



# Single-knob beam line for transverse emittance partitioning

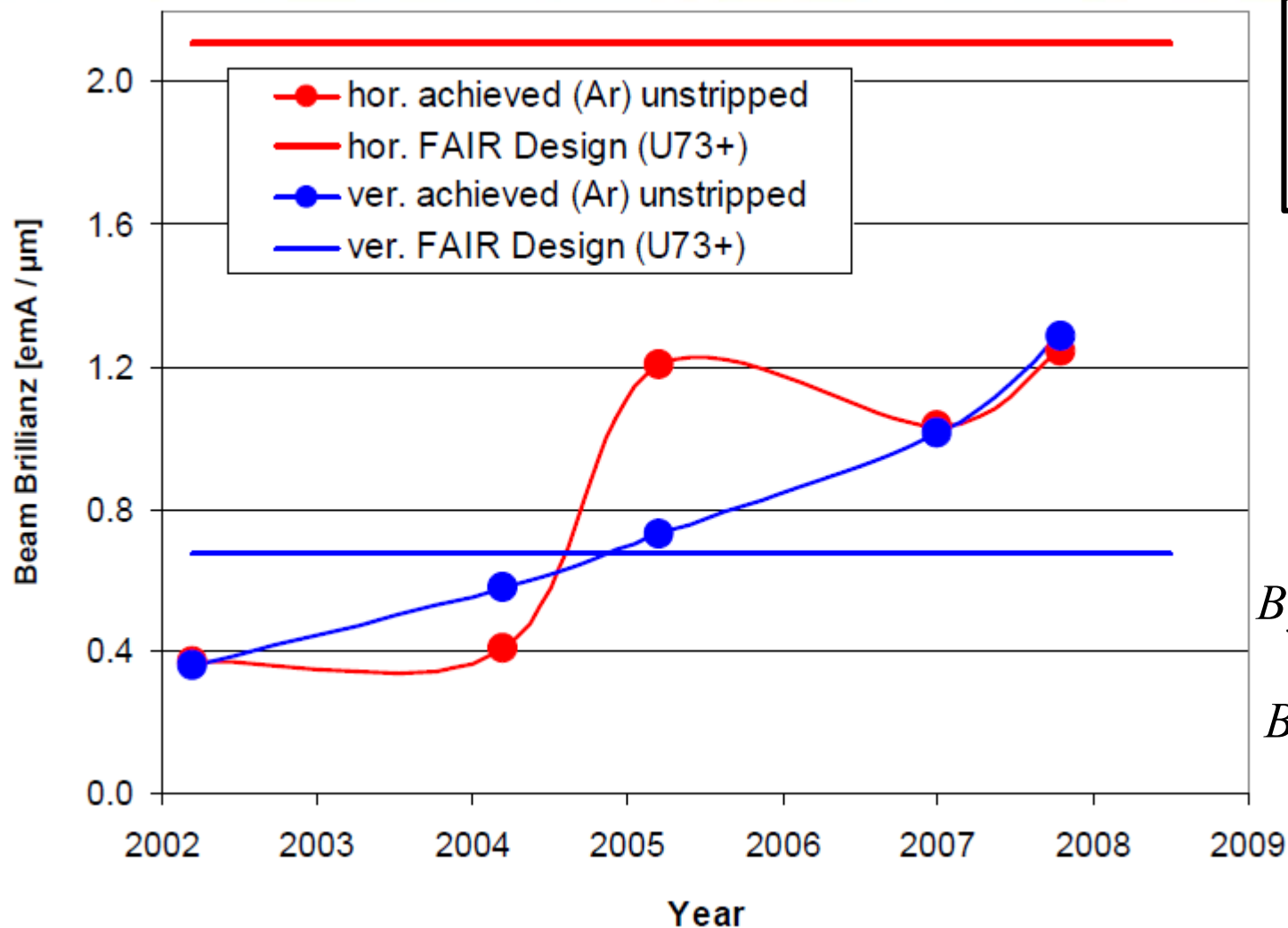
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Darmstadt, Germany

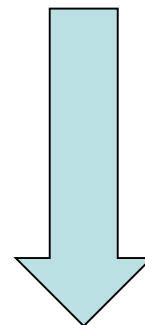
# Outline

- Motivation
- Basic terms
- **E**mittance **t**ransfer **e**xperiment (**E**mt**e**x)
- Decoupling and matching capability

