

Implementational Aspects of Eigenmode Computation based on Perturbation Theory

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Overview

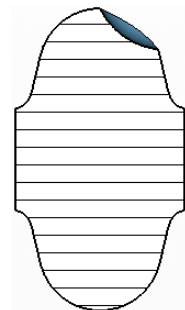
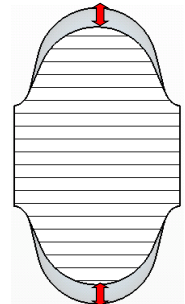
- Introduction
- Perturbation Theory
 - Basics
 - Introduction of two different Perturbative Methods
 - Analytical Proof of Principle
- Numerical Implementation
- Results
- Conclusion & Outlook



Introduction

Geometrical Perturbations of a Cavity

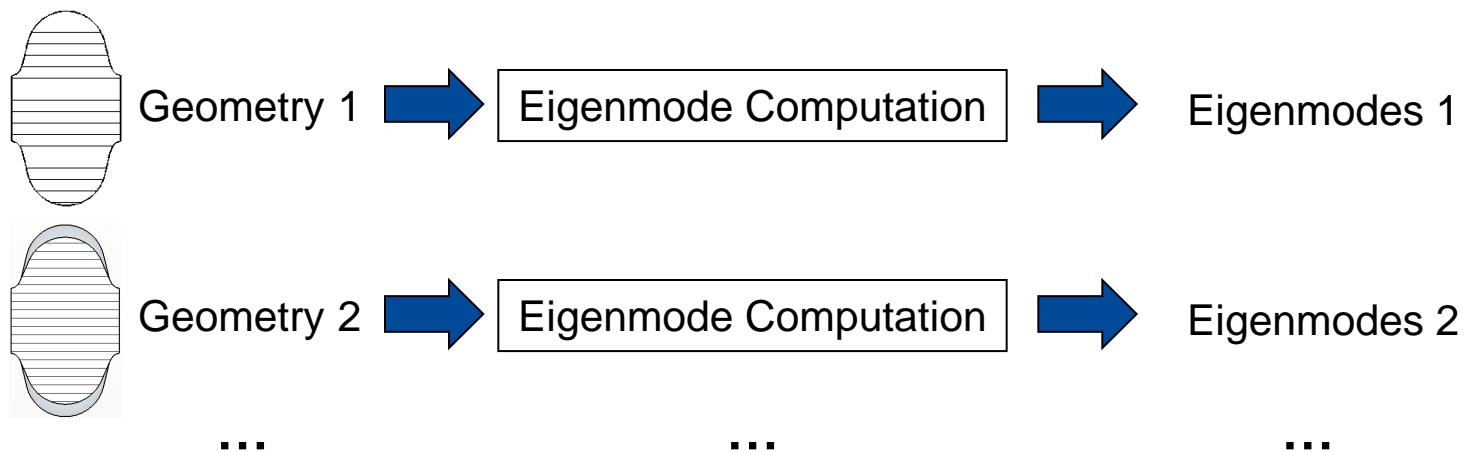
- Geometrical shape of a cavity determines
 - Eigenmodes inside the cavity
 - Cavity characteristics
- Perturbation of the cavity shape changes its characteristics
- Forms of perturbations:
 - Desired modification
 - Optimization of cavity characteristics ($E_{acc}/E_{peak}, Q, \dots$)
 - Undesired perturbation: Deviation of desired geometry due to manufacturing tolerances and operational demands
 - Impairment of accelerating performance (π -Mode)
 - Beam deflection / wakefield excitation



➔ Need to assess perturbation effects

Motivation for Using Perturbation Theory

- Parameter studies to investigate perturbation effects
→ Computation of eigenmodes for numerous different cavity geometries
- Common numerical solvers:
Perform a full computation even if geometry is only slightly changed
→ **Computationally extensive and inefficient**

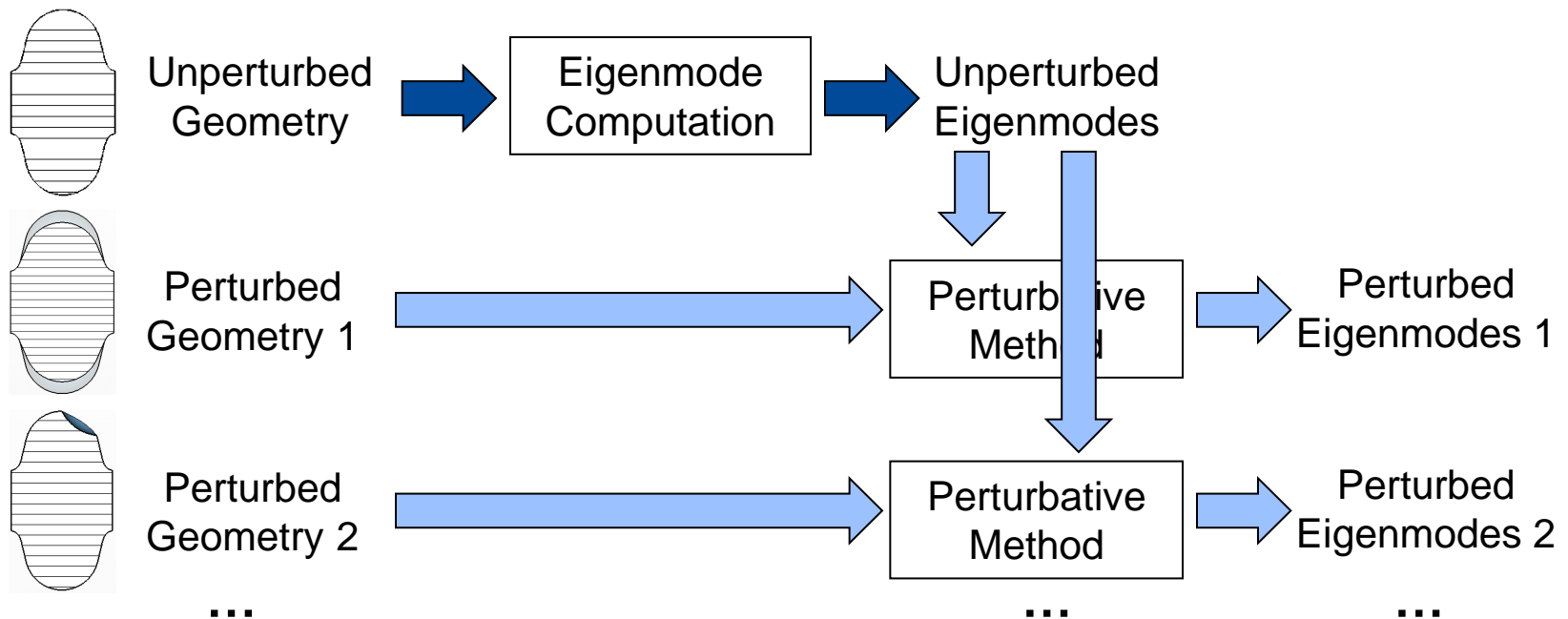


Motivation for Using Perturbation Theory

- Perturbative methods:

- Perform full eigenmode computation solely for one geometry
- Derive eigenmodes of every other geometry from these eigenmodes

