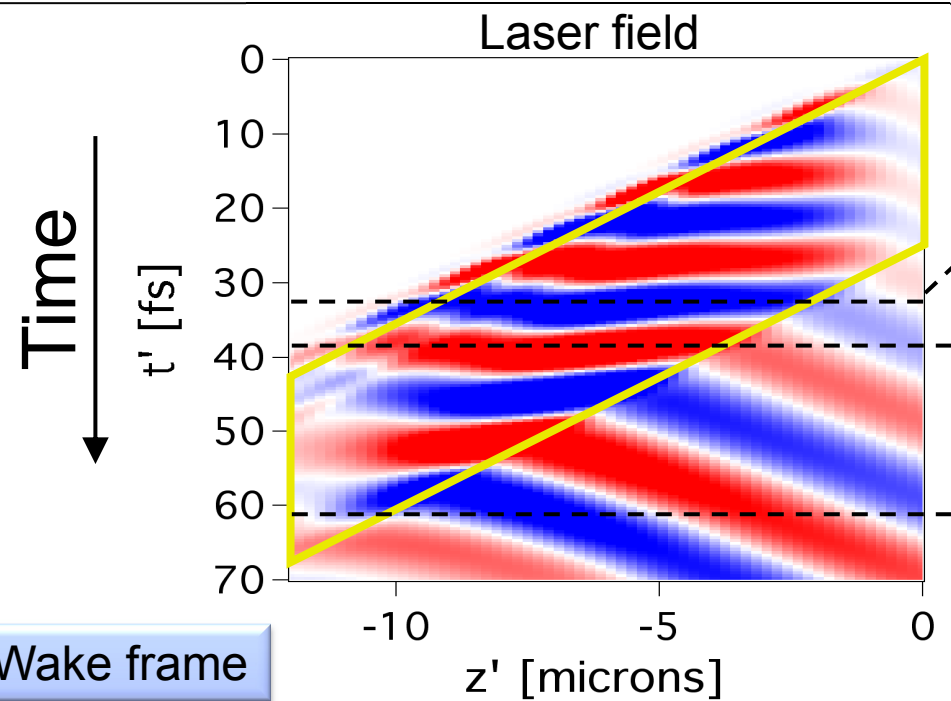
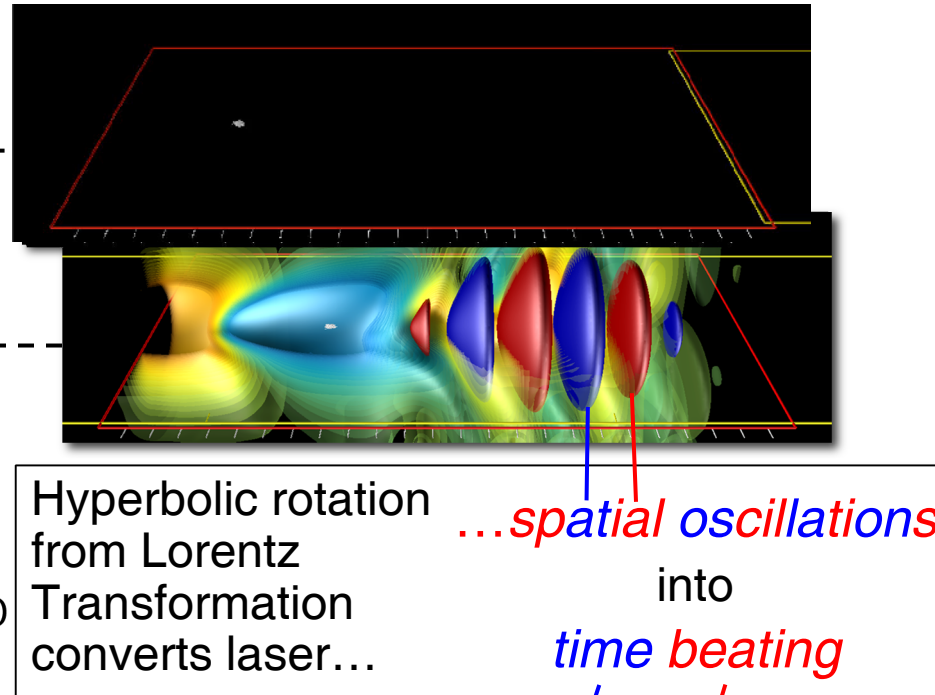
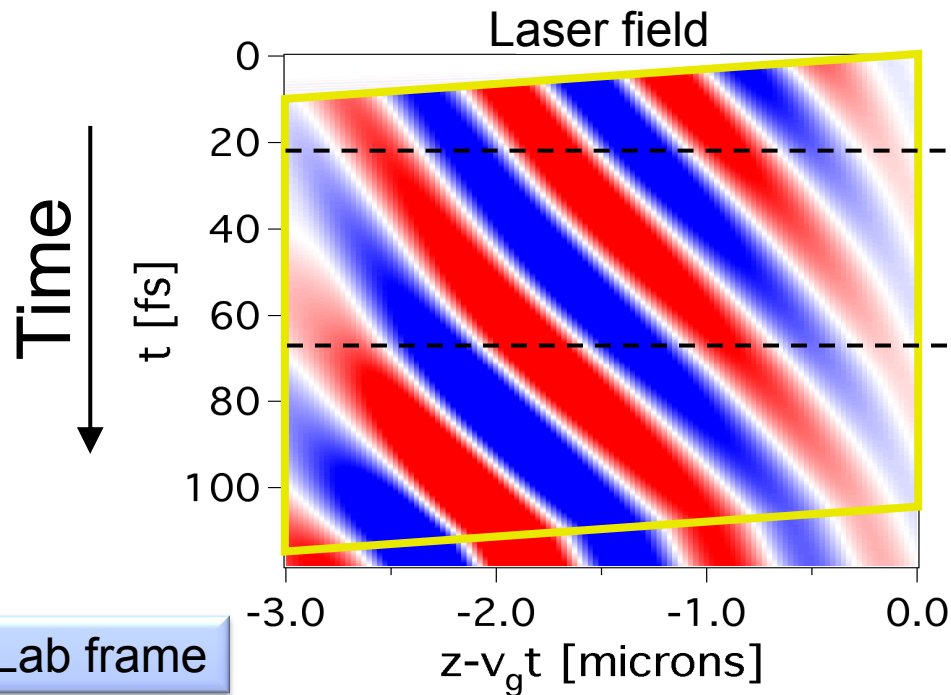


Instability at short wavelength: can it be filtered?