# FAIR Injector Linac Design Requirements



$$B_{x,y} = \frac{\frac{q}{A} \times I_{beam}}{\varepsilon_{x,y}}$$



$B_x$  is not enough

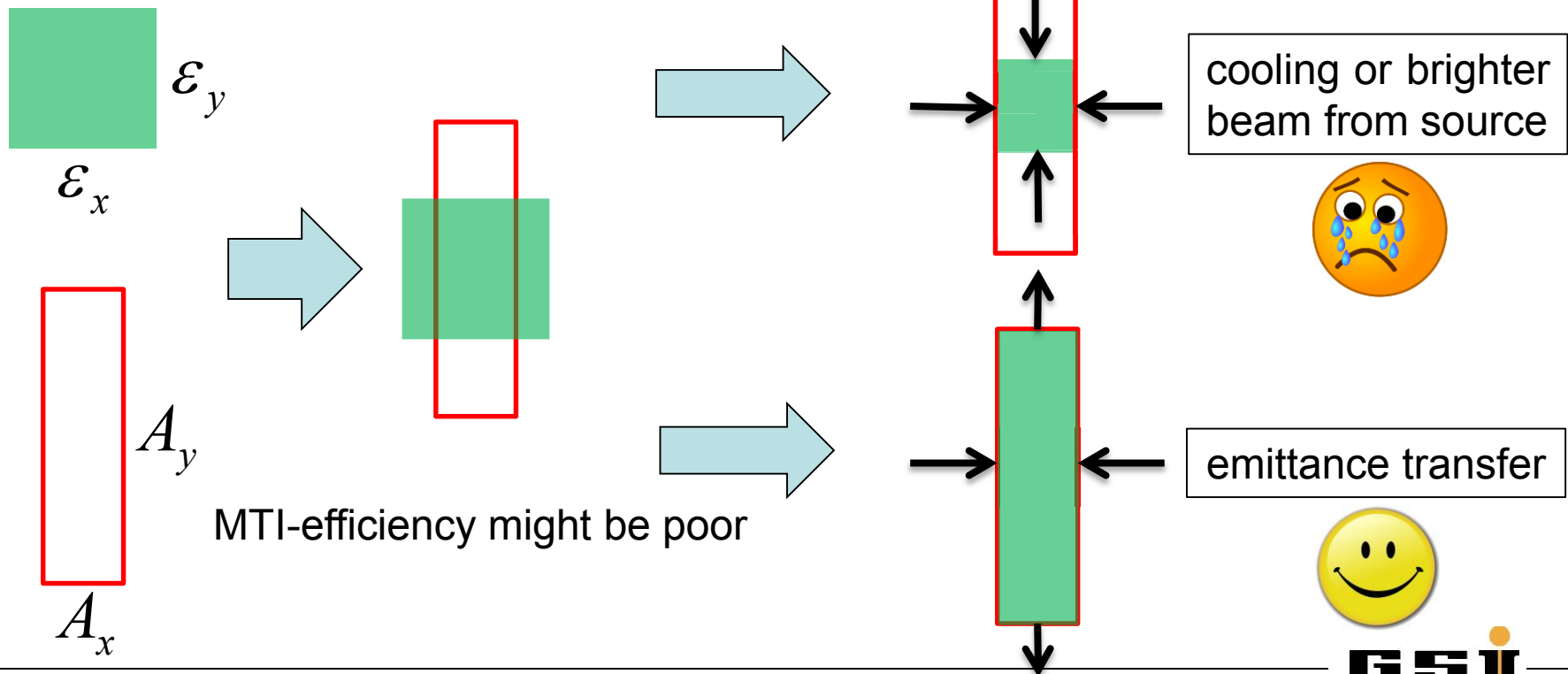
$B_y$  is enough

$$\varepsilon_x \approx \varepsilon_y$$

# Injection into Synchrotron: Acceptance versus Emittance

- Round beam is provided by the injector  $\varepsilon_x \approx \varepsilon_y$

- Flat beam is required for synchrotron  $A_x \approx \frac{1}{3} A_y$



# Second Beam Moments & Emittances

$$C = \begin{bmatrix} \langle xx \rangle & \langle xx' \rangle & \langle xy \rangle & \langle xy' \rangle \\ \langle x'x \rangle & \langle x'x' \rangle & \langle x'y \rangle & \langle x'y' \rangle \\ \langle yx \rangle & \langle yx' \rangle & \langle yy \rangle & \langle yy' \rangle \\ \langle y'x \rangle & \langle y'x' \rangle & \langle y'y \rangle & \langle y'y' \rangle \end{bmatrix}.$$

If at least one of the elements of the off-diagonal sub-matrix is non-zero, the beam is x-y coupled.

$$\varepsilon_1 = \frac{1}{2} \sqrt{-\text{tr}[(CJ)^2] - \sqrt{\text{tr}^2[(CJ)^2] - 16|C|}},$$
$$\varepsilon_2 = \frac{1}{2} \sqrt{-\text{tr}[(CJ)^2] + \sqrt{\text{tr}^2[(CJ)^2] - 16|C|}}.$$

Transverse eigen-emittances are equal to the transverse rms emittances when inter-plane correlations are zero.

$$J = \begin{bmatrix} 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & -1 & 0 \end{bmatrix} \quad M^T J M = J.$$

Eigen-emittances are invariant under symplectic transformations.

# Beam through thin Solenoid with Foil in its Center

$$\begin{bmatrix} \varepsilon\beta & 0 & 0 & 0 \\ 0 & \frac{\varepsilon}{\beta} & 0 & 0 \\ 0 & 0 & \varepsilon\beta & 0 \\ 0 & 0 & 0 & \frac{\varepsilon}{\beta} \end{bmatrix}$$

$$R_{in} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & k_{in} & 0 \\ 0 & 0 & 1 & 0 \\ -k_{in} & 0 & 0 & 1 \end{bmatrix}$$

$$R_{out} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & -k_{out} & 0 \\ 0 & 0 & 1 & 0 \\ k_{out} & 0 & 0 & 1 \end{bmatrix}$$

$$k = k_{in} \neq k_{out}$$

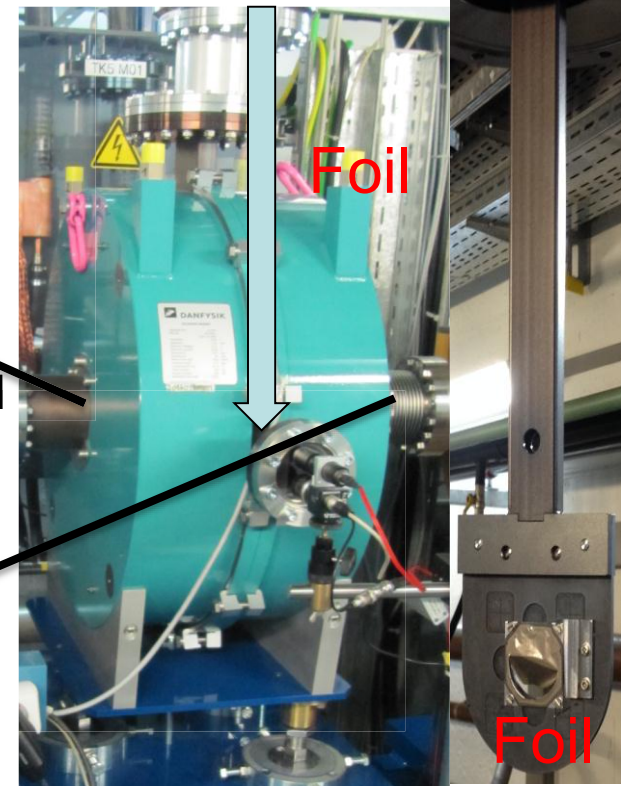
$$k_{in} = k, k_{out} = \delta q k$$

$$\delta q = \frac{(B\rho)_{in}}{(B\rho)_{out}}, a = \delta q - 1$$

fringe field matrices

$$\begin{bmatrix} \varepsilon\beta & 0 & 0 & -k\varepsilon\beta \\ 0 & \frac{\varepsilon}{\beta} + k^2\varepsilon\beta & k\varepsilon\beta & 0 \\ 0 & k\varepsilon\beta & \varepsilon\beta & 0 \\ -k\varepsilon\beta & 0 & 0 & \frac{\varepsilon}{\beta} + k^2\varepsilon\beta \end{bmatrix}$$

