

SYNCHROTRON OSCILLATION DAMPING DUE TO BEAM-BEAM COLLISIONS

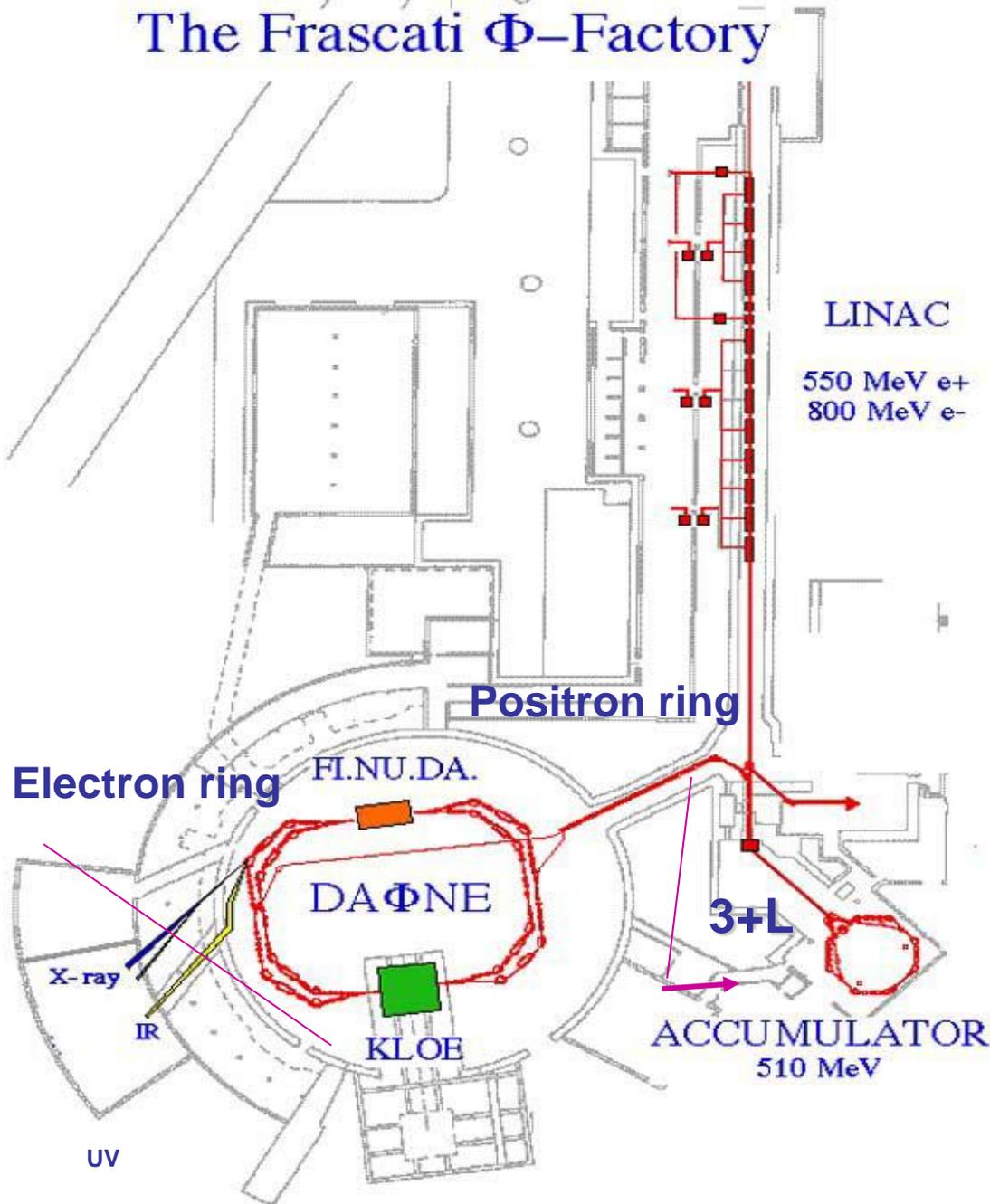
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Introduction

- In DAFNE, the Frascati e⁺/e⁻ collider, the *crab waist* collision scheme has been successfully implemented in 2008 and 2009, for the Siddharta experiment.
- During the collision operations, an unusual synchrotron damping effect has been observed.
- Switching off the longitudinal feedback and having beam currents in the order of 200-300 mA, the positron beam of course becomes unstable
- Nevertheless the longitudinal instability is damped by bringing the positron beam in collision with a high current electron beam (~2A).
- Besides, doing this, we have observed a shift of $\approx -600\text{Hz}$ in the residual synchrotron sidebands.

The Frascati Φ -Factory



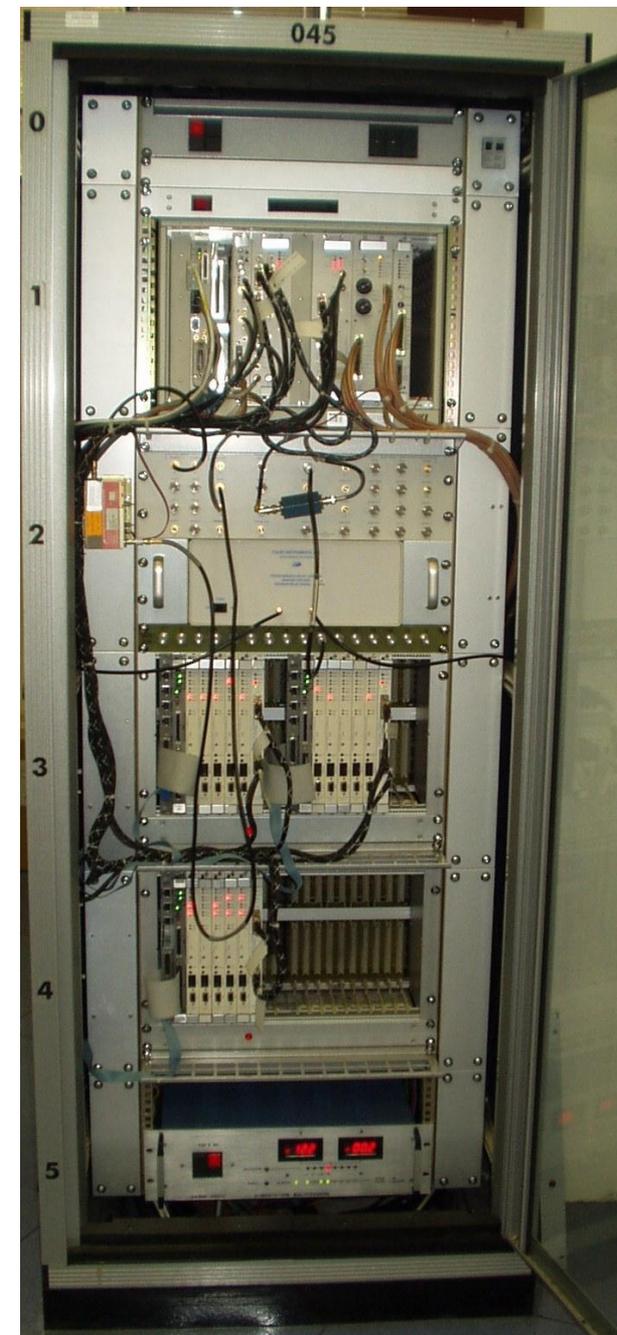
- LINAC
- Transfer lines
- Accumulator ring
- 2 main rings
- 2 interaction points
- Multi-detectors:
in the last run,
Siddhartha

Instrumentation used

- Precise measurements on this effect have been performed by using two different ways:



- a commercial Real-time Spectrum Analyzer RSA 3303 by Tektronix
- the diagnostics capabilities of the DAFNE longitudinal bunch-by-bunch feedback (developed in collaboration with SLAC and LBNL in 1993-96)



Acquisitions from the Real-time Spectrum Analyzer RSA 3303

- Transverse and longitudinal pickup made by 4 high frequency buttons
- H9 (MA-COM) hybrids to have difference in H & V
- ~30 m. low attenuation cable
- Bandpass filter (@ 360MHz, +/- 5MHz)
- Amplifier
- Spectrum analyzer