→ Significant reduction of computational effort





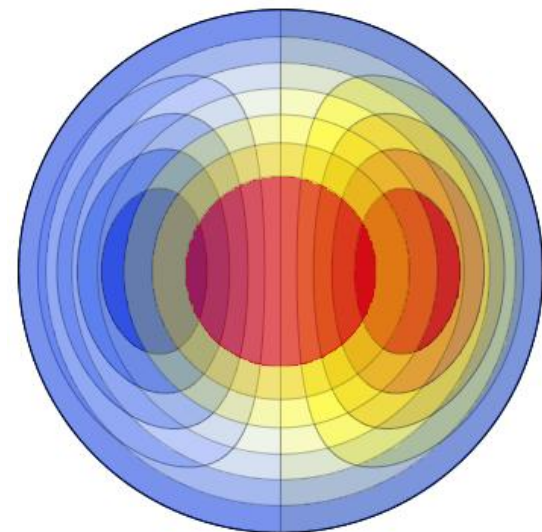
Basics of Perturbation Theory

Orthogonality of Eigenmodes

The integral of the product of the stationary fields of two different modes over the complete cavity volume is always zero

$$\frac{\epsilon}{2U_i} \iiint_V \mathbf{E}_i(\mathbf{r}) \cdot \mathbf{E}_k(\mathbf{r}) dV = \delta_{ik} = \begin{cases} 1 & i = k \\ 0 & i \neq k \end{cases}$$

$$\frac{\mu}{2U_i} \iiint_V \mathbf{H}_i(\mathbf{r}) \cdot \mathbf{H}_k(\mathbf{r}) dV = \delta_{ik} = \begin{cases} 1 & i = k \\ 0 & i \neq k \end{cases}$$



Cross section of cylindrical cavity
 E_z of TM_{01} and TM_{11} mode

V : Volume of the unperturbed cavity

U_i : Energy stored in unperturbed mode i

Orthogonality of Eigenmodes → Series Expansion

Fields of unperturbed eigenmodes

=

System of mutually orthogonal functions



Field pattern of perturbed eigenmodes:
Expansion as a series of the unperturbed modes

$$\begin{aligned}\tilde{\mathbf{E}}_i(\mathbf{r}) &= \sum_{k=1}^N e_{ik} \cdot \mathbf{E}_k(\mathbf{r}) \\ \tilde{\mathbf{H}}_i(\mathbf{r}) &= \sum_{k=1}^N h_{ik} \cdot \mathbf{H}_k(\mathbf{r})\end{aligned}$$

Condition: Perturbed volume has to be part of unperturbed volume

How to determine the perturbed eigenmodes

- Key aspect:

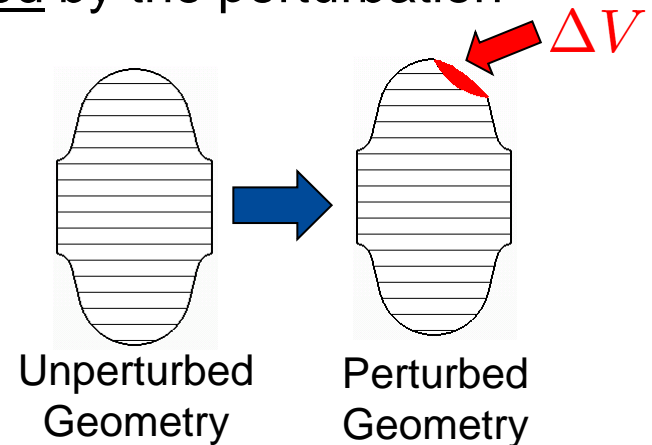
Interaction of each unperturbed mode with every other unperturbed mode inside the volume ΔV that is removed by the perturbation

$$IT_{E(ik)} = \iiint_{\Delta V} \mathbf{E}_i(\mathbf{r}) \cdot \mathbf{E}_k(\mathbf{r}) dV$$

$$IT_{H(ik)} = \iiint_{\Delta V} \mathbf{H}_i(\mathbf{r}) \cdot \mathbf{H}_k(\mathbf{r}) dV$$



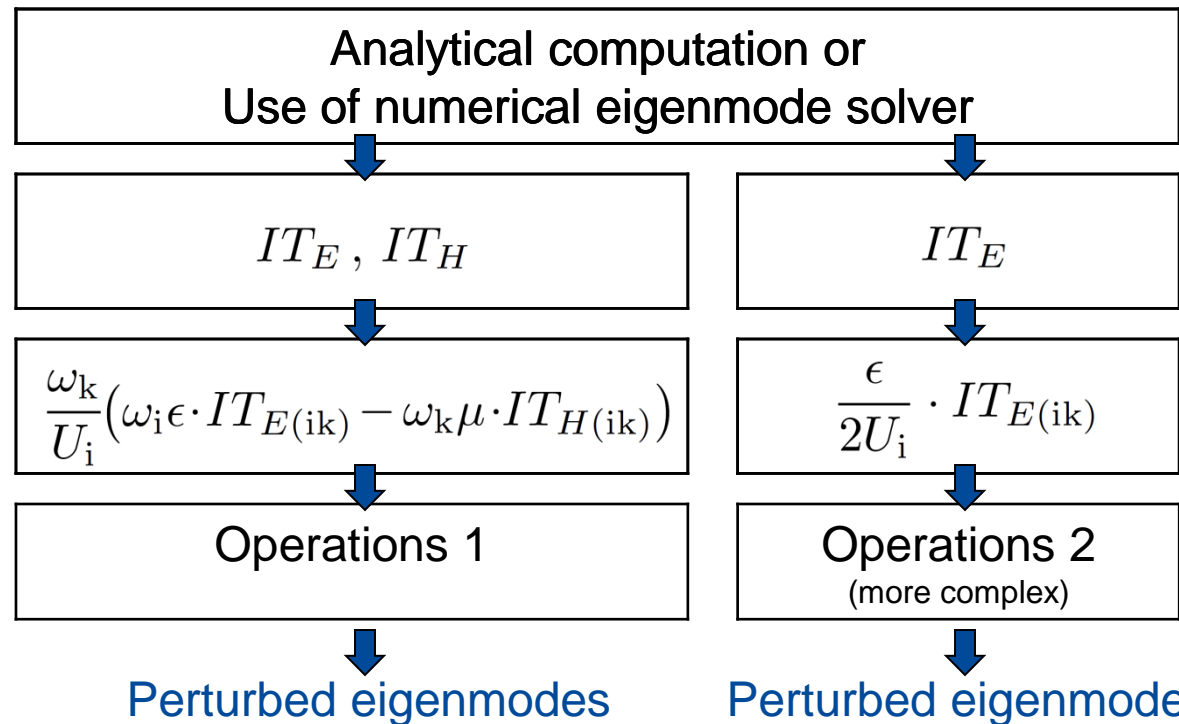
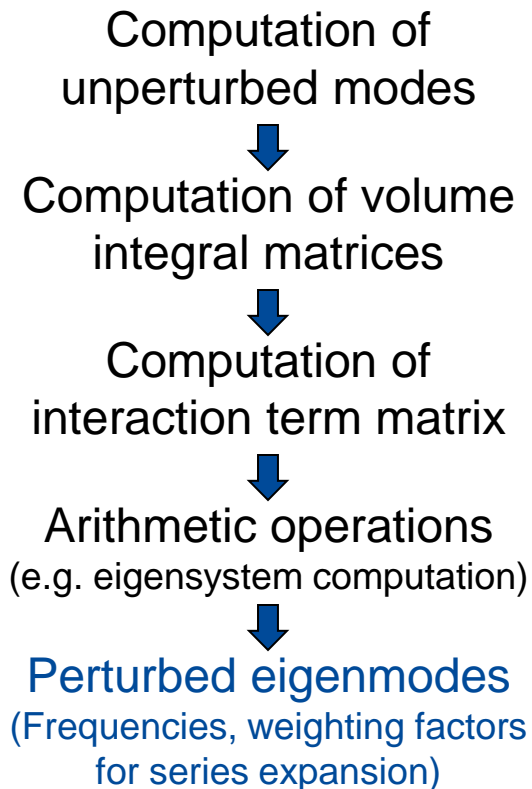
Volume integral
matrices
 IT_E, IT_H



