

Generation of Monoenergetic Protons by Laser Acceleration of Multi-Ion Foils with Polarization Switch

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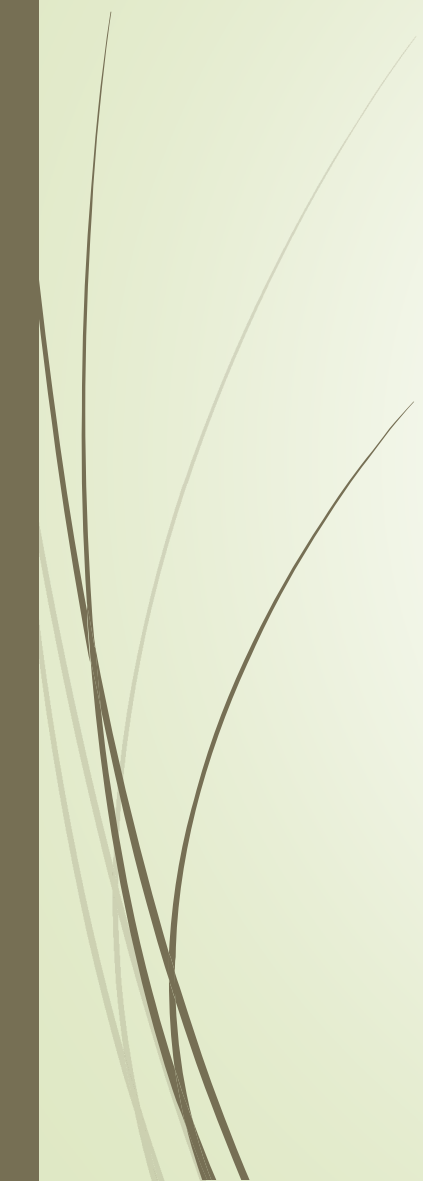
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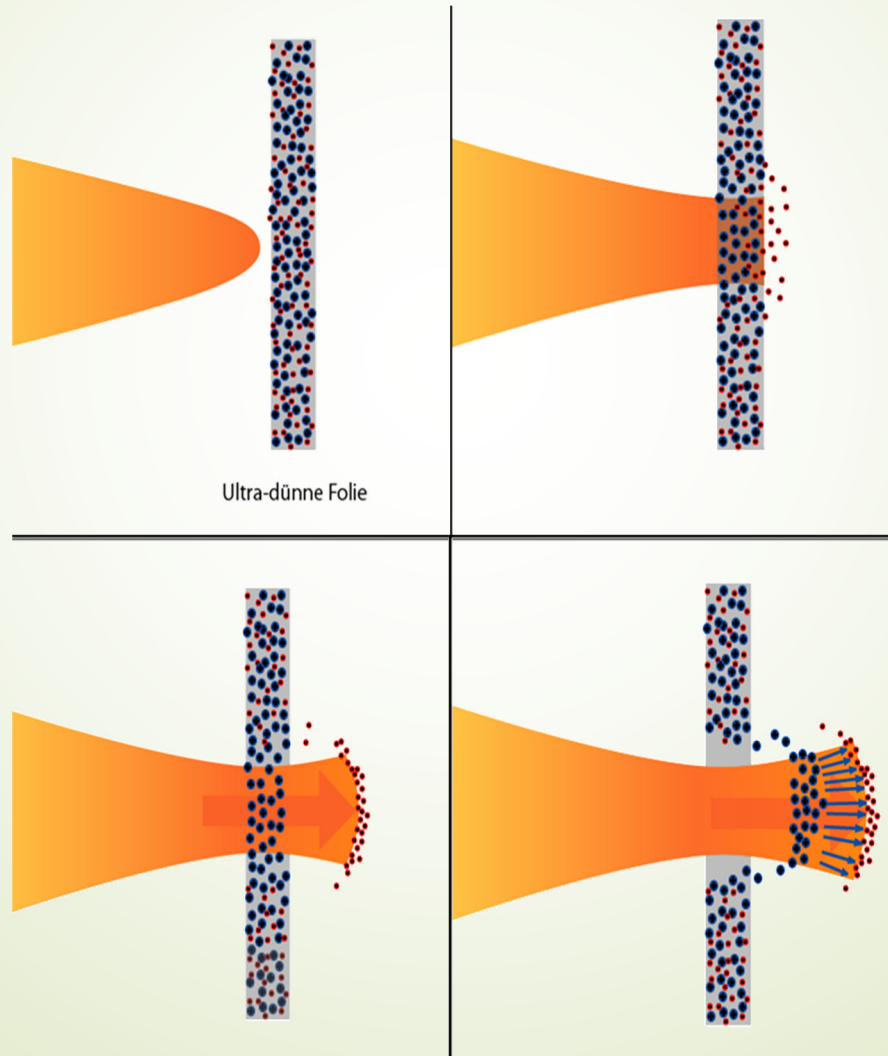