after entrance fringe field

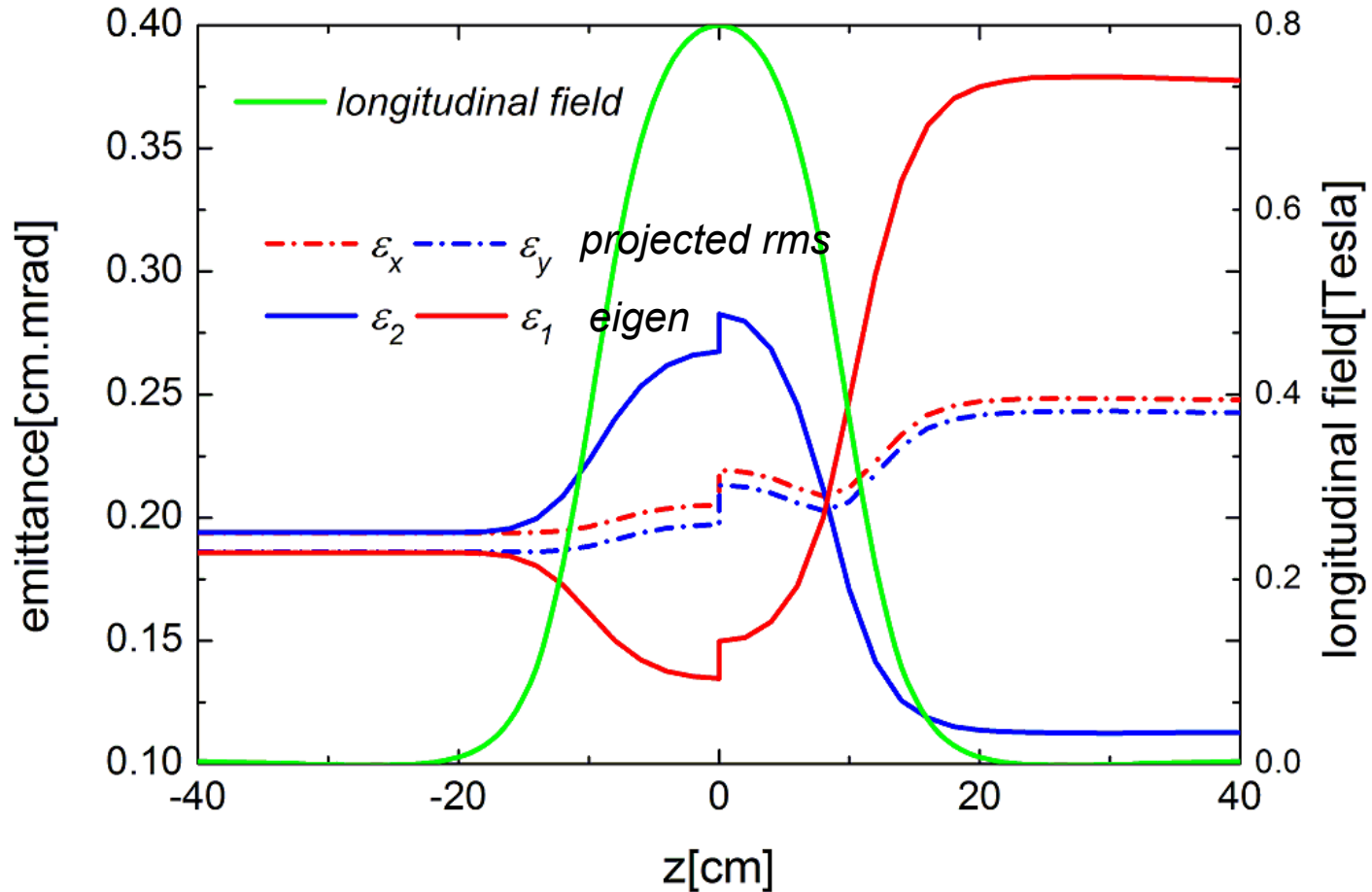


$$\begin{bmatrix} \varepsilon\beta & 0 & 0 & a k \varepsilon \beta \\ 0 & \frac{\varepsilon}{\beta} + a^2 k^2 \varepsilon \beta + \Delta \varphi^2 & -a k \varepsilon \beta & 0 \\ 0 & -a k \varepsilon \beta & \varepsilon \beta & 0 \\ a k \varepsilon \beta & 0 & 0 & \frac{\varepsilon}{\beta} + a^2 k^2 \varepsilon \beta + \Delta \varphi^2 \end{bmatrix}$$

after exit fringe field



# To Create Inter-Plane Correlations & Change Eigen-Emittances



$$t = \frac{\varepsilon_x \varepsilon_y}{\varepsilon_1 \varepsilon_2} - 1 \geq 0$$

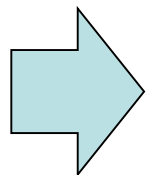
Stripping is non-symplectic action

To quantify the inter-plane coupling

# Then Remove Inter-Plane Correlations & Preserve Eigen-emittances

$$R_{\text{triplet}} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & \mu \\ 0 & 0 & -\frac{1}{\mu} & 0 \end{bmatrix}$$