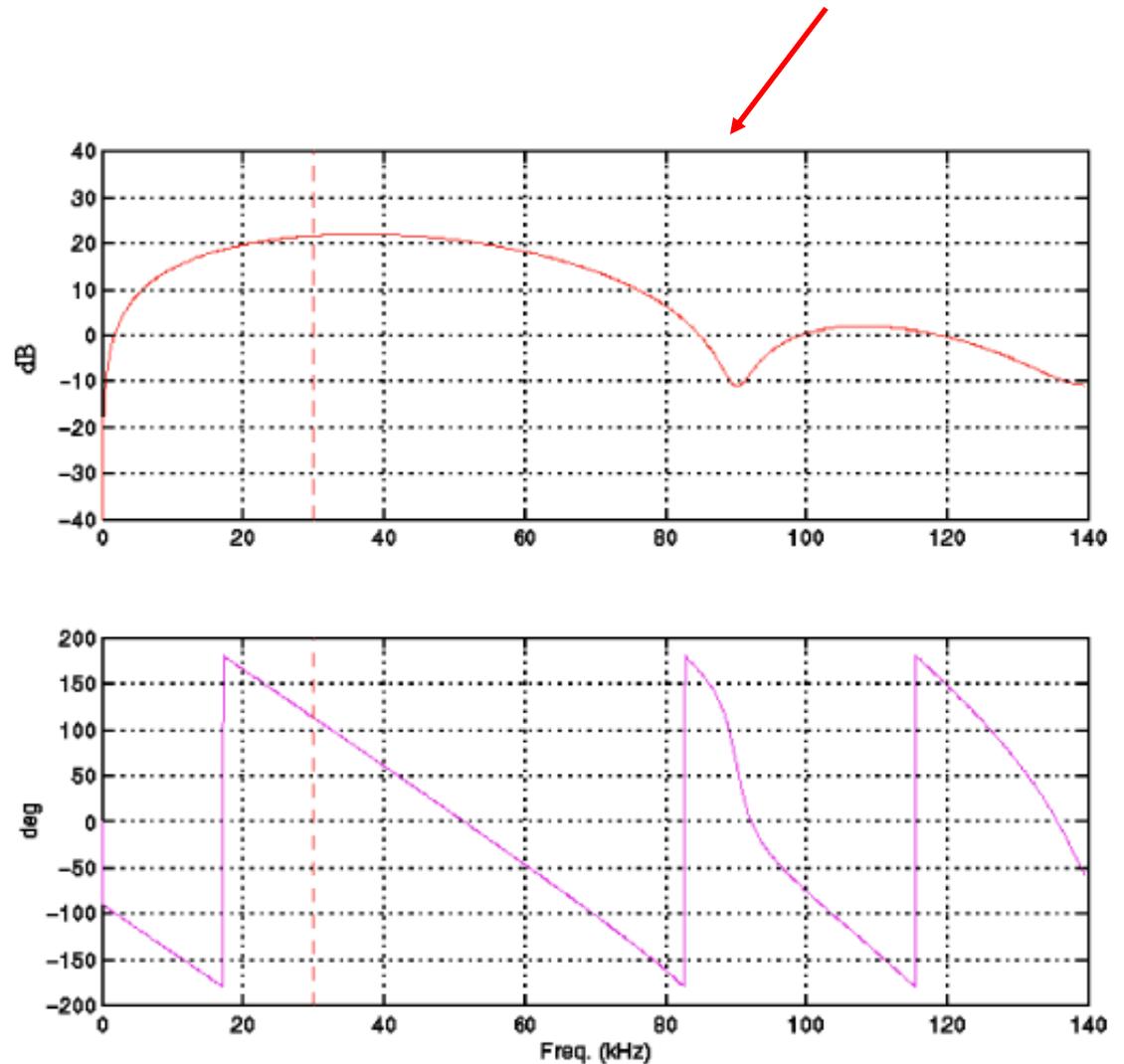




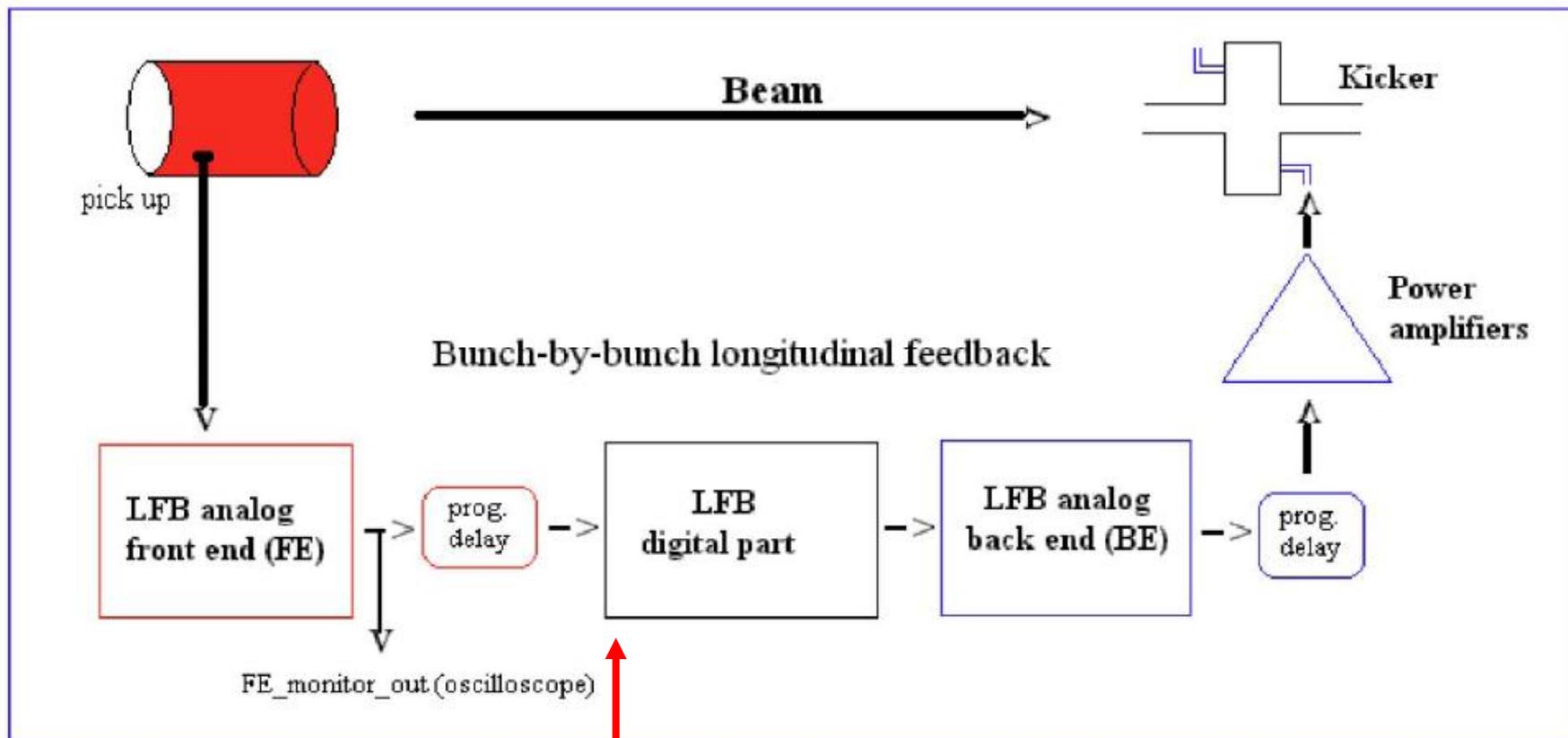
Long.Feedback used at ALS, PEP-II, DAFNE, Bessy, Pohang