Basic Operations of Perturbation Theory

Method 1: GST Generalisation of Slater's Theorem

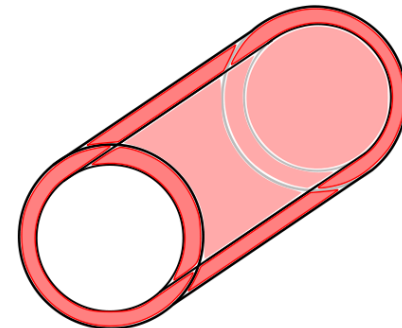
Method 2



Proof of Principle I

- Analytically evaluable cavity shape
 - Use of analytically computed unperturbed and eigenmodes
 - Any desired number of unperturbed modes usable for series expansion
 - Very high precision for implementation
 - Very low effort

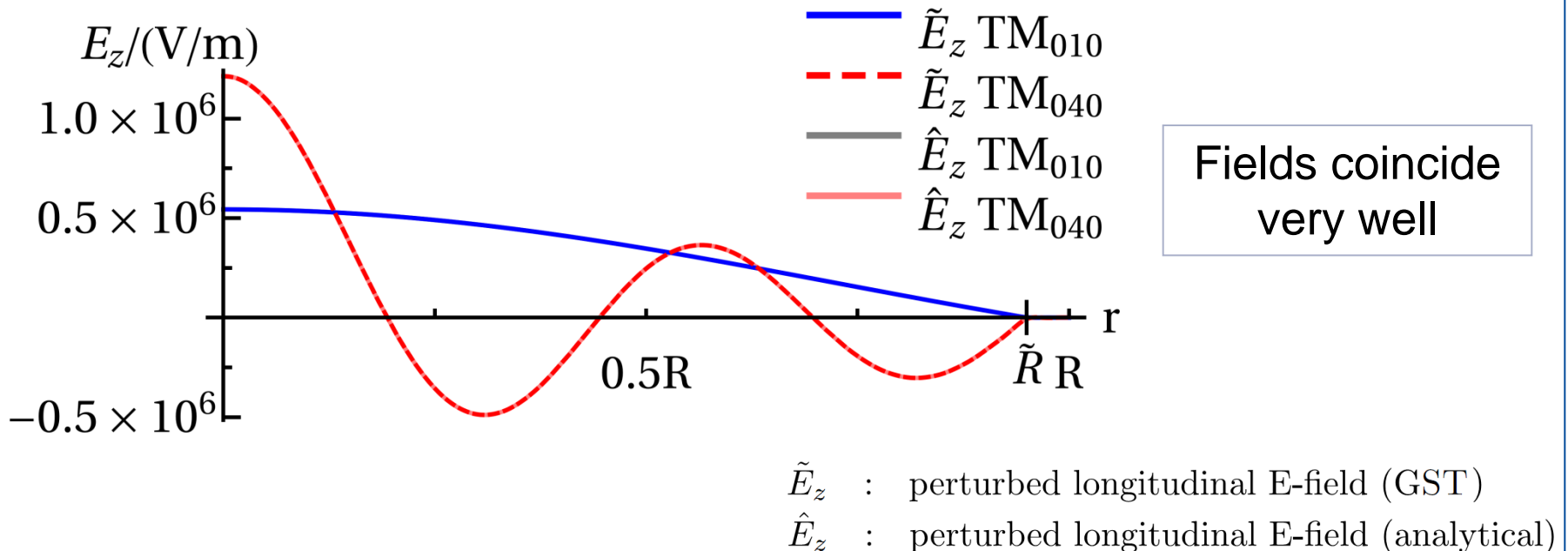
- Computation example
 - cylindrical cavity
 - subject to radial perturbation of 5%
 - Investigation of TM_{0n0} modes (n: radial mode index)



Proof of Principle II:

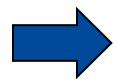
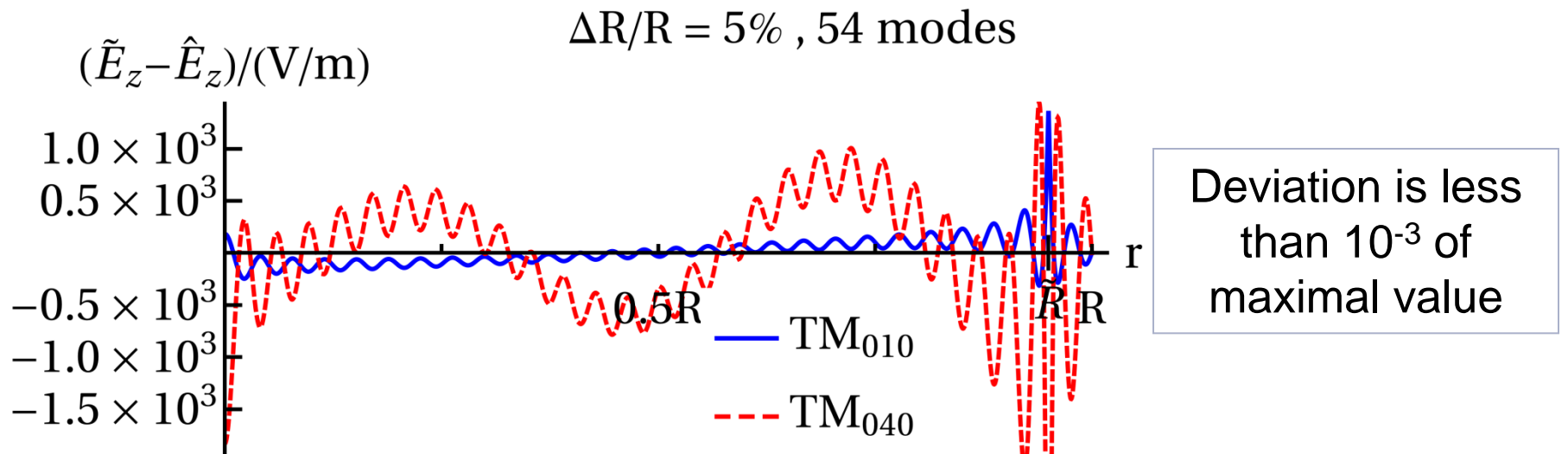
- Using only a small number of unperturbed modes $TM_{0.1.0}$ to $TM_{0.54.0}$
- Relative error of perturbed frequencies $< 10^{-3}$
- Longitudinal electric field along the radius r

$\Delta R/R = 5\%$, 54 modes



Proof of Principle II:

- Using only a small number of unperturbed modes $TM_{0.1.0}$ to $TM_{0.54.0}$
- Relative error of perturbed frequencies $< 10^{-3}$
- Longitudinal electric field along the radius r



Very accurate results

\tilde{E}_z : perturbed longitudinal E-field (GST)
 \hat{E}_z : perturbed longitudinal E-field (analytical)



