

## A Tutorial on Accelerator-Based Light Sources

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## Outline

- Discovery and nature of synchrotron radiation
- Applications of synchrotron radiation
- Radiation-producing devices
- Basics of electron accelerators for light sources
- Storage ring light sources
- Energy recovery linac light sources
- Free-electron lasers
- Conclusion
- Slides mostly contain references to learning resources, not necessarily to seminal papers



## **Discovery of Synchrotron Radiation**

- Lienard predicted in 1898 that a charged particle moving on a circular trajectory would radiate
  - Analysis shows that acceleration causes the radiation
- Accidental direct observation of visible radiation in 1947
  - General Electric 70 MeV electron synchrotron
  - Hence the name "synchrotron radiation"



H. C. Pollock, Am. J. Phys. 51 (1983). G. C. Baldwin, Physics Today, Jan. 1975.



## **Properties of Synchrotron Radiation**

Radiation power is related to  $\gamma$  and the acceleration



 Deflection of relativistic particles is easily accomplished with magnetic fields

$$\vec{F} = q\vec{E} + q\vec{v} \times \vec{B}$$

See, e.g., J. D. Jackson, Classical Electrodynamics, Ch. 14 and K. J. Kim, AIP Conf. Proc. 184 (1989).

#### Radiation2D Demo (K. Shirasawa, T. Shintake)



Available from www-xfel.spring8.or.jp/cband/e/index.html

## Interaction of Radiation with Matter

Radiation energy, frequency, and wavelength are related

$$e(\text{keV}) = 4.14 f(\text{EHz}) = \frac{12.39}{\lambda(A)}$$

- Low-frequency radiation interacts with vibrational modes of molecules
- High-frequency radiation interacts with light, tightly-bound electrons
- Absorption spectra are molecule and element-specific



## Some Applications of Synchrotron Radiation

- Imaging
  - E.g., Diagnostic x-ray takes advantage of differences in absorption for bone and flesh
- Bragg diffraction
  - Typical inter-atomic distances in solids are a few Å
  - X-ray scattering from crystal planes can result in interference effects
  - Used to understand structure of ordered solids



See, e.g., Yale Course "X-Ray Diffraction" on Youtube







Images: Wikipedia

## **Radiation Producing Devices: Bending Magnets**

- Synchrotron radiation originally observed for beam circulating in a "bend" or "dipole" magnet
- Modern rings have many bends to force the beam into a closed path
  - The radiation they produce may or may not be used
- Bend magnet radiation is like a "moving searchlight"





#### **Properties of Bend Magnet Radiation**

- Broad fan of radiation can serve a series of users
- Broad spectrum
- "Critical energy" divides power spectrum in half



## Radiation Producing Devices: Wigglers

- A wiggler is a series of N<sub>w</sub> strong bends of alternating sign
- The beam "wiggles" in a sinusoidal trajectory



Trajectory angle and amplitude characterized by K parameter

$$K = \frac{\theta_{\max}}{1/\gamma}$$

$$x_{\max} = \frac{K\lambda_w}{2\pi\gamma}$$

See, e.g., K. J. Kim, AIP Conf. Proc. 184 (1989).

## **Properties of Wiggler Radiation**

In wigglers, the deflection angle is large



- Only the radiation from the cusps of the trajectory is pointed at the observer
  - **Dipole-like radiation**
  - N \_ times the flux
- Can make field stronger, since there is no net deflection
- Dipole-like radiation with higher flux and higher energy

$$e_c(\text{keV}) = 0.66E(\text{GeV})^2 B(\text{T}) = 2.22 \frac{E(\text{GeV})^3}{\rho(\text{m})}$$

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## **Radiation Producing Devices: Undulators**

- If K≤3 radiation from successive periods interferes constructively
- Average forward velocity of a wiggling electron is

$$\bar{v}_z = c \left( 1 - \frac{1 + \frac{1}{2}K^2}{2\gamma^2} \right)$$

 Radiation has v=c, so after each period it slips ahead by

$$\Delta l = \frac{\lambda_w \left(1 + \frac{1}{2}K^2\right)}{2\gamma^2}$$

 Coherent addition between poles for certain radiation wavelengths

$$\lambda_r n = \Delta l, \qquad n = 1, 2, 3, \dots$$

Undulator radiation peaked harmonics

$$\lambda_{r,n} = rac{\lambda_w \left(1 + rac{1}{2}K^2\right)}{2n\gamma^2}$$

Even harmonics nominally absent on-axis due to symmetry of trajectory

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Diagram after B. McNeil, "How an FEL works FEL2010"

## **Comparison of Radiation Spectra**



- With realistic electron beam, undulator shows even harmonics as well as odd harmonics
- Undulators preferred for high-brightness applications
- Wigglers provide high energy and flux

Undulator: K=1.3, 3.3cm period, 2.4m length Wiggler: B=1.0T, 8.5 cm period, 2.4m length Dipole: B=0.6T

Flux computed through  $5x5 \text{ mm}^2$  pinhole at 30 m for 100 mA APS beam.