Quad triplet



$$\bar{R} = R_r R_{\text{triplet}} R_r^T \quad \text{Skew quad triplet}$$

See K.J. Kim PRST-AB 6, 104002 (2003)

$$\mu = \beta_n = \sqrt{\frac{\varepsilon\beta}{\frac{\varepsilon}{\beta} + a^2 k^2 \varepsilon\beta + \Delta\varphi^2}}$$

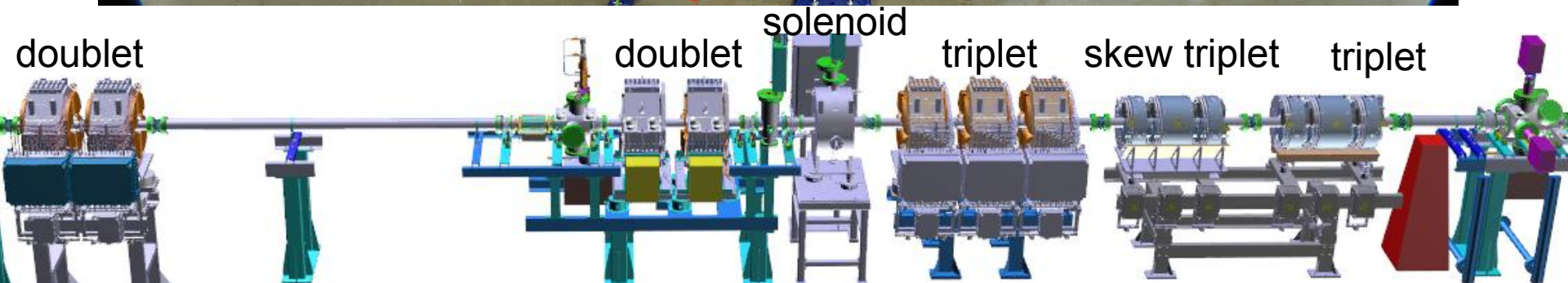
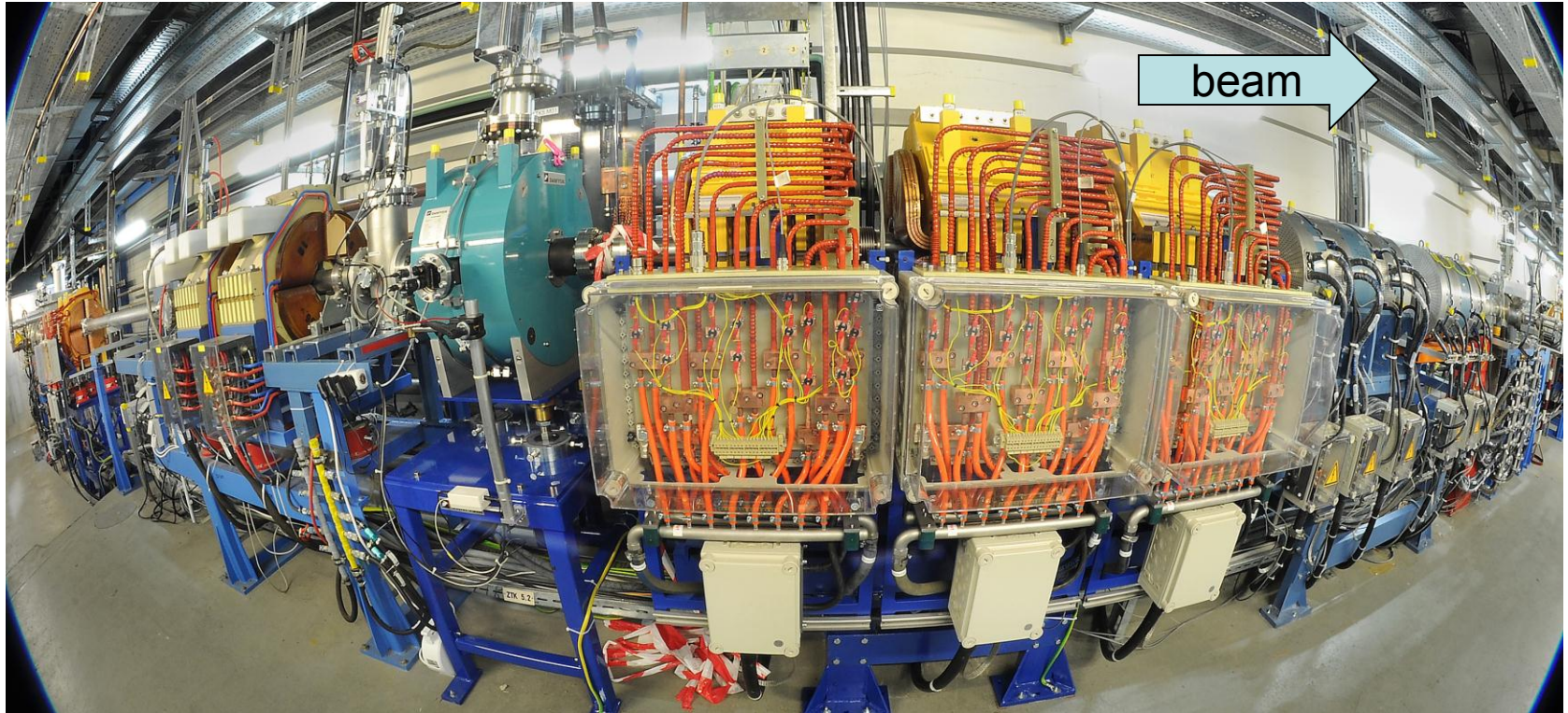
Correlations are removed!!!

angular scattering in foil

Symplectic action, the projected rms emittances are minimized to the eigen-emittance

