

# A COMPACT THERMIONIC RF INJECTOR WITH RF BUNCH COMPRESSION FED BY A QUADRUPOLE-FREE MODE LAUNCHER\*

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## Abstract

We present a design for a compact X-Band RF thermionic injector consisting of two iris-loaded accelerator structures. Both structures are fed by a single quadrupole-free  $TM_{01}$  mode launcher. In the first structure the electron bunches are extracted from a thermionic cathode. The second structure creates an energy chirp in the bunch for its further ballistic compression. This injector can produce short electron bunches without the need for a magnetic bunch compressor. We are developing this injector as part of a linac-based 91.392 GHz RF power source, which further comprises a booster linac and a mm-wave decelerator structure that extracts 91.392 GHz RF power from the electron beam. This source will be used to power a short-period RF undulator with 1.75 mm period.

## INTRODUCTION

We propose to build a 13.5 nm extreme ultra violet light source, based on a novel mm-wave undulator operating at 91.392 GHz [1, 2]. For this source we will reuse most of the existing SLAC X-Band (11.424 GHz) test accelerator infrastructure [3]. We will replace the existing photo-gun with a thermionic RF injector. This injector, which is the topic of this paper, consists of a thermionic RF gun, an accelerating structure, and a quadrupole-free  $TM_{01}$  mode launcher feeding both the gun and the accelerator structure. The electron bunches are accelerated in a booster linac downstream of the injector, and RF power at 91.392 GHz is extracted from the beam in a W-Band decelerator structure. This power is then fed into the RF undulator.

The mm-wave undulator is cavity that has a loaded quality factor of approximately 12,700 [1], which corresponds to a filling time of 22 ns. The undulator requires an input RF pulse length that is several times its filling time in order to reach its steady state fields. A decelerator structure at the W-Band has a loaded quality factor of only 1,700, which corresponds to a filling time of approximately 3 ns. Therefore, when driven by a single bunch, the RF pulse generated by a decelerator structure will not have sufficient length to fill the undulator. We therefore decided to develop a multibunch thermionic RF injector, in order to produce a W-Band RF pulse that is long enough to fill the undulator. Additionally we need bunches with a small 6-D emittance because we will compress them to a small fraction of the wavelength at 91.392 GHz and focus them transversely into the small

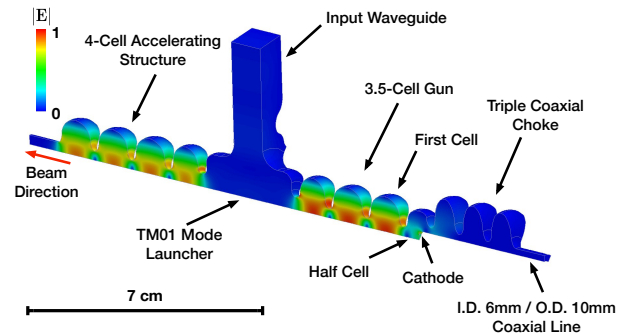


Figure 1: RF injector surface electric field on one eighth of the injector. The length of the structure shown is 18 cm.

aperture of the decelerator structure. We also considered DC injectors with subharmonic bunching. A high-voltage low emittance thermionic DC injector using a small  $CeB_6$  cathode has been developed for a Free Electron Laser [4–6]; however, the fact that it requires a DC voltage of approximately 500 kV and a complex bunching system makes it large and expensive. A low voltage DC injector with subharmonic bunching cannot satisfy the small emittance requirement.

Thermionic RF injectors at the S-Band have been employed in synchrotron light sources [7–10]. Most of these injectors used an alpha magnet for bunch compression. A higher frequency X-Band thermionic RF injector was built by the University of Tokyo for a Compton scattering source [11–17]. This injector also uses an alpha magnet for bunch compression. In this work we report a design for a compact X-Band thermionic RF injector with only RF bunch compression. We decided to use RF bunch compression instead of an alpha magnet due to space constraints in our system. A similar approach has been employed in a photo-gun [18–20].