Implemented by 60 real DSP's at DAFNE

Every DSP implement a FIR function for each bunch

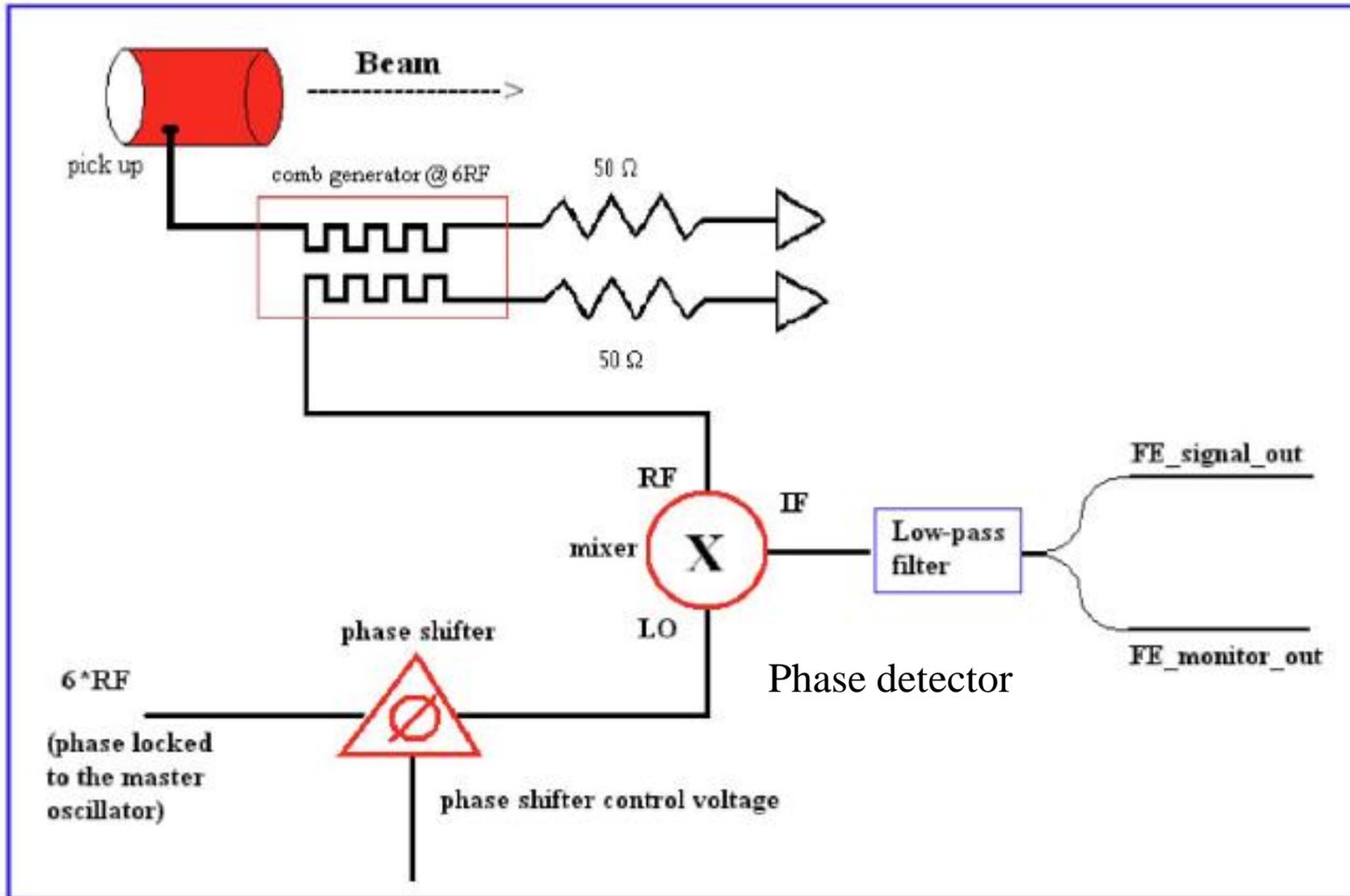


Block diagram of the longitudinal feedback system



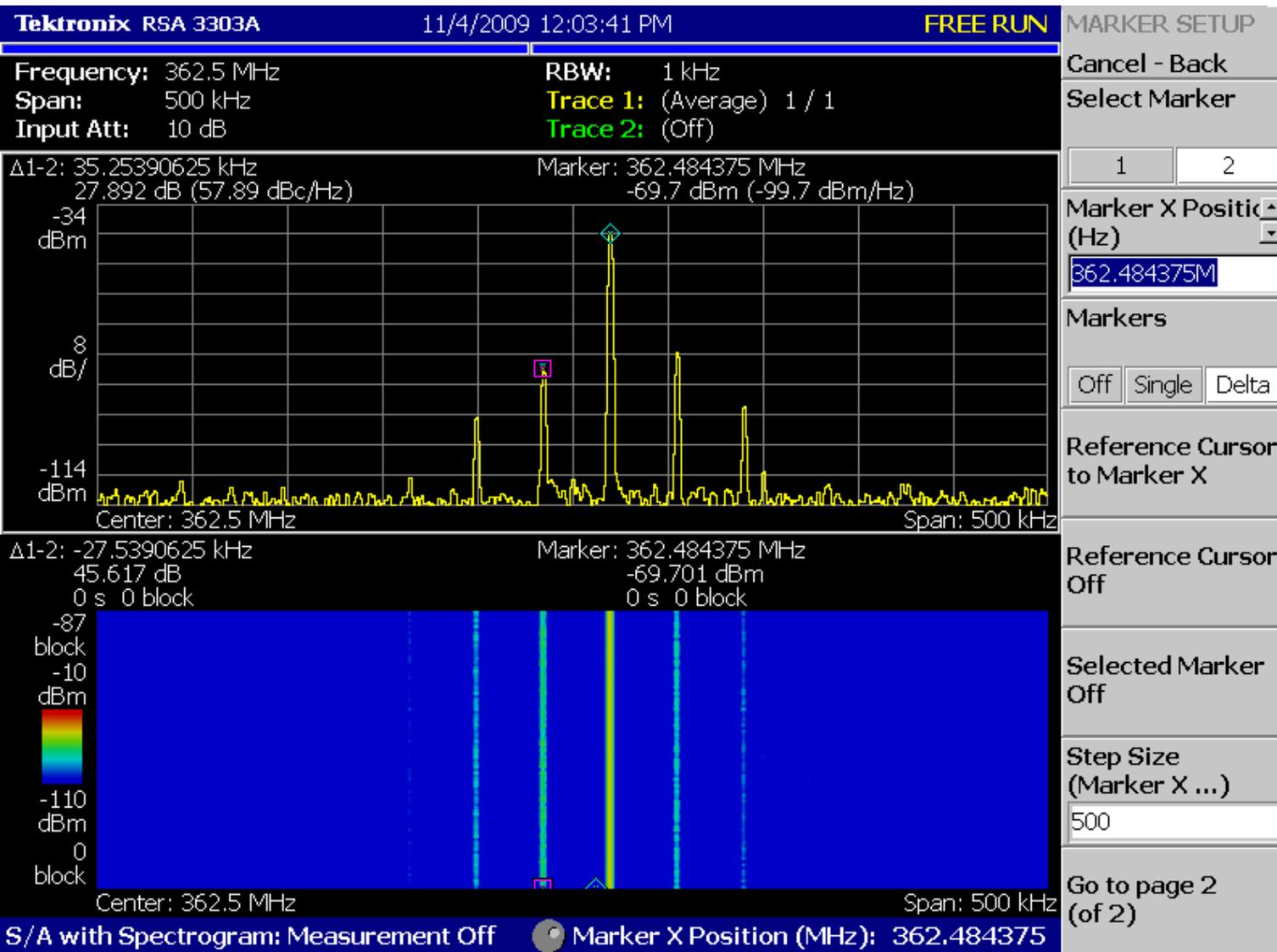
Bunch-by-bunch record

longitudinal feedback front-end (simplified scheme)



Longitudinal sidebands in e+ beam

(set of measurements recorded at DAFNE in 11/2009)



I+ = 200mA

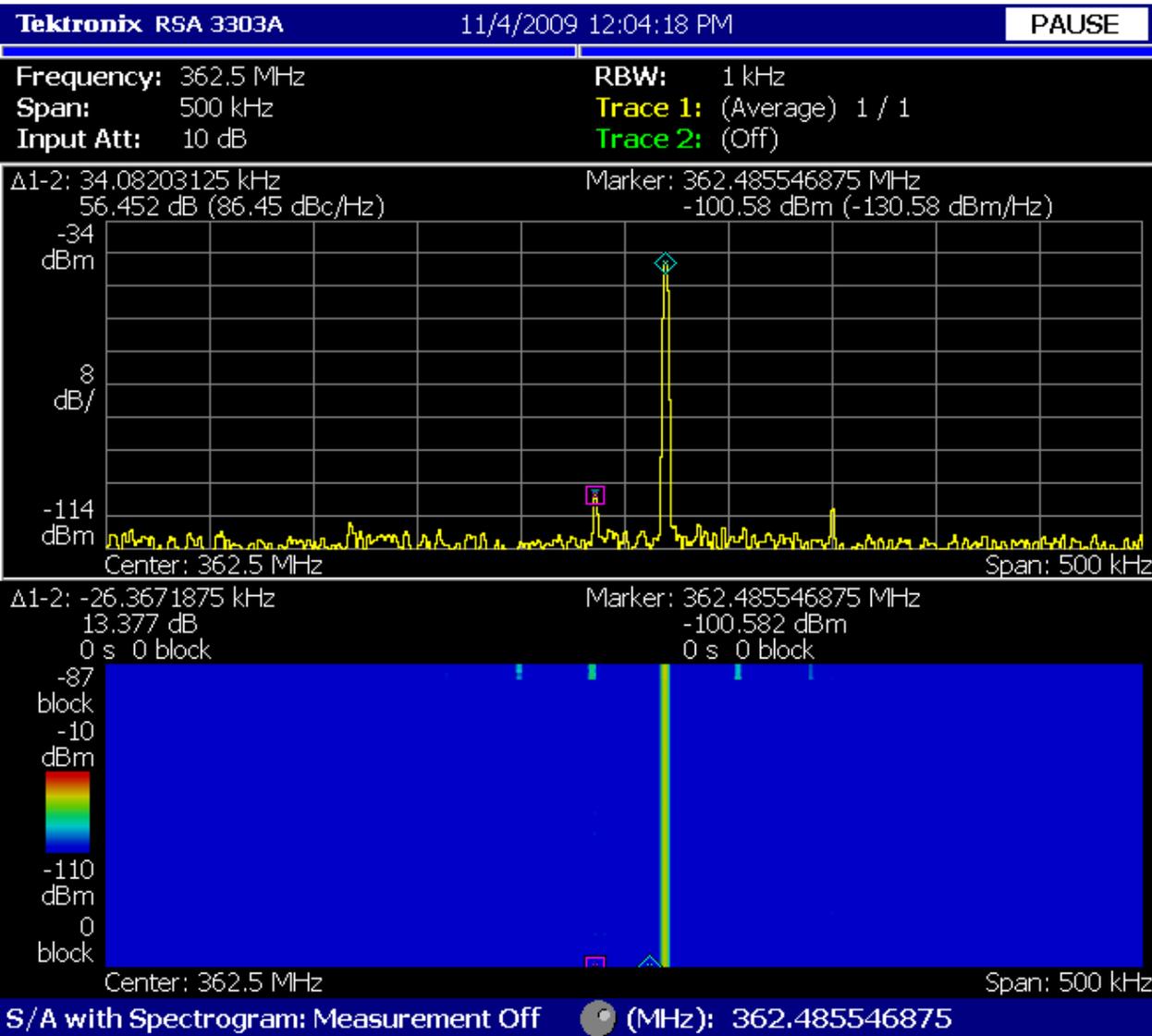
e+ longitudinal feedback off

No collisions at the IP

The highest peak is the 118th harmonic

Other peaks: longitudinal sidebands

Longitudinal sidebands in e+ beam in collision with e- beam



MARKER SETUP

Cancel - Back

Select Marker

1 2

Marker X Position (Hz)
 362.485546875M

Markers
 Off Single Delta

Reference Cursor to Marker X

Reference Cursor Off

Selected Marker Off

Step Size (Marker X ...)
 500

Go to page 2 (of 2)

