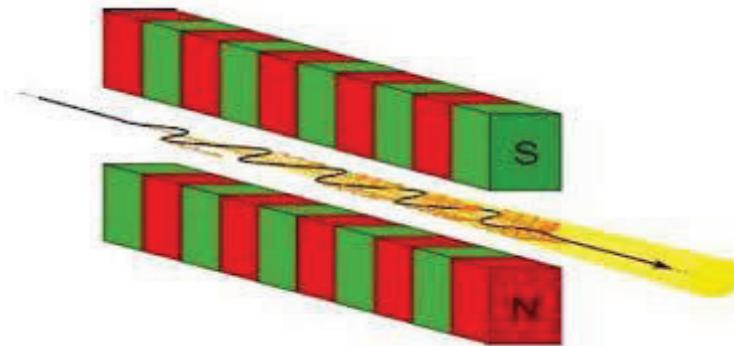


# The beamlet method

Methods to improve FEL output from “dirty” beams



James Henderson, Lawrence Campbell and Brian McNeil

# Outline

- Motivation Plasma accelerators
- Potential beamlet solution modulate and disperse
- Chicanes to maintain resonant interaction
- Simplified Model
- The beamlet method
- Conclusion

# Motivation

- Plasma accelerator have accelerating gradients of  $10^3$  greater than conventional accelerators, potential for “table top” FELs.

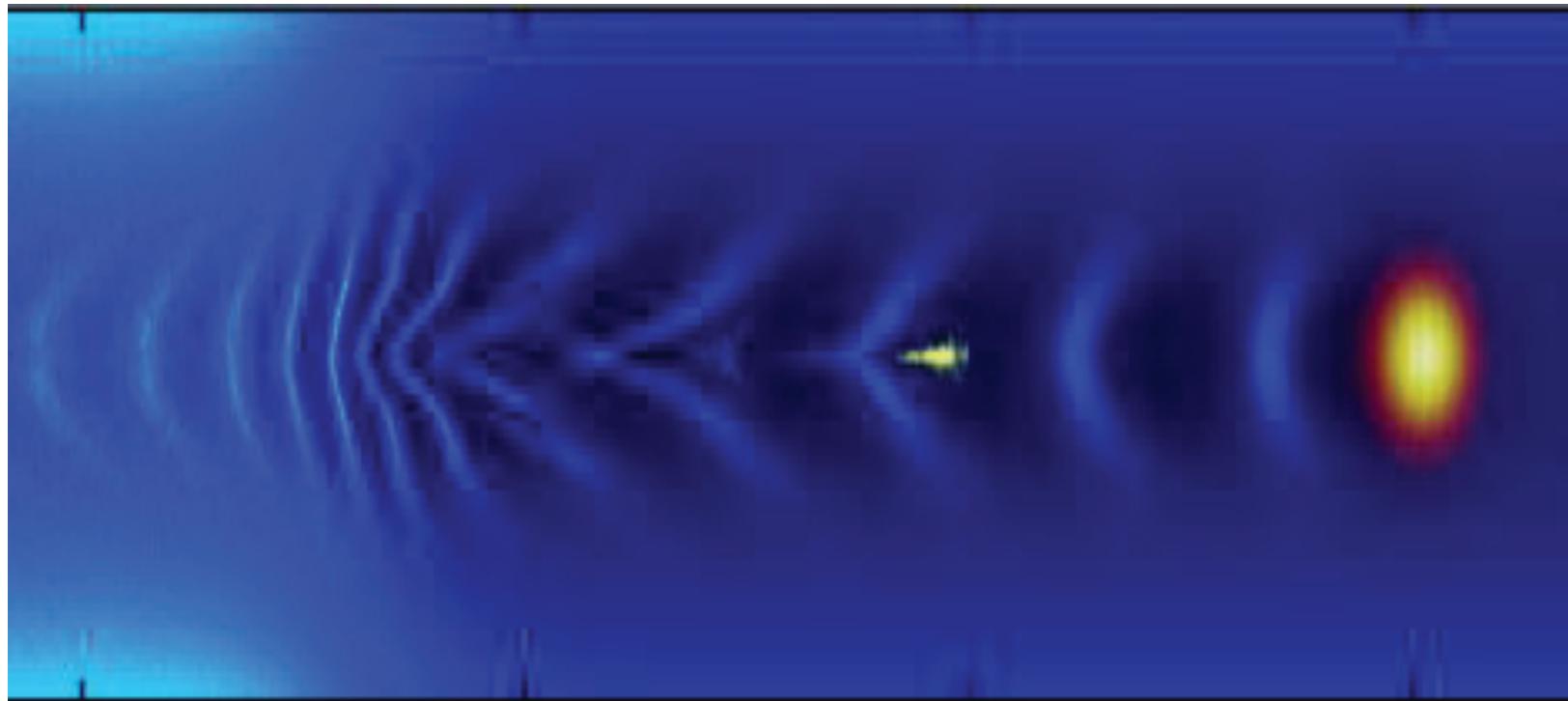
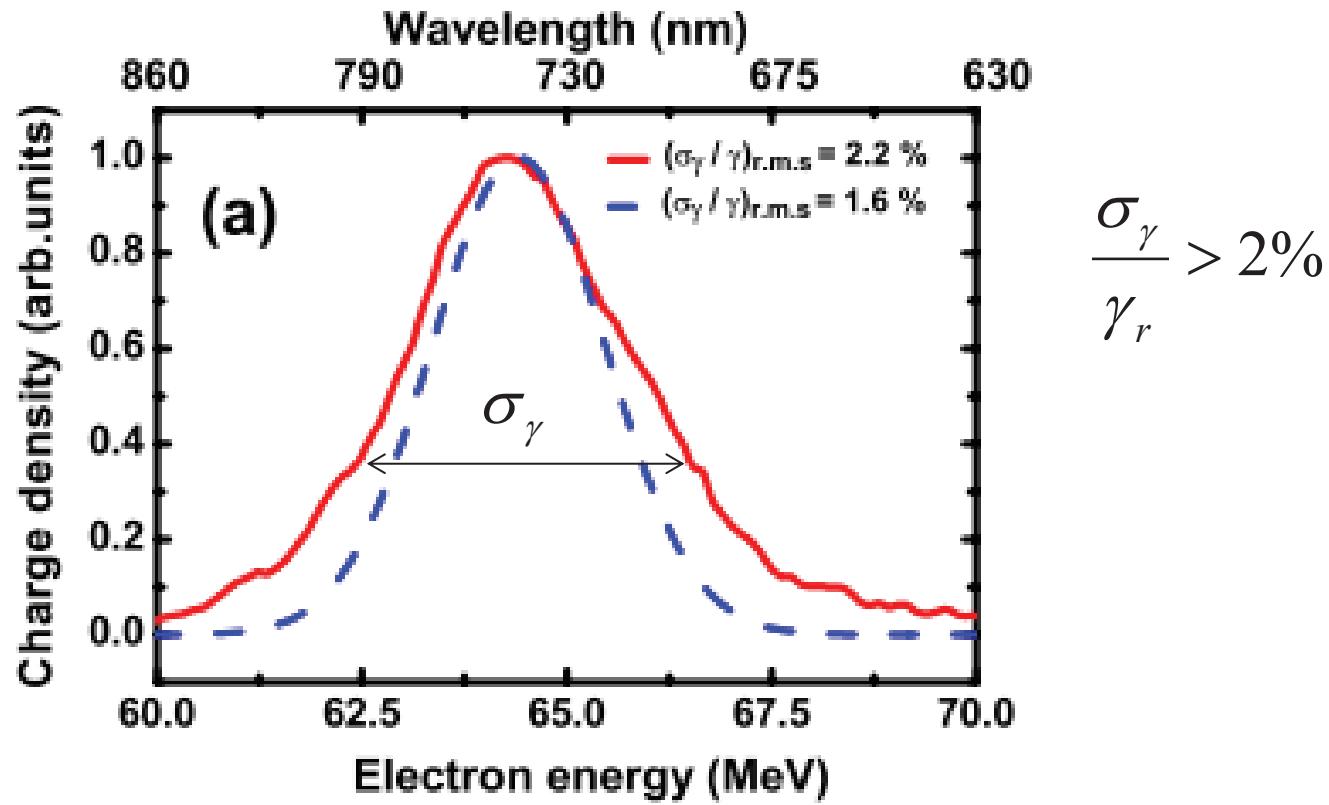


Image credit - <https://www2.physics.ox.ac.uk/research/plasma-accelerators>  
Front page image – <http://lpap.epfl.ch/page--en.html>

