

Status of collective effects at GSI

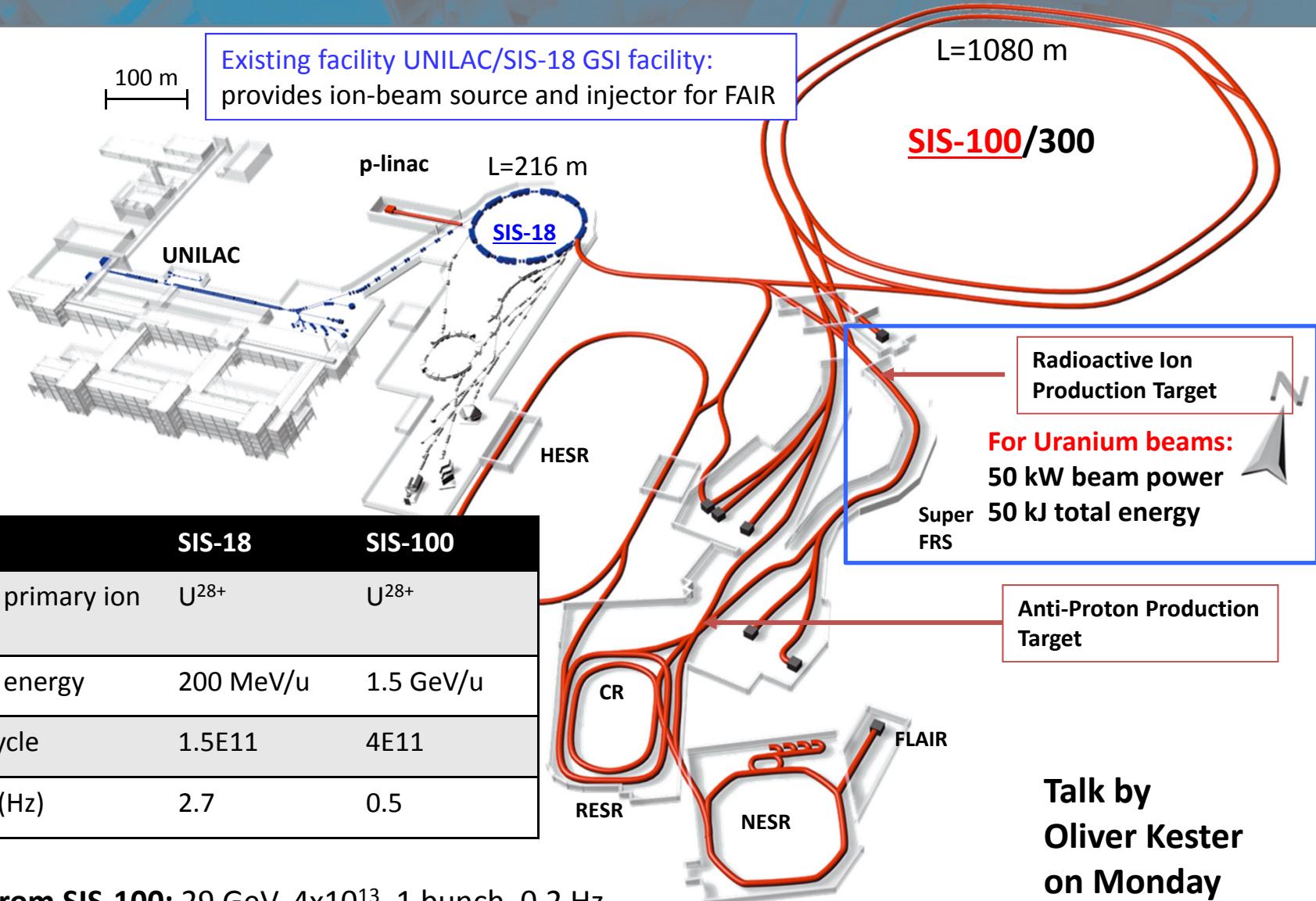
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Contents

- Collective effects at FAIR
- Simulation codes for collective effects in rings at GSI
- Impedances and expected transverse coherent instabilities
- Beam diagnostics: Interpretation of tune spectra using head-tail modes
- Conclusions

The FAIR accelerator facility



Collective effects in the FAIR rings

Incoherent space charge:

$$\epsilon_0 \nabla \cdot \vec{E} = \rho \quad (\text{in the rest system of the beam})$$

-> tune shift: $\Delta Q_y^{sc} \propto -\frac{q^2}{m} \frac{N}{B_f} \frac{4}{\epsilon_y \beta_0^2 \gamma_0^3} \frac{1}{1 + \sqrt{\epsilon_y/\epsilon_x}}$

-> $\Delta Q^{sc} \lesssim 0.4$

-> beam loss and modification of coherent effects

Impedances:

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad \nabla \times \vec{B} = \mu_0 \vec{j} + \frac{1}{c^2} \frac{\partial \vec{E}}{\partial t} \quad (\text{laboratory system})$$