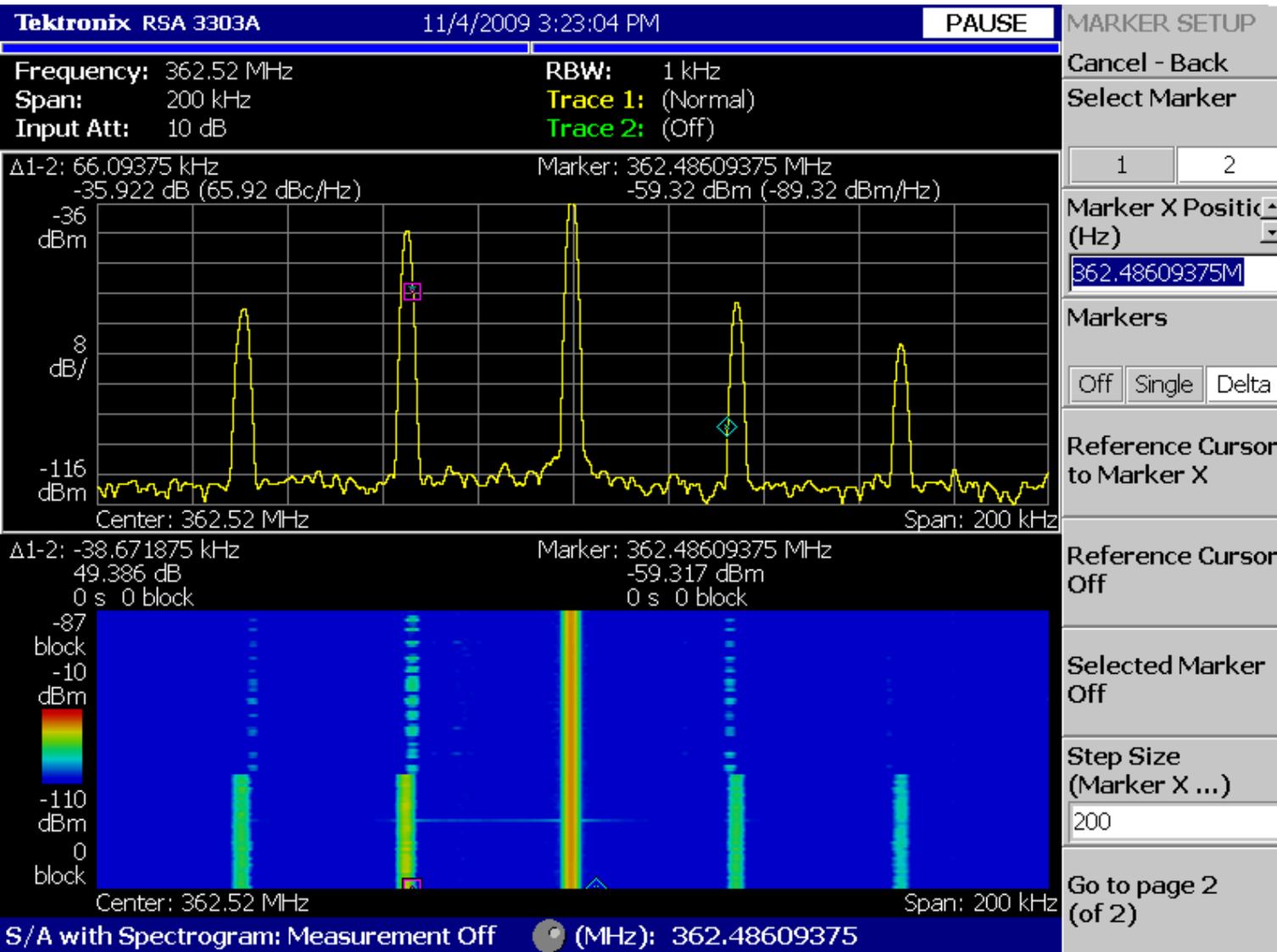
I+beam = 200mA

Longitudinal feedback off

collisions at the IP with the e- beam

Recorded by Tektronix S.A. 3303

In collision $\leftarrow \rightarrow$ out of collision

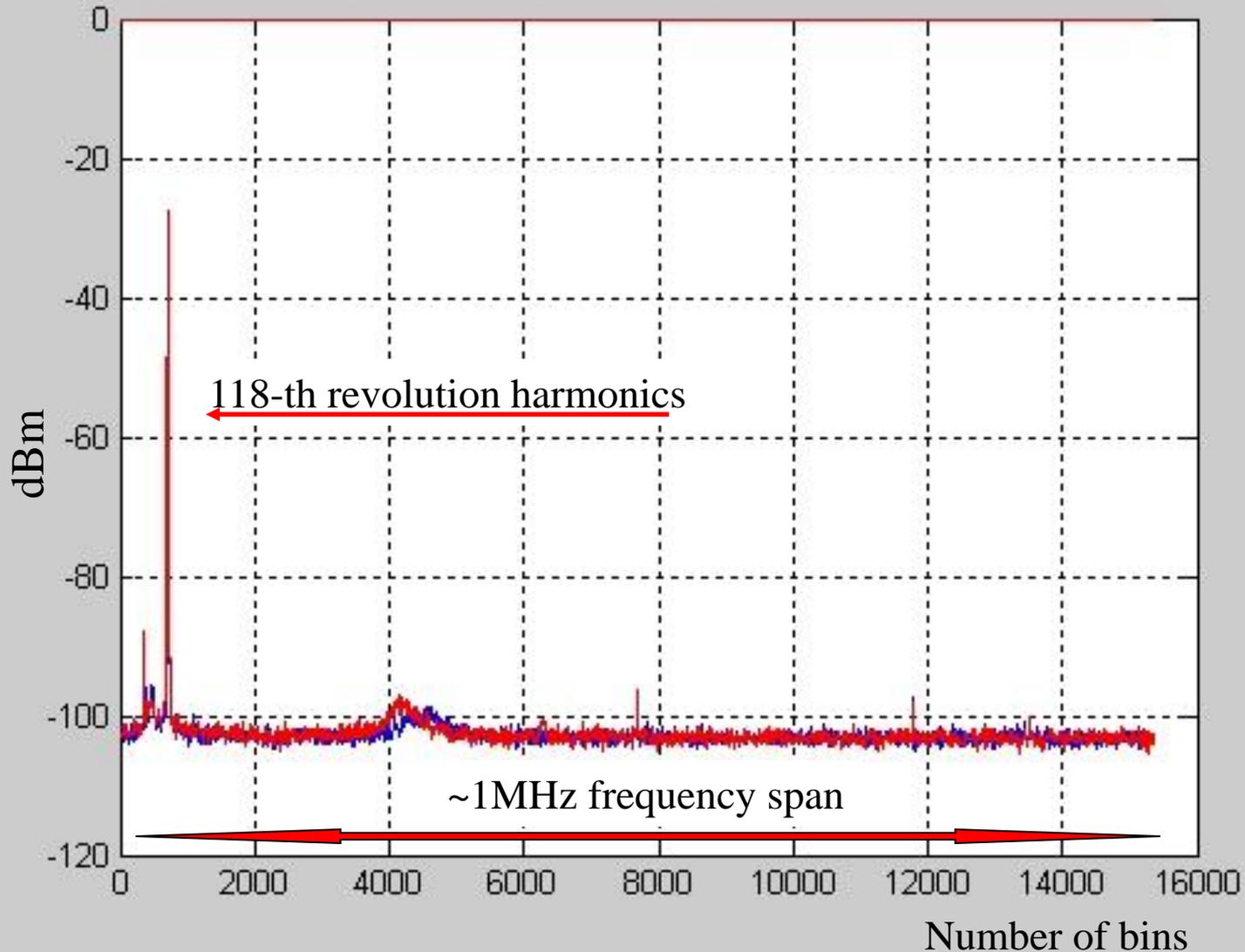


← Last record

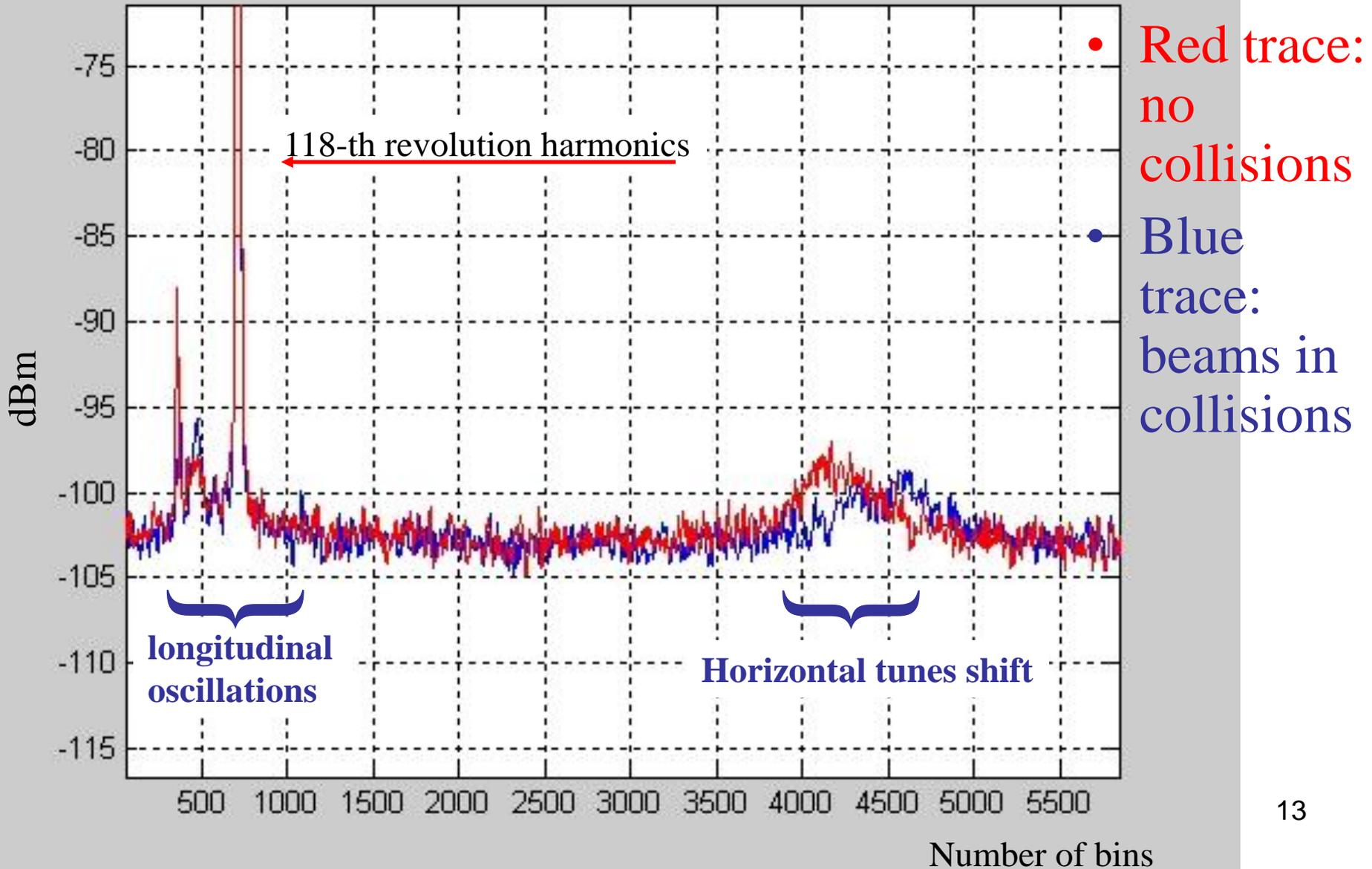
← Last 60 sec. story

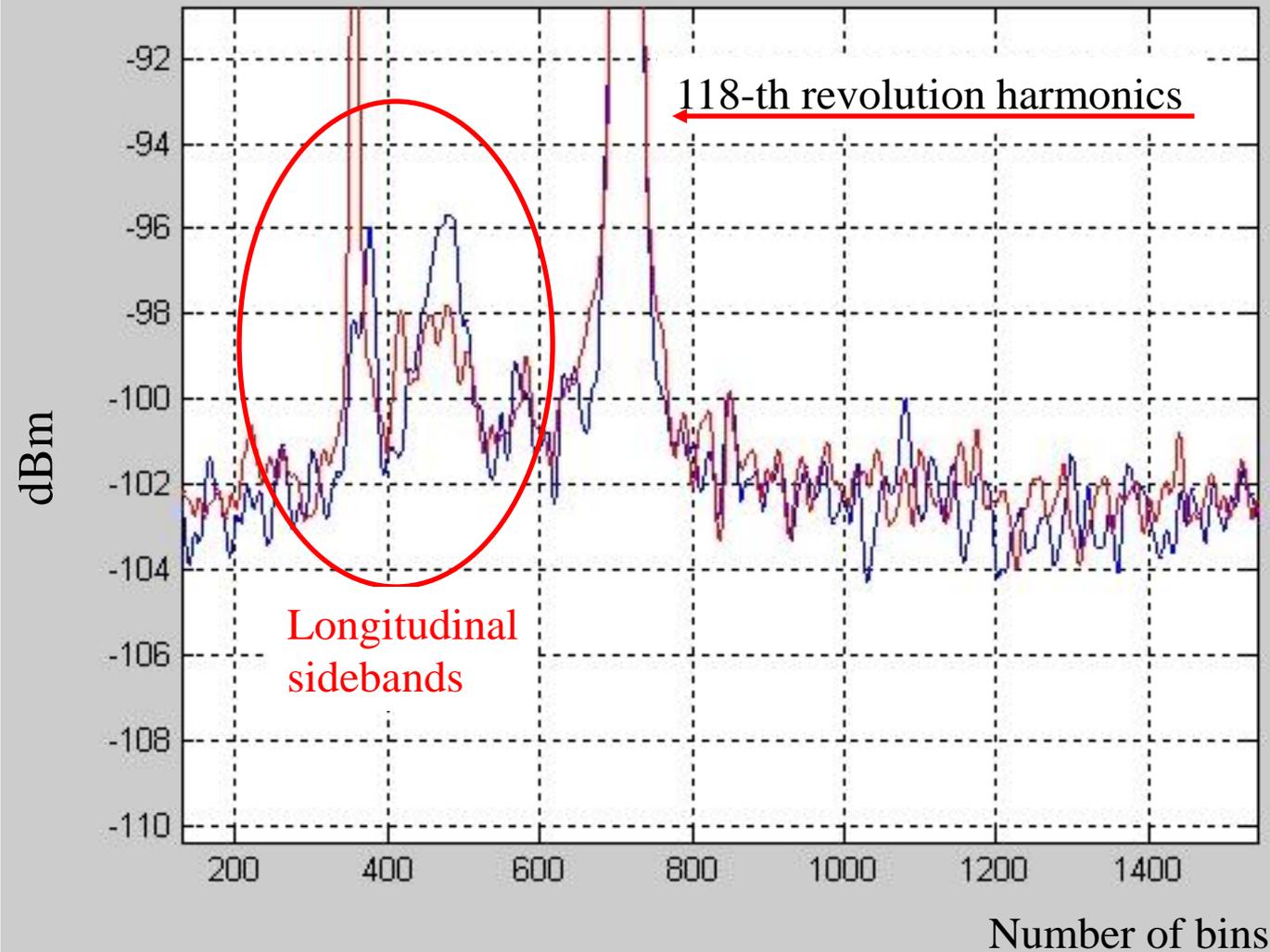
← Tune shift

Data downloaded from the spectrum analyzer RSA3303 to a PC.