# Motivation

- Plasma accelerators have accelerating gradients of  $10^3$  greater than conventional accelerators, potential for “table top” FELs.
- However electron pulses from plasma accelerators typically have a large energy spread



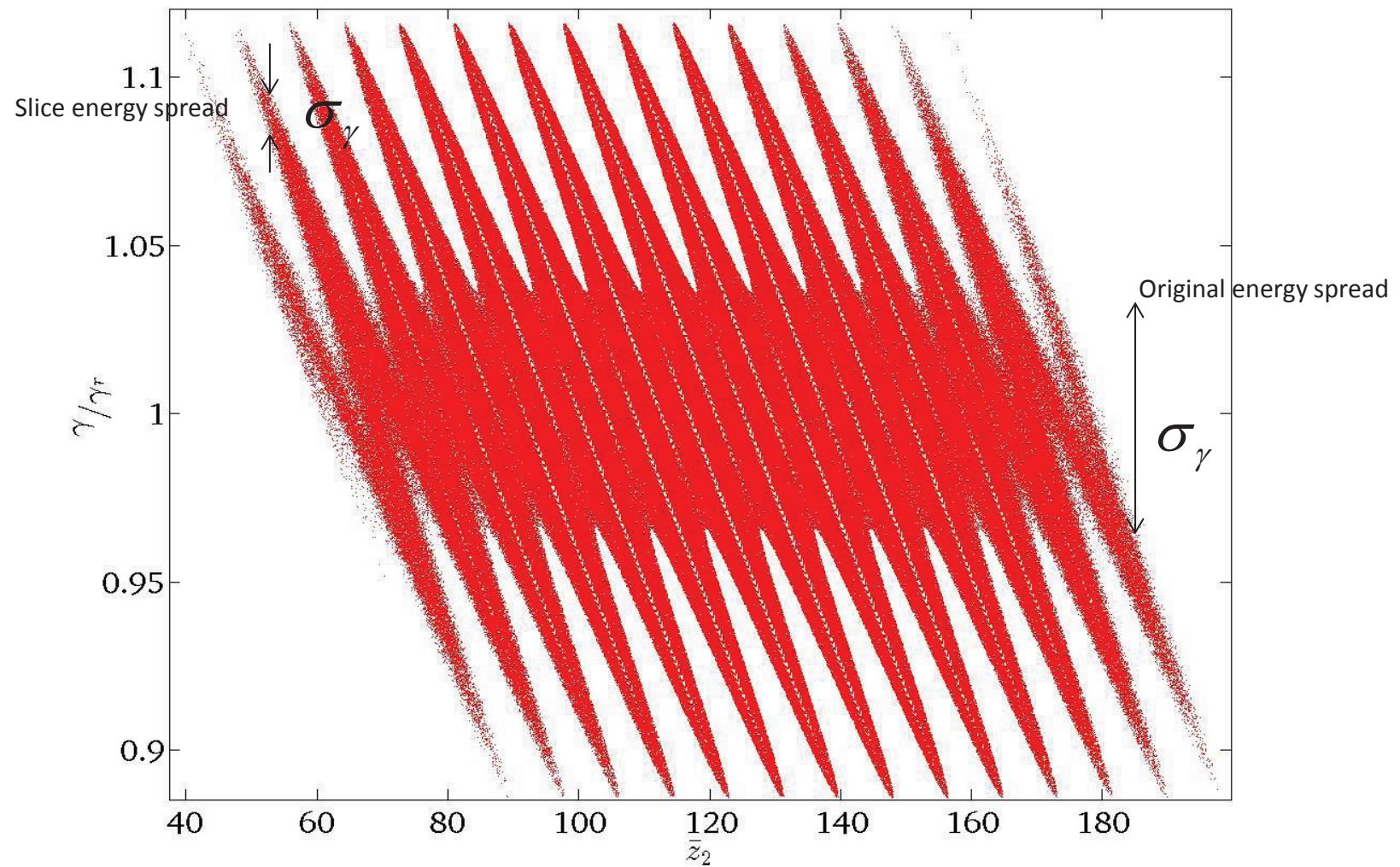
For an FEL we require

$$\frac{\sigma_\gamma}{\gamma_r} < \rho \quad \text{where} \quad \rho = 10^{-2}$$

A method of determining narrow energy spread electron beams  
from a laser plasma wakefield accelerator using undulator radiation  
J. G. Gallacher, et. al. PHYSICS OF PLASMAS **16**, 093102 2009

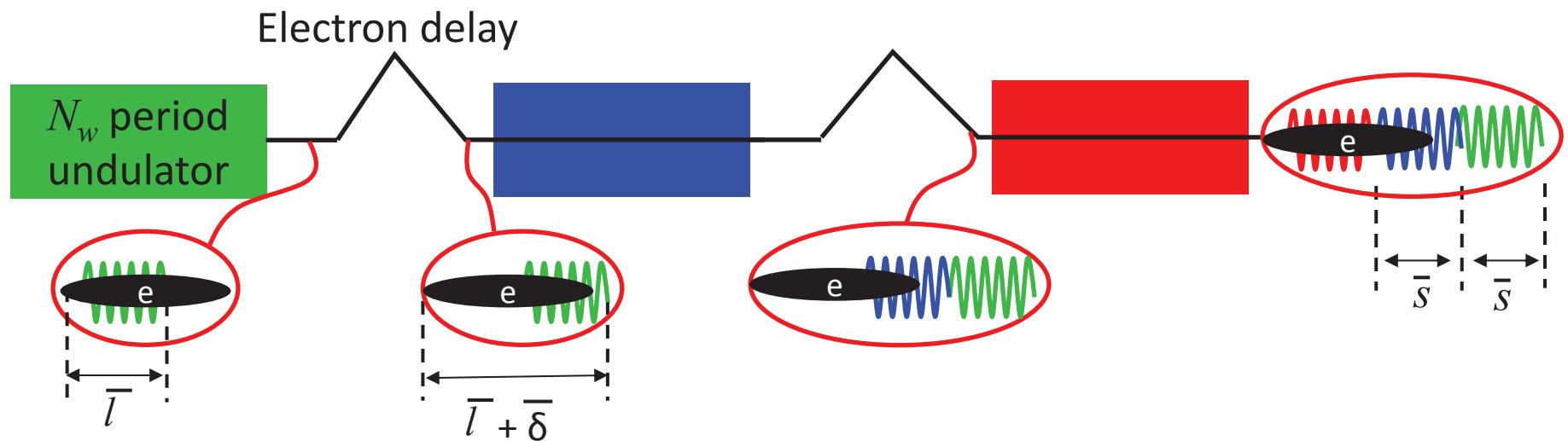
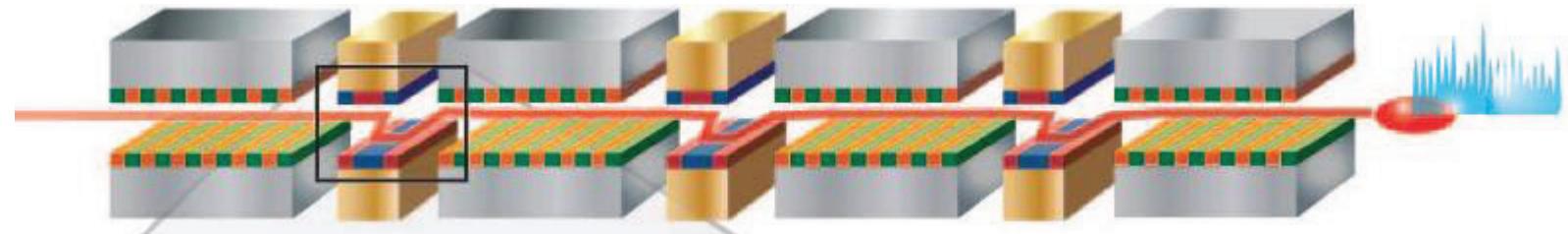
# Potential Solution

- Beamlet method – based on EEHG technique
- Electron Pulse is modulated and then dispersed