Computed with SPECTRA (T. Tanaka, H. Kitamura).

## **Tuning Curve for Undulator**

 Undulators have adjustable gaps that allow changing the field

 $K = 93.4B(T)\lambda_w(m)$ 

- Hence, users can move the location of the maximum brightness to correspond to experimental needs
- Maximum brightness occurs for first harmonic when K=1.3



#### **Basic Electron Beam Properties**

- Accelerators don't just provide energetic electrons, but energetic beams of electrons
  - Electron beam properties strongly affect the properties of the radiation
- Important measures of beam quality
  - Low energy spread (0.02~0.1%)
  - Brief time duration (20 fs ~ 50 ps)
  - Small transverse size and divergence (10~100  $\mu$ m by 1~10  $\mu$ rad)
- The quality of a beam is expressed by the brightness

$$B \propto \frac{N_e}{\sigma_E \sigma_t \sigma_x \sigma'_x \sigma_y \sigma'_y}$$

(simplified form)

Commonly combine transverse quantities into "emittances"

$$\epsilon_x = \sigma_x \sigma_{x'}$$

$$\epsilon_{x,n} = \gamma \sigma_x \sigma_{x'}$$

Geometric emittance (simplified) Constant in absence of radiation

Normalized emittance (simplified) Constant under acceleration.

## **Two Basic Types of Light Source Facility**

- Single-pass light source
  - Beam used once and then discarded



- Beam quality determined primarily by the electron gun
  - See D. Dowell's tutorial on Friday morning
- Operating examples: LCLS, FLASH, SCSS, JLab ERL, ...

- Circulating-beam light source
  - Beam circulates for hours in a "storage ring"
  - Beam quality determined by storage ring optical design
  - Operating examples: ALS, APS, NSLS, SPEAR, ...



## **Quantum Excitation of Electron Beams**

- Radiation emission has a random component
  - Different electrons emit differently
  - Quantum-mechanical effect
- Bending will diminish beam brightness
  - Directly by increasing energy spread
  - Indirectly by increasing bend-plane emittance
- Hence, when electron beam brightness requirements are very demanding, use system with very minimal bending
- Effect can be reduced to some extent (more later)



#### **Radiation Damping**

Geometric emittance is the product of size and divergence

$$\epsilon_x = \sigma_x \sigma_{x'}$$

Divergence decreases when beam is accelerated

$$x' = \frac{p_x}{p_z} \to \frac{p_x}{p_z + \Delta p_z}$$

- "Adiabatic damping" of emittance in a linac
- Results in damping of emittance in rings
- In storage rings, an equilibrium is reached between QE and damping

$$\epsilon_0 \sim \frac{E^2}{N_d^3} \qquad \left(\frac{\sigma_E}{E}\right)_0 \sim \frac{E}{\sqrt{\rho}}$$

See, e.g., M. Sands, op. cit., and J. Murphy, Light Source Data Book.

## **Contemporary Storage Ring Light Sources**

- Most rings are highly periodic and symmetric
  - APS cell is a typical Chasman-Green configuration



## **Methods of Decreasing Emittance**



## Beam Lifetime and Top-Up Injection

- Beam circulating in ring is gradually lost
  - Scattering on residual gas molecules
  - Electron-electron collisions within a bunch (Touschek effect)
- Touschek effect dominates in low-emittance rings

$$\frac{1}{I}\frac{dI}{dt} \sim \frac{I}{N_b\sigma_x\sigma_y\sigma_t}$$

- Prior to 2000, all rings operated in "decay mode"
- In June 2000, APS began "top-up" operation
  - Add current every few minutes
- Widely used, supporting
  - Lower emittance
  - Higher bunch current
  - Greater stability



## **Dynamic and Momentum Acceptance**

- Problem: must get beam into the ring and keep it there for a long time (billions of turns!)
- "Dynamic Acceptance" is the region within which particles can stably exist
  - Determines the aperture available for injection of beam
- "Local momentum acceptance" describes how large a momentum kick a particle can tolerate before being lost
  - Determines Touschek lifetime
- Designer must adjust quadrupoles and sextupoles to maximize these apertures
  - Avoid resonances
  - Taylor amplitude- and momentumdependent focusing ("tune") variation