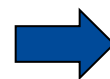
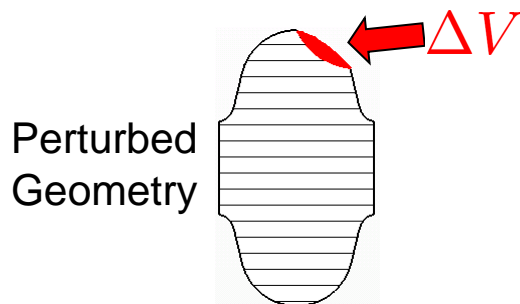
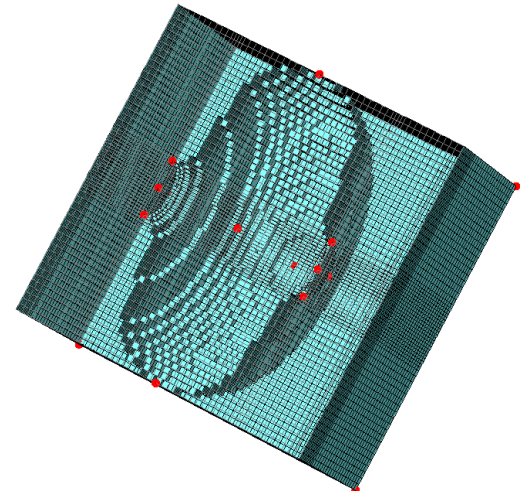
Numerical Implementation

Basic Operations of Numerical Implementation

- Computation of unperturbed eigenmodes
→ Unperturbed frequencies and fields
 - Computation of volume integrals
 - Computation of interaction terms
 - Evaluation of arithmetic operations
→ Perturbed frequencies and weighting factors
 - Series expansion in terms of unperturbed fields
→ Perturbed fields
- } • Simple & Low effort
• Equal for analytical and numerical implementation

Numerical computation of unperturbed eigenmodes

- Result relevant parameters
 - Mesh density
 - Exactness of boundary discretization
→ Accuracy of unperturbed frequencies and fields
- Depends on
 - Cavity geometry
 - Frequency range / minimal wavelength
→ Computation method

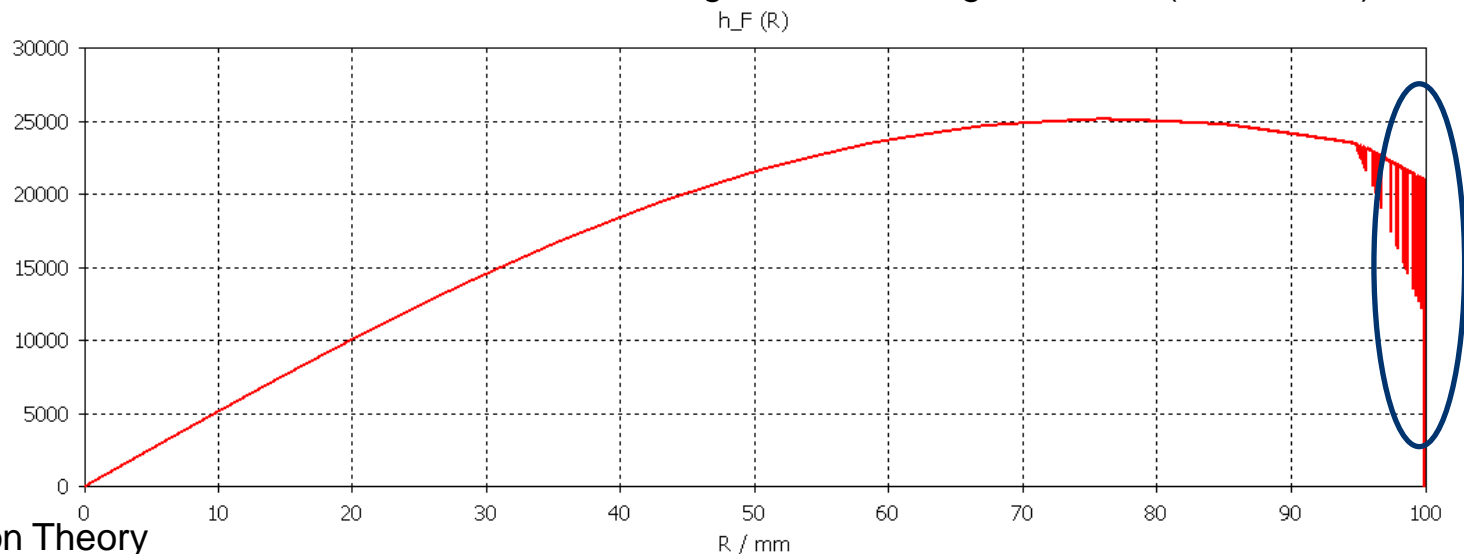


Accuracy of boundary fields
very important

FIT*-based Eigenmode-Computation (CST-MicrowaveStudio©)

- Dual grid with hexahedral cells and PBA (partially filled cells)
 - Inner fields: Good approximation
 - Fields near to the boundary: abrupt transition to zero

Magnetic field along the radius (R=100 mm)



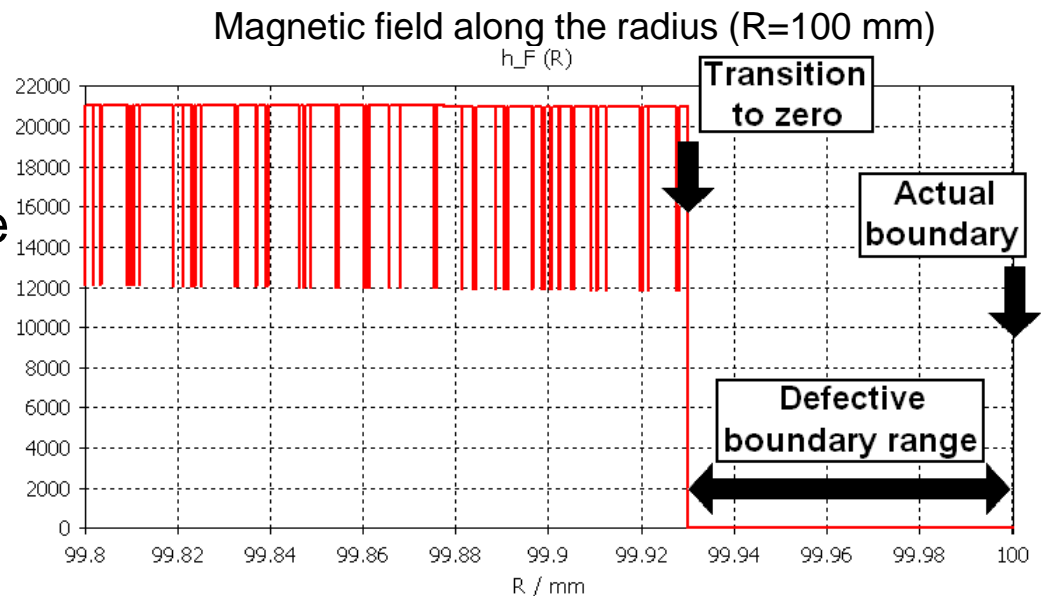
* Finite Integration Theory

FIT*-based Eigenmode-Computation (CST-MicrowaveStudio©)

- Dual grid with hexahedral cells and PBA (partially filled cells)
 - Inner fields: Good approximation
 - Fields near to the boundary: abrupt transition to zero
- Boundary fields are of crucial importance for volume integrals