# Chicanes

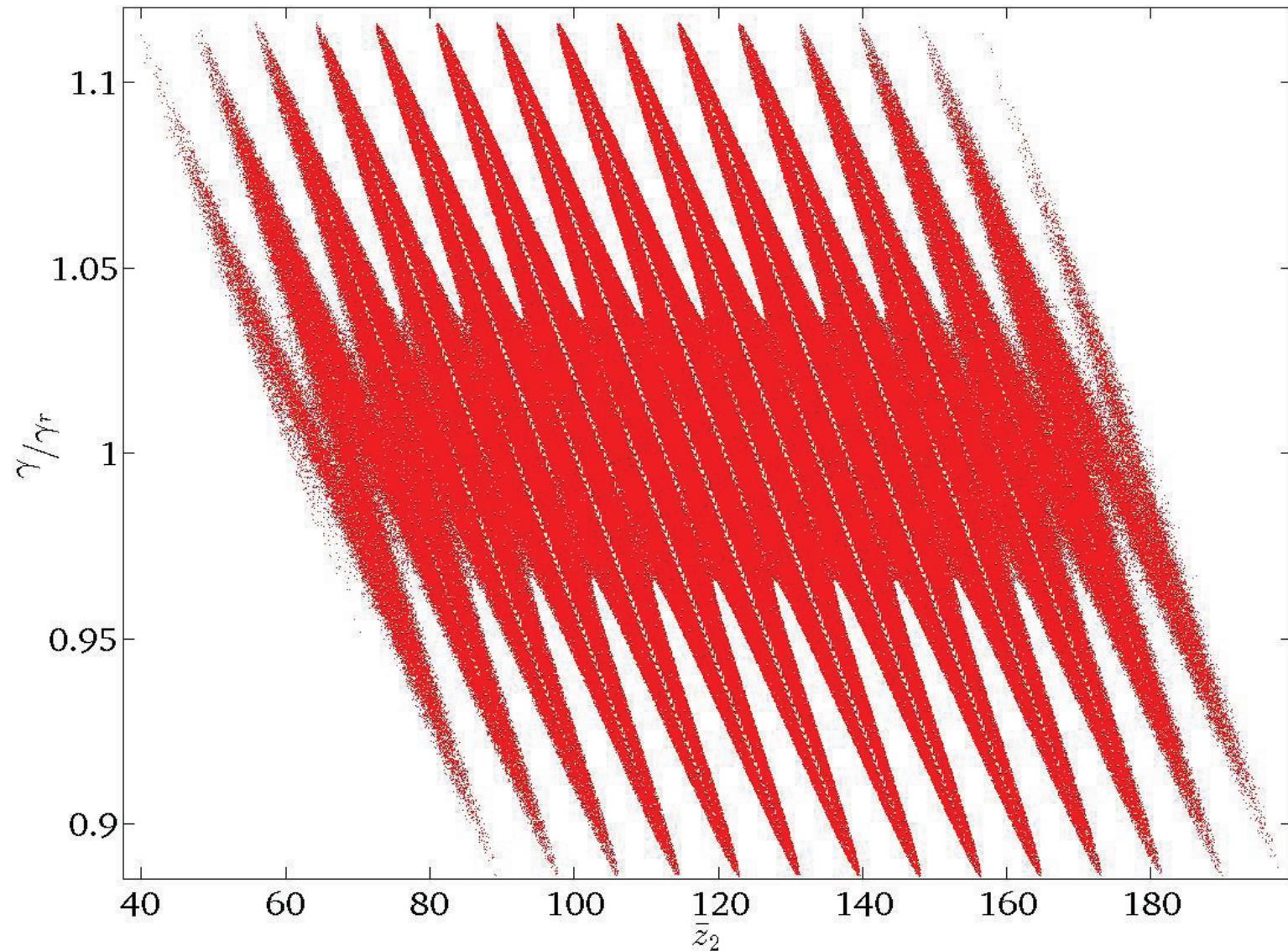
Pass radiation between beamlets using a undulator-chicane lattice



$$\bar{s} = \bar{l} + \bar{\delta}$$

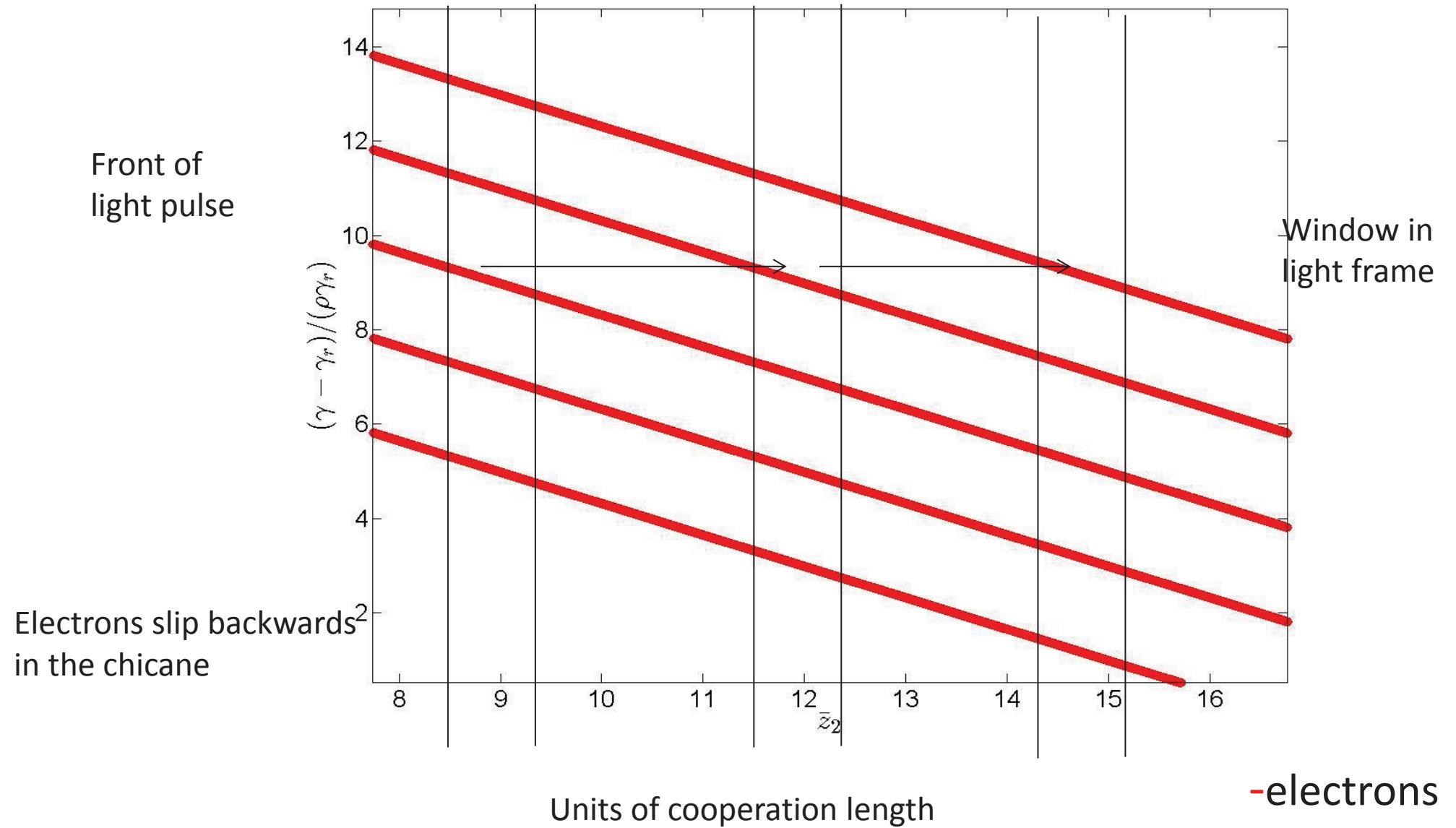
$$\frac{\Delta\omega}{\omega_r} = \frac{4\pi\rho}{\bar{s}}$$