## DESIGN REQUIREMENTS

We designed the gun for a  $70 \text{ MV m}^{-1}$  peak electric field on the cathode and a  $175 \text{ MV m}^{-1}$  peak on-axis electric field of the other cells of the gun. The gun should produce appreciable current at cathode peak fields ranging from  $60 \text{ MV m}^{-1}$  to  $90 \text{ MV m}^{-1}$ . Due to beamline length constraints, the point of maximum bunch compression has to be within 0.6 m from the cathode. The following requirements must be met at that point: beam energy greater than 4 MeV, RMS bunch length  $\sigma_z$  less than 0.5 ps, transverse RMS bunch size  $\sigma_x, \sigma_y$  less than 1 mm, and bunch charge  $Q$  on the order of 7 pC. This charge corresponds to 80 mA of average current during the RF pulse. The RF power budget for the injector is 10 MW, and the RF pulse length is 250 ns.

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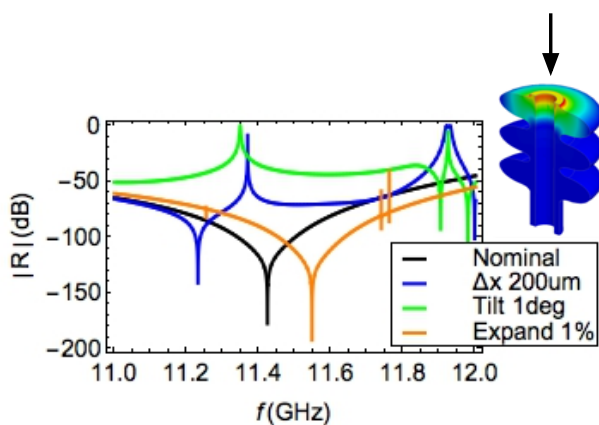


Figure 2: Triple choke reflection coefficient vs frequency for different perturbations of the cathode stem position.

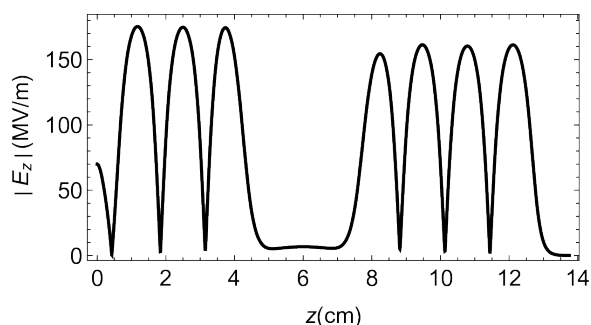


Figure 3: On-axis electric field amplitude profile for 8 MW of input power.

The cathode current density is conservatively assumed to be  $10 \text{ A cm}^{-2}$ .

## RF DESIGN

The injector comprises a 3.5-cell thermionic RF gun, a 4-cell accelerating structure, and a quadrupole free mode launcher feeding both structures through the overmoded beam pipe. Figure 1 shows the circuit configuration and the surface electric field on one eighth of the injector. We plan to use a commercial 0.134 inch diameter flat face dispenser cathode. The cathode is mounted on a 6 mm diameter cathode stem. The gap between the cathode stem and the RF circuit of the gun is chosen to be 2 mm in order to reduce

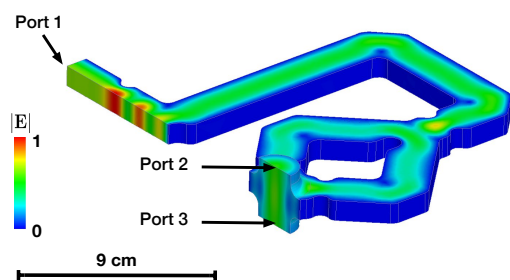


Figure 4: Quadrupole-free mode launcher half.

the possibility of RF breakdowns. The cathode half cell is formed using coaxial chokes with a 10 mm inner diameter. The cathode stem forms a coaxial line. Potential misalignment of the cathode stem could couple power to the  $\text{TE}_{11}$  in that coaxial line and result in power leakage. Therefore the dimensions of the coaxial line were chosen such that the  $\text{TE}_{11}$  cutoff frequency (11.9281 GHz) is above the operating frequency 11.424 GHz.

The shapes of the chokes and the accelerator cells are elliptical to reduce the probability of multipacting. We found that one or two chokes are not sufficient to stop RF leakage in the event of the cathode stem being misaligned. We therefore used three chokes in succession. Figure 2 shows the reflection coefficient of the triple choke when the stem is at its nominal position, when the stem is misplaced by  $200 \mu\text{m}$ , when it is tilted by  $1^\circ$ , and when it is expanded by 1% or  $30 \mu\text{m}$  in radius. We found that the choke performance is most sensitive to cathode stem tilt.