Outline



- Generation of monoenergetic ions to MeV level – laser radiation pressure acceleration
- Extended acceleration using multi-ion foils – a two-stage acceleration scheme combining radiation pressure and Coulomb repulsion
- Further acceleration with laser profile design – switch of laser polarization during acceleration

Radiation Pressure Acceleration



Robinson et al. 2008
Klimo et al. 2008
Yan et al. 2008
M. Schnürer *et al.*, 2012

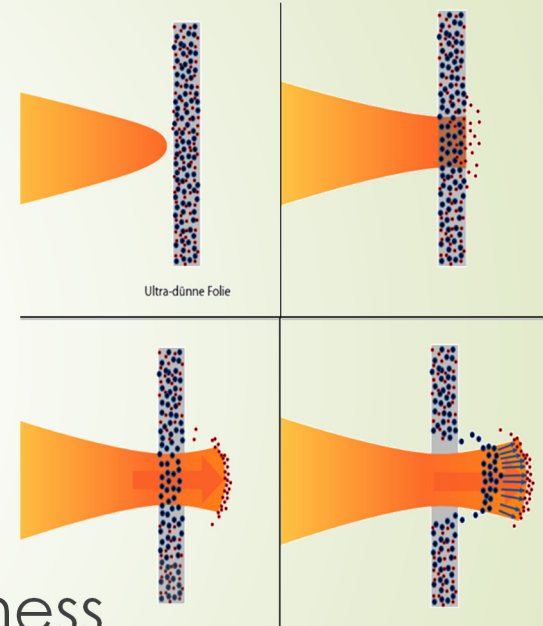
Radiation Pressure Acceleration of Quasi-monoenergetic Hadrons

Requirements:

- Circularly polarized laser
 - Avoiding electric oscillation and 2nd harmonics
- High contrast ratio ($> 10^9$)
 - Minimizing pre-pulse
- Ultra-thin optimal target thickness
 - Stable ion trapping of double-layer structure

Advantage

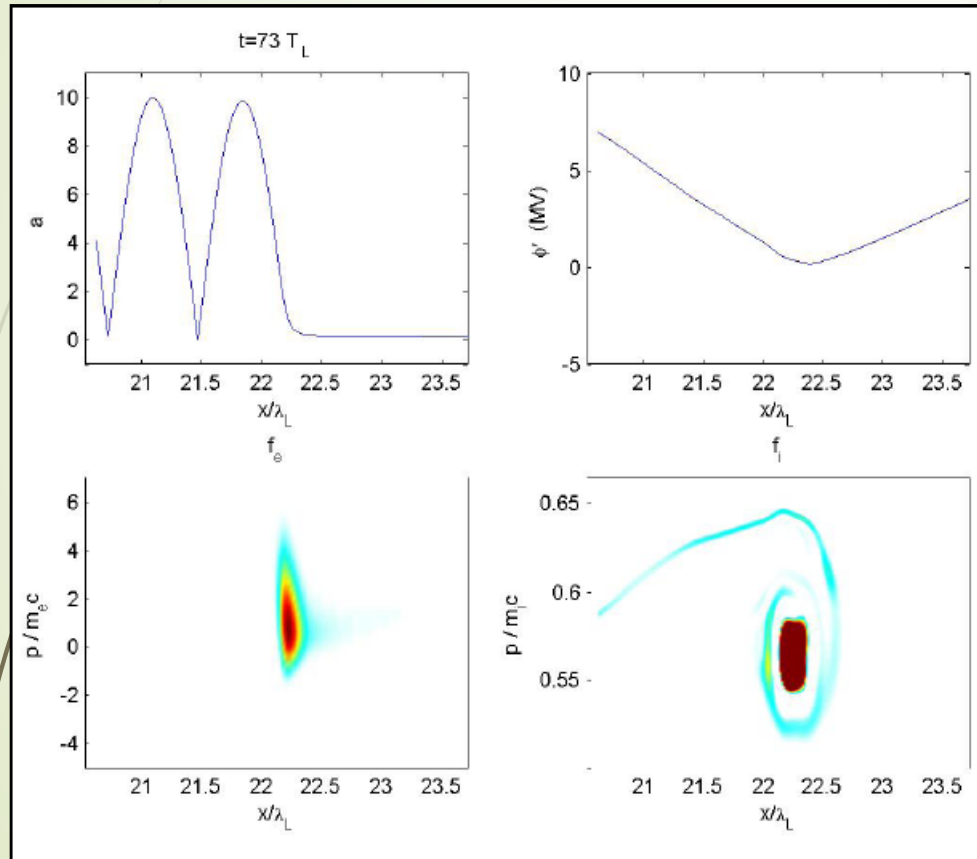
- Monoenergetic ion spectra with high energy