# Real Beamlets

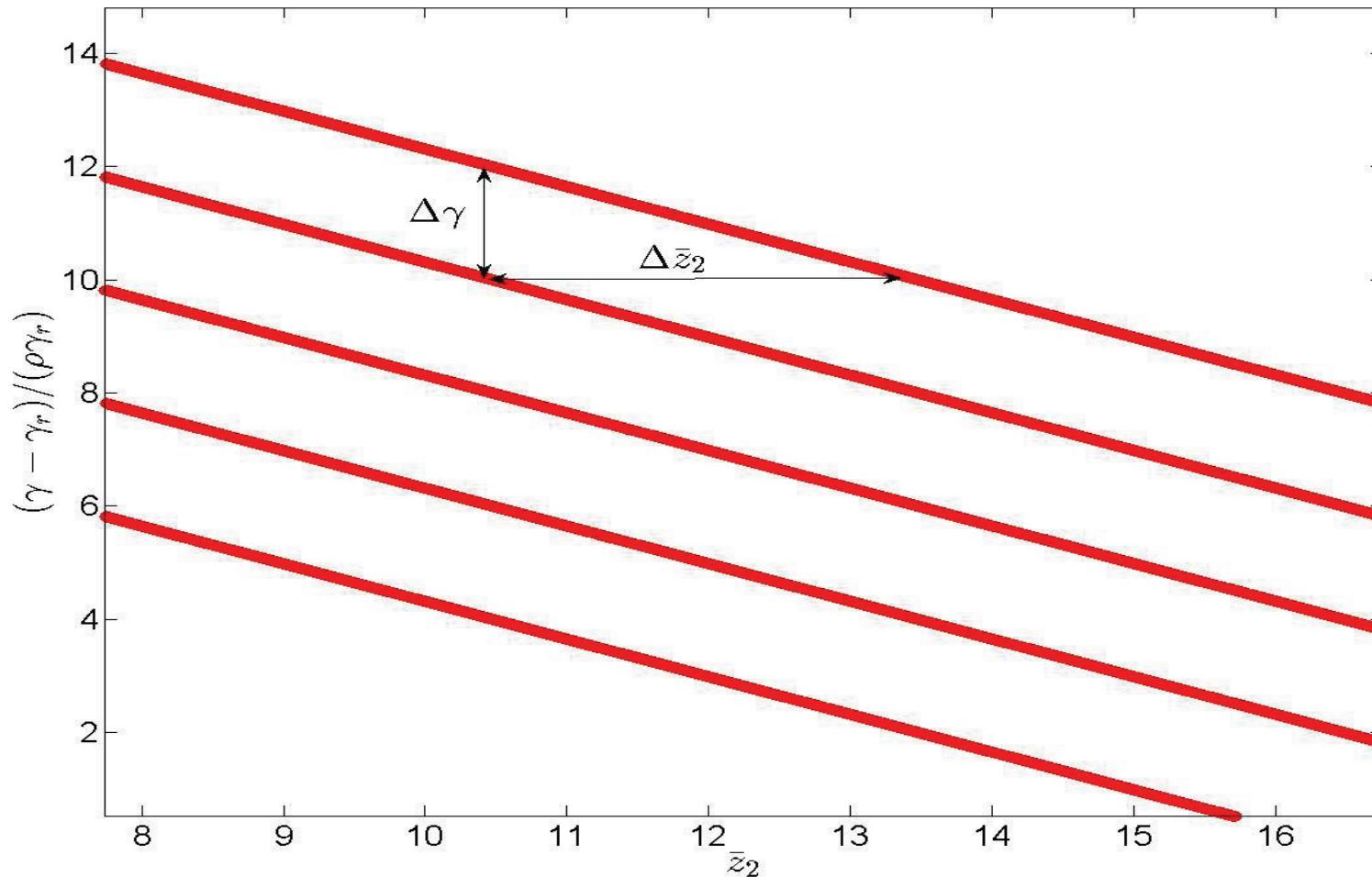


# Simplified Model

Approximating the beamlets



# Simplified Model



lattice and the resonant frequency difference of the beamlets

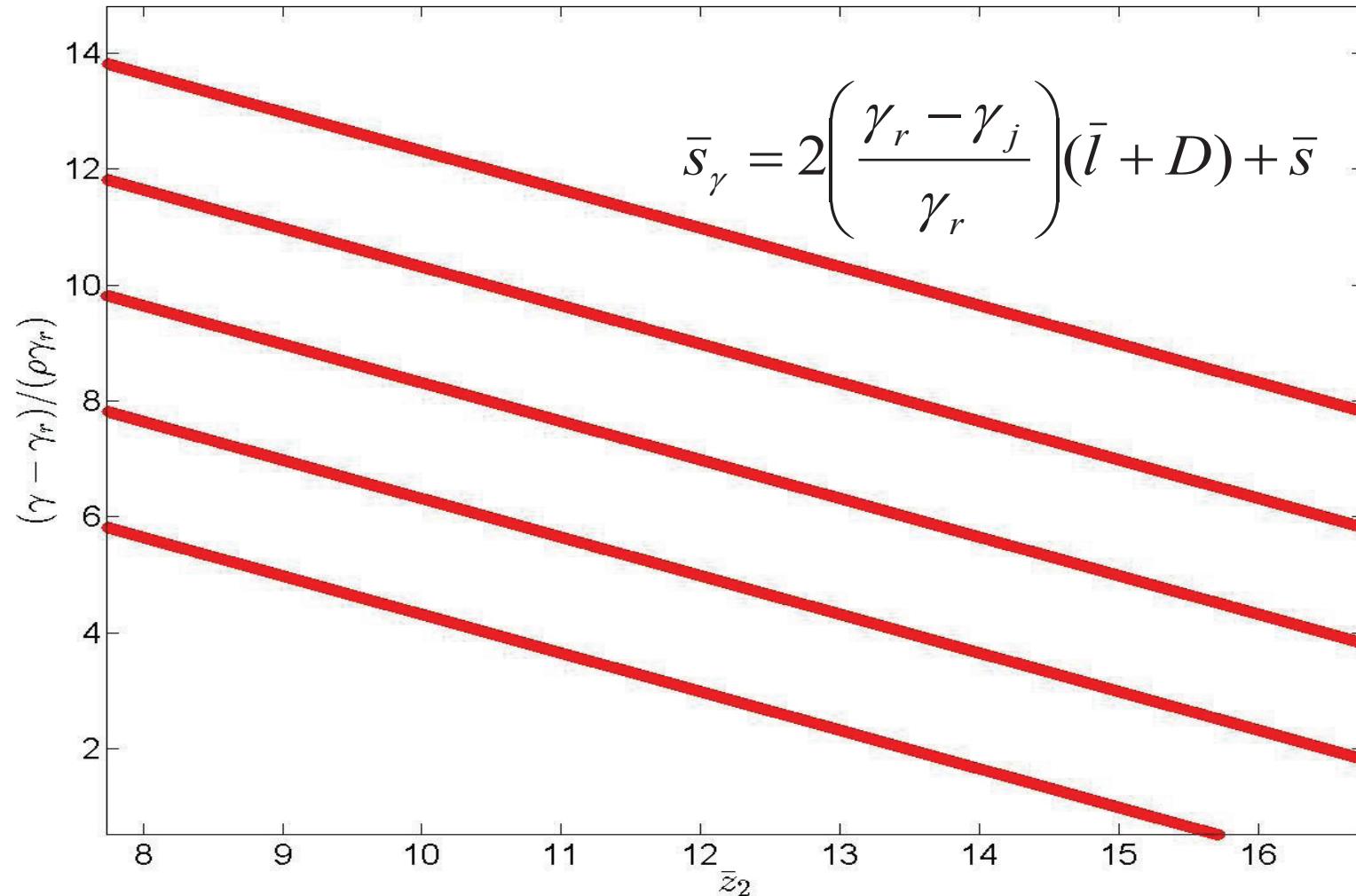
$$\frac{\Delta\omega_{\text{beamlet}}}{\omega_r} = 2 \frac{\Delta\gamma}{\gamma_r}$$