Frequency-map analysis of an APS upgrade

#### **Ultimate Storage Rings**

- Present rings have asymmetric emittance, e.g., 1 nm x 10 pm
- Some users want small, equal emittances, e.g,. 10 pm x 10 pm
- Combining several ideas gives one possible approach
  - Large (~3 km circum.) ring with  $N_d \sim 10$ /cell
    - Makes  $\varepsilon_0 \approx 30$  pm possible
  - Run on coupling resonance to get  $\varepsilon_x = \varepsilon_0/2 = 15$  pm
  - Fill thousands of buckets to get reasonable lifetime
  - Use "swap-out" injection to replace depleted bunch trains
  - 2~3 orders of magnitude brighter present rings
- Challenges are similar to present rings



![](_page_22_Picture_12.jpeg)

Tutorial on Accelerator-Based Light Sources, M. Borland, March 29, 2011 See, e.g., M. Bei et al., NIM A 622 (2010).

## **Energy Recovery Linac X-ray Sources**

- Linac emittances can be very small compared to rings
  - State-of-art linac: 20 pm x 20 pm (LCLS at 14.1 GeV)
  - State-of-art ring: 1000 pm x 10 pm (PETRA-III at 6 GeV)
- However
  - PETRA-III has 100 mA average current, 14 beamlines
  - LCLS has 30 nA average current, 1 beamline
- ERL is attempt to have some of the best of both worlds
  - High brightness gun
  - Beam not stored, modest emittance dilution
  - "Round beams," like USR
  - High average current made possible by energy recovery
- Operating example at JLab (135 MeV)

![](_page_23_Figure_13.jpeg)

See, e.g., I. Bazarov et al., PAC 2001, 230 (2001).

## Challenges for Hard X-ray ERLs

- Obtaining required emittance to beat storage rings
  - High-voltage DC gun design
  - Limited choice of cathode material
- Need ~24 hour cathode lifetime for an operating facility
  - Difficult when average current must be several 10's of mA
  - Difficult when laser spot is very small (to get small emittance)
- Minimization of beam losses
  - May require transport losses at parts-per-billion level
  - Beam halo is a difficult problem to predict and manage
- Minimization of cryogenic power load
- Unclear if ERLs can compete with USRs

![](_page_24_Figure_12.jpeg)

See, e.g., M. Borland, Ring-based sources overview, FLS2010.

## From Undulators to Free Electron Lasers

- Undulator radiation is bright because *individual* electrons emit inphase at each period
- However, any two electrons have random relative phases, so no coherent addition

$$I_{\rm rad} \propto N_e$$

![](_page_25_Figure_4.jpeg)

 If electrons were locked at positions separated by the wavelength, they'd all emit coherently

$$I_{\rm rad} \propto N_e^2$$

![](_page_25_Figure_7.jpeg)

- Since  $N_{a} \sim 10^{9}$ , this is a huge <u>potential</u> improvement
  - FELs are devices that capitalize on this potential
- The wavelength is Angstroms
  - How can we arrange for electrons to bunch at that scale?

Graphics courtesy A. Zholents. For more, Google, e.g., B. McNeil, "How an FEL works FEL2010"

## Microbunching by Radiation in an Undulator

- Recall the undulator "resonance condition":
  - Radiation slips ahead by one wavelength for each period
    - Electron sees the radiation phase oscillate through 360 deg
  - In the same time, the electron makes a full oscillation in the undulator field
  - Depending on relative phase of oscillations, electron can gain or lose energy to the radiation

$$\Delta U = \vec{F} \cdot \Delta \vec{x} = -e\vec{E} \cdot \Delta \vec{x} \qquad \text{Work=Force*Distance}$$
$$\frac{dU}{dt} = -e\vec{v} \cdot \vec{E} \qquad \text{Power=Force*Velocity}$$

 $\left(\frac{dU}{dt}\right)_1 = -e(-v_x \sin kz)(E_x \sin kz)$  Energy gain after 1 period  $\left(\frac{dU}{dt}\right)_2 = -e(v_x \cos kz)(E_x \sin kz)$  Energy constant  $\left(\frac{dU}{dt}\right)_3 = -e(v_x \sin kz)(E_x \sin kz)$  Energy loss after 1 period

#### **Energy Exchange Leads to Microbunching**

![](_page_27_Figure_1.jpeg)

#### **1D FEL Simulation**

![](_page_28_Figure_1.jpeg)

Program by B. McNeil, available from http://phys.strath.ac.uk/eurofel/rebs/rebs.htm

## **Amplifier FELs**

- The process just described will amplify radiation at the firstharmonic wavelength
- Radiation power grows exponentially with characteristic length L<sub>a</sub>
- Initial radiation may come from several sources
  - Ordinary laser: "seeded FEL"
  - Spontaneous radiation in the undulator: SASE FEL
  - Upstream FEL with spectral filtering: self-seeded FEL
- In SASE (Self-Amplified Spontaneous Emission) case, saturation occurs in about 16 gain lengths
  - About 10<sup>7</sup> gain over spontaneous

![](_page_29_Figure_9.jpeg)

See, e.g., S. Schreiber, Rev. Accel. Sci. Tech., Vol. 3 (2010); J. Feldhaus, et al., Opt. Commun. 140 (1997).