Physics insight

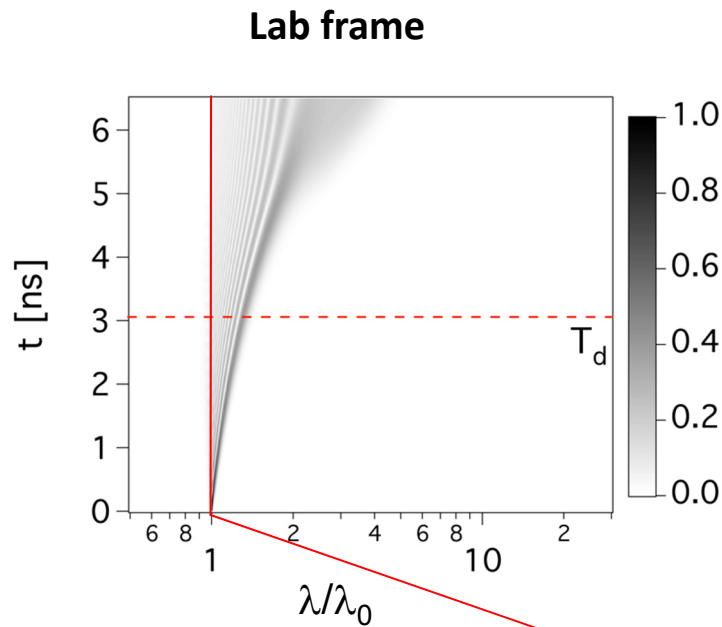
Hyperbolic rotation → more filtering





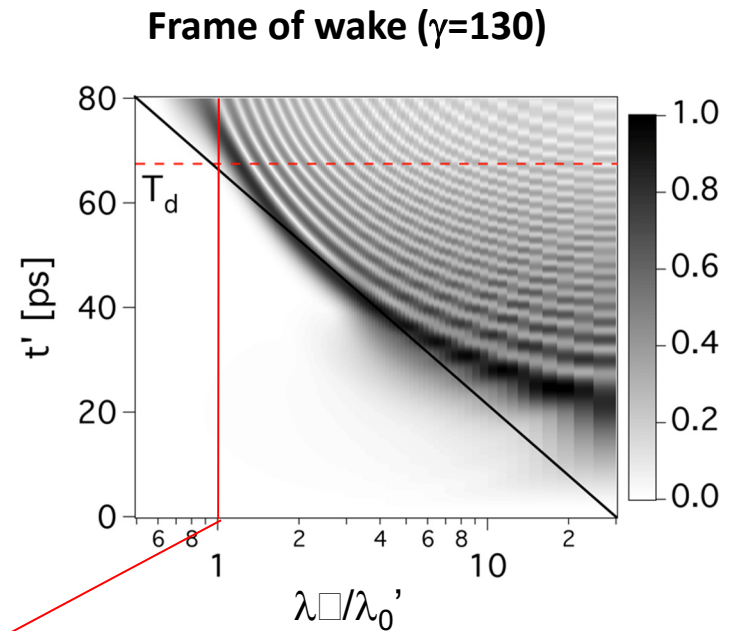
More filtering is possible without altering physics

Laser spectrum on axis



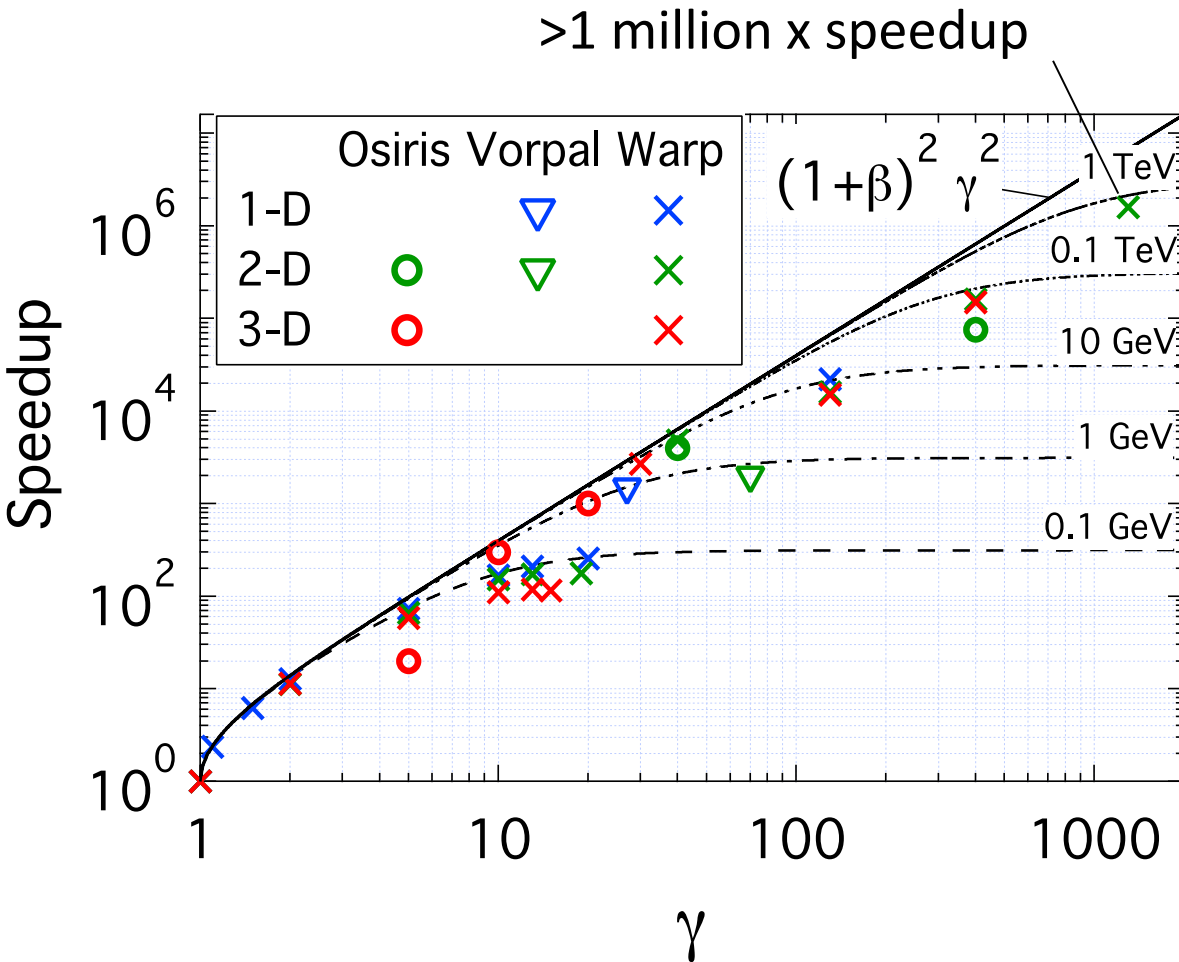
Laser in vacuum

Content concentrated around λ_0



Content concentrated at much larger λ

Speedup verified by us and others to over a million



Warp:

1. J.-L. Vay, et al., *Phys. Plasmas* **18** 123103 (2011)
2. J.-L. Vay, et al., *Phys. Plasmas (letter)* **18** 030701 (2011)
3. J.-L. Vay, et al., *J. Comput. Phys.* **230** 5908 (2011)
4. J.-L. Vay et al, PAC Proc. (2009)

Osiris:

1. S. Martins, et al., *Nat. Phys.* **6** 311 (2010)
2. S. Martins, et al., *Comput. Phys. Comm.* **181** 869 (2010)
3. S. Martins, et al., *Phys. Plasmas* **17** 056705 (2010)
4. S. Martins et al, PAC Proc. (2009)

Vorpil:

1. D. Bruhwiler, et al., *AIP Conf. Proc* **1086** 29 (2009)

Enabled simulations previously untractable: e.g. 10 GeV stage

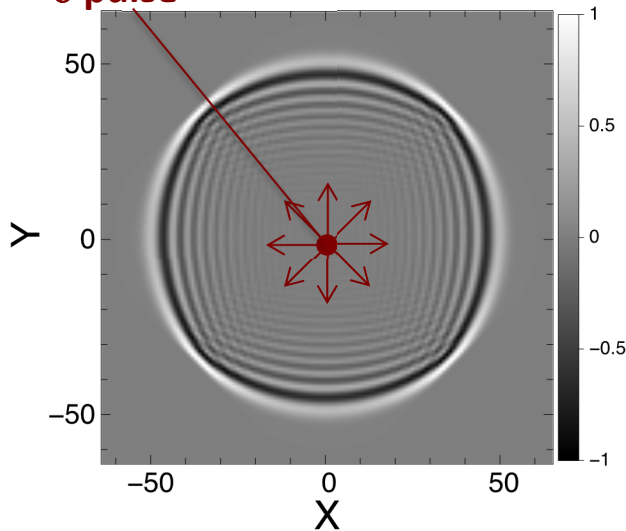
2006 (lab): in 1D ~ 5k CPU-hours → 2011 (boost): in 3D ~ 1k CPU-hours