$$\frac{\Delta\omega_{\text{modal}}}{\omega_r} = \frac{4\pi\rho}{\bar{s}}$$

$$\Delta\omega_{\text{modal}} = \Delta\omega_{\text{beamlet}}$$

# Potential Solution

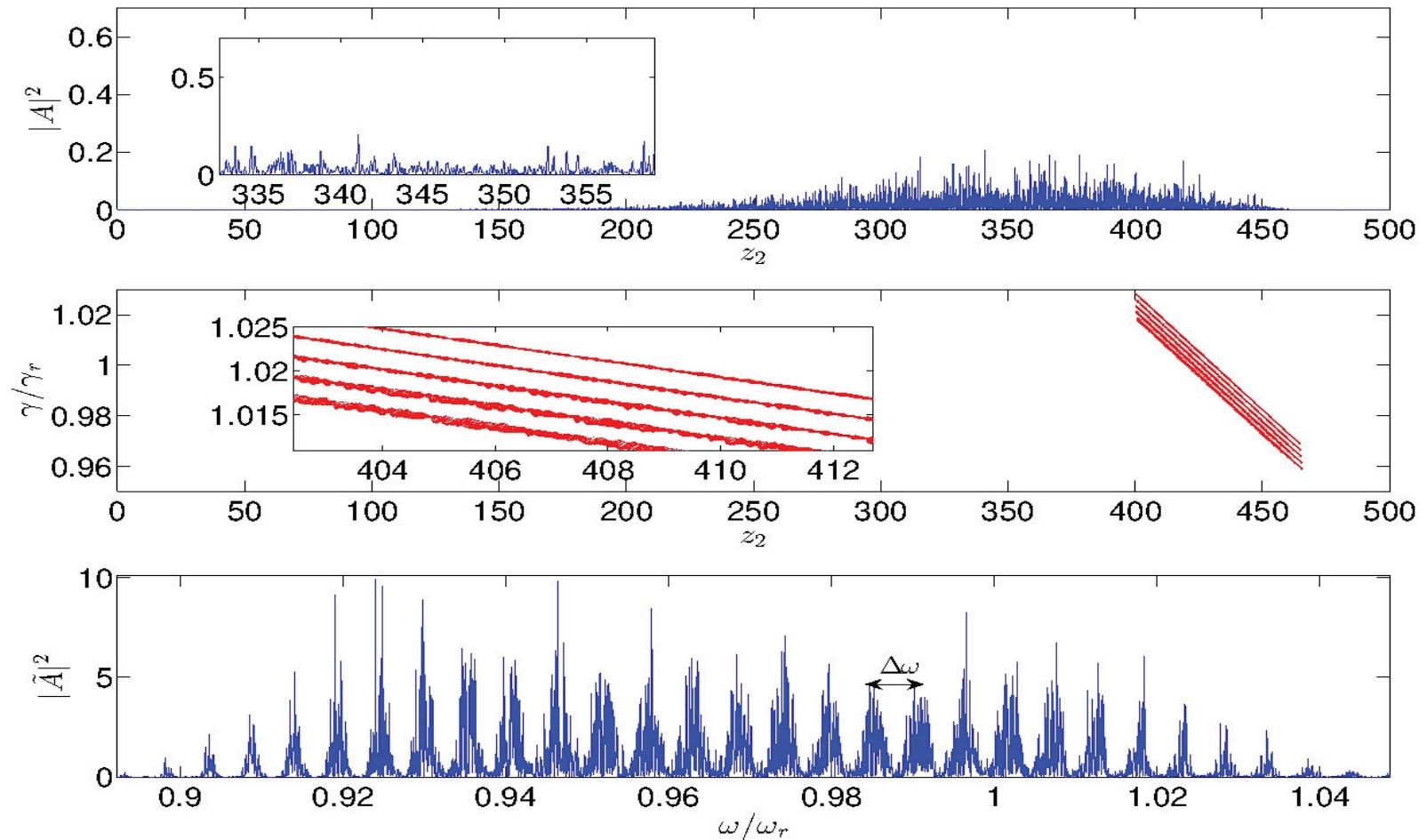
- Electron pulses with a large energy chirp exhibit an energy dependent slippage



- Electron of different speeds will take different paths from undulator and chicane
- Can lead to mismatching when passing radiation from beamlet to beamlet