The length of the half cell and of the first full cell of the gun – shown in Fig. 1 – were optimized for maximum *brightness times bunch charge* for 20% of the bunch particles. This parameter was calculated using beam dynamics simulation in Astra [21]. For the optimization we run Astra without space charge. The beam dynamics simulation setup and post-processing are described in the next section. The resulting cell lengths were  $0.5\lambda/2$  for the half cell, and  $0.96\lambda/2$  for the first cell, where  $\lambda = 2.62 \text{ cm}$  is the wavelength. The gun and the accelerating structure were then assembled with the mode launcher, as shown in Fig. 1. The distance between the structures was chosen using Astra. Figure 3 shows the on-axis electric field profile in the injector for 8 MW of input power. The peak surface electric field on the copper is  $190 \text{ MV m}^{-1}$ , and the peak surface magnetic field is  $310 \text{ kA m}^{-1}$ . From [22] the peak pulsed surface heating for  $1 \mu\text{s}$  RF pulses is  $41.3^\circ\text{C}$ , which is considered safe for copper [23].

Both the gun and the accelerator structure are fed through a  $\text{TM}_{01}$  mode launcher. This mode launcher comprises four rectangular  $\text{TE}_{10}$  to cylindrical  $\text{TM}_{01}$  power dividers and a power combining network. Figure 4 shows one half of

Table 1: Bunch Parameters at the Point of Maximum Compression  $z = 65.78 \text{ cm}$

	80%	50%	20%	Spec.
$E$ (MeV)	4.596	4.506	4.486	$> 4$
$Q$ (pC)	6.322	3.900	1.485	$\sim 7$
$I_{\text{bunch}}$ (A)	4.895	41.14	56.43	
$B_n$ ( $\frac{\text{A}}{\text{mm}^2 \text{mrad}^2}$ )	0.302	17.44	204.9	
$\sigma_x$ ( $\mu\text{m}$ )	576.8	189.6	69.38	$< 1000$
$\sigma_y$ ( $\mu\text{m}$ )	578.5	189.8	69.13	$< 1000$
$\sigma_z$ (fs)	322.8	23.70	6.578	$< 500$
$\epsilon_{x,n}$ (mm · mrad)	5.674	2.167	0.747	
$\epsilon_{y,n}$ (mm · mrad)	5.707	2.176	0.737	
$\Delta E_{\text{rms}}$ (%)	3.498	1.012	0.4398	