-> image currents in the beam pipe

-> magnetic/resistive materials: ferrite, magnetic alloy

-> coherent instabilities and feedback requirements

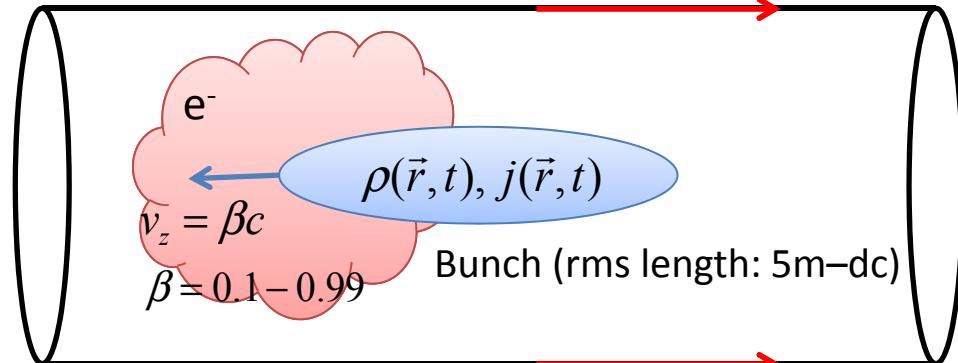
Secondary particles:

electron clouds created by residual gas ionization and SEY.

-> trapping of electron during slow extraction, two-stream instability.

Thin beam pipe (0.3 mm stainless steel)

Image current



Intrabeam scattering:

-> Beam cooling equilibrium in the HESR

In the FAIR synchrotrons SIS-18 and SIS-100 different incoherent/coherent effects occur simultaneously.

Beam loss in SIS-100 has be limited below 5 % (injection energy) and 1-2 % (extraction energy)

-> Computer modeling in combination with dedicated experiments (model validation) is essential.

Simulation codes for collective effects in rings at GSI

MICROMAP (G. Franchetti):

- 3D particle tracking with error multipoles,
- 2D self-consistent, 3D adaptable space charge.
- **see presentation on Wednesday**

G. Franchetti et al., *Experiment on space charge driven nonlinear resonance crossing in an ion synchrotron*, PRST-AB 2010

External tracking codes with collective effects:

HEADTAIL, PTC/Orbit (CERN) + **pyorbit** (SNS)

LOBO (O. Boine-F.):

- Longitudinal beam dynamics with space charge, impedances, cooling, IBS, internal targets
- direct Vlasov-Fokker-Planck solver or PIC

O. Boine-F., *rf barrier compression with space charge*, PRST-AB 2010

O. Boine-F., A. Lehrach, et al., *Cooling equilibrium and beam loss with internal targets in the HESR*, NIM A 2006

PATRIC (O. Boine-F., V. Kornilov, et al.):

- 3D particle tracking with self-consistent 2.5D space charge solver and wake fields
- MADX maps, arbitrary rf bucket forms
- Implemented for multi-core CPUs using MPI.

V. Kornilov, O. Boine-F., *Head-tail instability and Landau damping in bunches with space charge*, PRST-AB 2010

CST EM and Particle Studio® (CST AG, Darmstadt)
+ 3D EM frequency domain solver for impedances
poster by U. Niedermayer (Monday)

U. Niedermayer, O. Boine-F., *Analytical and numerical calculations of resistive wall impedances for thin beam pipe structures at low frequencies*, NIM A 2012

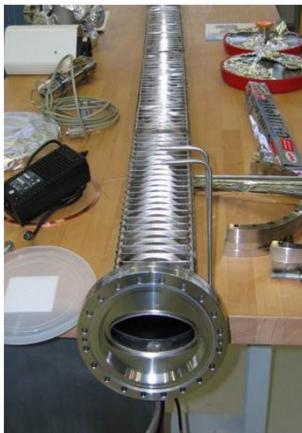
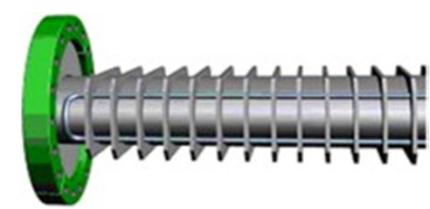
VORPAL (Tech-X) 3D EM PIC + 2D e-cloud codes

O. Boine-F., E. Gjonaj, F. Petrov, F. Yaman, T. Weiland, G. Rumolo, *Energy loss and longitudinal wakefield of short proton bunches in electron clouds*, PRST-AB 2012

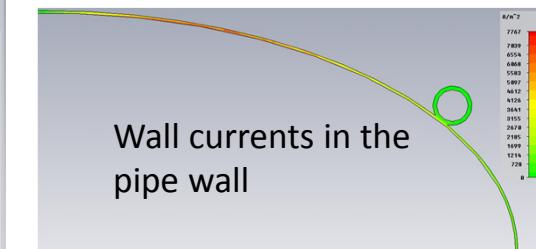
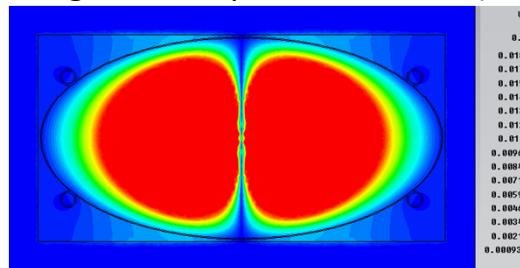
Transverse SIS-100 beam pipe impedance