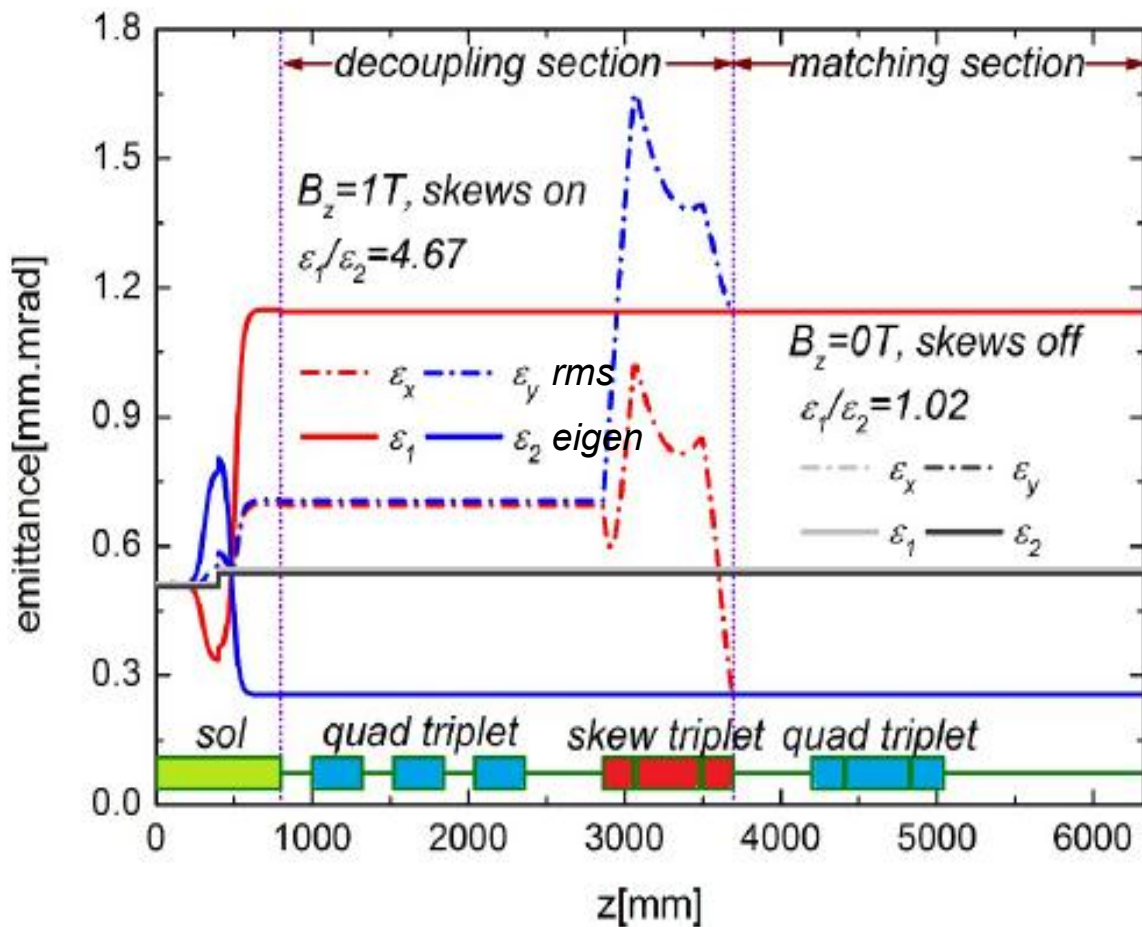


# Emittance Transfer Experiment at GSI



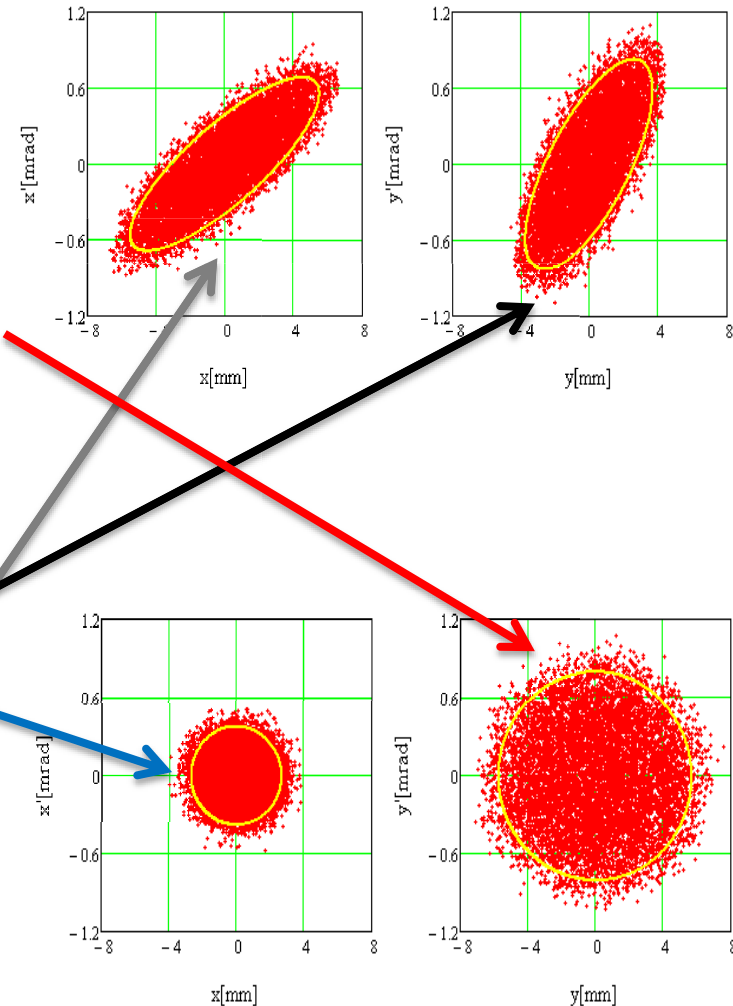


# Beam Dynamics Simulations