Defective boundary range
must be as small as possible

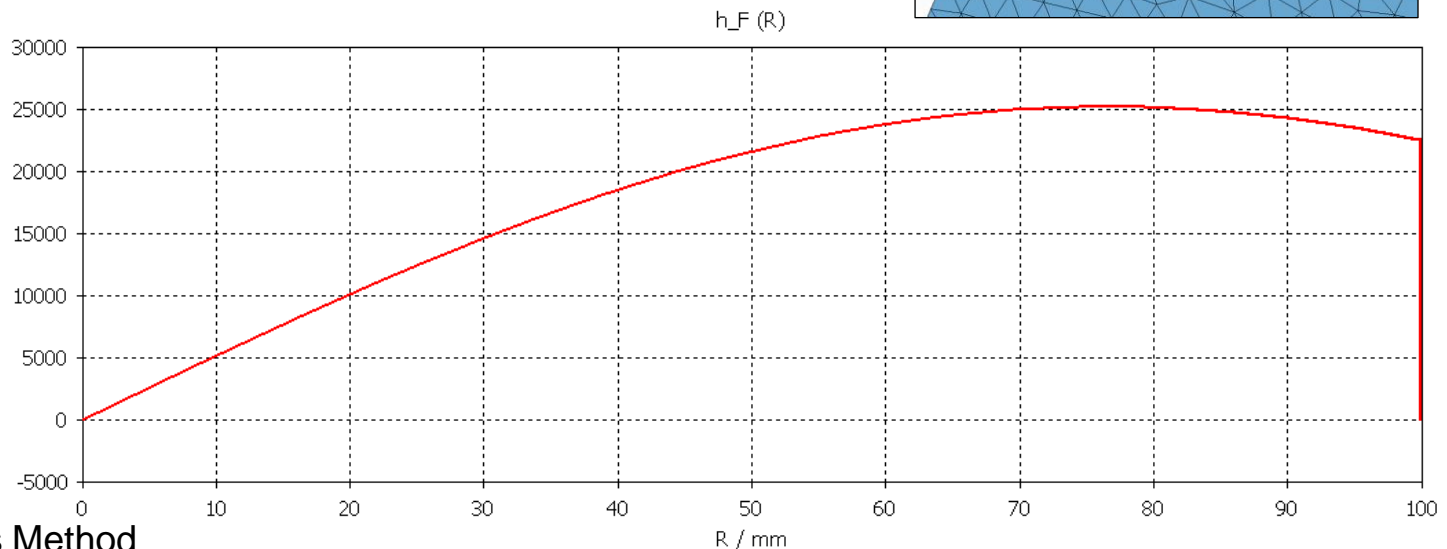
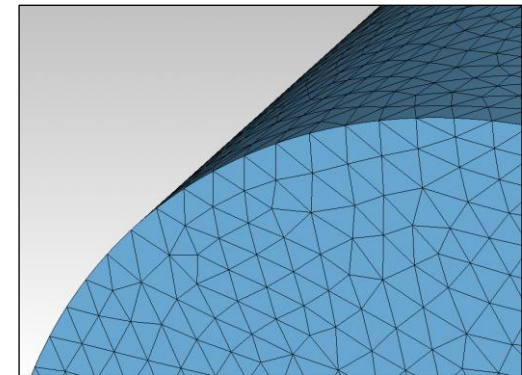
➔ FIT not appropriate for
Perturbation Theory



* Finite Integration Theory

FEM*-based Eigenmode-Computation (CST-MicrowaveStudio©)

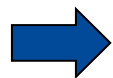
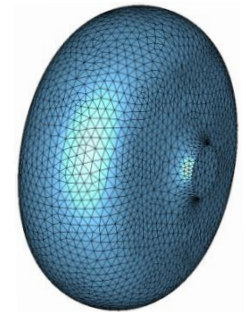
- Tetraeder-Grid with Curved Elements
 - Better approximation of boundary curve
 - No oscillations



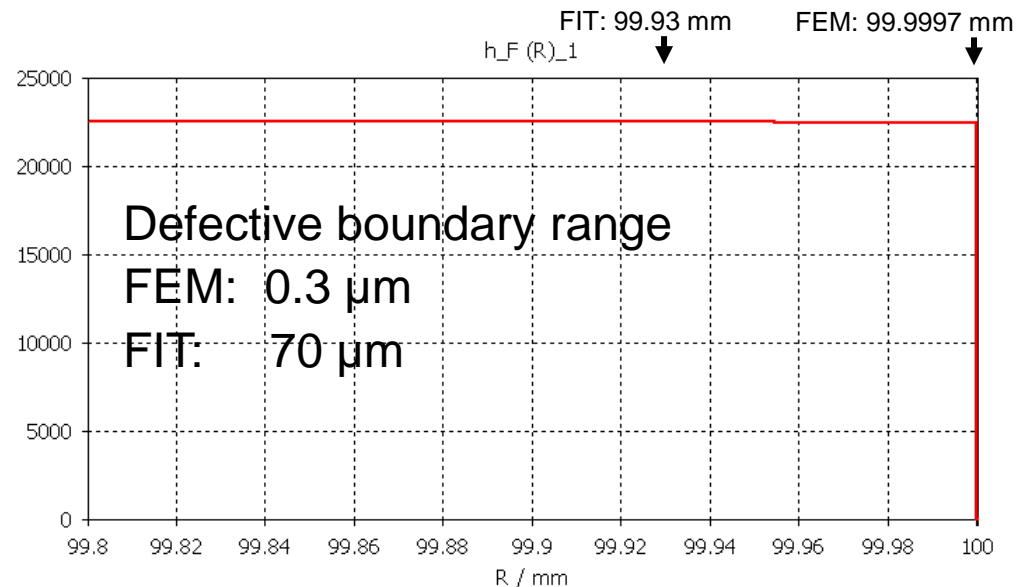
* Finite Elements Method

FEM*-based Eigenmode-Computation (CST-MicrowaveStudio©)

- Tetraeder-Grid with Curved Elements
 - Better approximation of boundary curve
 - No oscillations
 - Defective boundary range essentially smaller
- Especially for elliptical cavities more accurate results



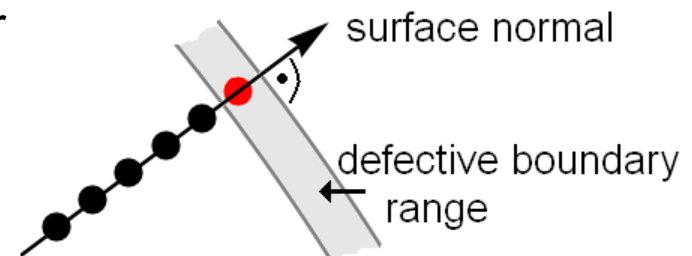
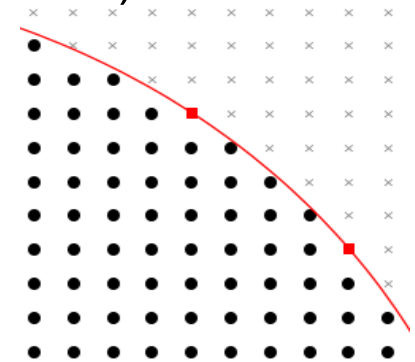
FEM more suitable
for boundary fields



* Finite Elements Method

Extrapolation of Boundary Field Values

- After computation of unperturbed eigenmodes (CST MWS)
 - Fields have to be exported as discrete field points
- But still: Field points inside defective boundary range
 - Incorrect field values of zero
 - Impairment of volume integrals
- Simple but effective solution
 - Extrapolation of defective field value from set of correct values along surface normal vector
 - 1D interpolation



- Interpolation values
- Extrapolated value

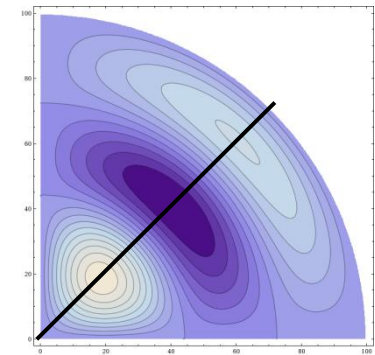
Computation of Volume Integrals: Approach 1