**Filtering the instability is good,
understanding the instability is better:
→ stability analysis.**

Analysis performed for various field solvers

Kronecker
 δ pulse

FDTD



- Numerical dispersion,
- anisotropy,
- Courant condition:

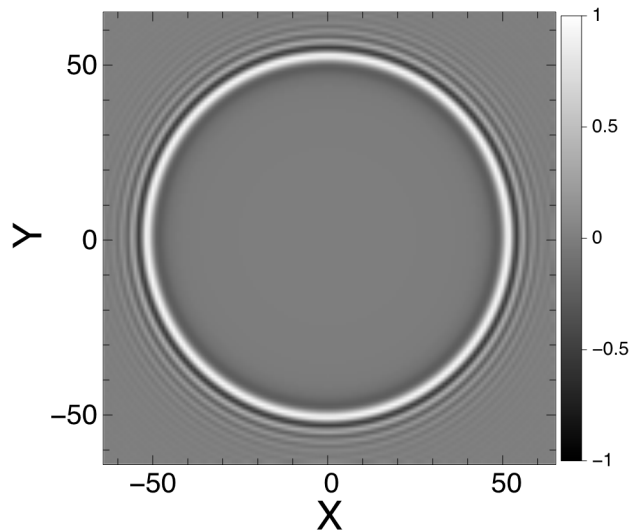
$$c\Delta t \leq 1 / \sqrt{\frac{1}{\Delta x^2} + \frac{1}{\Delta y^2} + \frac{1}{\Delta z^2}}$$

FDTD = Finite-Difference Time Domain: Warp¹ (LBNL/LLNL/U. Maryland), Osiris² (UCLA)

PSTD = Pseudo-Spectral Time Domain: UPIC-EMMA² (UCLA)

PSATD = Pseudo-Spectral Analytical Time Domain: Warp³ (LBNL/LLNL/U. Maryland)

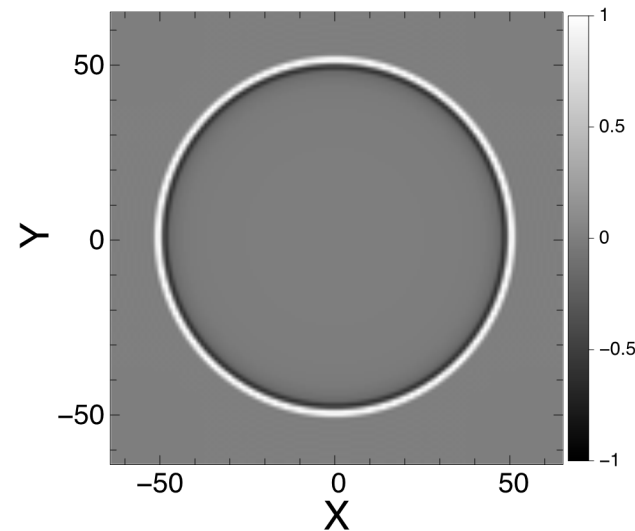
PSTD



- Numerical dispersion,
- isotropy,
- Courant condition:

$$c\Delta t \leq 2/\pi \sqrt{\frac{1}{\Delta x^2} + \frac{1}{\Delta y^2} + \frac{1}{\Delta z^2}}$$

PSATD



- Exact dispersion,
- isotropy,
- Courant condition:

None

¹B. B. Godfrey, J.-L. Vay, "Numerical stability of relativistic beam multidimensional PIC simulations employing the Esirkepov algorithm" J. Comp. Phys. 248 (2013)

²X. Xu, et al, "Numerical instability due to relativistic plasma drift in EM-PIC simulations", Comp. Comm. Phys. 184 (2013)

³B. B. Godfrey, J. L. Vay, I. Haber, "Numerical stability analysis of the PSATD-PIC algorithm", *submitted*

Numerical dispersion relation of spectral-PIC algorithm*

2-D numerical dispersion relation (Fourier space):

$$\begin{pmatrix} \xi_{z,z} + [\omega] & \xi_{z,x} & \xi_{z,y} + [k_x] \\ \xi_{x,z} & \xi_{x,x} + [\omega] & \xi_{x,y} - [k_z] \\ [k_x] & -[k_z] & [\omega] \end{pmatrix} \begin{pmatrix} E_z \\ E_x \\ B_y \end{pmatrix} = 0.$$

$$[\omega] = \sin\left(\omega \frac{\Delta t}{2}\right) / \left(\frac{\Delta t}{2}\right) \quad [k_z] = k_z \sin\left(k \frac{\Delta t}{2}\right) / \left(k \frac{\Delta t}{2}\right) \quad [k_x] = k_x \sin\left(k \frac{\Delta t}{2}\right) / \left(k \frac{\Delta t}{2}\right)$$

$$S^J = \left[\sin\left(k'_z \frac{\Delta z}{2}\right) / \left(k'_z \frac{\Delta z}{2}\right) \right]^{\ell_z+1} \left[\sin\left(k'_x \frac{\Delta x}{2}\right) / \left(k'_x \frac{\Delta x}{2}\right) \right]^{\ell_x+1},$$

$$S^{E_z} = \left[\sin\left(k'_z \frac{\Delta z}{2}\right) / \left(k'_z \frac{\Delta z}{2}\right) \right]^{\ell_z} \left[\sin\left(k'_x \frac{\Delta x}{2}\right) / \left(k'_x \frac{\Delta x}{2}\right) \right]^{\ell_x+1} (-1)^{m_z},$$

$$S^{E_x} = \left[\sin\left(k'_z \frac{\Delta z}{2}\right) / \left(k'_z \frac{\Delta z}{2}\right) \right]^{\ell_z+1} \left[\sin\left(k'_x \frac{\Delta x}{2}\right) / \left(k'_x \frac{\Delta x}{2}\right) \right]^{\ell_x} (-1)^{m_x},$$

$$S^{B_y} = \cos\left(\omega \frac{\Delta t}{2}\right) \left[\sin\left(k'_z \frac{\Delta z}{2}\right) / \left(k'_z \frac{\Delta z}{2}\right) \right]^{\ell_z} \left[\sin\left(k'_x \frac{\Delta x}{2}\right) / \left(k'_x \frac{\Delta x}{2}\right) \right]^{\ell_x} (-1)^{m_z+m_x}.$$

*B. B. Godfrey, J. L. Vay, I. Haber, "Numerical stability analysis of the Pseudo-Spectral Analytical Time-Domain PIC algorithm", *submitted*

Numerical dispersion relation of spectral PIC algorithm (II)

$$\xi_{z,z} = -n\gamma^{-2} \sum_m S^J S^{E_z} \csc^2 \left[(\omega - k'_z v) \frac{\Delta t}{2} \right] \\ (kk_z^2 \Delta t + \zeta_z k_x^2 \sin(k\Delta t)) \Delta t [\omega] k'_z / 4k^3 k_z,$$

$$\xi_{z,x} = -n \sum_m S^J S^{E_x} \csc \left[(\omega - k'_z v) \frac{\Delta t}{2} \right] \eta_z k'_x / 2k^3 k_z,$$

$$\xi_{z,y} = nv \sum_m S^J S^{B_y} \csc \left[(\omega - k'_z v) \frac{\Delta t}{2} \right] \eta_z k'_x / 2k^3 k_z,$$

$$\xi_{x,z} = -n\gamma^{-2} \sum_m S^J S^{E_z} \csc^2 \left[(\omega - k'_z v) \frac{\Delta t}{2} \right] \\ (k\Delta t - \zeta_z \sin(k\Delta t)) \Delta t [\omega] k_x k'_z / 4k^3,$$

$$\xi_{x,x} = -n \sum_m S^J S^{E_x} \csc \left[(\omega - k'_z v) \frac{\Delta t}{2} \right] \eta_x k'_x / 2k^3 k_x,$$

$$\xi_{x,y} = nv \sum_m S^J S^{B_y} \csc \left[(\omega - k'_z v) \frac{\Delta t}{2} \right] \eta_x k'_x / 2k^3 k_x,$$