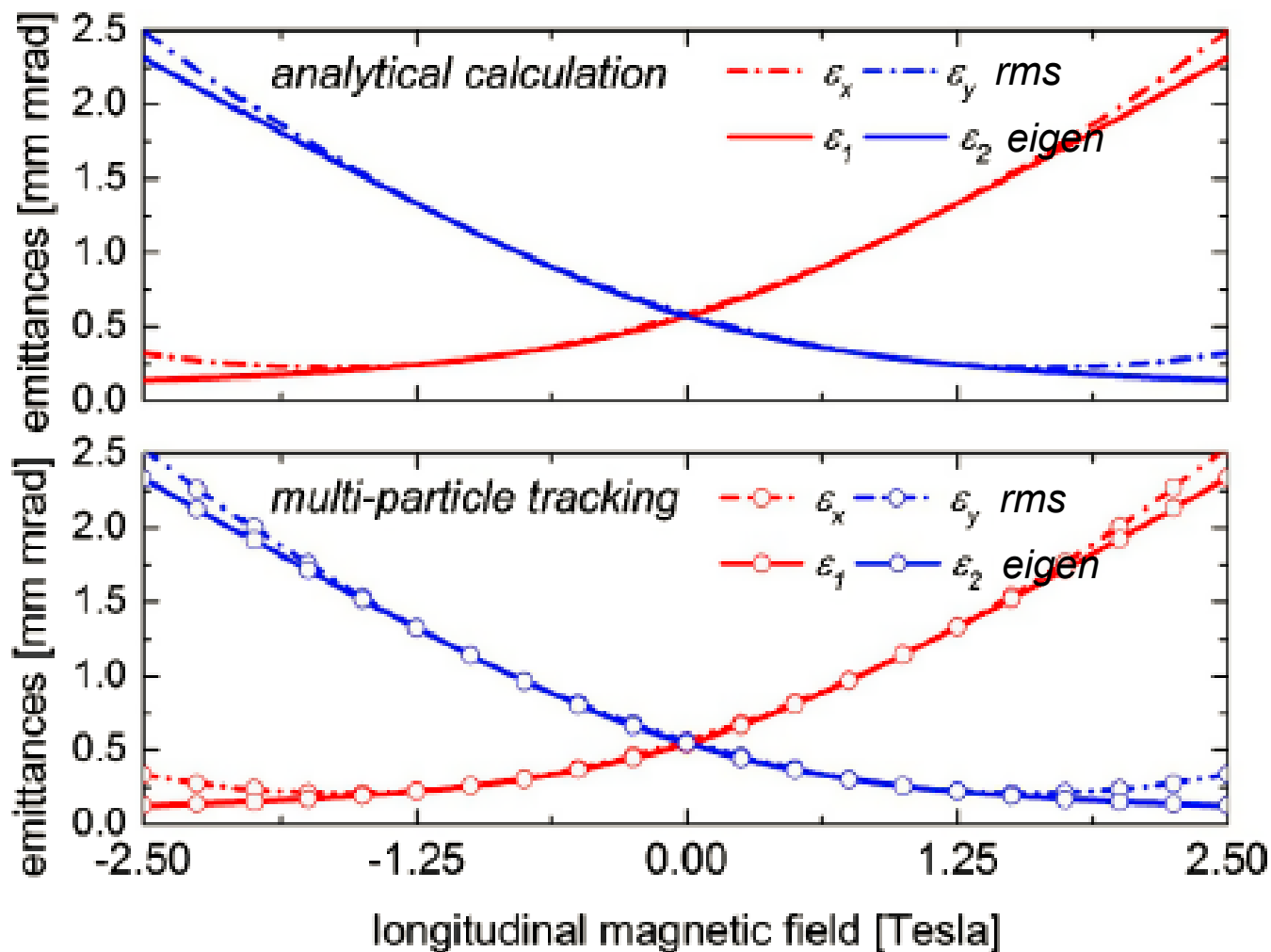


11.4 MeV/u

$D_3^+ \Rightarrow 3D^+$

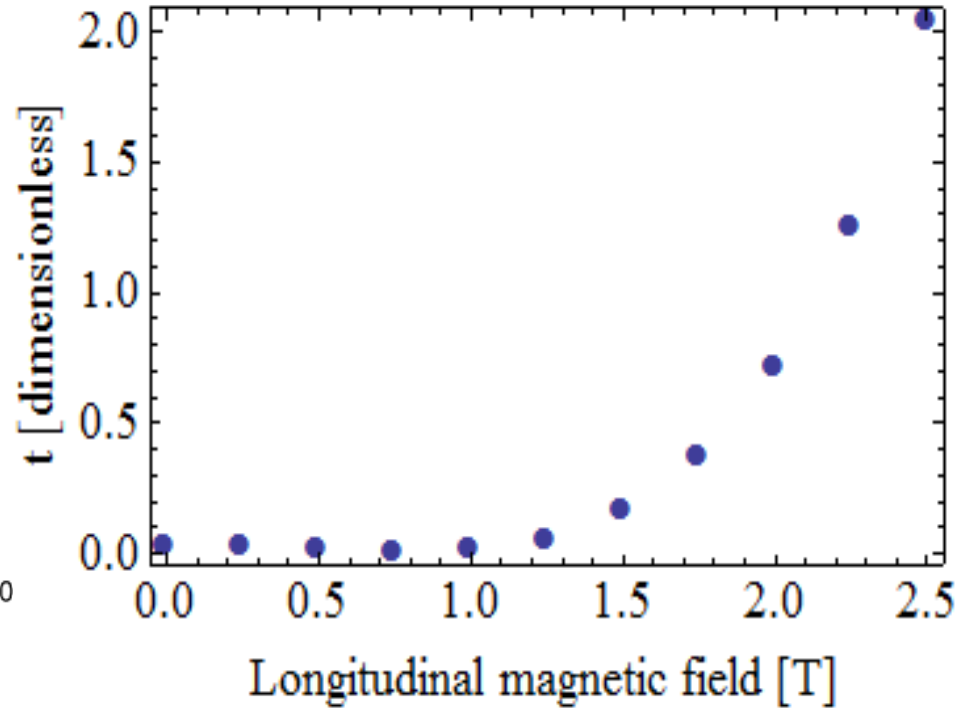
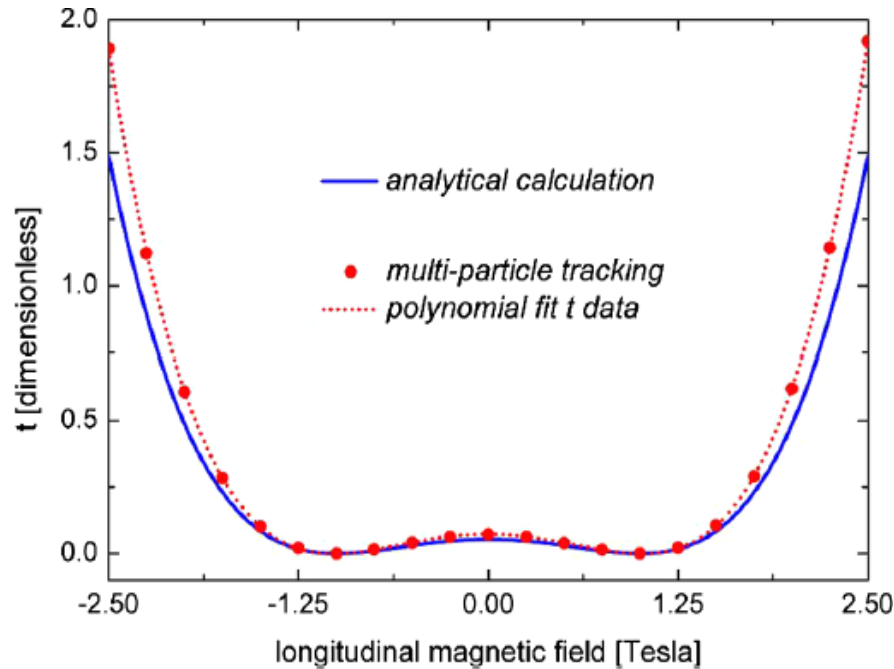