- Interpolation of discrete field data
→ Continuous 3D IP-functions for integration

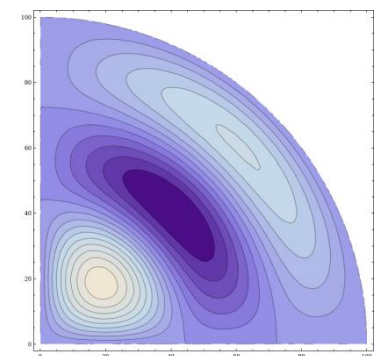
$$IT_{E(i,k)} = \iiint_{\Delta V} \mathbf{E}_i(\mathbf{r}) \cdot \mathbf{E}_k(\mathbf{r}) dV$$

- For very accurate results
 - IP-functions with polynomial degree > 1 needed

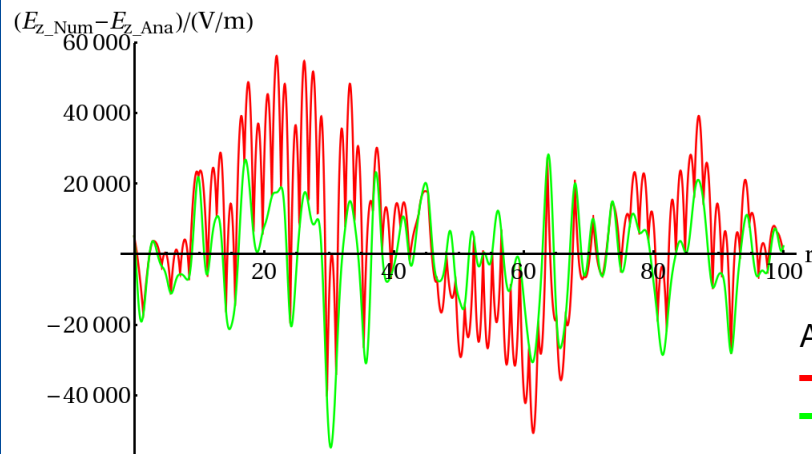
Cylindrical cavity
 E_z -Field of TM_{232}



Analytical



Interpolation



Absolute error of field along radius:

— Interpolation (degree of 1)

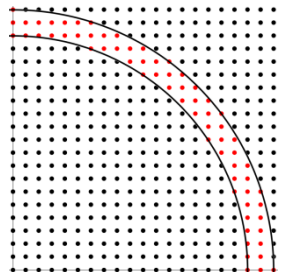
— Interpolation (degree of 3)

Computation of Volume Integrals: Approach 1

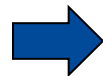
- Interpolation of discrete field data
→ Continuous 3D IP-functions for integration

$$IT_{E(i,k)} = \iiint_{\Delta V} \mathbf{E}_i(\mathbf{r}) \cdot \mathbf{E}_k(\mathbf{r}) dV$$

- For very accurate results
 - IP-functions with polynomial degree > 1 needed
 - Only realizable on a structured grid
 - Very large number of field points



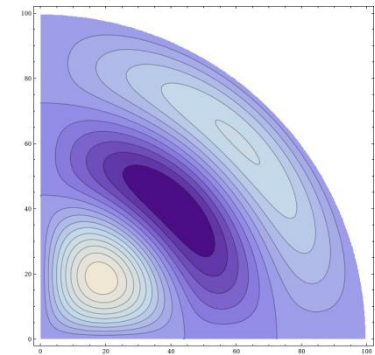
Radial perturbation



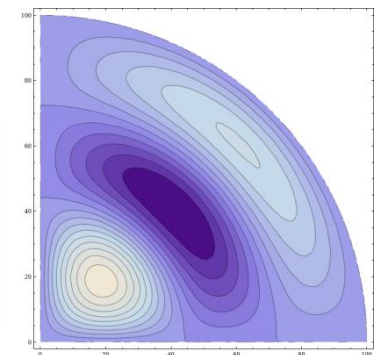
Enormous effort

- Export of field values
- 3D Interpolation of degree > 1

Cylindrical cavity
 E_z -Field of TM_{232}



Analytical



Interpolation

Computation of Volume Integrals: Approach 2

- Numerical integration: Summation of products of discrete fields and discrete volume

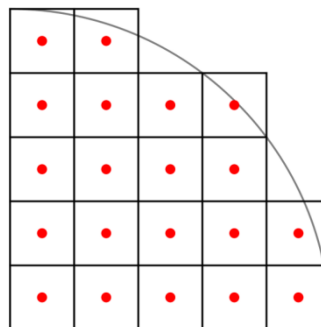
$$\iiint_{\Delta V} \mathbf{E}_i(\mathbf{r}) \cdot \mathbf{E}_k(\mathbf{r}) \, dV \rightarrow \sum_{m=1}^M \mathbf{E}_i(\mathbf{r}_m) \cdot \mathbf{E}_k(\mathbf{r}_m) \cdot \Delta V_m$$

- Partitioning of ΔV into volume elements

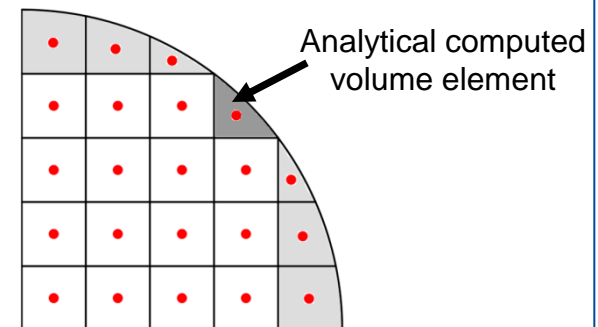
Commonly used cubic elements: very inaccurate for boundary elements

- Solution: Analytical volume elements

If boundary of ΔV and element intersect: analytical computation of volume and center



Cubic elements



Cubic elements (analytical)

Computation of Volume Integrals: Approach 2

- Numerical integration: Summation of products of discrete fields and discrete volume

$$\iiint_{\Delta V} \mathbf{E}_i(\mathbf{r}) \cdot \mathbf{E}_k(\mathbf{r}) \, dV \rightarrow \sum_{m=1}^M \mathbf{E}_i(\mathbf{r}_m) \cdot \mathbf{E}_k(\mathbf{r}_m) \cdot \Delta V_m$$

- Partitioning of ΔV into volume elements

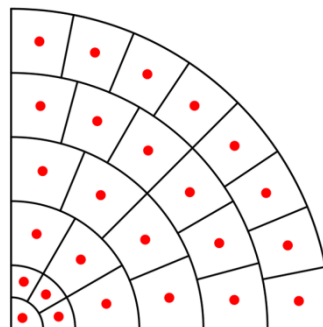
Commonly used cubic elements: very inaccurate for boundary elements

- Solution: Analytical volume elements

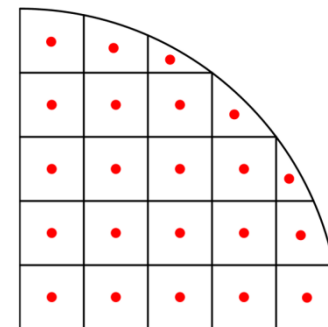
If boundary of ΔV and element intersect: analytical computation of volume and center



- Precise discretization
- Low effort



Cylindrical elements (analytical)



Cubic elements (analytical)

Computation of perturbed fields

- Final arithmetic operations yield
 - Perturbed frequencies
 - Weighting factors for series expansion
- Series expansion of perturbed fields

$$\tilde{\mathbf{E}}_i(\mathbf{r}) = \sum_{k=1}^N e_{ik} \cdot \mathbf{E}_k(\mathbf{r})$$

$$\tilde{\mathbf{H}}_i(\mathbf{r}) = \sum_{k=1}^N h_{ik} \cdot \mathbf{H}_k(\mathbf{r})$$