Numerical dispersion relation of spectral PIC algorithm (III)

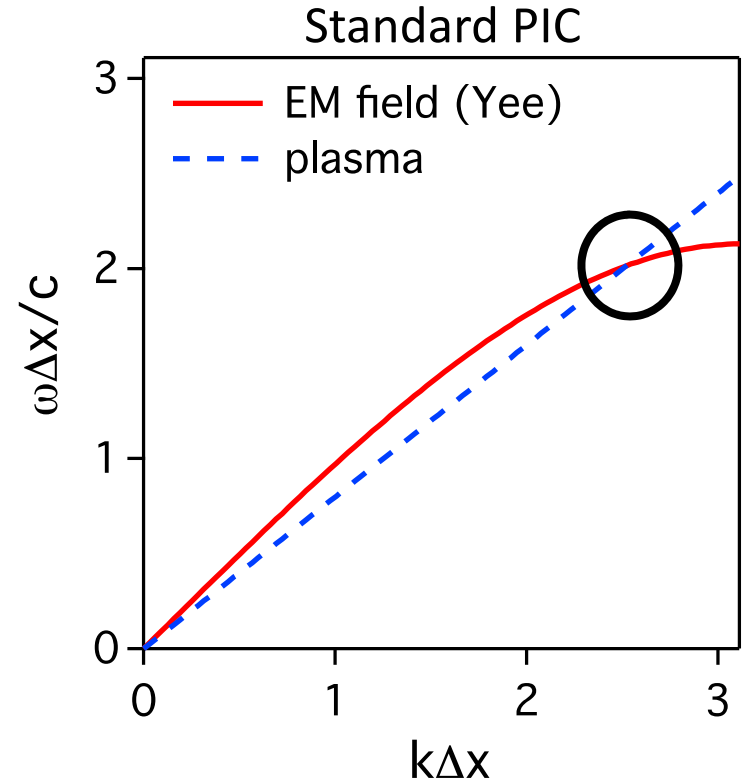
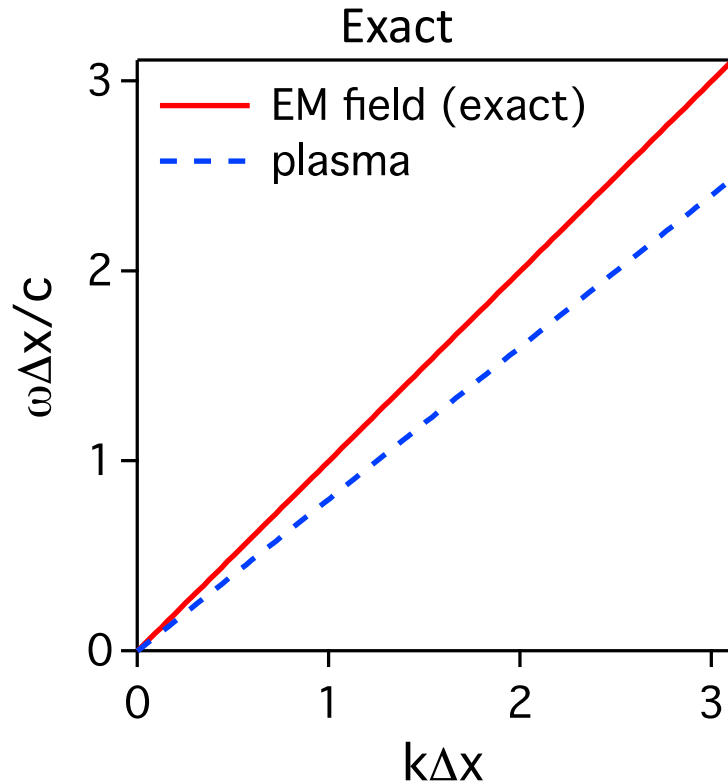
$$\eta_z = \cot \left[(\omega - k'_z v) \frac{\Delta t}{2} \right] (k k_z^2 \Delta t + \zeta_z k_x^2 \sin(k \Delta t)) \sin \left(k'_z v \frac{\Delta t}{2} \right) \\ + (k \Delta t - \zeta_x \sin(k \Delta t)) k_z^2 \cos \left(k'_z v \frac{\Delta t}{2} \right),$$

$$\eta_x = \cot \left[(\omega - k'_z v) \frac{\Delta t}{2} \right] (k \Delta t - \zeta_z \sin(k \Delta t)) k_x^2 \sin \left(k'_z v \frac{\Delta t}{2} \right) \\ + (k k_x^2 \Delta t + \zeta_x k_z^2 \sin(k \Delta t)) \cos \left(k'_z v \frac{\Delta t}{2} \right).$$

Then simplify and solve with Mathematica...

Numerical Cherenkov in 1-D*

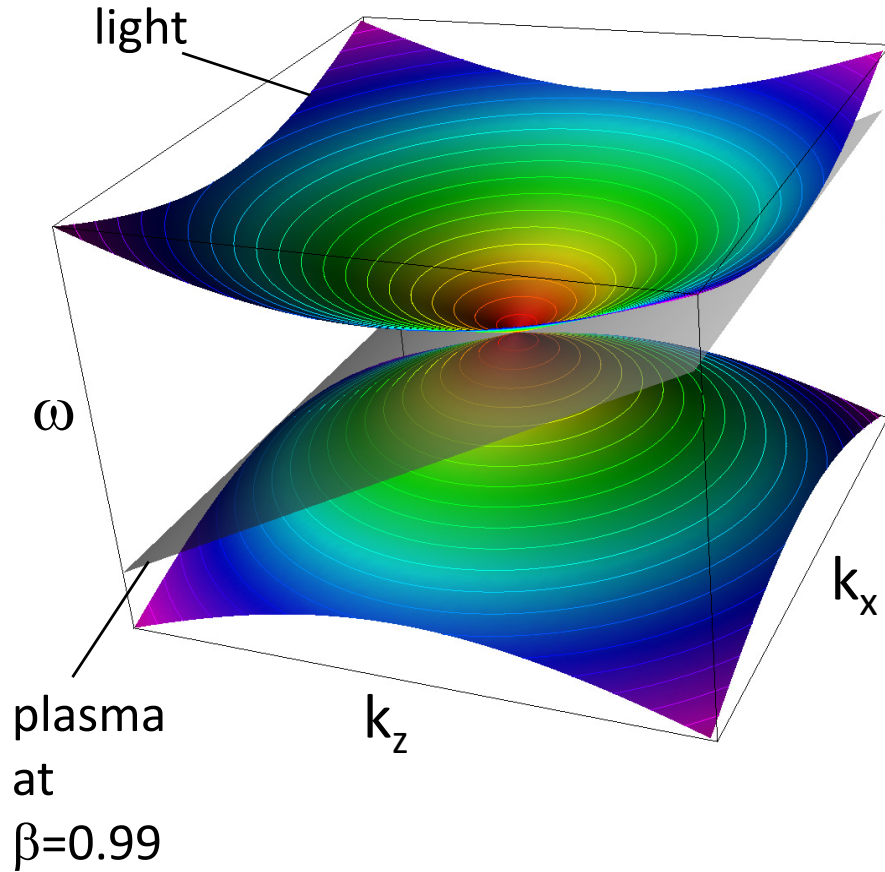
Numerical dispersion leads to crossing of EM field and plasma modes $\rightarrow$ instability.