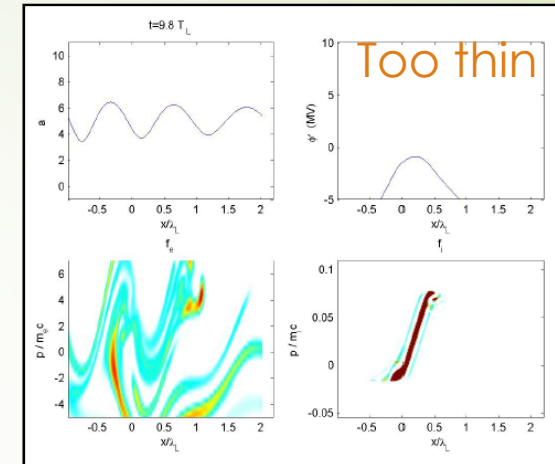
1D Vlasov-Maxwell Simulation

Simulation half the optimal width $l_0 = 0.1 \lambda_L$

Simulation optimal width $l_0 = 0.2 \lambda_L$

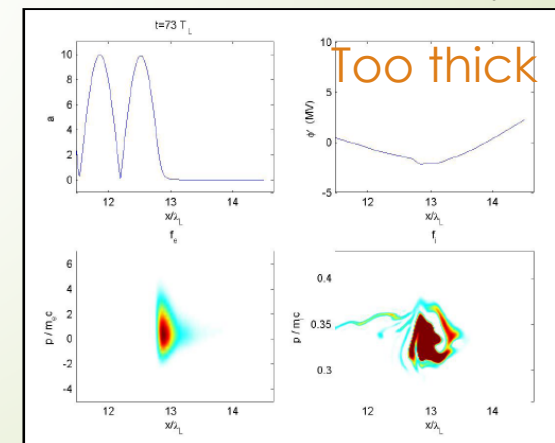


Optimal thickness, stable ion trapping



Laser burns through. Poor acceleration

Simulation twice the optimal width $l_0 = 0.4 \lambda_L$