# The beamlet method (D=0)\*

## Mode-locked beamlet modes

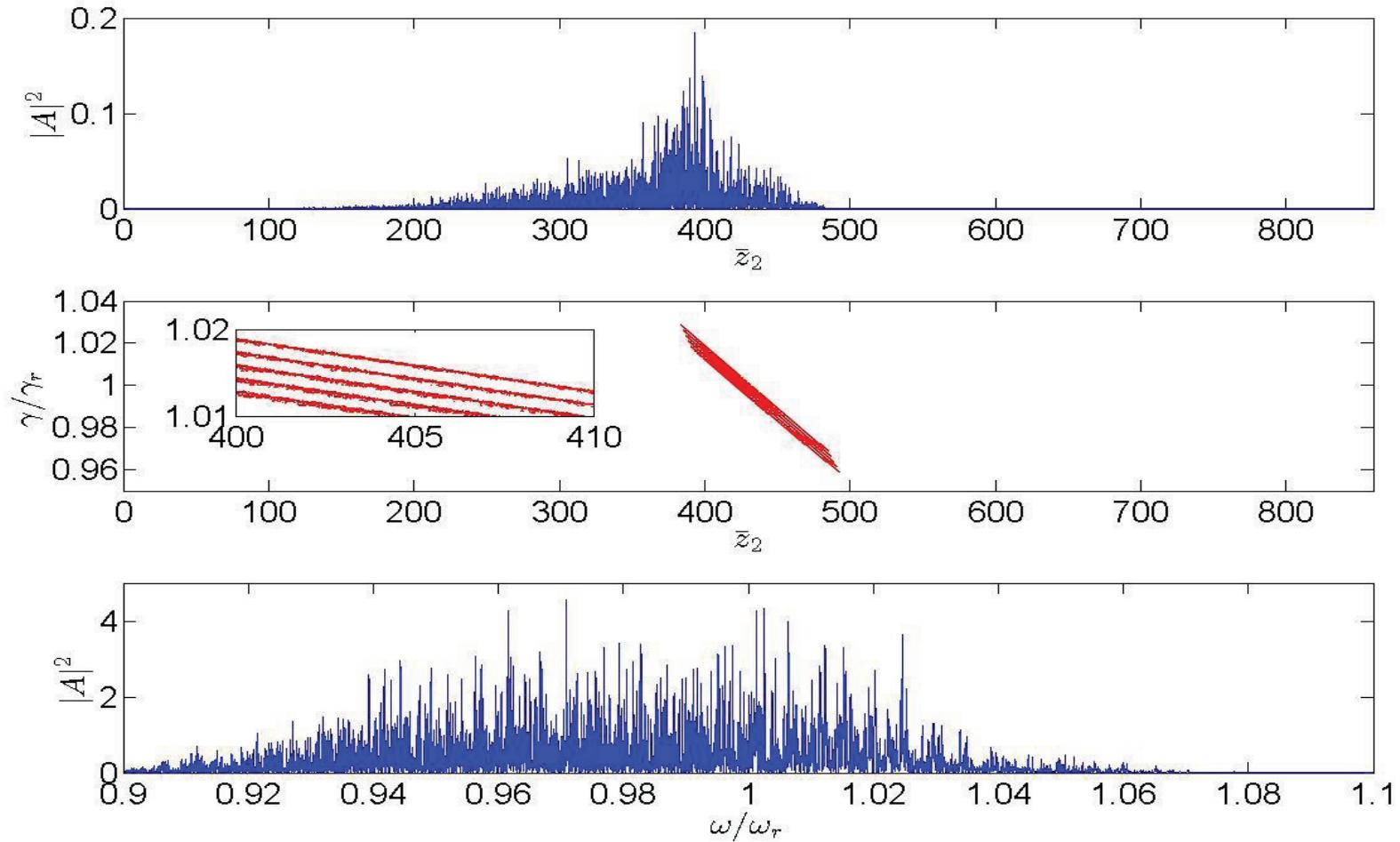


$$\bar{s} = 2.51, \bar{l} = 0.25, \bar{\delta} = 2.26, D = 0$$

$$\frac{\Delta\omega}{\omega_r} = \frac{4\pi\rho}{\bar{s}}$$

\* James Jones' Ref

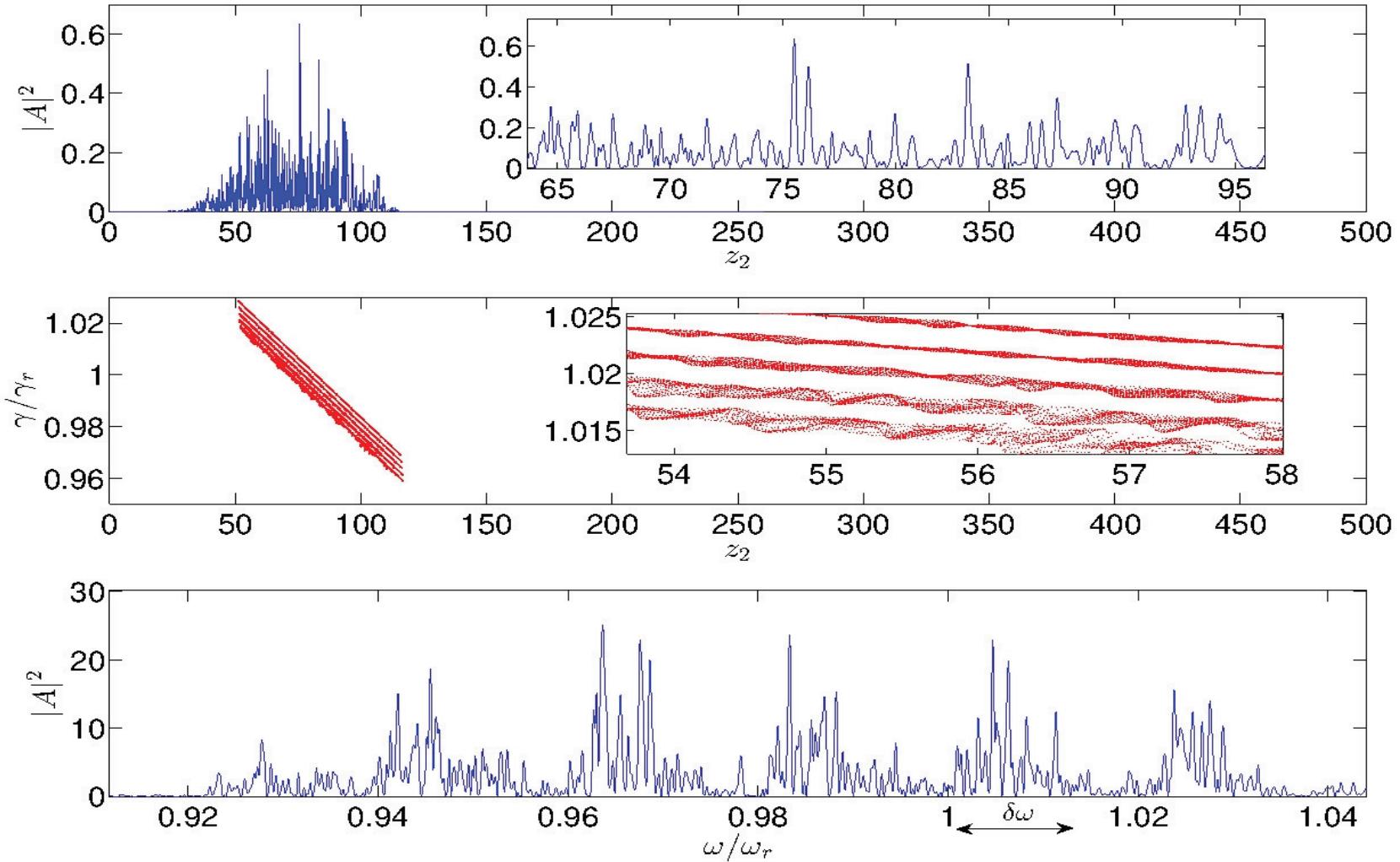
# The beamlet method ( $D>0$ )



$$\bar{s} = 2.51, \bar{l} = 0.25, \bar{\delta} = 2.26, D = 0.3$$

$$\bar{s}_\gamma = 2 \left( \frac{\gamma_r - \gamma_j}{\gamma_r} \right) (\bar{l} + D) + \bar{s}$$

# The beamlet method (D=0)



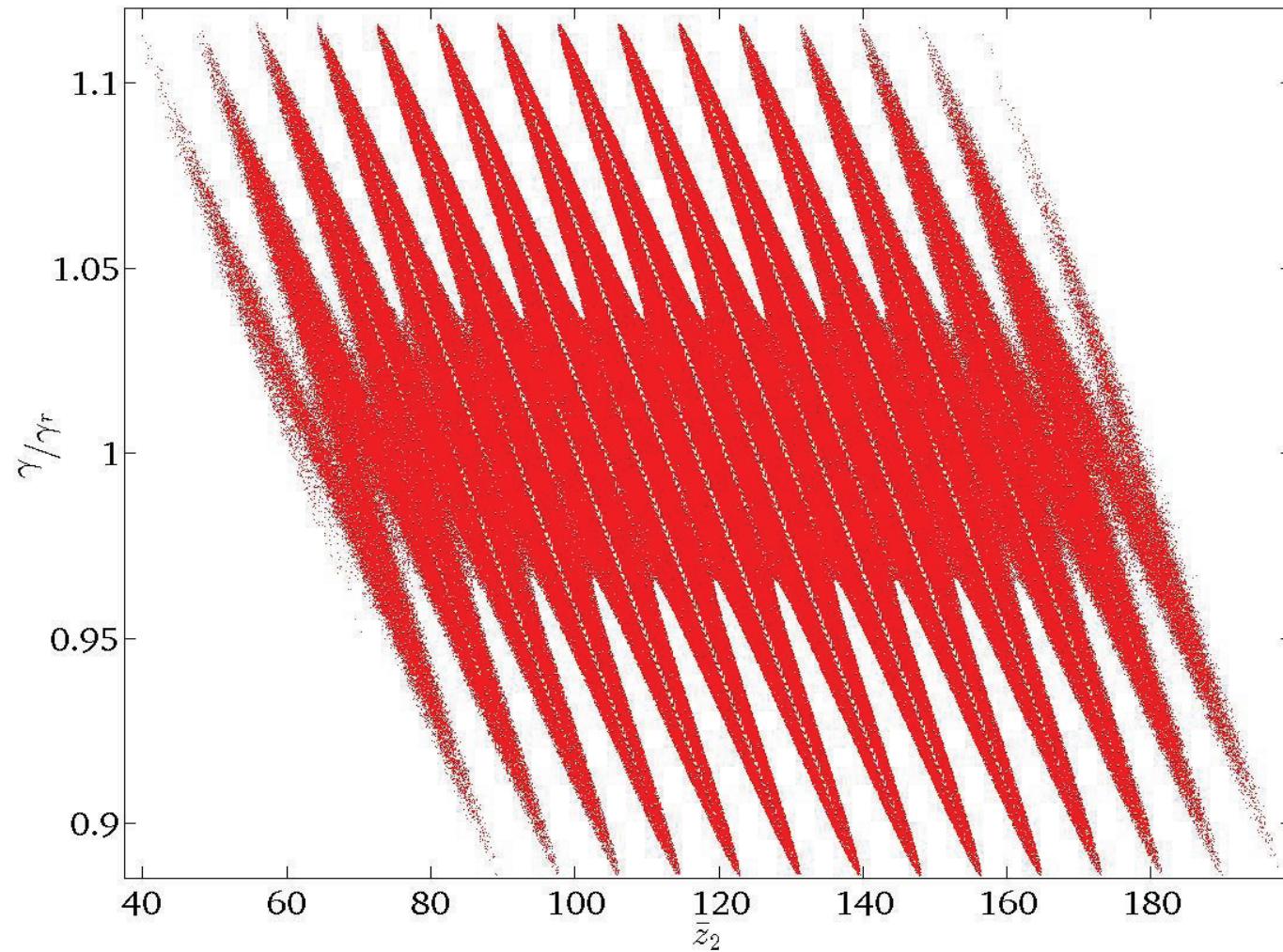
$$\frac{\Delta\omega}{\omega_r} = \frac{4\pi\rho}{\bar{\delta}}$$

$$\bar{s} = 2.51, \boxed{\bar{l} = 1.88}, \bar{\delta} = 0.63, \boxed{D = 0}$$

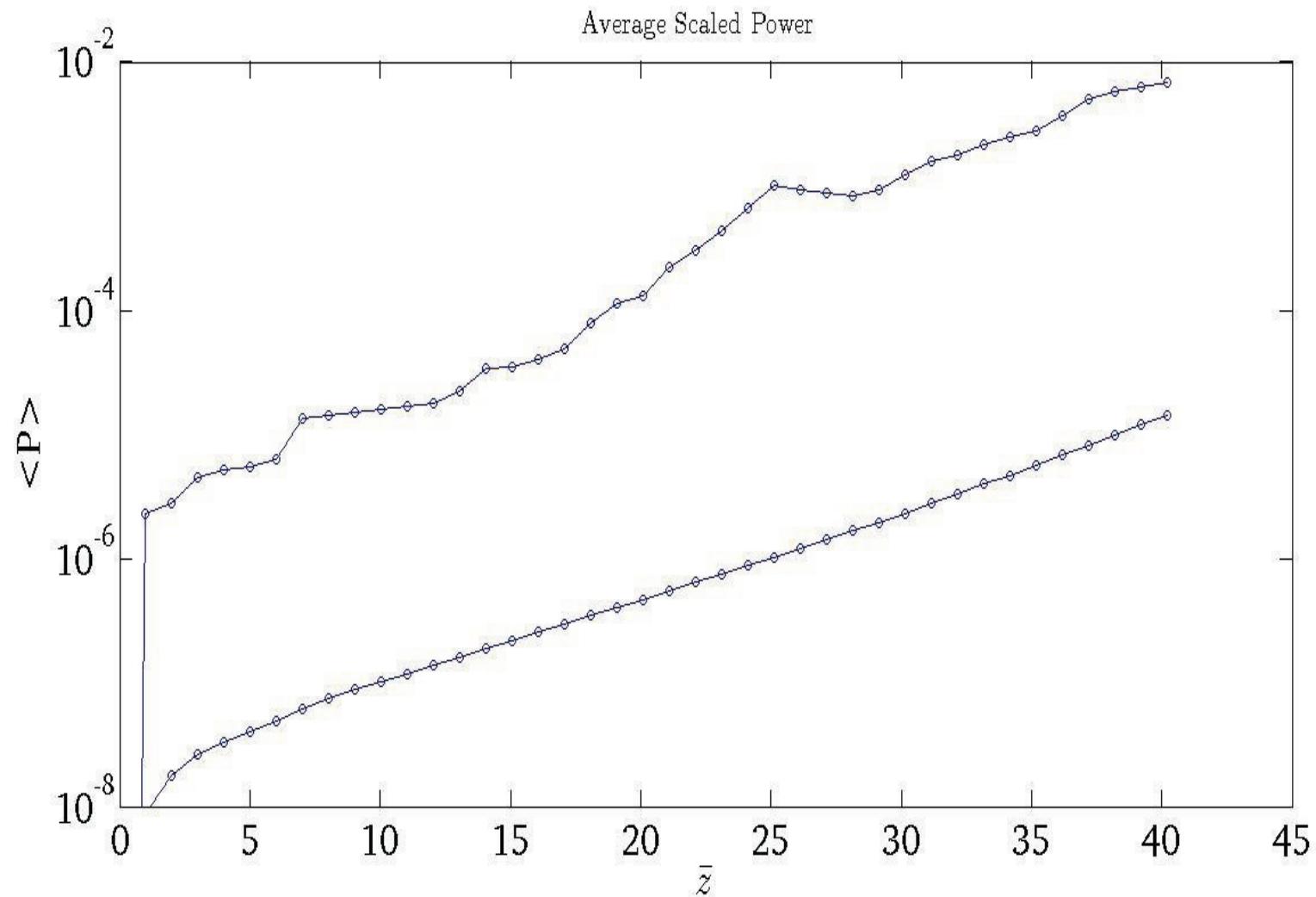
$$\bar{s}_\gamma = 2 \left( \frac{\gamma_r - \gamma_j}{\gamma_r} \right) (\bar{l} + D) + \bar{s}$$