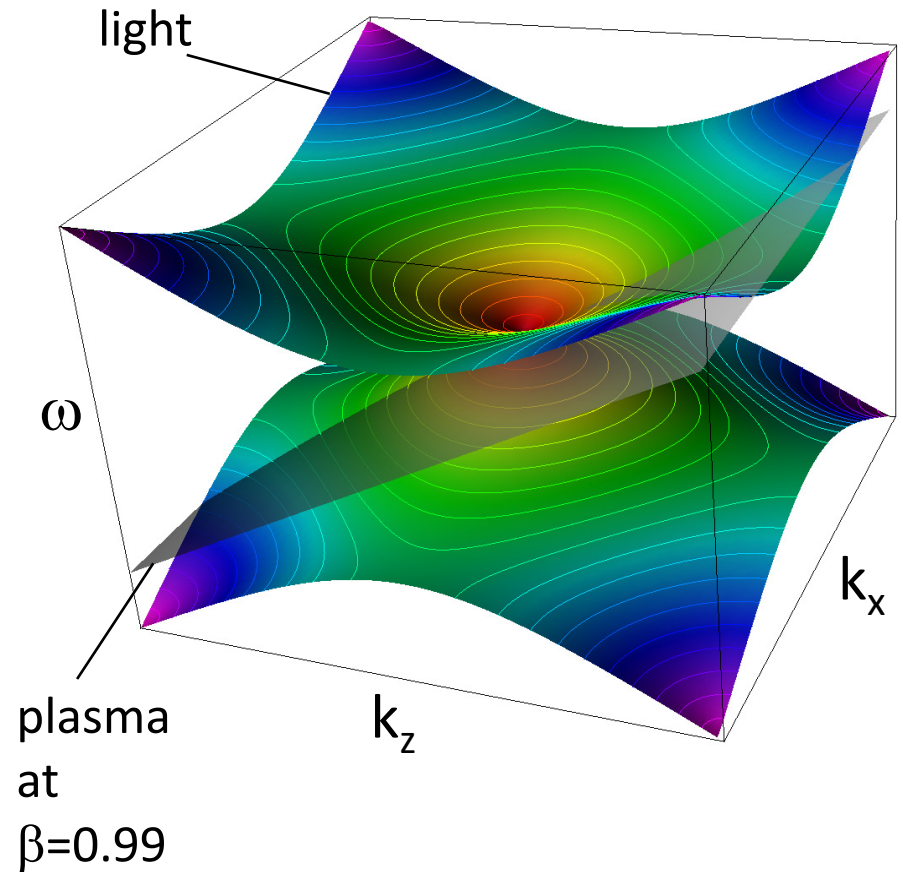
*B. B. Godfrey, "Numerical Cherenkov instabilities in electromagnetic particle codes", *J. Comput. Phys.* **15** (1974)

Numerical Cherenkov in 2-D (for $c\Delta t=\Delta x/\text{sqrt}(2)$)