Wider energy spread, less acceleration

Rayleigh-Taylor Instability (RTI) and Induced Transparency

$$a_0 = 25$$

$$n_{e0} = 41.7n_{cr}$$

$$I_0 = 0.2\lambda_L$$

Proton
density

Electron
density

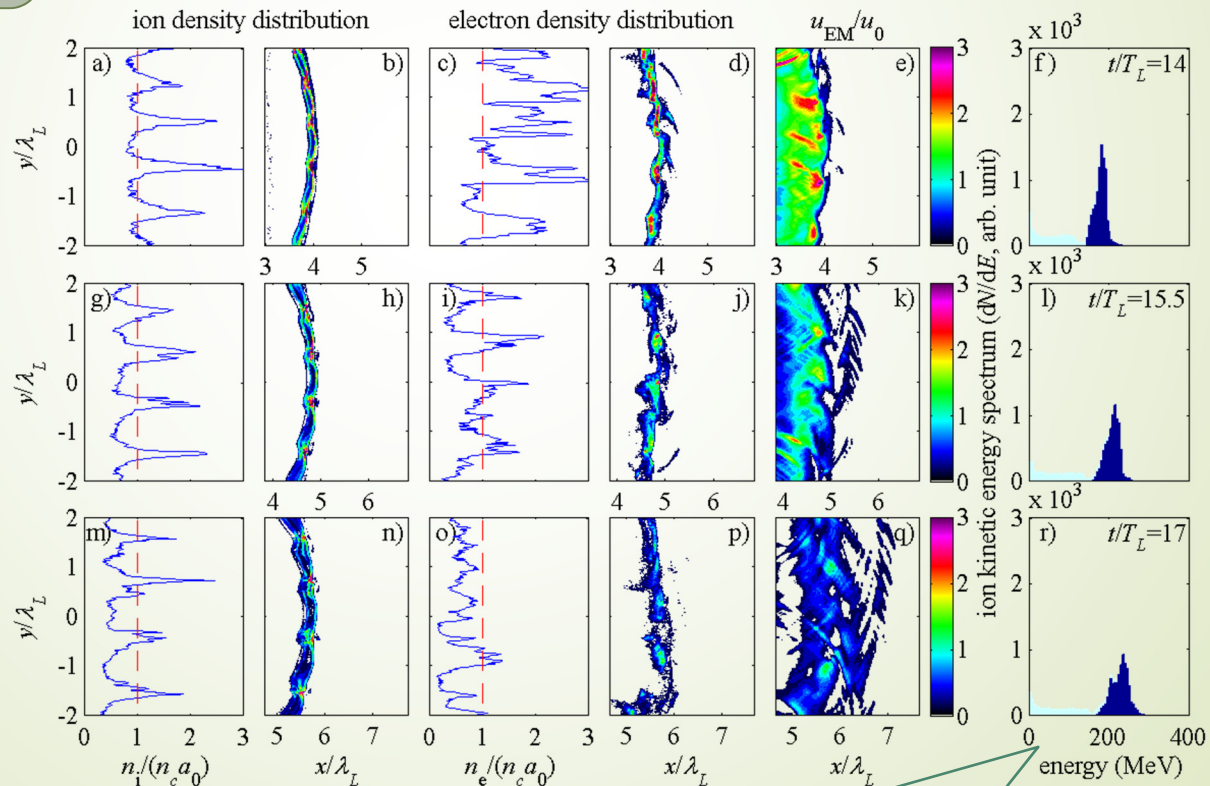
EM
energy
density

Proton
energy
spectra

$t = 14T_L$

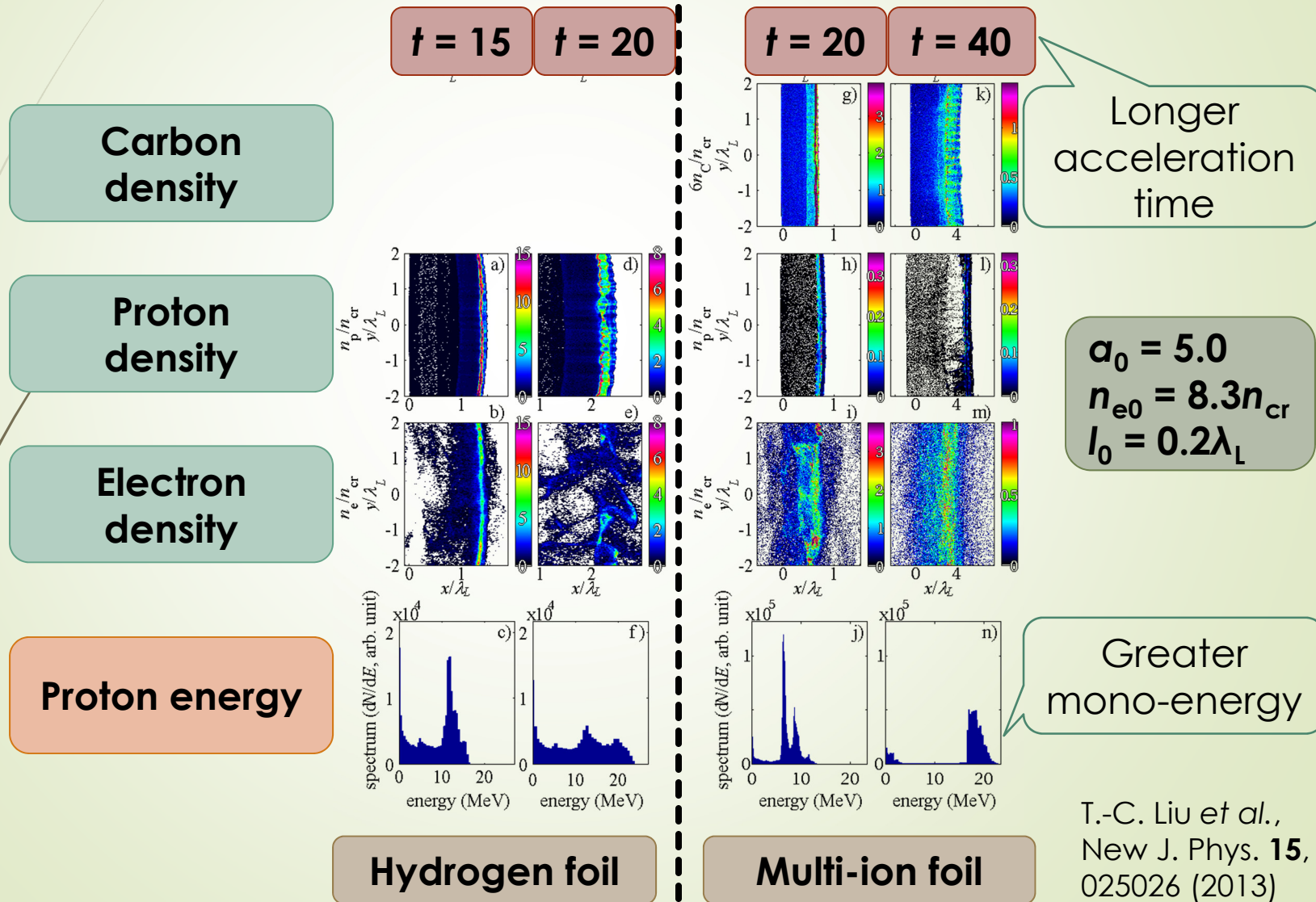
$t = 15.5T_L$

$t = 17T_L$

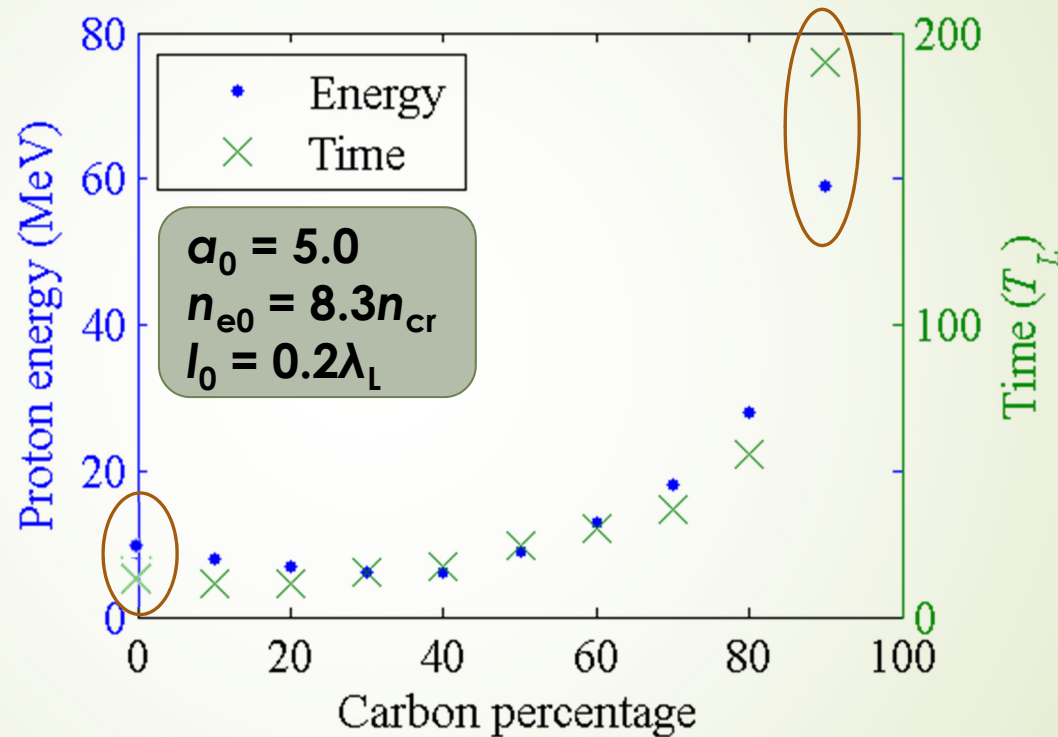