The highest peak is the e+ 118-th harmonics. Data are elaborated by MATLAB:
in red beams out of collision, **in blue** beams in collision.

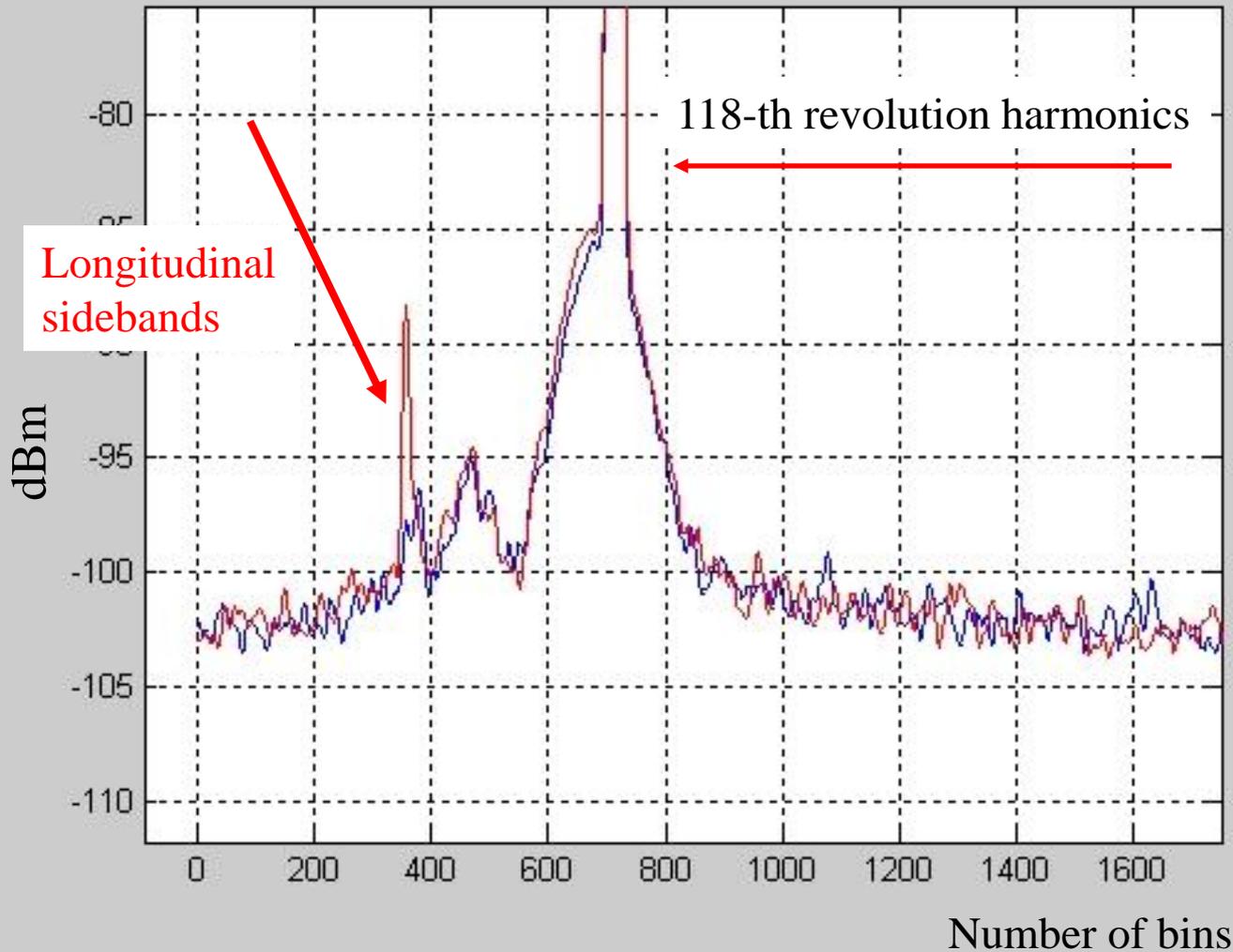


Zoom of the previous figure, showing the longitudinal and horizontal tunes





Zoom of the previous figure, particular of the longitudinal sidebands



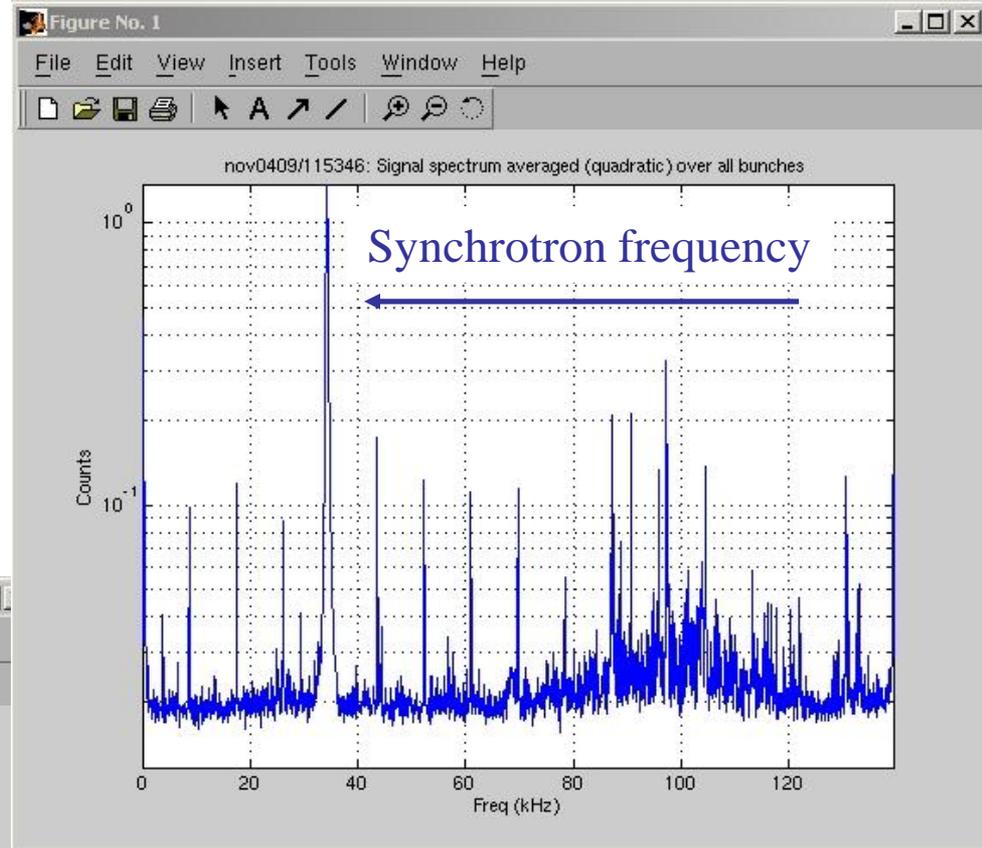
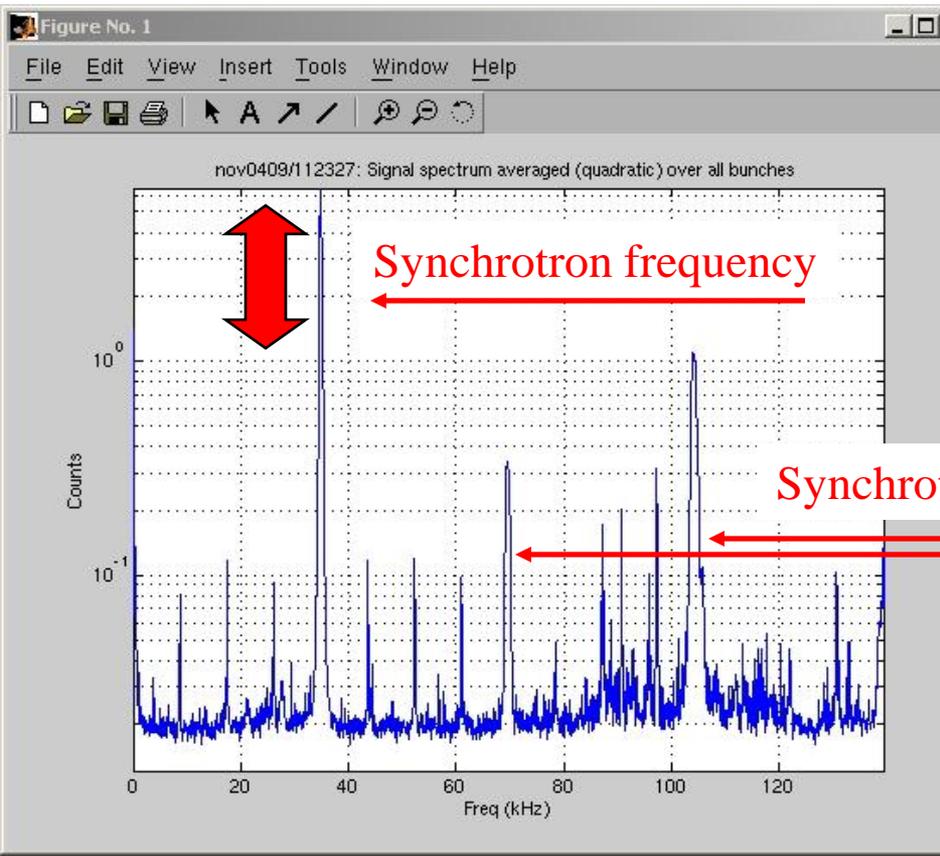
Another case, similar to the previous figure, particular of the longitudinal sidebands.