# The beamlet method

- Electron Pulse is modulated and dispersed

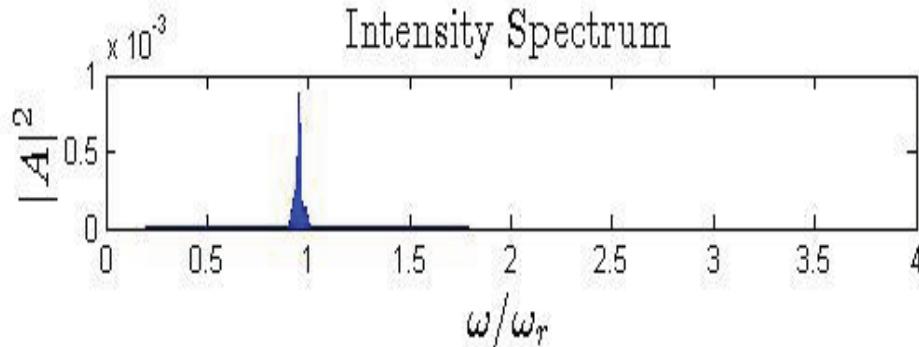
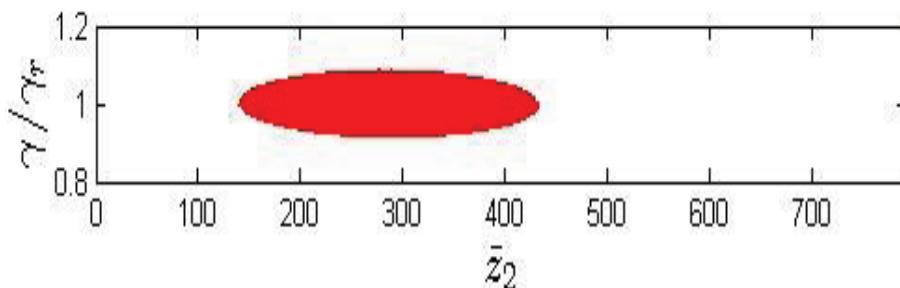
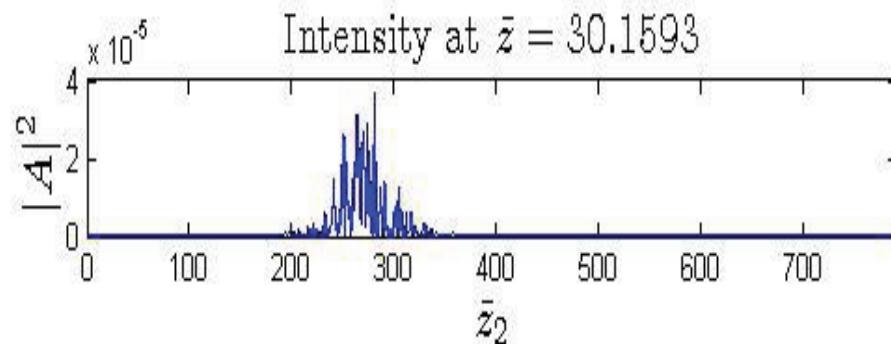


# The beamlet method

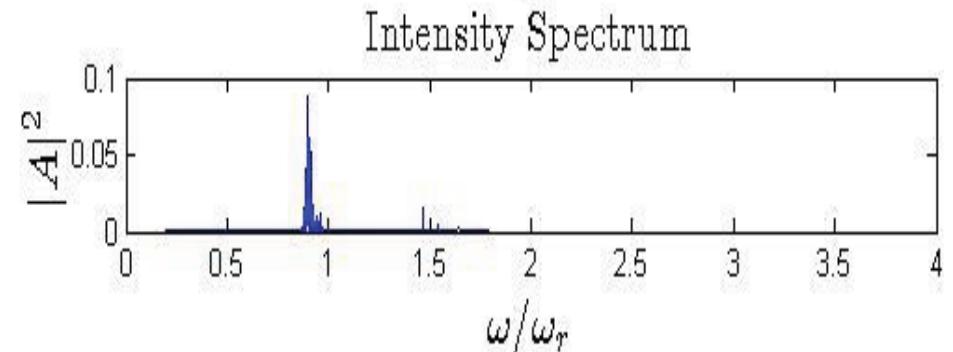
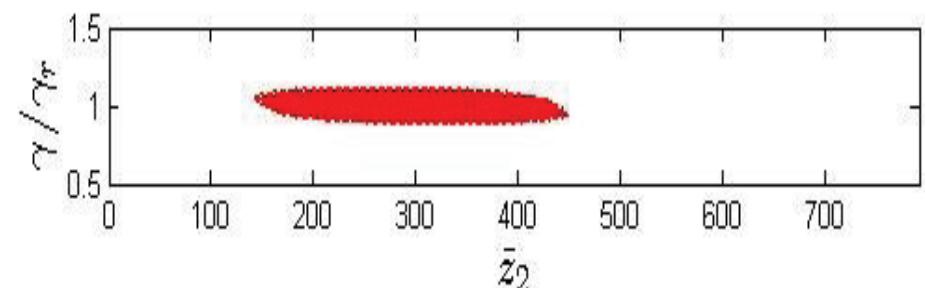
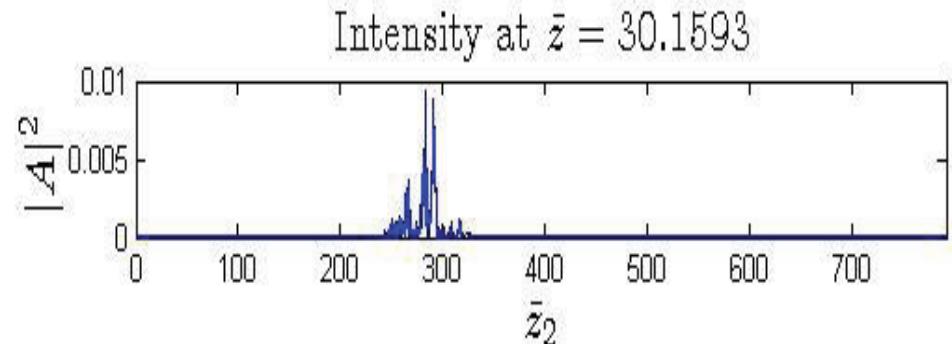


# The beamlet method

No Beamlets



Beamlets



# Conclusion

- Method to improve the FEL output was presented
- This is on-going research
- Results so far are promising with two-three order magnitude improvement