Losing Monoenergetic Property

Multi-Ion Foils – Significantly Enhanced Proton Energy due to Longer Acceleration time



Scaling of Monoenergetic Proton with Different Carbon Concentration



➡ Compared with pure hydrogen foil:

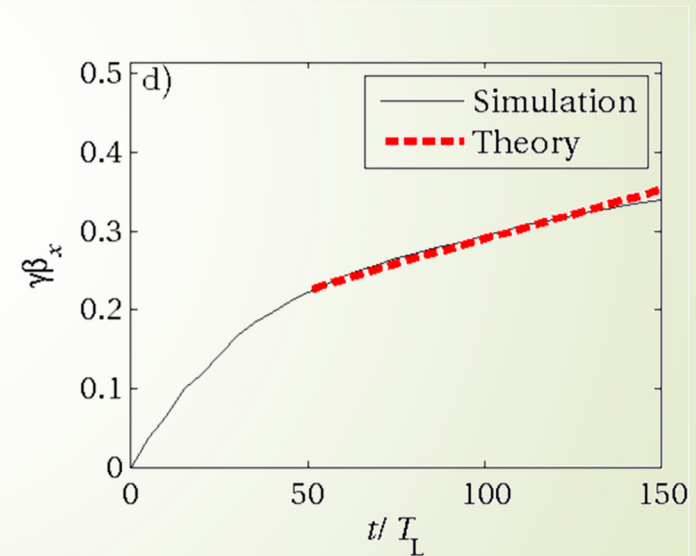
- ➡ Proton energy is **6x** as much.
- ➡ Acceleration time is **10x** extended.