CST® EM Studio + frequency domain solvers

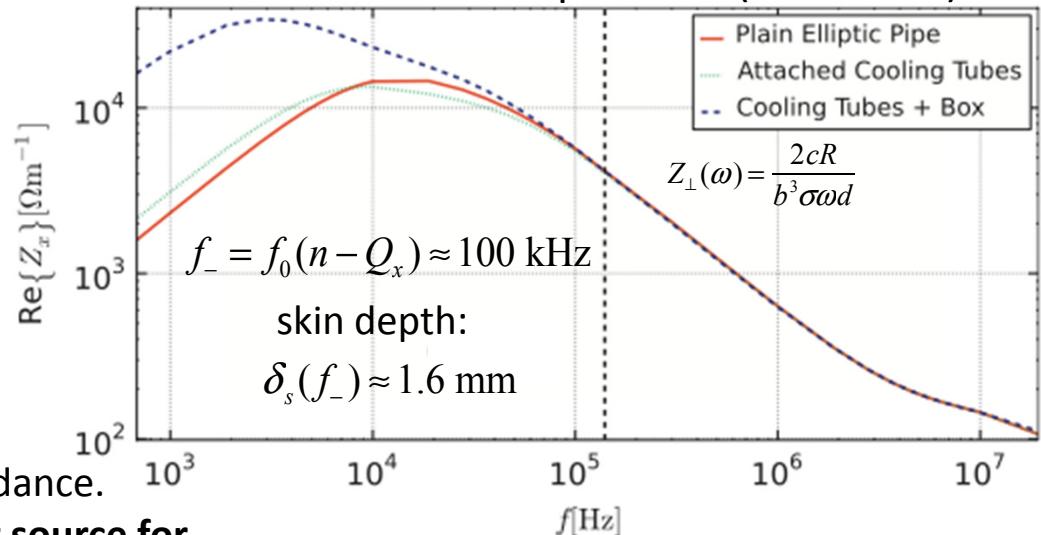
SIS-100 thin (**0.3 mm**) stainless steel beam pipewith cooling tubes attached.



Longitudinal electric field in the SIS100 pipe structure resulting from a dipolar excitation (300 kHz).



Transverse impedance (horizontal)



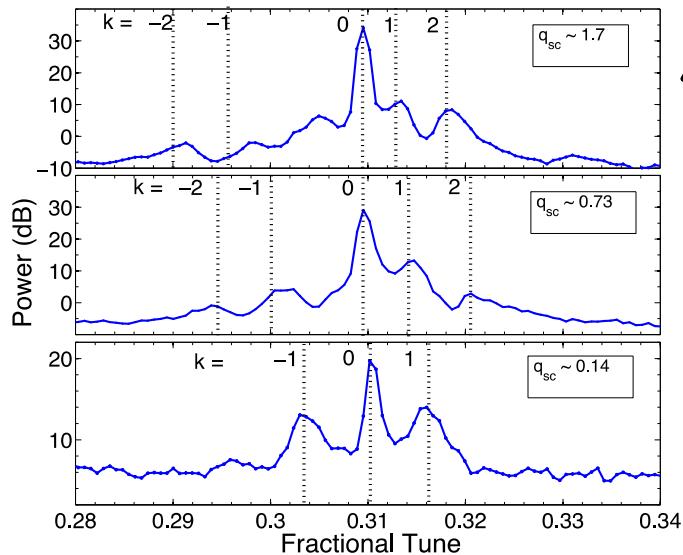
U. Niedermayer, O. Boine-F., *Analytical and numerical calculations of resistive wall impedances for thin beam pipe structures at low frequencies*, NIM A 2012

- In the frequency range of interest outside structures do not contribute to the impedance.
- **The thin resistive beam pipe is the major source for the expected head-tail instabilities in SIS-100.**