Courtesy P. Emma (SLAC).

#### Beam Quality Requirements for SASE FELs

- X-ray FELs are only now becoming available
  - Beam quality requirements very demanding, hard to meet
- Geometric emittance must be less than radiation emittance

$$\frac{\epsilon_n}{\gamma} \lesssim \frac{\lambda_r}{4\pi}$$

Shorter wavelength requires higher beam quality and/or energy

Energy spread must be less than Pierce parameter

$$\sigma_{\delta} \lesssim \rho = \frac{1}{4} \left( \frac{1}{2\pi^2} \frac{I_{pk}}{I_A} \frac{\lambda_w^2 K^2}{\beta_x \epsilon_n \gamma^2} \right)^{\frac{1}{3}}$$

Relative energy spread must be small, peak current must be high, emittance must be low.

Want gain length as short as possible

$$L_g \approx \frac{\lambda_w}{4\pi\sqrt{3}\rho}$$

Want short undulator period, high peak current, low emittance

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See, e.g., P. Emma, SLAC Summer School Lecture, June 2010.

## LCLS Hard X-ray FEL (1.5 to 15 Å)

![](_page_31_Figure_1.jpeg)

- Rf photocathode gun: 250 pC/bunch,  $\varepsilon_n = 0.5 \mu m$
- Energy up to 14 GeV
- ~130m long undulator
- Two magnetic bunch compressors (250 MeV and 4.3 GeV)
  - Compress bunch from 2 ps to  $\sim$ 20 fs rms

See P. Emma, et al., Nature Photonics, Vol. 4 (2010).

# **Magnetic Bunch Compression**

![](_page_32_Figure_1.jpeg)

#### Many Challenges Overcome Already for SASE FELs

- Microbunching can be suppressed if beam quality is poor
  - Emittance must be very good (0.5  $\sim$  1 um normalized)
  - Energy spread must be low (e.g,. well below 0.1% rms)
- To get workable gain length, peak current must be very high
  - To get gain length of a few meters need a few kA
  - Bunch compression must be used, but has side effects
    - Emittance and energy spread growth due to coherent radiation
    - Possible microbunching instability
- These challenges have been overcome by several operating facilities
  - E.g., LCLS, FLASH, SCSS
- Several others will come on line in the next few years
  - E.g., SPRing8 XFEL, Euro XFEL, LCLS-II

## **Challenges for SASE FELs**

- SASE builds from spontaneous radiation, so output is noisy
  - Pulse-to-pulse jitter
  - Spectral jitter
  - Several approaches to improving this
    - Seeding with external laser
    - Self-seeding with spectrally-filtered upstream FEL radiation
- Hard to get high repetition rate from affordable linac, e.g., LCLS is 120 Hz, 1~100 kHz ideal
- Similarly, hard to support multiple users (one at a time for now)
  - Pulse sharing less palatable when linac rate is low
- Expensive to operate (10x cost per experiment compared to rings)
- Challenging to get very hard (>30 keV) x-rays
  - Higher energy, higher beam quality needed
- In spite of these issues, FELs have a decisive advantage for certain demanding experiments, e.g.,
  - Time-resolved studies requiring sub-picosecond resolution
  - Structure determination for very small crystals

## Conclusions

- The basic physics behind accelerator light sources was discovered well over 100 years ago
- Storage ring sources are in high demand and ideal for many experiments
  - Mature, flexible, reliable, cost-effective technology
  - Very wide tuning range with dozens of independent users
  - Appears 10<sup>2</sup> to 10<sup>3</sup> brightness improvement possible, but expensive
- SASE FELs offer  $\sim 10^{10}$  higher peak brightness than rings
  - Full transverse coherence
  - Intense femtosecond pulses
  - Expensive but indispensable for some research
  - Very active research to improve characteristics
- ERLs serve a niche for coherent IR and THz radiation
  - Much R&D needed to prove potential for wider spectral range
- This talk has only scratched the surface
  - Exciting new ideas and opportunities abound