Eigenmode Computation For Ferrite-Loaded Cavity Resonators Klaus Klopfer*, Wolfgang Ackermann, Thomas Weiland Institut für Theorie Elektromagnetischer Felder, TU Darmstadt



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GSI Helmholtzzentrum für Schwerionenforschung



Motivation



GSI SIS18 Ferrite Cavity



Main benefits of ferrite cavities:

- reduction of wavelength
 more compact cavity
- modification of resonance frequency in a wide range SIS 18 cavity:
 - $\sim 0.6\,\text{MHz}-5\,\text{MHz}$

Motivation



GSI SIS18 Ferrite Cavity: Main components





Assumptions: $|\vec{H}_d| \ll |\vec{H}_{\text{bias}}|$, effect of hysteresis negligible

$$\vec{B}(t) = \mu_0 \mu_{\text{bias}} \vec{H}_{\text{bias}} + \mu_0 \overleftrightarrow{\mu}_d \vec{H}_d(t) = \mu_0 \mu_{\text{bias}} \vec{H}_{\text{bias}} + \mu_0 \overleftrightarrow{\mu}_d \operatorname{Re}\left(\vec{H}_d \cdot e^{-i\omega t}\right)$$

Linearization at working point





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Linearization at working point



- Modification of bias current
 - ⇒ Modification of differential permeability
 - ⇒ Adjustment of eigenfrequency



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$$\vec{B}(t) = \mu_0 \mu_{\text{bias}} \vec{H}_{\text{bias}} + \mu_0 \overleftrightarrow{\mu_d} \vec{H}_d(t) = \mu_0 \mu_{\text{bias}} \vec{H}_{\text{bias}} + \mu_0 \overleftrightarrow{\mu_d} \operatorname{Re}\left(\vec{H}_d \cdot e^{-i\omega t}\right)$$

Linearization at working point



Eigensolutions determined by:

$$\epsilon^{-1} \nabla \times \left(\mu_0^{-1} \overset{\leftrightarrow}{\mu_d}{}^{-1}_d \nabla \times \vec{E} \right) = \omega^2 \vec{E}$$

Boundary condition: $\vec{n} \times \vec{E} = 0$ on cavity boundary



[D. Polder, Phil. Mag., 40 (1949)]

Properties of the differential permeability tensor $\overleftrightarrow{\mu}_{d}$:

Fully occupied (3×3)-tensor, for $\vec{H} = H_{\text{bias}} \cdot \vec{e}_z$ reduces to the Polder tensor

$$\stackrel{\leftrightarrow}{\mu}_{d} = \begin{pmatrix} \mu_{1} & i\mu_{2} & 0\\ -i\mu_{2} & \mu_{1} & 0\\ 0 & 0 & 1 \end{pmatrix}$$

with
$$\mu_{1,2} = \mu_{1,2} \Big(\vec{H}_{\text{bias}}, \omega \Big)$$

If magnetic losses are included: Im(µ_{1,2}) ≠ 0 ⇒ Non-Hermitian



$$\epsilon^{-1} \nabla \times \left(\mu_0^{-1} \overset{\leftrightarrow}{\mu}_d^{-1} \nabla \times \vec{E} \right) = \omega^2 \vec{E}$$



$$\epsilon^{-1} \nabla \times \left(\mu_0^{-1} \overset{\leftrightarrow}{\mu_d} ^{-1} \nabla \times \vec{E} \right) = \omega^2 \vec{E} - -$$

Discretization by Finite Integration Technique (FIT):

$$M_{\epsilon}^{-1}\widetilde{\mathbf{C}}M_{d,\mu}^{-1}\mathbf{C}\widehat{\mathbf{e}} = \omega^{2}\widehat{\mathbf{e}}$$



$$\epsilon^{-1} \nabla \times \left(\mu_0^{-1} \overset{\leftrightarrow}{\mu_d} ^{-1} \nabla \times \vec{E} \right) = \omega^2 \vec{E} - -$$

Discretization by Finite Integration Technique (FIT):

$$M_{\epsilon}^{-1}\widetilde{\mathbf{C}}M_{d,\mu}^{-1}\mathbf{C}\widehat{\mathbf{e}} = \omega^{2}\widehat{\mathbf{e}} \qquad \checkmark$$