Diagnostics by the DAFNE longitudinal bunch-by-bunch

feedback:

front-end data plotted in frequency domain (all bunches)

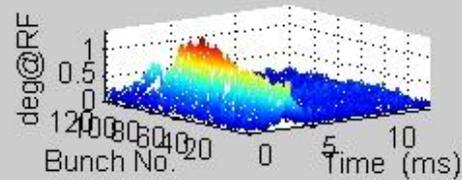
Out of collision



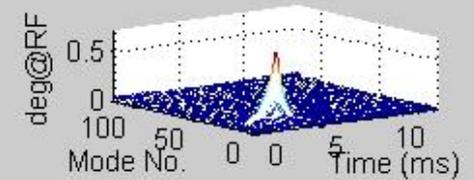
In collision

In collision e+ long. modal analysis

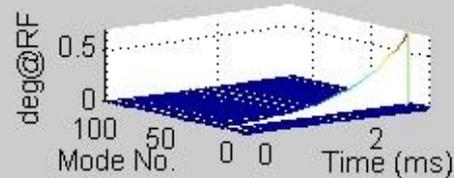
a) Osc. Envelopes in Time Domain



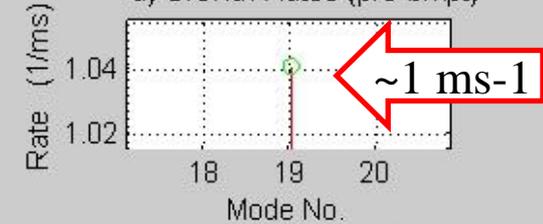
b) Evolution of Modes



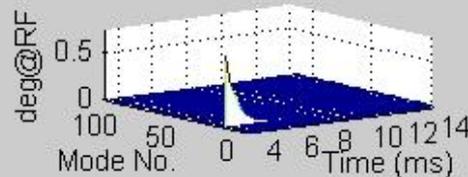
c) Exp. Fit to Modes (pre-brkpt)



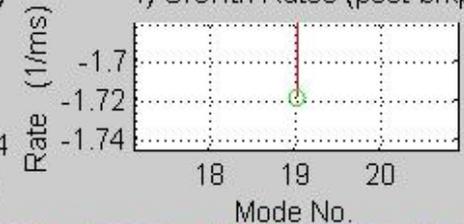
d) Growth Rates (pre-brkpt)



e) Exp. Fit to Modes (post-brkpt)



f) Growth Rates (post-brkpt)

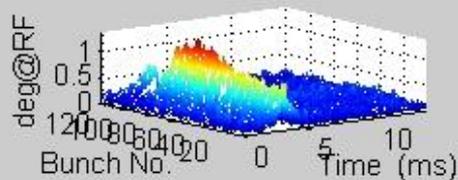


DAFNE E+: nov0509\145636: Io= 350mA, Dsamp= 11, ShifGain= 3, Nbn= 12
Gain1= 1 Gain2= 0 Phase1= 220 Phase2= 40 Brkt= 945 Calib= 0.6362

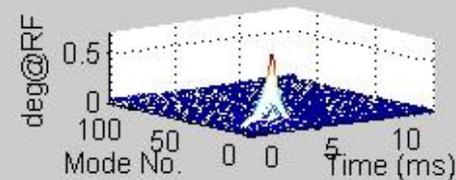
In collision e+ long. modal analysis

Out of collision longitudinal modal analysis

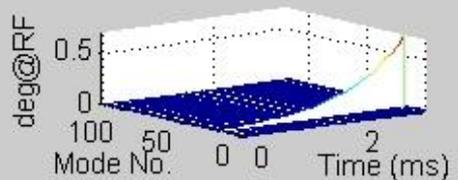
a) Osc. Envelopes in Time Domain



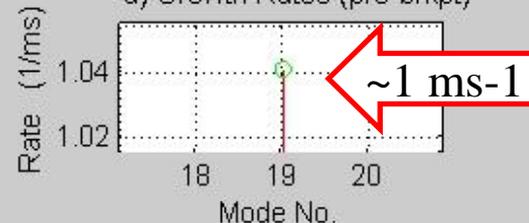
b) Evolution of Modes



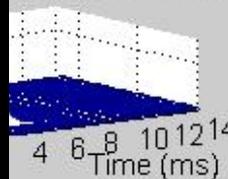
c) Exp. Fit to Modes (pre-brkpt)



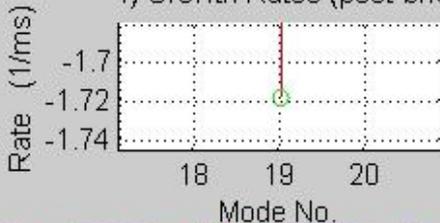
d) Growth Rates (pre-brkpt)



e) Exp. Fit to Modes (post-brkpt)

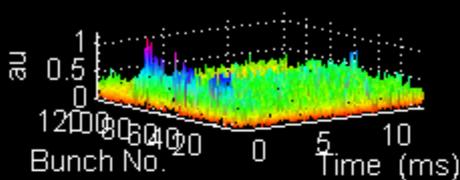


f) Growth Rates (post-brkpt)

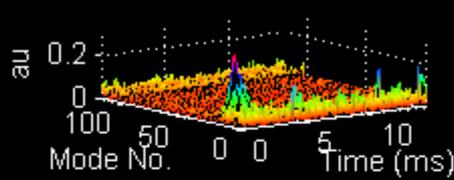


45636: Io= 350mA, Dsamp= 11, ShifGain= 3, Nbn= 12
Phase1= 220 Phase2= 40 Brkpt= 945 Calib= 0.6362

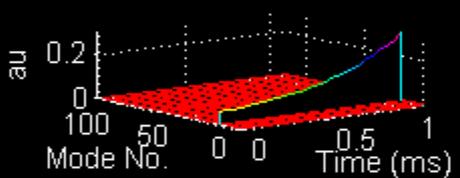
a) Osc. Envelopes in Time Domain



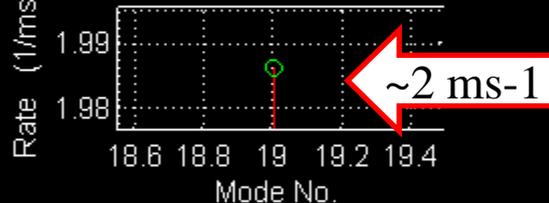
b) Evolution of Modes



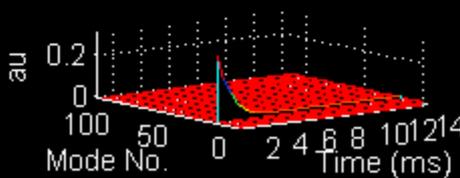
c) Exp. Fit to Modes (pre-brkpt)



d) Growth Rates (pre-brkpt)



e) Exp. Fit to Modes (post-brkpt)



f) Growth Rates (post-brkpt)