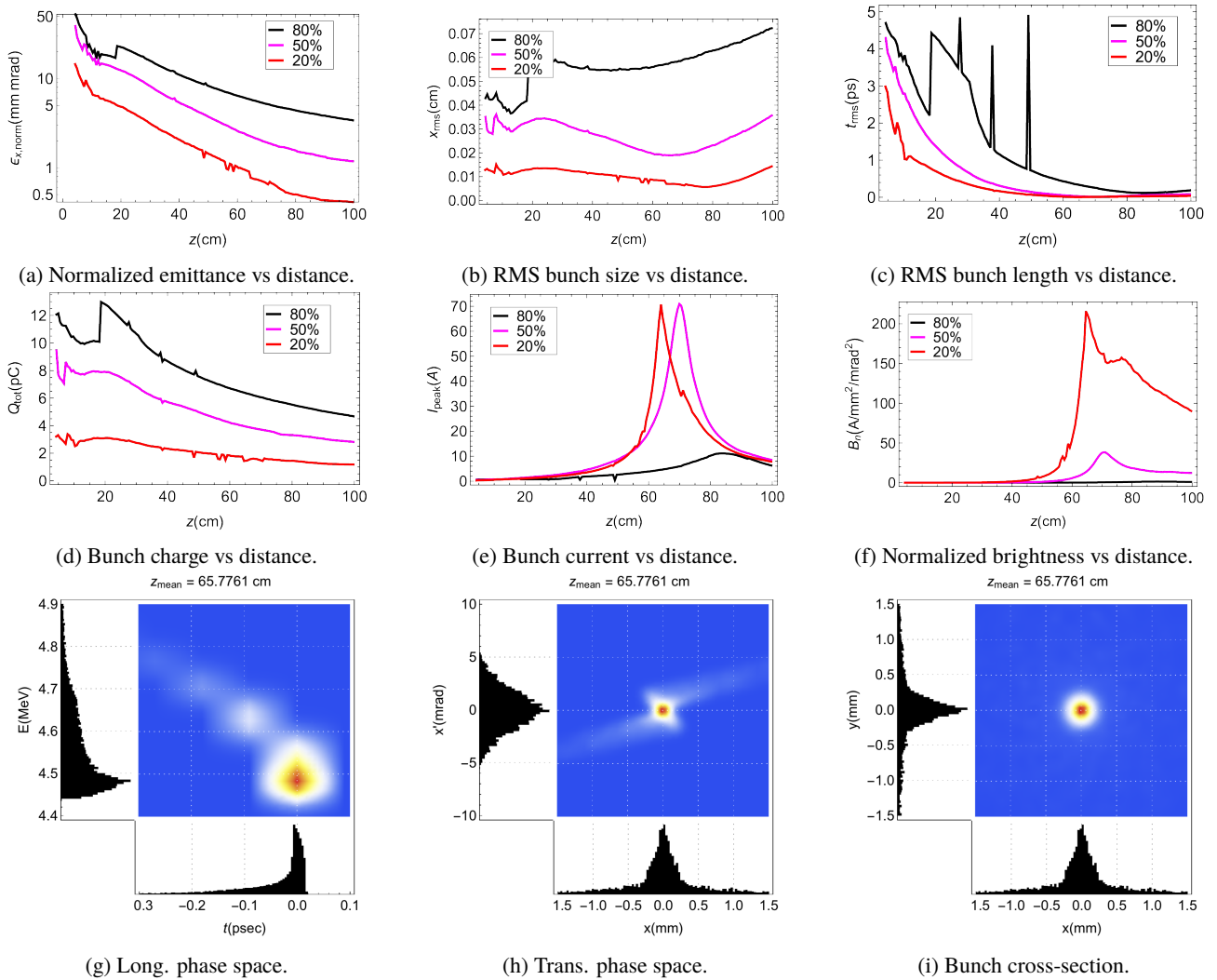


Figure 5: Beam dynamics simulation results and phase space at the point of maximum compression.

the mode launcher. The four feeds around the cylindrical waveguide do not excite either dipole or quadrupole modes.

## BEAM DYNAMICS

The beam input to Astra had a uniform in time distribution, equivalent to three RF buckets. We used the cylindrical space charge model in Astra. The beam pipe radius of the injector is 3 mm. Particles beyond that radius were discarded. To calculate single bunch parameters out of a multi-bunch phase space, the phase space was exported every 40 ps and was post-processed in Mathematica. Bunches were extracted from the phase space, and the parameters of the middle bunch were calculated for different fractions of the bunch. That fraction was calculated as follows: particles were iteratively filtered out of the particle array if they were more than two standard deviations away in energy or position from the particle array mean, until the number of remaining particles in the array was the required percentage of the original. The bunch current is calculated as  $I_{bunch} = Q_{tot} / (4\sigma_z)$ . The normalized brightness is defined as  $B_n = 2I_{bunch} / (\epsilon_{x,n} \cdot \epsilon_{y,n})$ ,

where  $\epsilon_{x,n}$  and  $\epsilon_{y,n}$  are the normalized transverse emittance in the  $x$  and  $y$  directions. Figure 5 shows the bunch parameters versus distance from the cathode and phase space at  $z = 65.78$  cm. Table 1 summarizes the bunch parameters at this  $z$  coordinate. As shown in Fig. 5d, the bunch charge drops with  $z$  because a significant part of the beam is diverging due to RF focusing and is lost in the beam pipe.

## CONCLUSION

We have designed a compact X-Band thermionic RF injector, consisting of a 3.5-cell gun and a 4-cell accelerating structure, fed by a single  $TM_{01}$  mode launcher. This injector generates a train of electron bunches spaced 87.5 ps apart. The RF gun is designed to operate with  $70 \text{ MV m}^{-1}$  peak electric field on the cathode and approximately  $175 \text{ MV m}^{-1}$  peak on-axis field for the rest of the cells. The parameters of 80% of the bunch are  $Q = 6.3 \text{ pC}$ ,  $\sigma_{x/y} = 578 \text{ }\mu\text{m}$ ,  $\sigma_z = 323 \text{ fs}$ , and  $\epsilon_{x/y,n} = 5.7 \text{ mm mrad}$ . The parameters of 20% of the bunch are  $Q = 1.5 \text{ pC}$ ,  $\sigma_{x/y} = 69 \text{ }\mu\text{m}$ ,  $\sigma_z = 6.6 \text{ fs}$ , and  $\epsilon_{x/y,n} = 0.75 \text{ mm mrad}$ .

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