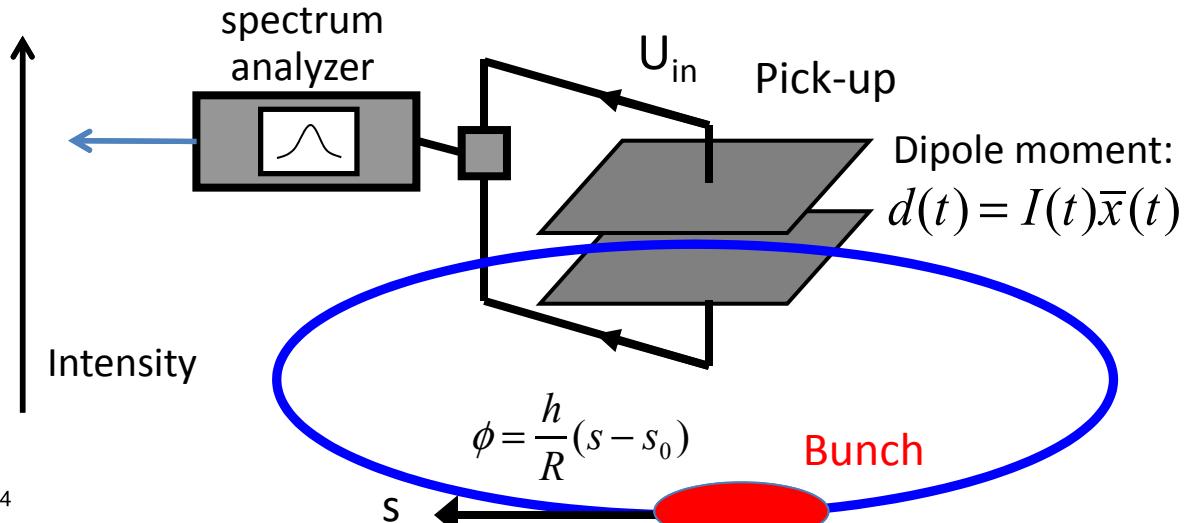
Next step: ferrite kickers and bench measurements (see poster by Uwe Niedermayer)

Tune spectra (low intensity bunches)

Tune spectrum: FFT[Dipole moment per turn]



Difference signal U_Δ : beam 'offset' fluctuations \bar{x}



local dipole moment: $I(\phi)\bar{x}_k(\phi) \propto e^{-i\chi\phi/\phi_b} e^{-iQ_k t}$

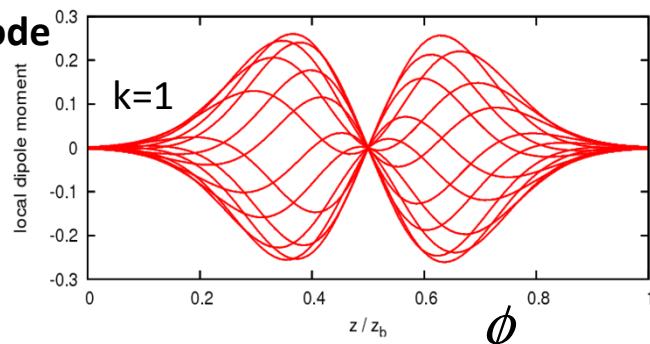
R. Singh, P. Forck, et al.

Synchrotron satellites (synchrotron tune Q_s):

$$Q_k = Q_0 + kQ_s \quad \text{or} \quad \Delta Q_k = kQ_s$$

Tune spread in a rf bucket (rms bunch length σ):

$$\delta Q \approx kQ_s \frac{\sigma_\phi^2}{4}$$



tune spectra -> measure and control the tune (e.g. to avoid resonances)

Tune spectra for high intensities: Theory

Simplified transverse particle equation of motion (e.g. D. Möhl, 1995): $\bar{x}(\phi, s) = \int x \rho(x, y, \phi, s) dx dy$

$$x''(\phi, s) + \frac{Q^2}{R^2} x - \frac{2Q_{x0}\Delta Q_x^{sc}}{R^2} (x - \bar{x}) = \frac{2Q_x \Delta Q_x^c}{R^2} \bar{x}$$

(bare betatron tune) (space charge) (image current/charges)

Longitudinal (synchrotron) oscillations: $\phi''(s) = \hat{\phi} \sin(Q_{s0}s / R)$

SIS-18/100: $q_{sc} = 10-20$
CERN PSB/PS: $q_{sc} \geq 100$

Coherent shift due to image currents: $\Delta Q^c \approx \frac{a^2}{b^2} \Delta Q^{sc}$
-> important for thick beams! (a: beam radius, b: pipe ra

space charge and image current parameters: $q_{sc} = \frac{\Delta Q^{sc}}{Q_{s0}}$ $q_c = \frac{\Delta Q^c}{Q_{s0}}$

head-tail tune shifts (longitudinal **airbag distribution**):

$$\Delta Q_k = -\frac{\Delta Q_{sc}}{2} \pm \sqrt{(\Delta Q_{sc}/2)^2 + (kQ_s)^2}$$

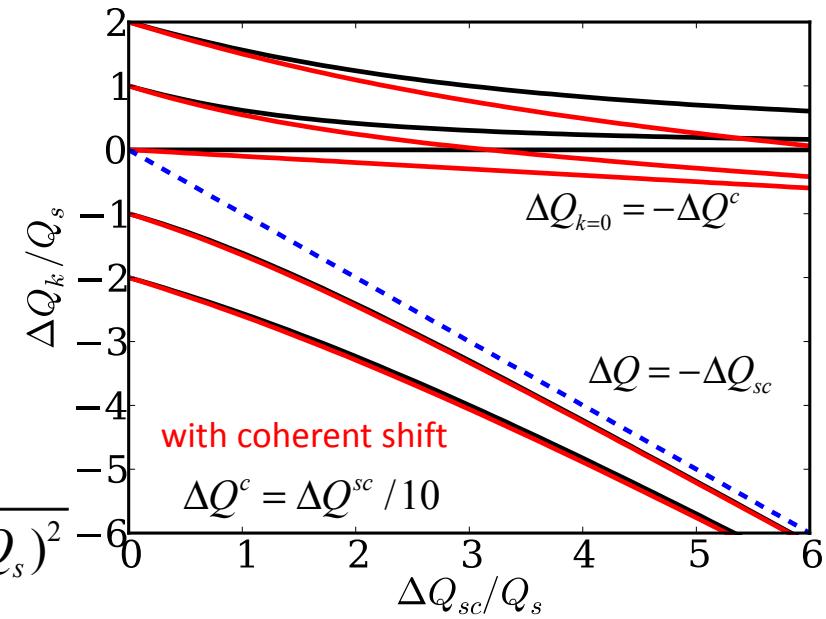
$$\Delta Q^{sc} = 0 : \quad \Delta Q_k = kQ_s^2$$