Analysis tools by S.Prabhakar,
PHD thesis, SLAC-R-554,
August 2001

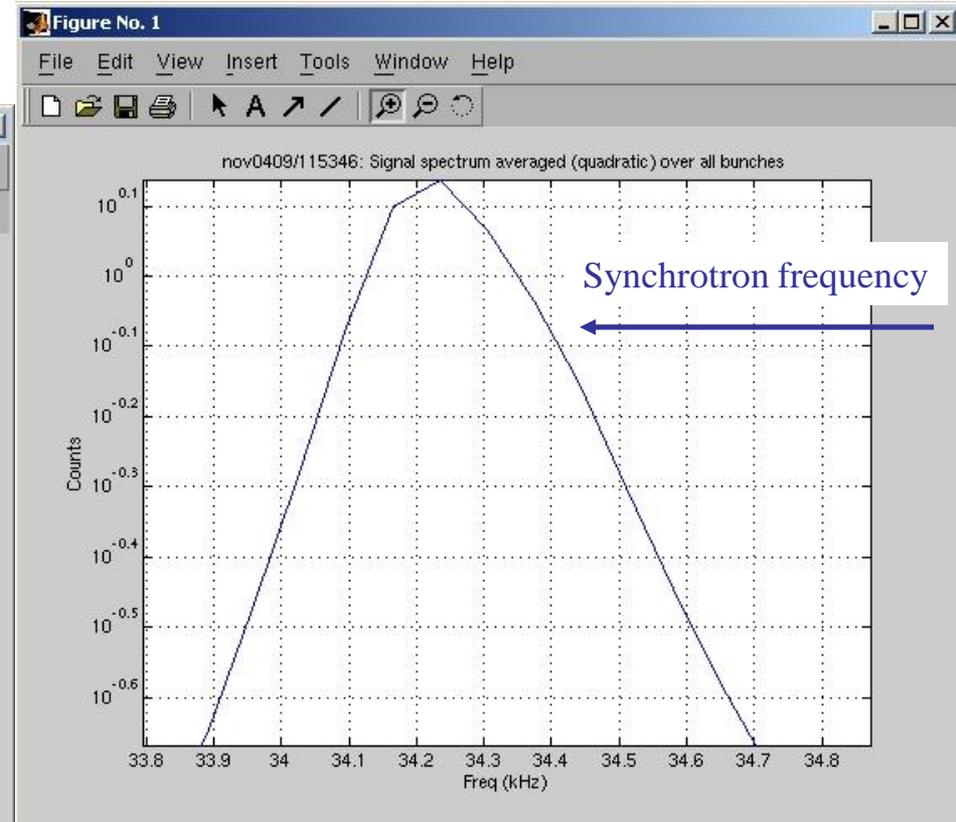
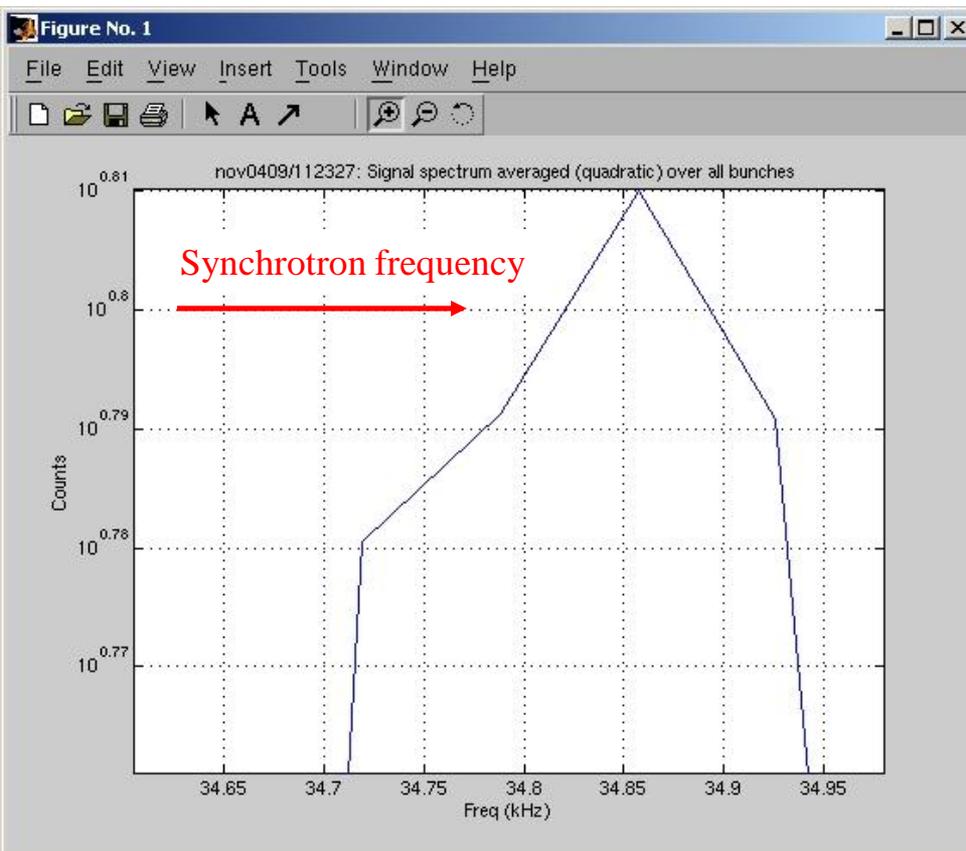
DAFNE E+: nov0409\151603: Io= 400mA, Dsamp= 11, ShifGain= 3, Nbn= 1
Gain1= 1, Gain2= 0, Phase1= 220, Phase2= 40, Brkpt= 305, Calib= 0.636

Signal spectrum averaged from all the bunch [data recorded by longitudinal feedback]

Difference= \sim -630Hz

Out of collision:

$f_{\text{sync}}=34.86$ kHz



In collision:

$f_{\text{sync}}=34.23$ kHz

Comments on the experimental data (1/2)

- This damping effect has been observed in DAFNE for the first time during collisions with the crab waist scheme.
- Our explanation is that beam collisions with a large crossing angle produce a longitudinal tune shift and a longitudinal tune spread, providing Landau damping of synchrotron oscillations.

Comments on the experimental data (2/2)

- Experimental observations and measurements at DAFNE have shown that beam-beam collisions can damp the longitudinal coupled bunch instability.
 - 1) Bringing into collisions a high current electron beam with an unstable positron one was stabilizing the synchrotron oscillations of the e+ beam, even with the longitudinal feedback system switched off.
 - 2) Besides, a negative frequency shift of positron beam synchrotron sidebands has been observed when colliding the beams.
- We attribute these two effects to a nonlinear longitudinal kick arising due to beam-beam interaction under a finite crossing angle.
- It is worthwhile to note here that we have observed this effect clearly only after implementation of the crab waist scheme of beam-beam collisions at DAFNE having twice larger horizontal crossing angle with respect to the previous operations with the standard collision scheme
- In the following, we show an analytical expression for the synchrotron tune shift, that is also a measure of the synchrotron tune spread, and compare the formula with numerical simulations.

Beam-beam kick formulae

$$x' = \frac{2r_e N}{\gamma} (x - z \operatorname{tg}(\theta/2)) \int_0^\infty dw \frac{\exp\left\{ -\frac{(x - z \operatorname{tg}(\theta/2))^2}{(2(\sigma_x^2 + \sigma_z^2 \operatorname{tg}^2(\theta/2)) + w)} - \frac{y^2}{(2\sigma_y^2 + w)} \right\}}{(2(\sigma_x^2 + \sigma_z^2 \operatorname{tg}^2(\theta/2)) + w)^{3/2} (2\sigma_y^2 + w)^{1/2}}$$

$$y' = \frac{2r_e N}{\gamma} y \int_0^\infty dw \frac{\exp\left\{ -\frac{(x - z \operatorname{tg}(\theta/2))^2}{(2(\sigma_x^2 + \sigma_z^2 \operatorname{tg}^2(\theta/2)) + w)} - \frac{y^2}{(2\sigma_y^2 + w)} \right\}}{(2(\sigma_x^2 + \sigma_z^2 \operatorname{tg}^2(\theta/2)) + w)^{1/2} (2\sigma_y^2 + w)^{3/2}}$$

$$z' = x' \operatorname{tg}(\theta/2)$$

In collisions with a crossing angle, the longitudinal kick of a particle is given by the projection of the transverse electromagnetic fields of the opposite beam onto the longitudinal axis of the particle itself. The kicks that the test particle receives while passing the strong beam with rms sizes σ_x , σ_y , σ_z under a horizontal crossing angle θ are in the above formulae. x , y , z are the horizontal, vertical and longitudinal deviations from the synchronous particle travelling on-axis, respectively.

N is the number of particles in the strong bunch, γ is the relativistic factor of weak beam

For the on-axis test particle ($x = y = 0$) the longitudinal kick is given by

$$z' = -\frac{2r_e N}{\gamma} z t g^2(\theta/2) \int_0^\infty dw \frac{\exp\left\{-\frac{(z t g(\theta/2))^2}{(2(\sigma_x^2 + \sigma_z^2 t g^2(\theta/2)) + w)}\right\}}{(2(\sigma_x^2 + \sigma_z^2 t g^2(\theta/2)) + w)^{3/2} (2\sigma_y^2 + w)^{1/2}}$$

For small synchrotron oscillations $z \ll \sigma_z$ the exponential factor in the integral can be approximated by 1 and at the end the formula becomes:

$$z' = -\frac{2r_e N}{\gamma} z t g^2(\theta/2) \frac{1}{(\sigma_x^2 + \sigma_z^2 t g^2(\theta/2)) + \sqrt{(\sigma_x^2 + \sigma_z^2 t g^2(\theta/2))\sigma_y^2}}$$

Then, analogously to the transverse cases, we can write the expression for the synchrotron tune shift:

$$\xi_z = -\frac{r_e N}{2\pi\gamma} \beta_z \frac{tg^2(\theta/2)}{\left((\sigma_x^2 + \sigma_z^2 tg^2(\theta/2)) + \sigma_y^2 \sqrt{(\sigma_x^2 + \sigma_z^2 tg^2(\theta/2))} \right)}$$

Remembering that the longitudinal beta function can be written as:

$$\beta_z = \frac{c|\eta|}{v_{z0}\omega_0} = \frac{\sigma_{z0}}{(\sigma_E/E)_0}$$

with c being the velocity of light; η the slippage factor, v_{z0} the unperturbed synchrotron frequency and ω_0 the angular revolution frequency, we obtain the final expression for the linear tune shift:

$$\xi_z = -\frac{r_e N^{strong}}{2\pi\gamma^{weak}} \frac{\left(\frac{\sigma_{z0}}{(\sigma_E/E)} \right)^{weak} tg^2(\theta/2)}{\left((\sigma_x^2 + \sigma_z^2 tg^2(\theta/2)) + \sigma_y^2 \sqrt{(\sigma_x^2 + \sigma_z^2 tg^2(\theta/2))} \right)^{strong}}$$