# Decoupling Capability



Gradients of quads (normal and skews) are optimized for B=1.0T

# Independent Confirmation



graph from C. Xiao, PRST-AB 2013

$$t = \frac{\varepsilon_x \varepsilon_y}{\varepsilon_1 \varepsilon_2} - 1 \geq 0$$

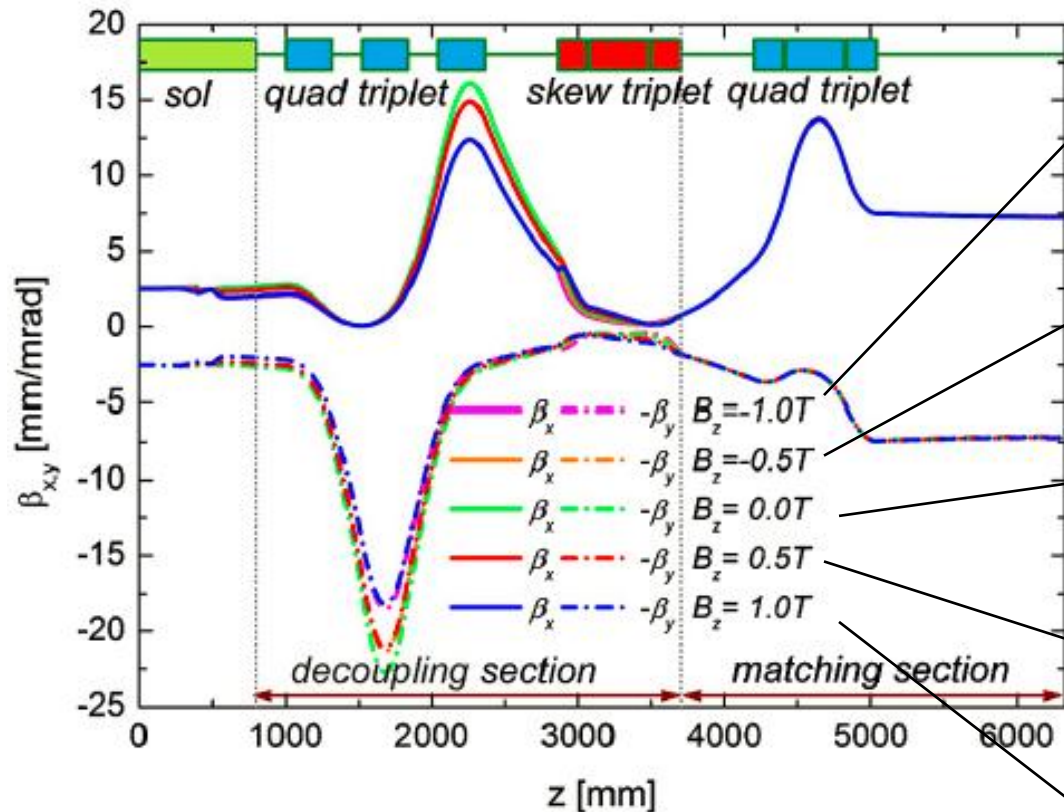
applying innovative 4d-envelope model first published in PRL 2013