PSATD

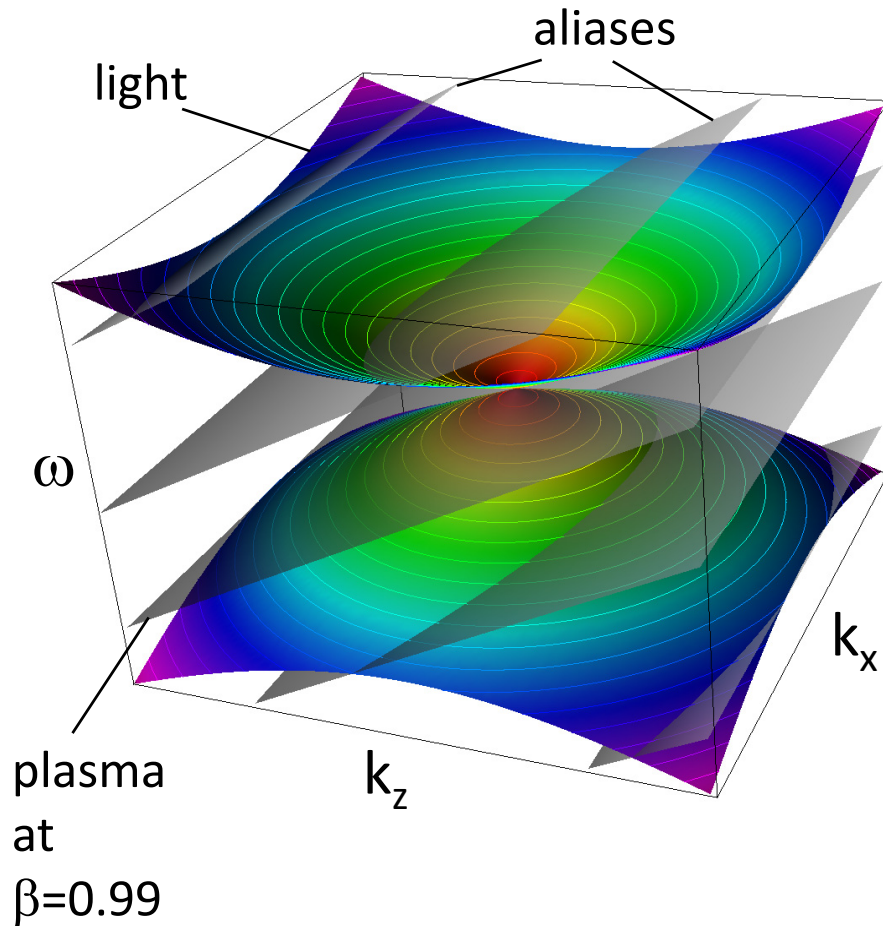


FDTD-Yee

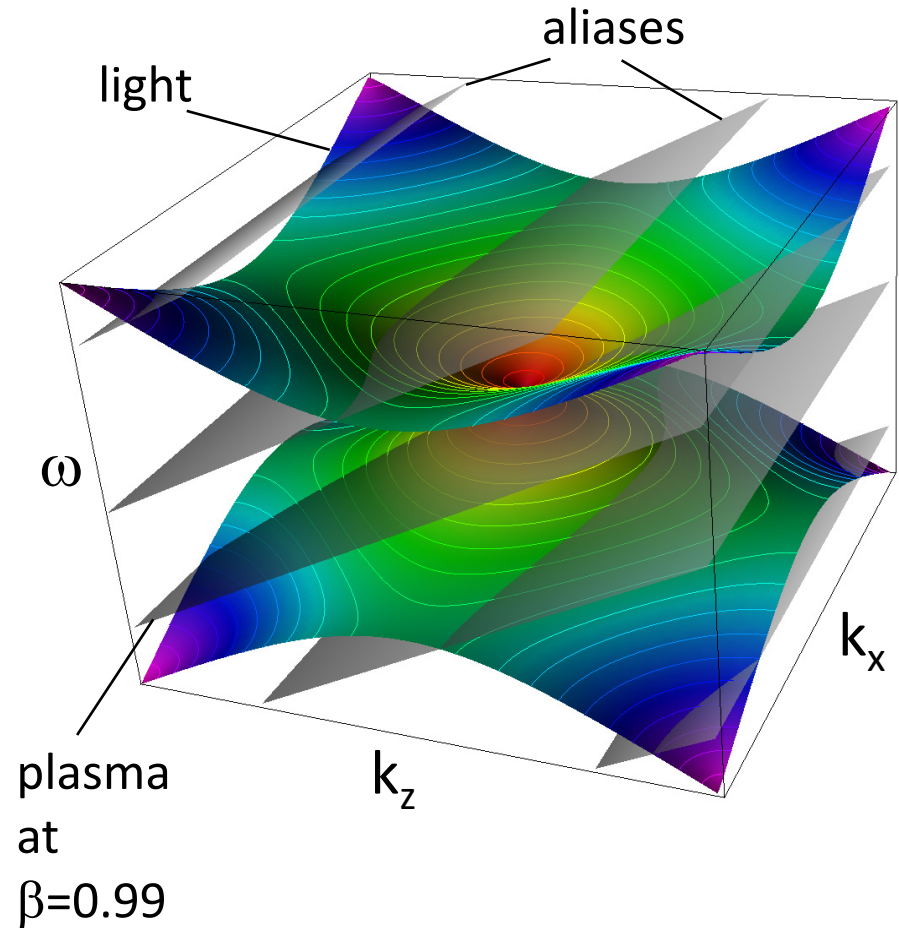


Need to take aliases into account in 2-D

PSATD



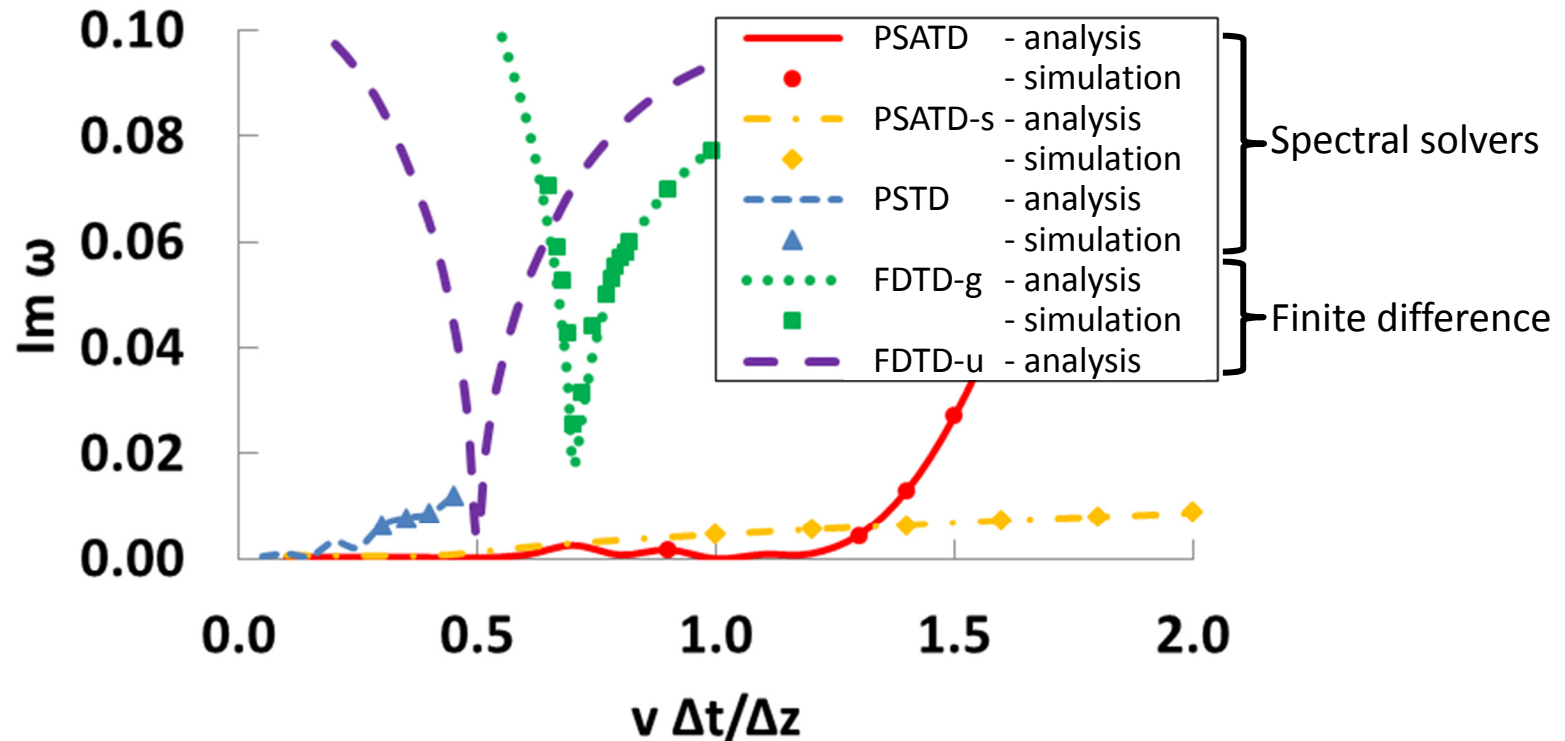
FDTD-Yee



Need to consider at least first aliases $m_x = \{-3 \dots +3\}$ to study stability.

Spectral solvers more stable than finite difference solvers

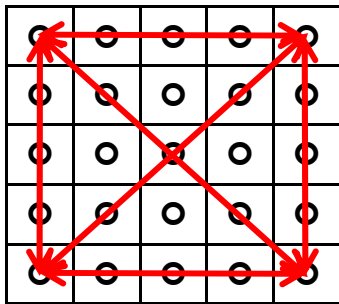
Growth rates of numerical Cerenkov instabilities