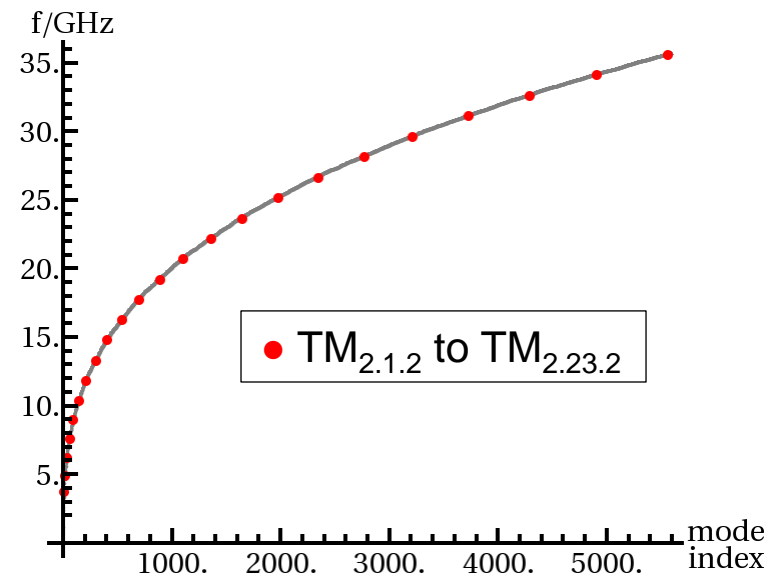
- CST MWS post processing: Summation of unperturbed fields multiplied by weighting factors
- External program: Export of unperturbed fields
→ Summation of discrete field values



Results

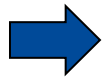
Numerical computation of unperturbed eigenmodes

- Computation example
 - Cylindrical cavity (R=100 mm, L=100 mm, fundamental mode: 1.15 GHz)
 - Investigation of TM_{2n2} modes subject to radial perturbations
- Eigenmodes computable in a very large frequency range
 - Up to 35.71 GHz

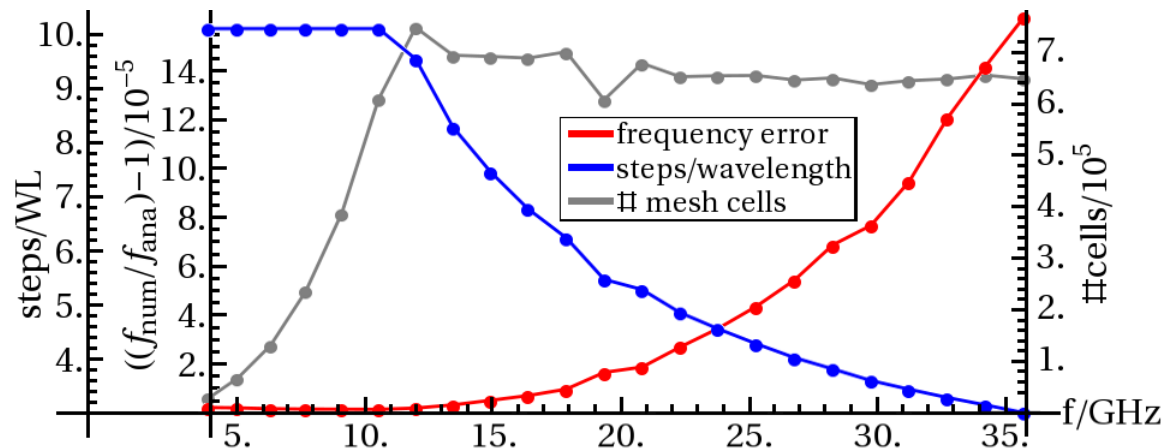


Numerical computation of unperturbed eigenmodes

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 - Investigation of TM_{2n2} modes subject to radial perturbations
- Eigenmodes computable in a very large frequency range
 - Up to 35.71 GHz
 - Frequency error $< 1.6 \cdot 10^{-4}$

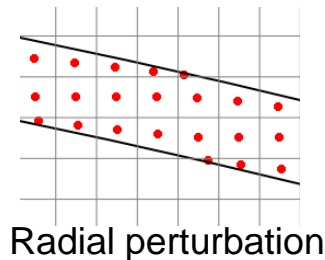
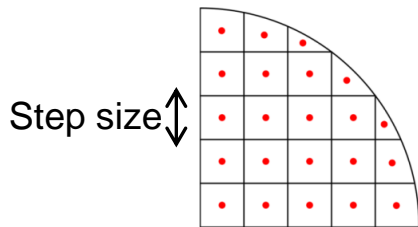


Very accurate in a
large frequency range



Computation of Volume Integrals: Size of elements

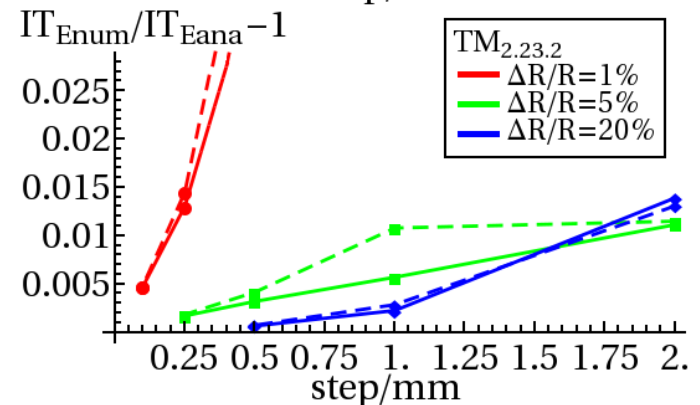
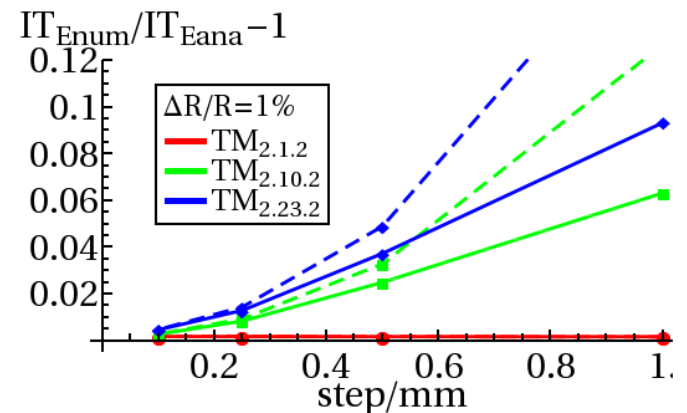
- Necessary step size depends on
 - Frequency / wave length
 - Extent of perturbation



$\Delta R/R$	ΔR	Step size	# elements
1 %	1 mm	0.25 mm	500000
5 %	5 mm	0.5 mm	306000
20 %	20 mm	1 mm	150000