1D Theoretical Model of Proton by Carbon Repulsion Shielded by Electrons

- Assuming protons as a test charge, the Poisson equation
$$\frac{d^2\phi}{dx^2} = \frac{en_0}{\epsilon_0} \exp\left(\frac{e\phi}{k_B T_e}\right)$$
 can be solved to obtain the equation of motion as

$$\begin{cases} \frac{dx_p}{dt} = v_p \\ \frac{d(\gamma_p v_p)}{dt} = \frac{eE_x}{m_p} = \frac{e\sigma_{\text{net}}}{2\epsilon_0 m_p} \coth\left(\frac{x_p - v_C t}{4\epsilon_0 k_B T}\right) \end{cases}$$

- Good agreement is achieved.





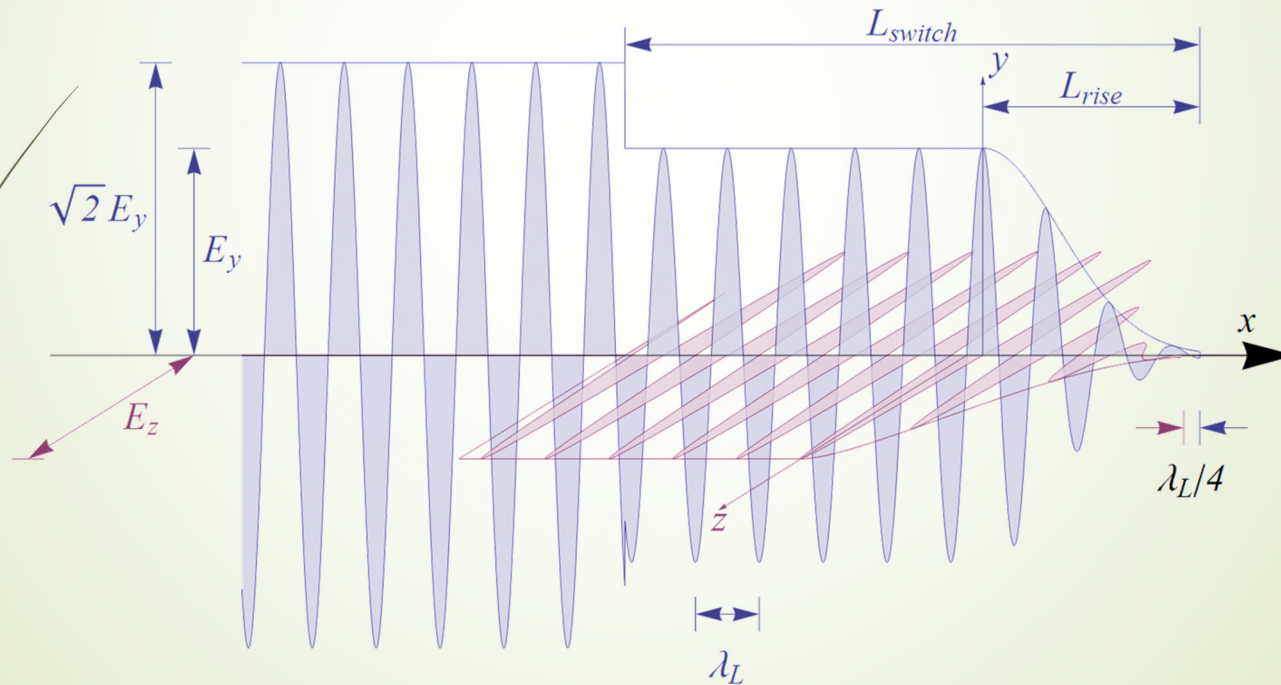
Further Improvement

- Since the **electron temperature** plays an important rule in deciding the shielding effect of Coulomb repulsion due to carbon ions, increasing it could slow down the returning electrons and boost the acceleration.
- **Linear polarized wave** – unwanted heating effect in radiation pressure acceleration now could be an advantage after transparency.