But spectral solvers harder to scale to large # of cores

Spectral

global “costly”
communications

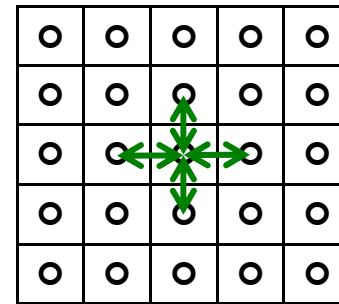


Harder to scale

VS

Finite Difference

local “cheap”
communications



Easier to scale

Physics insight

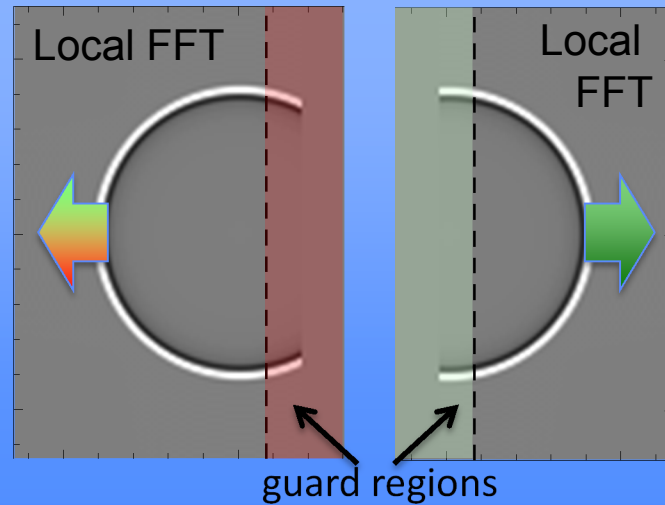
Finite speed of light → local FFTs

New concept on single pulse – part 1

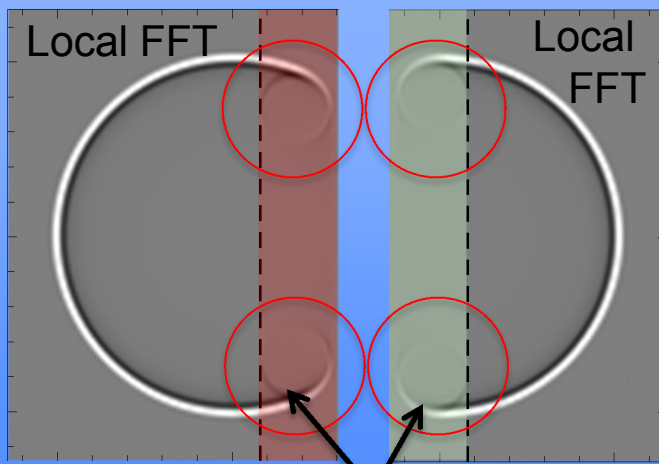
Example: unit pulse expansion at time T



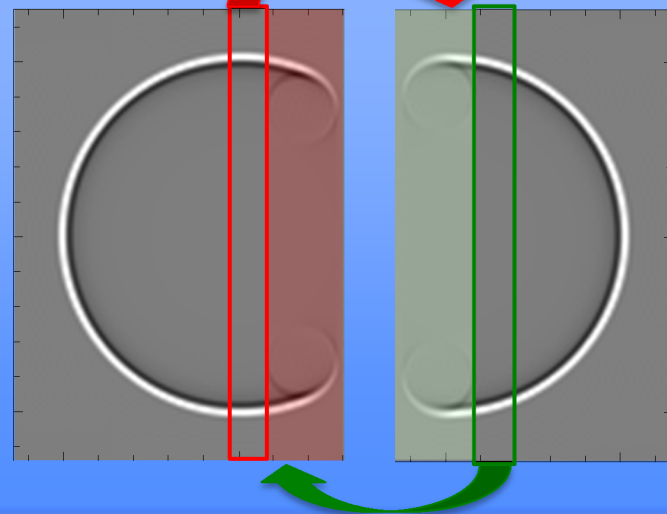
Separate calculation in two domains



Advance to time $T+DT$

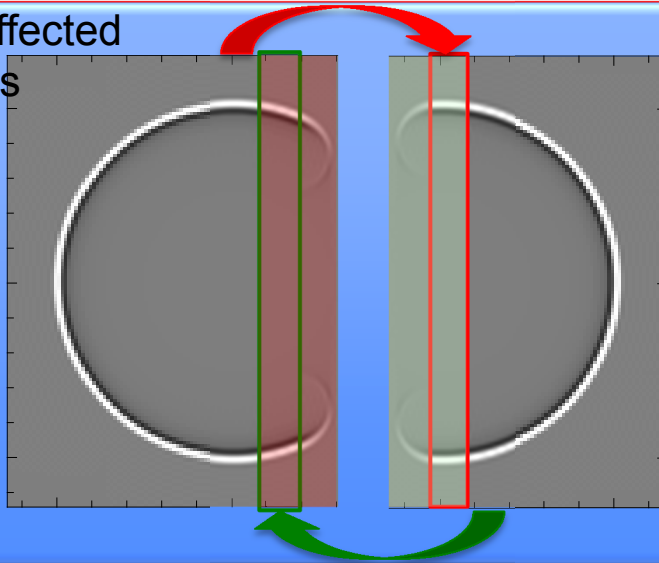


Copy unaffected data

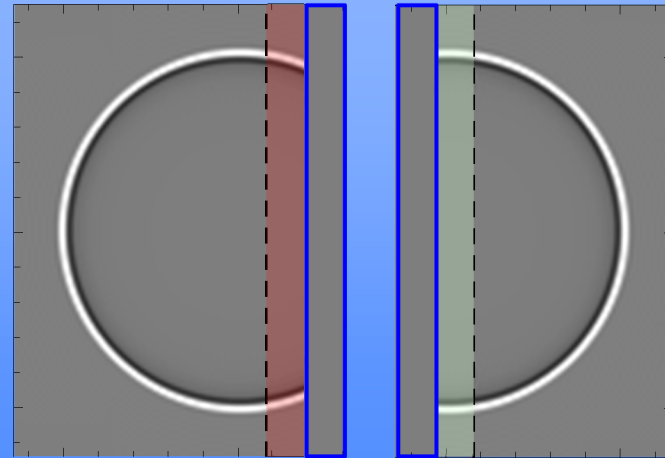


New concept on single pulse – part 2

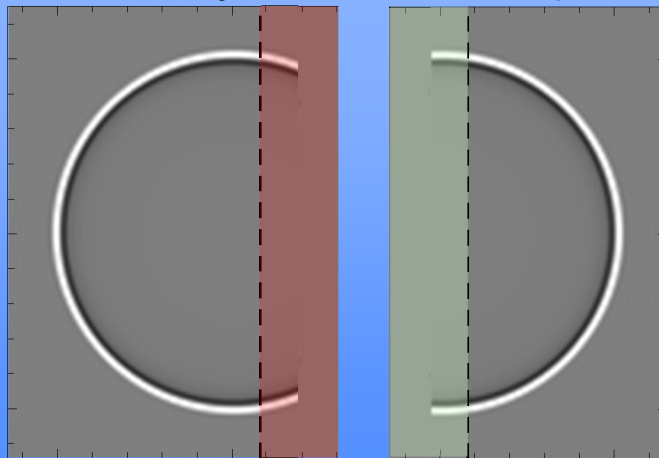
To affected areas



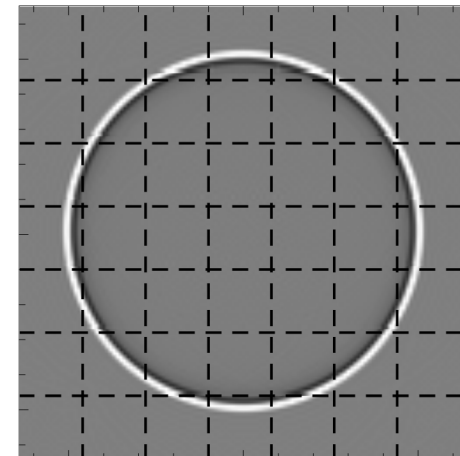
zero out remaining areas



Ready for next time step

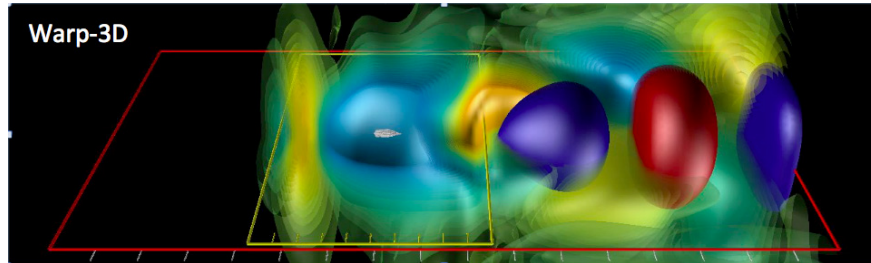


Successfully tested on 7x7 domain

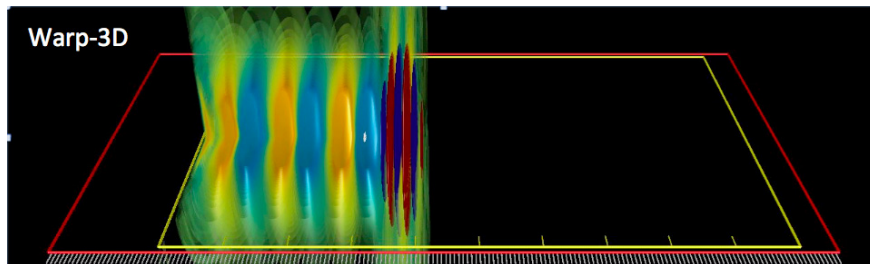


Successfully tested on 2-D modeling of short LPA stages

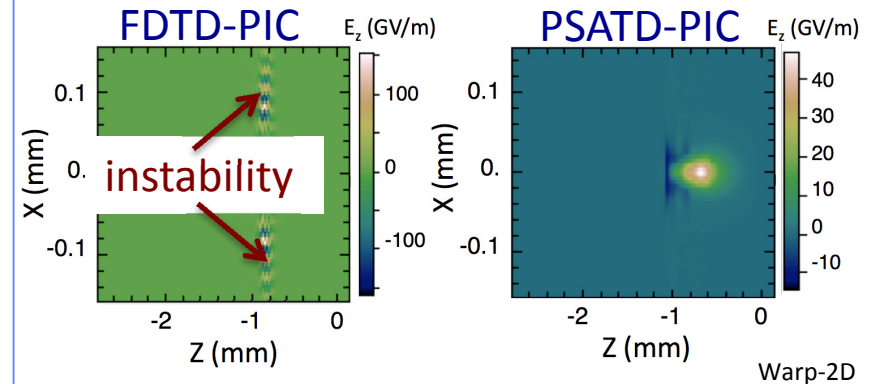
Lorentz boosted frame (wake)