H. Qin, Princeton University, USA

M. Chung, UIST, Korea

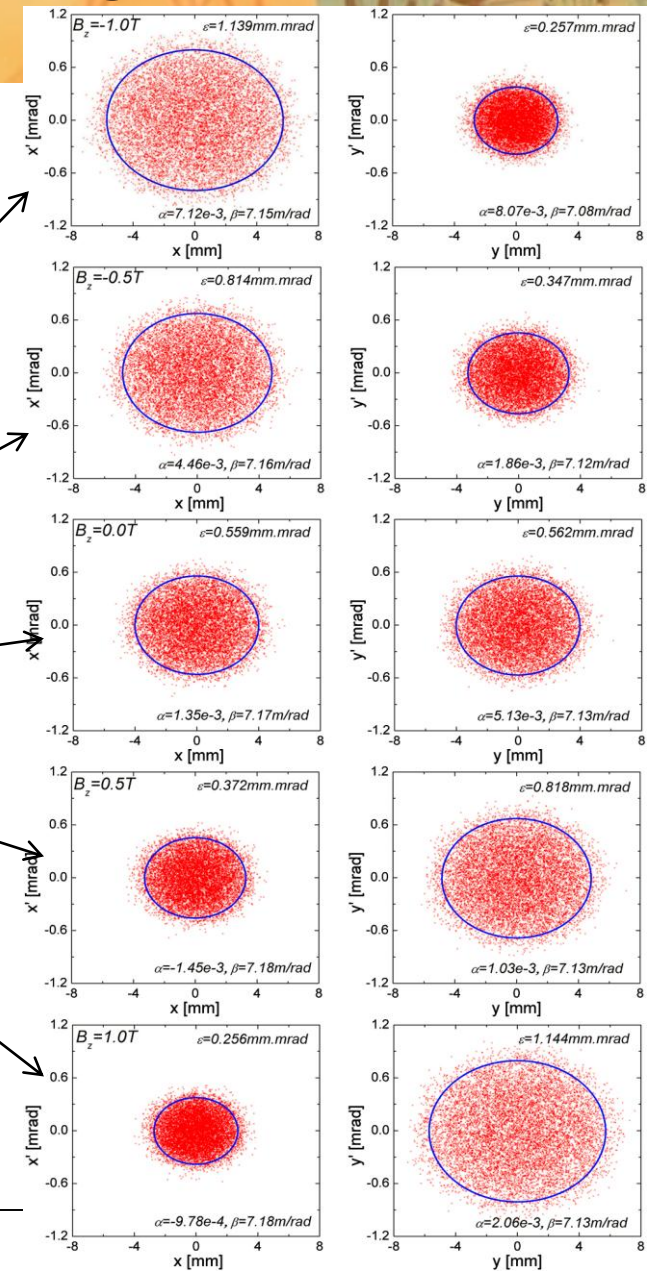


# Matching Capability



- Only the longitudinal magnetic field is varied
- all (skew) quads remain constant

See C.Xiao, PRST-AB 16, 044201 (2013)

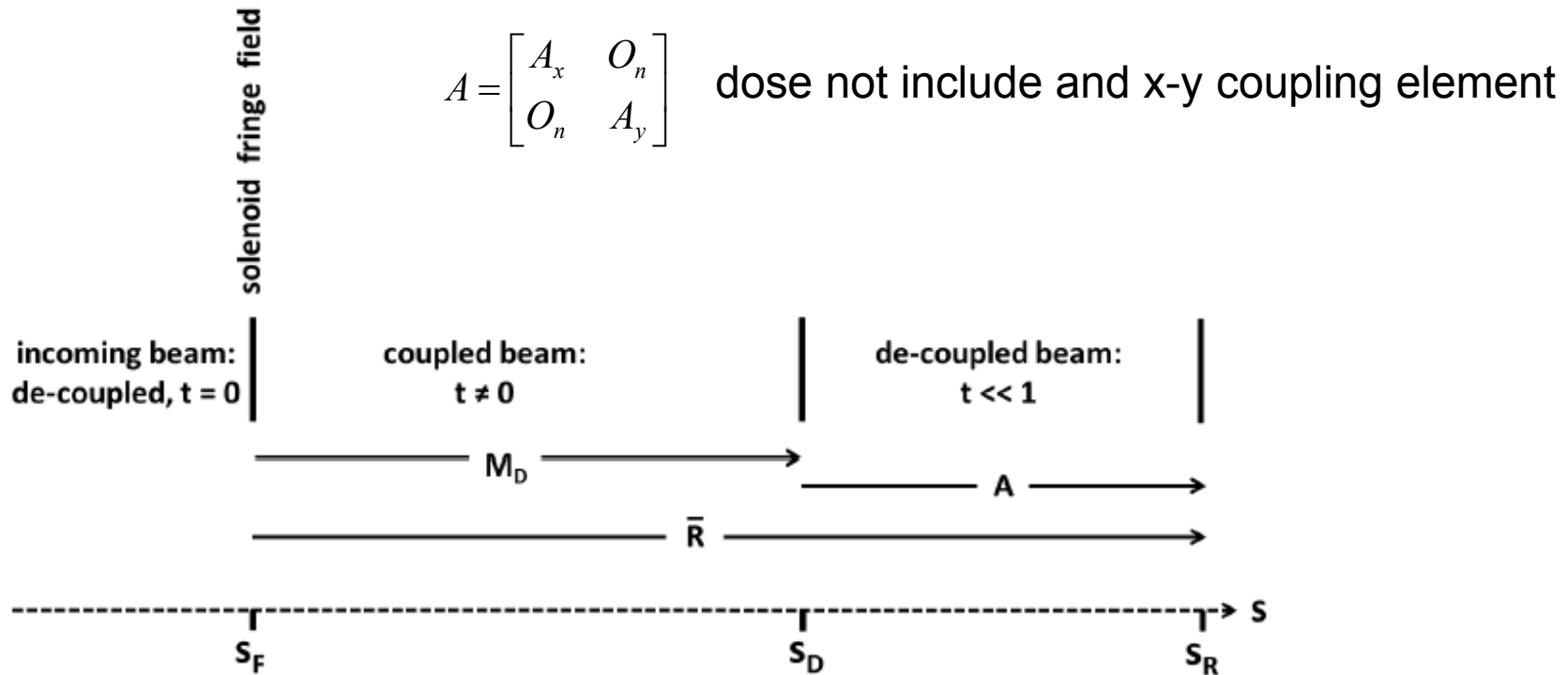


# Properties of the Emtex

- $k_0$  is the assumed fringe strength,  $k_1$  is the fringe strength actually applied for the coupling.  $t \ll 1$  holds over a wide range of  $k_1$ .
- The exit Twiss parameters (beta & alpha) do not depend on the actual fringe strength  $k_1$ .
- The only quantity considerably changed through the fringe strength is the transverse rms emittance partitioning.



# Decoupling in General Case



$\bar{R}$  has the properties  
 $M_D = A^{-1}\bar{R}$   
 $A^{-1}$  does not change them since it is non-coupling

the properties are also intrinsic properties of  $M_D$

See L. Groening, <http://arxiv.org/abs/1403.6962v1>

**My special gratitude to my colleagues at GSI**

**Thank you for attention!**