$$\Delta Q^c > 0 :$$

$$\Delta Q_k = -\frac{\Delta Q_{sc} + \Delta Q_c}{2} \pm \sqrt{(\Delta Q_{sc} - \Delta Q_c)^2 / 4 + (kQ_s)^2}$$

Blaskiewicz, Phys. Rev. ST Accel. Beams (1998)

Boine-F., Kornilov, Phys. Rev. ST Accel. Beams (2009), Burov (2009), Balbekov (2009)



Tune spectra from real bunches

The airbag model shows good agreement with tune spectra for Gaussian bunches (PATRIC simulations and SIS-18 measurements) -> V. Kornilov, O. Boine-F., arXiv (2012)

Modified airbag model:

$$\Delta Q_k = -\frac{\Delta Q_{sc}}{2} \pm \sqrt{(\Delta Q_{sc}/2)^2 + (kQ_{s1})^2}$$

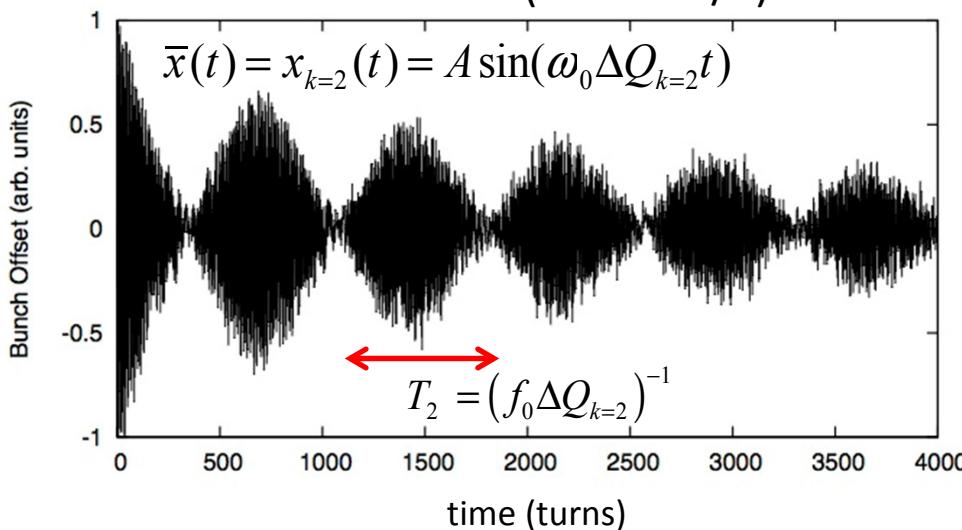
$$\Delta Q_{sc} = \frac{qI_p R}{4\pi\epsilon_0 c E_0 \gamma_0^2 \beta_0^3 \epsilon_x}$$

(sc tune shift in the bunch center)

$$\frac{Q_{s1}}{Q_{s0}} = \sqrt{1 - \frac{\sigma_\phi^2}{2}}$$

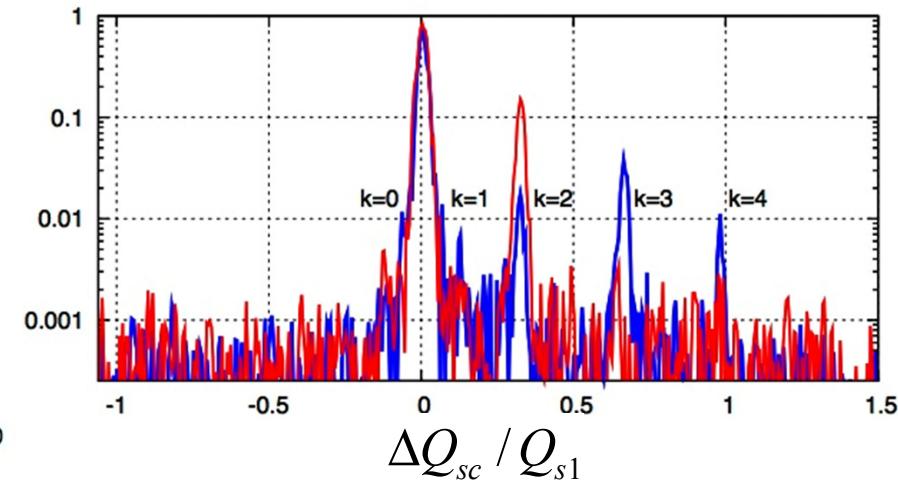
(longitudinal dipole tune)