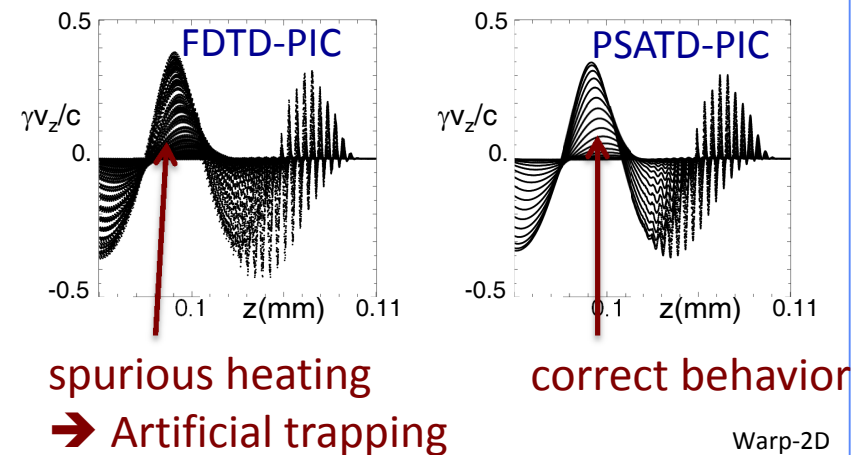
Lab frame



Improved stability

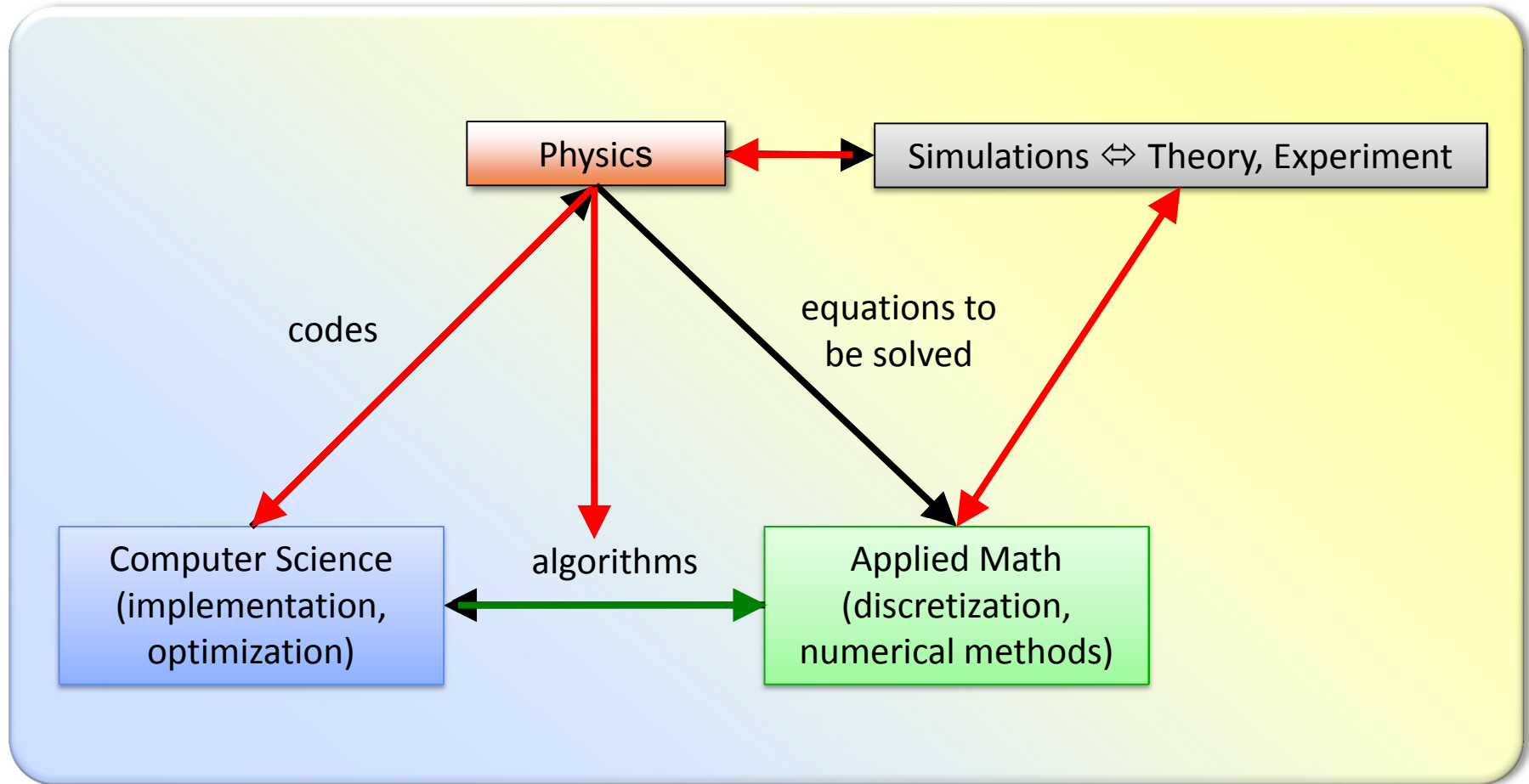


Improved phase space accuracy



Modeling of beams and accelerators is complex

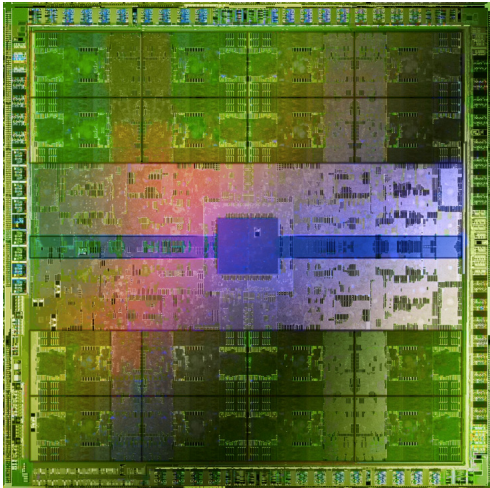
Physics impacts all aspects of modeling of beams and accelerators.



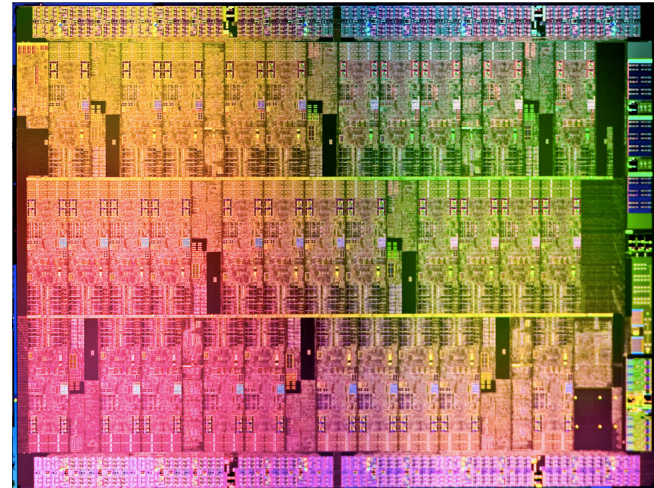
Summary & Outlook

- Importance of computer modeling is **on the rise**.
- Modeling of B&A are **very complex** and **specialized tasks**.
- Expect to see emergence of **local** and **global (virtual) centers**:
 - teams of **physicists + applied math + computer scientists**,
 - **dedicated** to the **science** and **technology** of **B&A modeling**,
 - develop codes **w/ more physics** on **more complex** machines.

NVIDIA GPU Fermi



Intel MIC



Thanks!

This was all possible thanks to the continuous support from:

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- LOASIS/BELLA: W. Leemans, C. Geddes, E. Cormier (now Tech-X).
- Center for Beam Physics: W. Fawley, M. Furman.
- Heavy Ion Fusion Virtual National Lab : Alex Friedman & Dave Grote for Warp.