Polarization Switch – Simulation Setup

Switch to Linear Polarization

Initially Circular Polarization



Comparison and Analysis – Electron Temperature

$$a_0 = 5.0$$

$$n_{e0} = 8.3n_{cr}$$

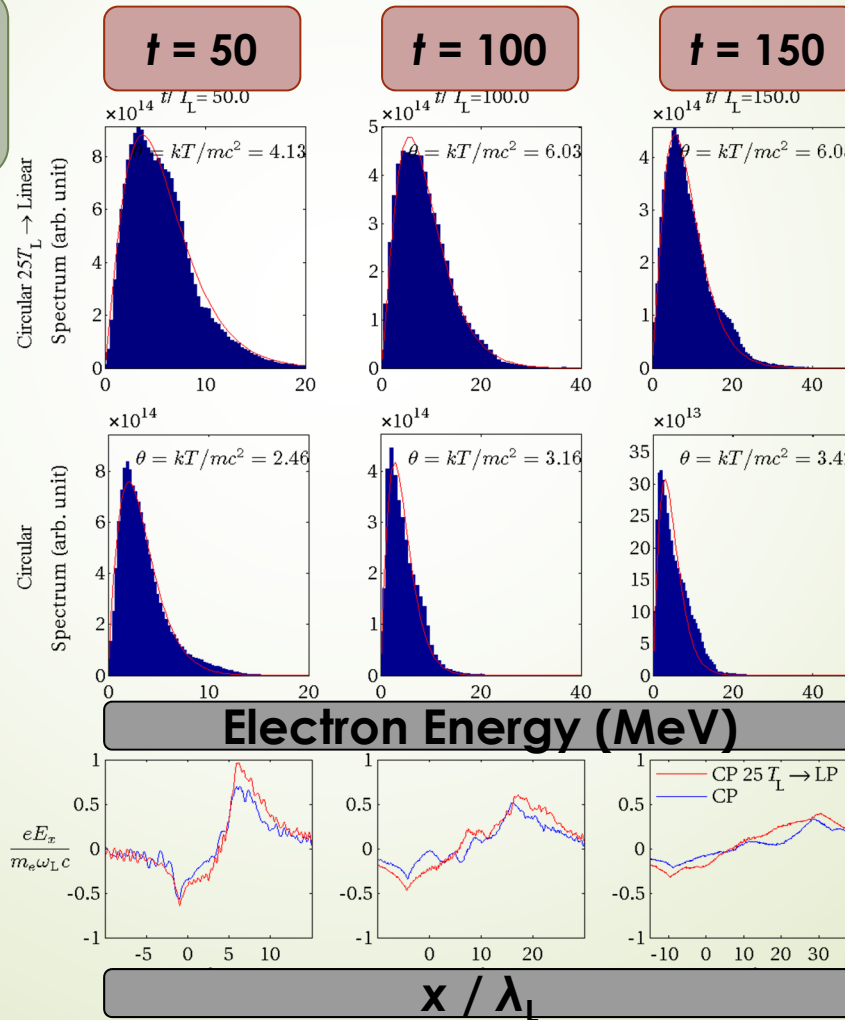
$$I_0 = 0.2\lambda_L$$

With
Switch

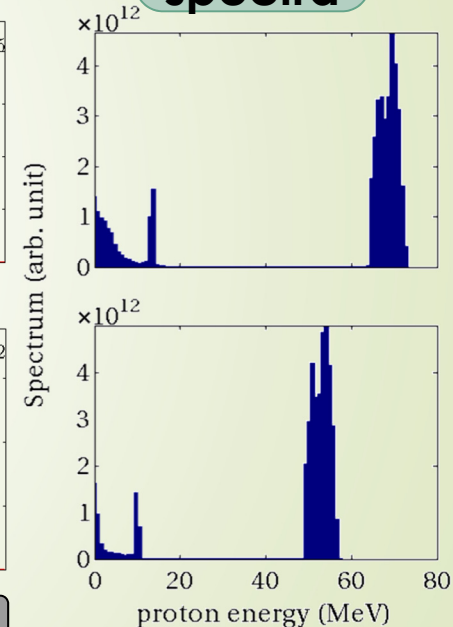
Electron Temperature
in COM Frame

Without
Switch

Electrostatic Field



Proton
energy
spectra

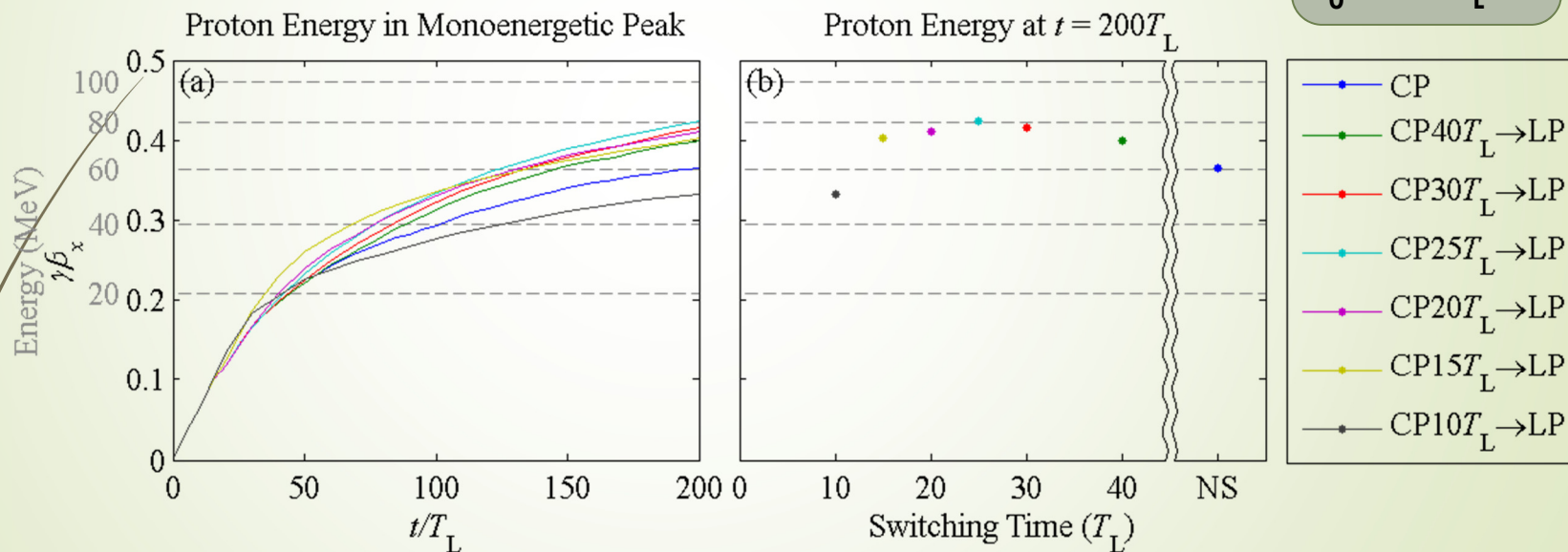


33% Boost in Monoenergy with Optimal Switching Time