Measured decoherence signal from a kicked Ar¹⁸⁺ bunch (100 MeV/u)



$$q_{sc} \approx 10, \quad q_c \ll 1$$

Resulting tune spectrum

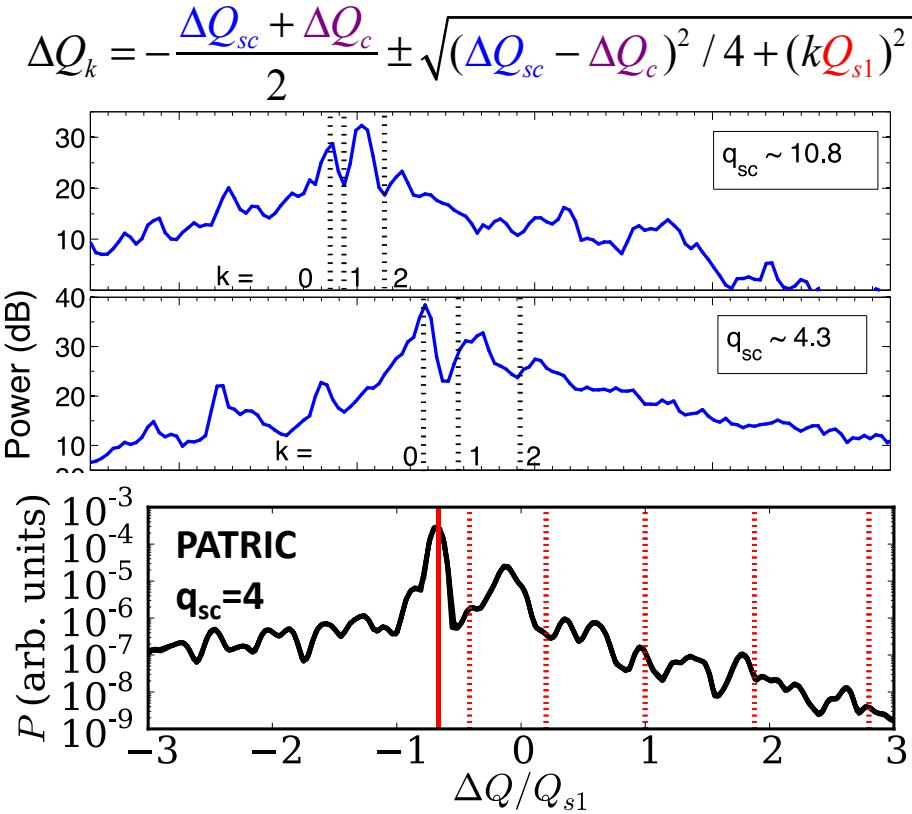


Thin beam ($a \ll b$): coherent tune shift can be neglected

Measured SIS-18 tune spectra (noise excitation)

11.4 MeV/u (injection energy), $0.3\text{-}1.0 \times 10^{10}$ N⁷⁺ bunches

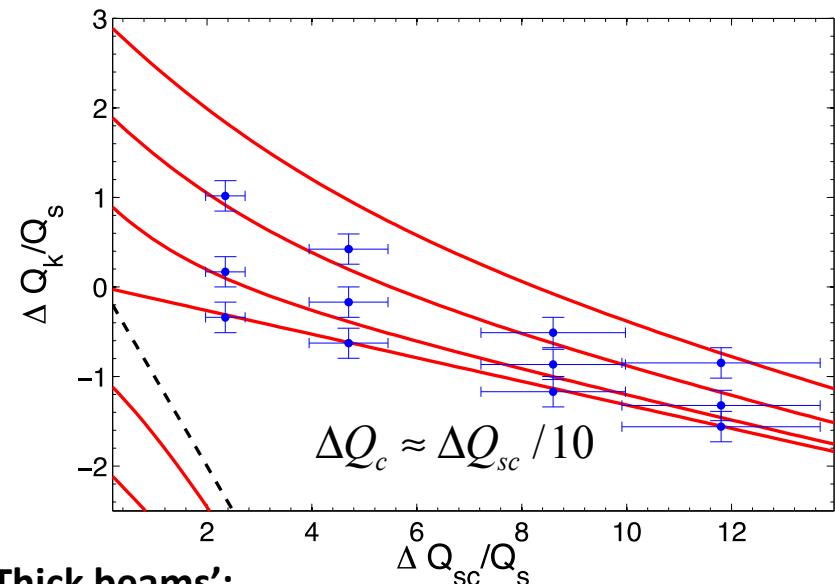
Thick beams ($a \approx b/3$): coherent tune shift is important



ΔQ_{sc} : from beam profile measurement (IPM)

Q_{s1} : from longitudinal Schottky signal

ΔQ_c : from the shift of the k=0 line.