➔ Reasonable number of mesh cells

Relative deviation of volume integrals

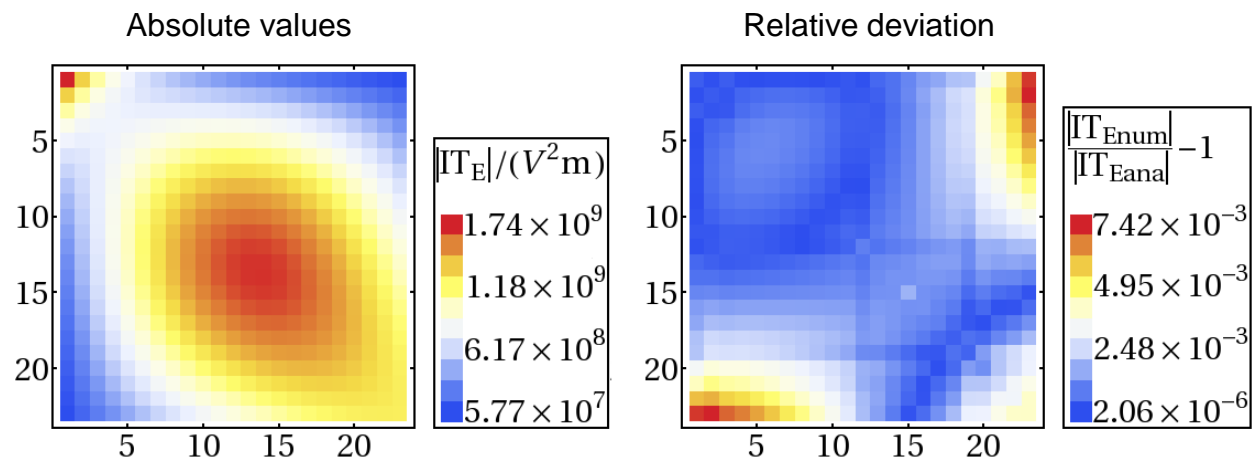


— Cubic elements
- - - Cylindrical elements

Computation of Volume Integrals: Accuracy

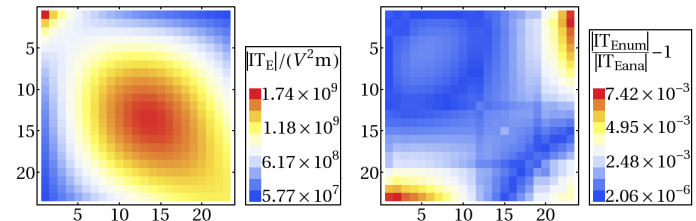
- IT_E (electric fields) : Very accurate
 - Relative error mainly $< 2 \cdot 10^{-3}$
 - Only for small values of IT_E larger error

Radial perturbation: 5% (5 mm)
Step size: 0.5 mm



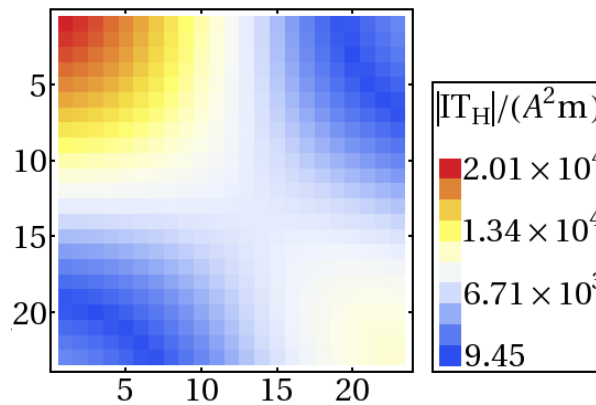
Computation of Volume Integrals: Accuracy

- IT_E (electric fields) : Very accurate
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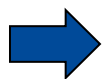
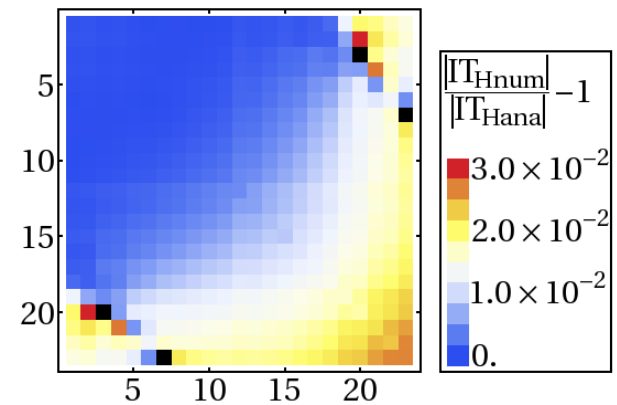


- IT_H (magnetic fields) : Very accurate for limited frequency range
 - Up to 21 GHz (13th mode): Relative error $< 7 \cdot 10^{-3}$
 - Increases up to $3 \cdot 10^{-2}$ for frequencies

Absolute values



Relative deviation



IT_E : very accurate
 IT_H : accurate

Radial perturbation: 5% (5 mm)
Step size: 0.5 mm

Interaction Terms & Final Results

- Accuracy depends on perturbative method

Method 1

$$\frac{\omega_k}{U_i} (\omega_i \epsilon \cdot IT_{E(ik)} - \omega_k \mu \cdot IT_{H(ik)})$$



$$\mu/\epsilon \approx 10^5$$

Absolute error of IT_H scaled up



Depending on ratio of $\omega_i \cdot \epsilon \cdot IT_{E(ik)}$ to $\omega_k \cdot \mu \cdot IT_{H(ik)}$
Error of resulting interaction term may increase

Method 2

$$\frac{\epsilon}{2U_i} \cdot IT_{E(ik)}$$

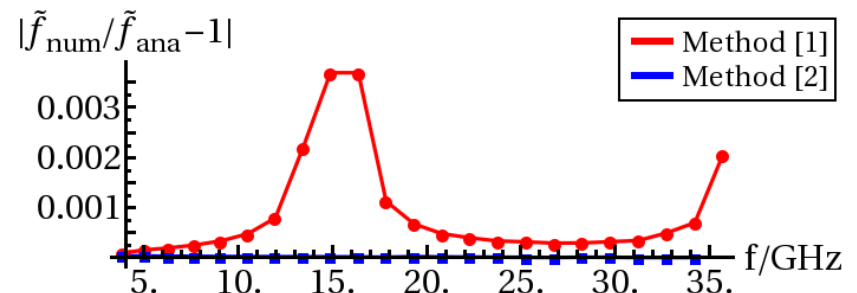


Same relative error as IT_E
Very accurate

Error of perturbed frequencies

- Method 2: Extremely small
- Method 1: Increased but still very small

Relative deviation between numerical
and analytical implementation





Conclusion & Outlook

Conclusion

- Perturbation Theory: Efficient method to compute perturbed eigenmodes of cavities in the context of parameter studies
 - Numerical implementation
 - Feasible with commonly used standard software
 - (Very) accurate results over a large frequency range
 - Reasonable computational effort
 - Differences in error propagation depending on perturbative method
- Application to arbitrary cavity geometries

Outlook

- Application to real cavities (elliptical cavities)
- Improvement of algorithm of perturbative methods

References

- [1] D. Meidlinger, *A General Perturbation Theory for Cavity Mode Field Pattern*, Proceedings of SRF2009, Berlin, Germany, THPPO005, 2009.
- [2] K. Brackebusch, U. van Rienen, *Eigenmode Computation for Cavities with Perturbed Geometry based on a Series Expansion of Unperturbed Eigenmodes*, Proceedings of 3rd IPAC, pp. 277-279, New Orleans, USA, 2012.
- [3] K. Brackebusch, H.-W. Glock, U. van Rienen, *Eigenmode Calculation of High Frequency Fields in Resonant Cavities based on Perturbation Theory*, Proceedings of 2nd IPAC, San Sebastian, Spain, pp.2235-2237, 2011.
- [4] Computer Simulation Technology AG, *Microwave Studio®*, 64289 Darmstadt, Germany.
- [5] Wolfram Research Inc., *Mathematica® Edition: Version 8.0*, Champaign, 2010.