$$a_0 = 5.0$$

$$n_{e0} = 8.3n_{cr}$$

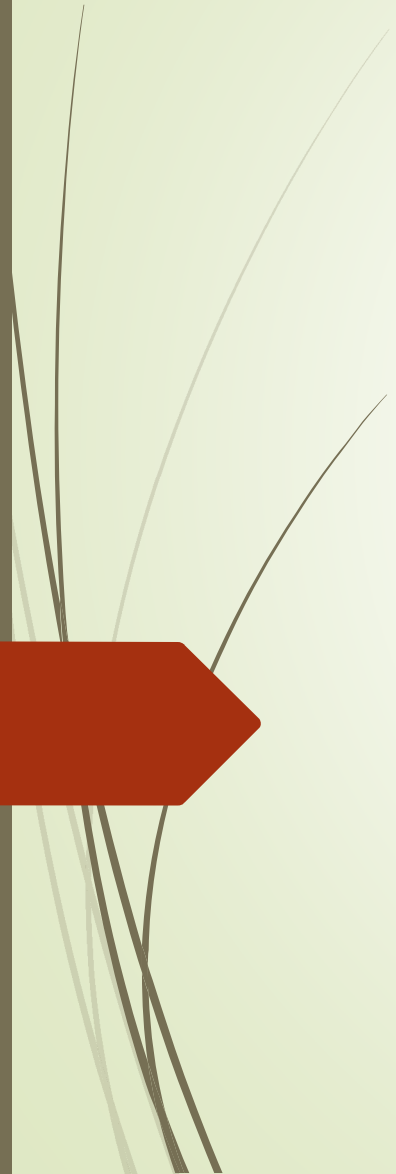
$$I_0 = 0.2\lambda_L$$





Conclusion

- Laser radiation pressure is useful in not only accelerating the proton mono-energetically, but also splitting the proton and carbon ion layers.
- Shielded repulsion can accelerate the protons stably and be modeled with good agreement.
- Switching the polarization increases the electron temperature after transparency and can significantly boost the obtainable energy.



Thank you!

Questions?