'Thick beams':

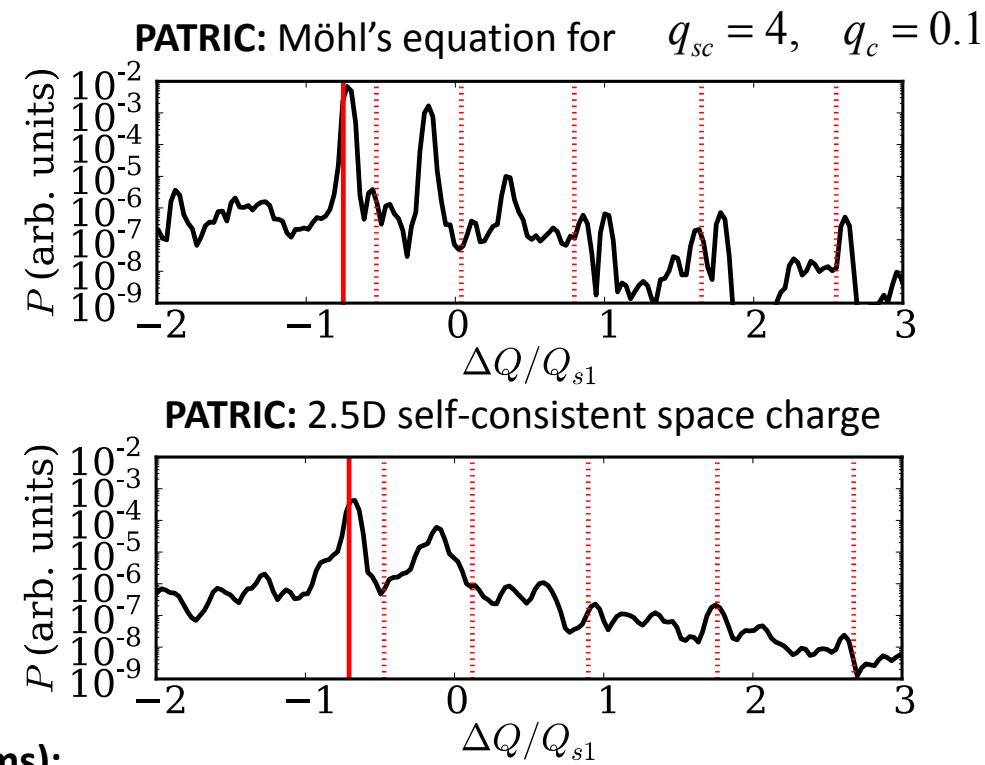
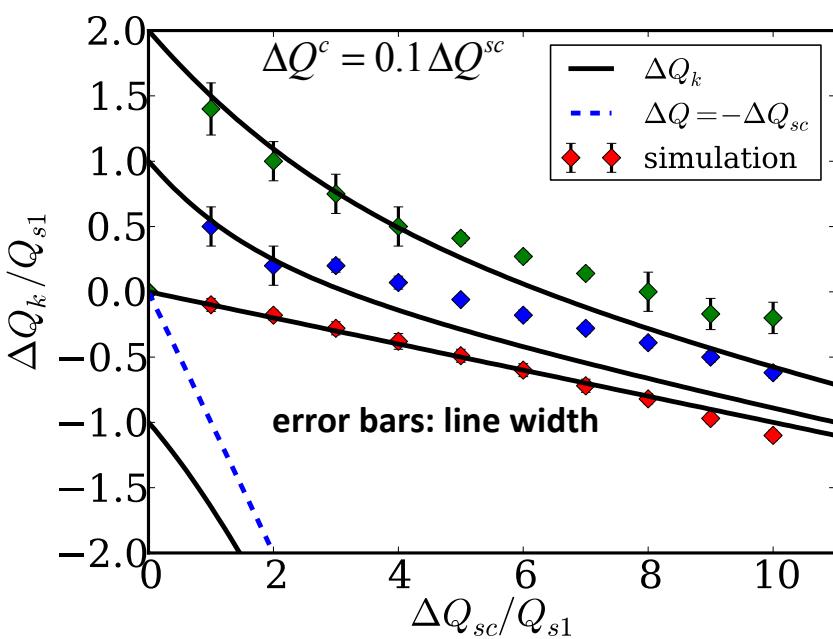
- deviation from the modified airbag model.
- broadening of the lines -> intrinsic Landau damping (Burov 2009, Balbekov 2009)

R. Singh, P. Forck, P. Kowina, O. Boine-F., et al., *Interpretation of tune spectra for high intensity*, to be published
-> see presentation by Oleksandr Chorniy

Computer tune spectra self-consistent space charge

PATRIC: 3D tracking with 2.5D self-consistent space charge solver, pipe with circular boundary.

$$\Delta Q_k = -\frac{\Delta Q_{sc}}{2} \pm \sqrt{(\Delta Q_{sc}/2)^2 + (kQ_{s1})^2}$$



In the presence of a coherent tune shift (thick beams):

- Agreement with the measurements (positions and widths of the lines)
- Similar deviations from the modified airbag model.
- Self-consistent space charge model is required !