- permeability tensor M_{d,µ}:
 - non-diagonal
 - dependend on \vec{H}_{bias} and ω



$$\epsilon^{-1}\nabla\times\left(\mu_{0}^{-1}\overset{\leftrightarrow}{\mu_{d}}^{-1}\nabla\times\vec{E}\right)=\omega^{2}\vec{E}$$

Discretization by Finite Integration Technique (FIT):

$$M_{\epsilon}^{-1}\widetilde{\mathbf{C}}M_{d,\mu}^{-1}\mathbf{C}\widehat{\mathbf{e}} = \omega^{2}\widehat{\mathbf{e}} \qquad \longleftarrow$$

- permeability tensor M_{d,µ}:
 - non-diagonal
 - dependend on \vec{H}_{bias} and ω
 - if magnetic losses included:

- requirements on eigensolver:
 - \Rightarrow nonlinear
 - \Rightarrow non-Hermitian



Concept



- calculation of bias magnetic field
- determination of permeability matrix M⁻¹_{d,µ}

General requirements:

- support of nonlinear material
- support of lossy material
- parallel computation with distributed memory (scalability)

calculation of eigenmodes





(*) Portable, Extensible Toolkit for Scientific Computation

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- ► *H_i*-algorithm Helmholtz decomposition $\vec{H} = \vec{H}_i + \vec{H}_h$ with $\nabla \times \vec{H}_i = \vec{J}$ and $\vec{H}_h = -\nabla \varphi$
- Solution of nonlinear equation: successive substitution or Newton method





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- Solution of nonlinear equation: successive substitution or Newton method

- Jacobi-Davidson algorithm harmonic Ritz-values for computation of interior eigenvalues
- Solution of nonlinear problem: successive substitution

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Aim: Efficient distributed computing





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- ► test model: $(|\vec{H}_{ext}| = 2750 \frac{A}{m}; \mu_r = 7)$ lossless, ferrite-filled cylindrical cavity resonator longitudinally biased by homogeneous magnetic field \rightarrow
- semi-analytical solution available [Chinn, Epp and Wilkins]













 test model: lossless, ferrite-filled cylindrical cavity resonator

longitudinally biased by homogeneous magnetic field





 test model: lossless, ferrite-filled cylindrical cavity resonator longitudinally biased by homogeneous magnetic field



$$\vec{\mathcal{H}}_{ext} \notin \vec{e}_{z}:$$

$$\vec{\mu}_{d} = \begin{pmatrix} \mu_{x,x} & \mu_{x,y} & \mu_{x,z} \\ \mu_{y,x} & \mu_{y,y} & \mu_{y,z} \\ \mu_{z,x} & \mu_{z,y} & \mu_{z,z} \end{pmatrix}$$



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Ferrite material is characterized by $B(H) = \mu_0 \ 2.5 \cdot 10^4 \tanh \left(H \cdot 10^{-2} \frac{\text{m}}{\text{A}}\right) \frac{\text{A}}{\text{m}} + \mu_0 H \text{ and } \epsilon_r = 1$







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- bias magnetic field excited by current winding (2 kA)



Results of fully nonlinear computation

Spectrum of the lowest nine eigenmodes





Results of fully nonlinear computation

Spectrum of the lowest nine eigenmodes





Results of fully nonlinear computation





Comparison nonlinear — linear computation





Comparison nonlinear — linear computation



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Goal:

calculation of eigenvectors for biased ferrite cavities

- 1. Magnetostatic solver (nonlinear material):
 - \Rightarrow permeability tensor $\overleftrightarrow{\mu}_{d}$
- 2. Eigensolver:
 - \Rightarrow nonlinear complex eigenvalue problem



Magnetostatic



Summary

Goal:

calculation of eigenvectors for biased ferrite cavities

1. Magnetostatic solver (nonlinear material):

 \Rightarrow permeability tensor $\stackrel{\leftrightarrow}{\mu}_{d}$

2. Eigensolver:

 \Rightarrow nonlinear complex eigenvalue problem

 Status: Functionality of fully nonlinear solver for Hermitian problems demonstrated