For the case of flat beams with $\left(\sigma_y \ll \sqrt{\sigma_x^2 + \sigma_z^2 \operatorname{tg}^2(\theta/2)} \right)$

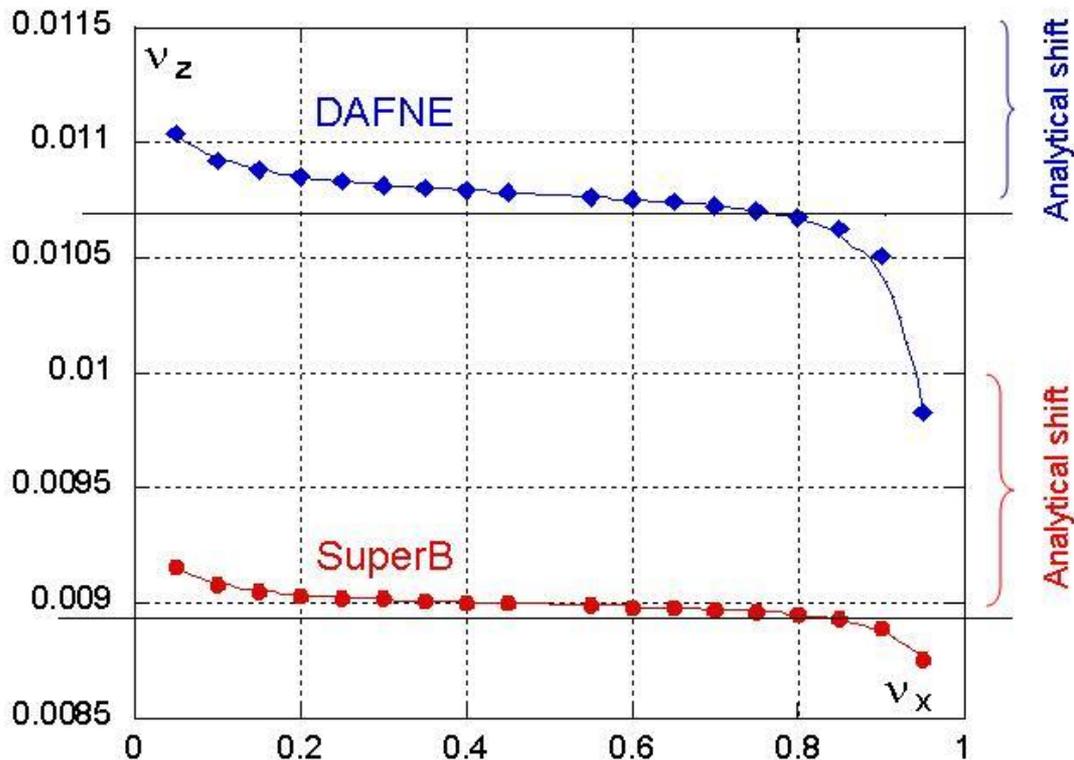
the tune shift expression can be further simplified to

$$\xi_z = -\frac{r_e N^{strong}}{2\pi\gamma^{weak}} \frac{\left(\frac{\sigma_{z0}}{(\sigma_E/E)} \right)^{weak}}{\left(\left(\frac{\sigma_x}{\operatorname{tg}(\theta/2)} \right)^2 + \sigma_z^2 \right)^{strong}}$$

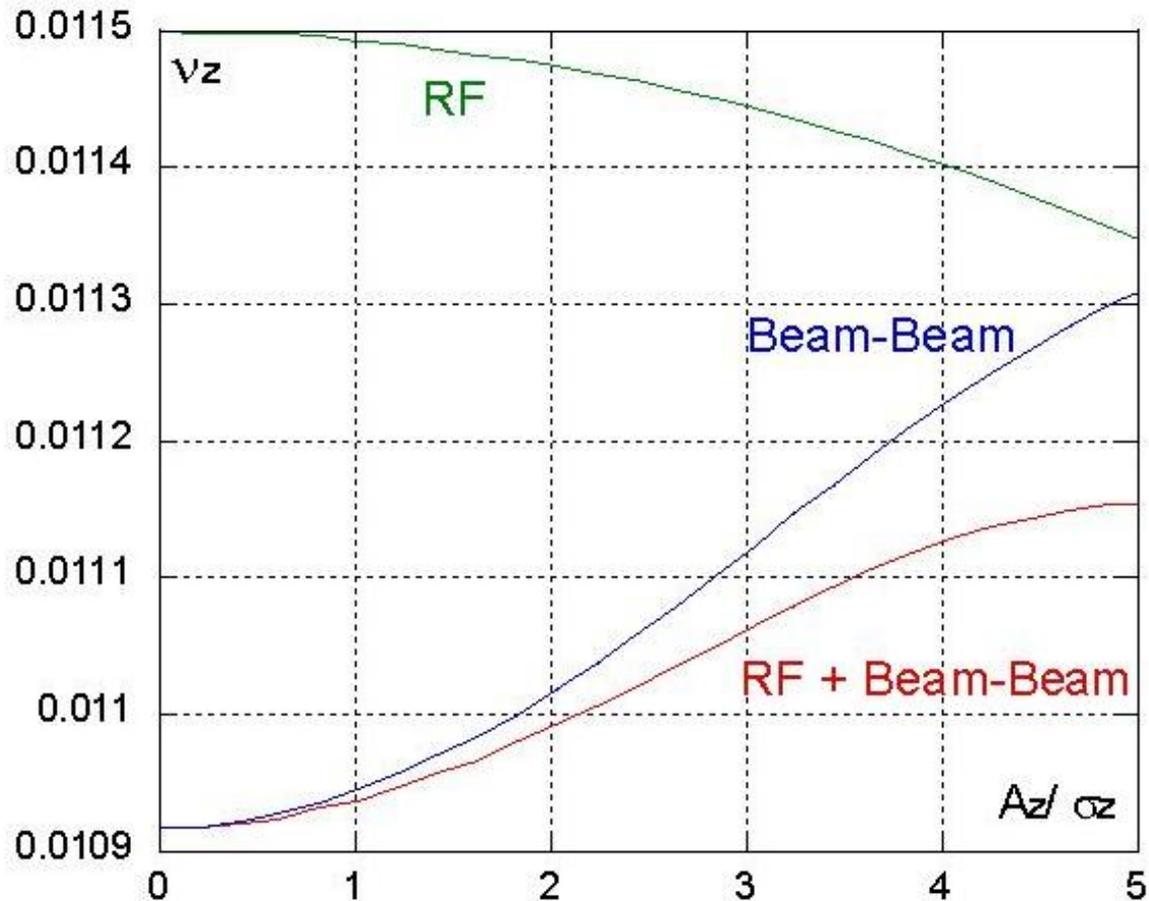
For the flat bunches the synchrotron tune shift practically does not depend on the vertical beam parameters. So, one should not expect any big variations due to crabbing and/or hour-glass effect. Since particles with very large synchrotron amplitudes practically do not “see” the opposite beam (except for a small fraction of synchrotron period) their synchrotron frequencies remain very close to the unperturbed value ν_{z0} . For this reason, like in the transverse cases, the linear tune shift can be used as a measure of the nonlinear tune spread

Numerical simulations

- In order to check validity of the previous formulae, we performed numerical simulations with the beam-beam code LIFETRAC.
- The synchrotron and betatron tunes in the presence of beam-beam effects are calculated by tracking and shown in the figure.



Synchrotron tune dependence on the horizontal tune. The solid straight lines correspond to the analytically predicted synchrotron tunes



Synchrotron tune dependence on normalized amplitude of synchrotron oscillations (blue curve – tune dependence created by beam-beam collisions alone, green – RF nonlinearity alone, red – both contributions).

Comments (1)

- First, our numerical simulations have confirmed that the synchrotron tune shift does not depend on parameters of the vertical motion, such as β_y and ν_y .
- Second, an agreement between the analytical and numerical estimates is quite reasonable for the horizontal tunes far from integers.
- Quite naturally, in a scheme with a horizontal crossing angle synchrotron oscillations are coupled with the horizontal betatron oscillations.
- One of the coupling's side effects is the ν_z dependence on ν_x , which becomes stronger in vicinity of the main coupling resonances.
- In order to make comparisons with the analytical formula we need to choose the horizontal betatron tune ν_x closer to half-integer, where its influence on ν_z is weaker.
- The coupling vanishes for very large Piwinski angles, that is why the ν_z dependence on ν_x is stronger for DAΦNE with respect to that of SuperB

Comments (2)

- Since ν_x for DAFNE is rather close to the coupling resonance, we will use numerical simulations in order to compare the calculated synchrotron tune shift with the measured one.
- In particular, when colliding the weak positron beam with 500mA electron beam, the measured synchrotron frequency shift was about -630 Hz (peak-to-peak).
- In our simulations we use the DAFNE beam parameters with respectively lower bunch current ($N = 0.9 \times 10^{10}$) and shorter bunch length ($\sigma_z = 1.6$ cm). This results in the synchrotron tune shift of -0.000232 corresponding to the frequency shift of -720 Hz.
- In our opinion the agreement is good considering experimental measurement errors and the finite width of the synchrotron sidebands

Conclusions

- Unexpected synchrotron oscillation damping due to beam-beam collisions experimental data have been collected by a commercial spectrum analyzer and by the bunch-by-bunch longitudinal feedback diagnostics
- Same result from two different diagnostic tools
- A simple analytical formula to explain synchrotron tune shift and tune spread due to beam-beam collisions with a crossing angle has been presented
- The formula agrees well with the simulations when the horizontal tune is far from the synchro-betatron resonances
- **The agreement is better for larger Piwinski angles.**
- **Measured and obtained by simulations synchrotron frequency shifts are in a good agreement**
- Calculations have shown that at high beam currents the synchrotron tune spread induced by the beam-beam interaction at DAFNE can be larger than the tune spread due to the nonlinearity of the RF voltage. This may result in additional Landau damping of the longitudinal coupled bunch oscillations.