Computer tune spectrum in dual rf buckets

Dual rf voltage: $V_{rf}(\phi) = V_0(\sin(\phi) + \frac{1}{2}\sin(2\phi))$

$$\lambda(\phi) = \frac{dN}{d\phi}$$

bunch in a single rf bucket

$$\Delta Q_{sc}(\phi) = \Delta Q_{sc,max} \frac{\lambda(\phi)}{\lambda(0)}$$

flattened bunch in a dual rf bucket

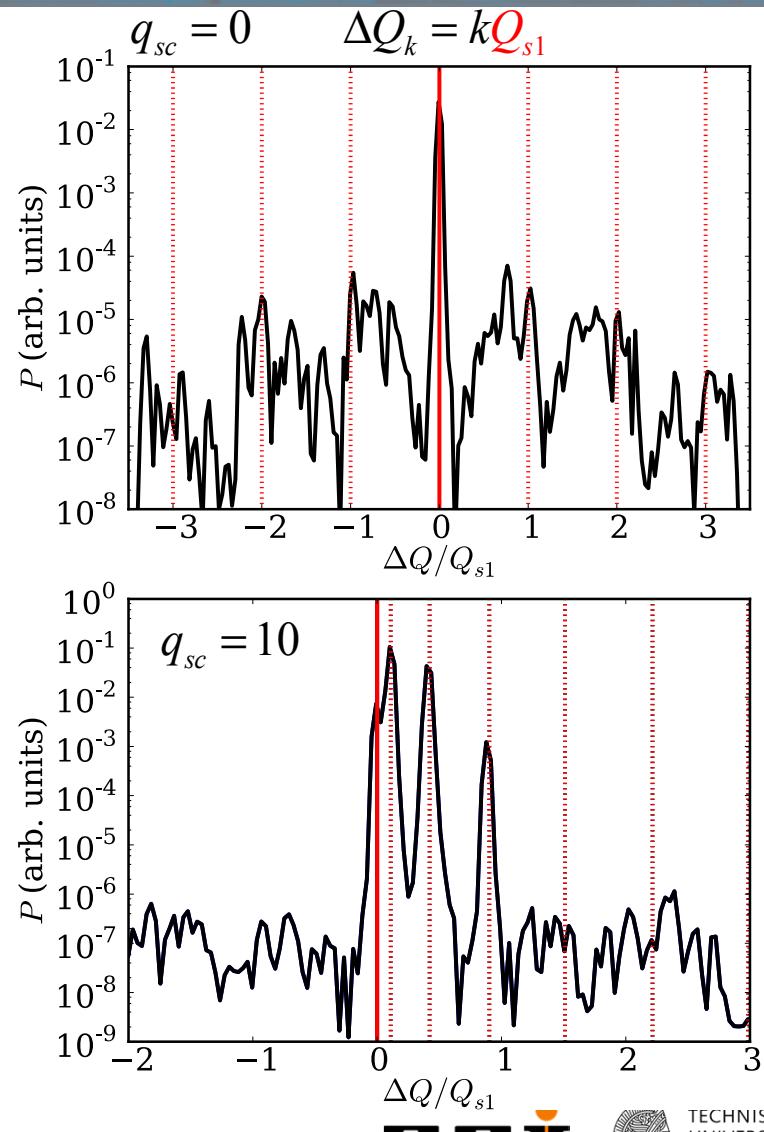
$$\frac{Q_{s1}}{Q_{s0}} \approx \sqrt{\frac{10}{7}} \sigma_\phi \quad (\text{synchrotron dipole mode})$$

-> poster by Monika Mehler

$$\Delta Q_k = -\frac{\Delta Q_{sc}}{2} \pm \sqrt{(\Delta Q_{sc}/2)^2 + (kQ_{s1})^2}$$

Preliminary results for dual rf:

- Good agreement with modified airbag model (thin beams).
 - No space charge: strongly damped head-tail modes
 - Space charge: undamped ($k < 4$) head-tail modes
- > Compare to observed head-tail instability thresholds in the CERN PSB (V. Kornilov, et al.)



Conclusions

- For the optimization of the GSI and FAIR rings, including **the important interplay of incoherent and coherent collective effects**, different **simulations codes** are employed, in combination with **dedicated beam physics experiments in the GSI SIS-18 and in the CERN PSB/PS**.
- For the interpretation of tune spectra in the presence of space charge we successfully used the PATRIC simulation code together with a modified 'airbag' model for head-tail oscillations.
-> direct measurements of the **incoherent space charge tune shift** from the **coherent tune spectrum** !
- Measurements, simulations (with self-consistent sc) indicate that **the coherent tune shift caused by the image currents (-> thick beams) in the pipe can strongly increase the damping of head-tail modes**.
- From simulations: **In a dual rf bucket the tune spectrum shows pronounced (low-order) head-tail modes** in the presence of space charge (similar indications from CERN PSB experiments)

For the FAIR synchrotrons (SIS-18/100) -> additional intrinsic damping of head-tail modes expected.

Next step (for the FAIR synchrotrons):

- Systematic **simulation+experimental studies** of head-tail modes, instability thresholds with space charge, including the pipe impedance, **dual rf buckets and dampers/octupoles**.