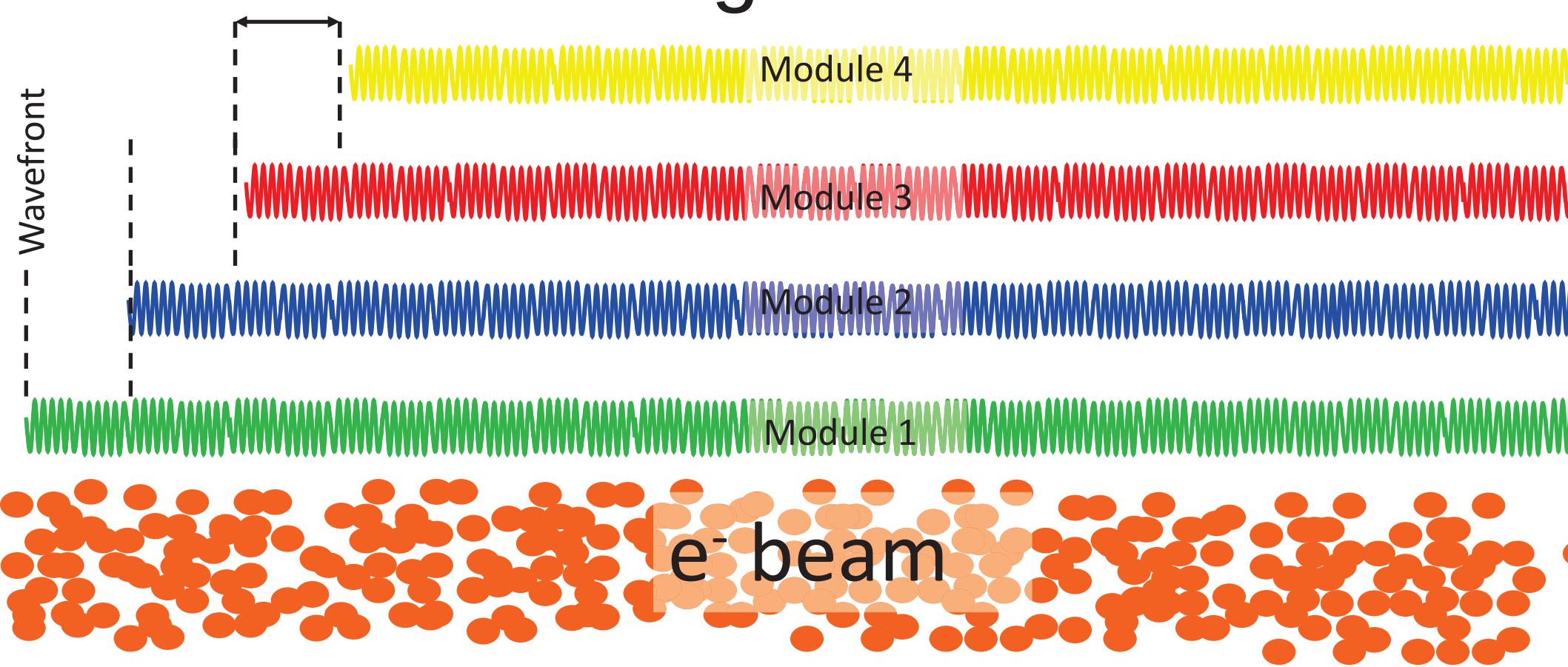
# Future development

- Optimization of the beamlet parameters
- Using a taper to cancel out the beamlet chirp
- Use of these techniques for multi electron pulse schemes

Thanks for listening.  
Any questions?

# •Mode generation

- Wavefront



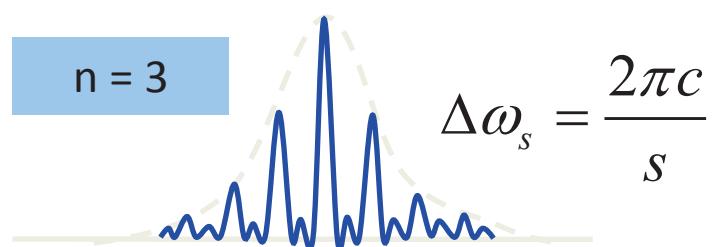
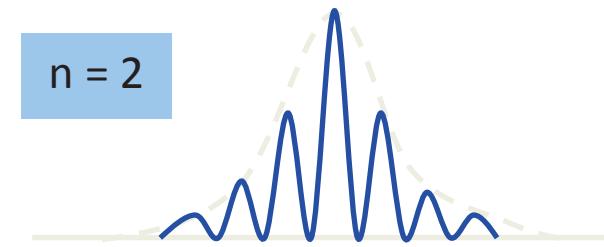
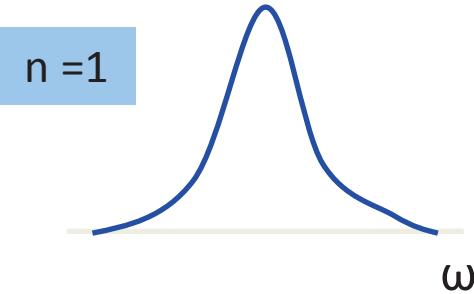
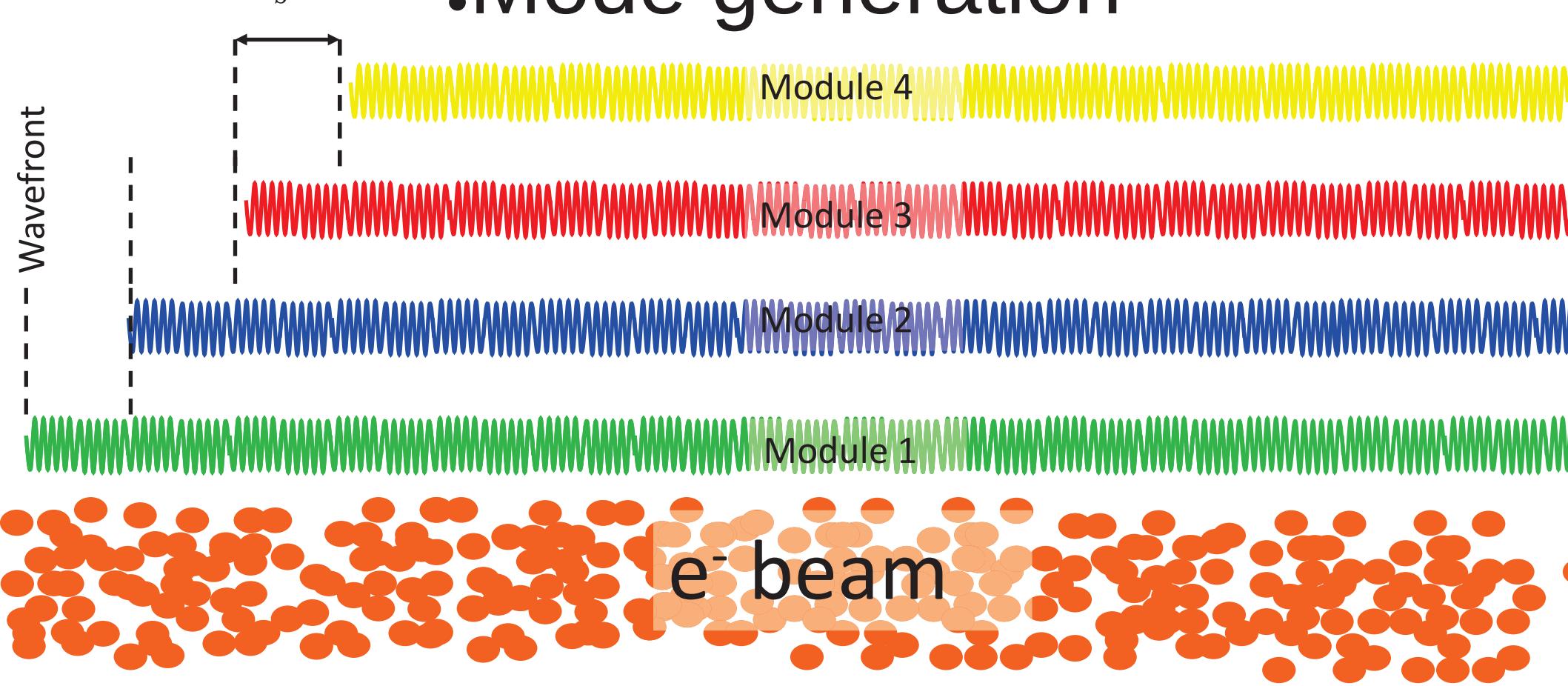
For continued slips of distance  $s$ , only those wavelengths with an integer number of periods in distance  $s$  will survive after many such slips. For  $s$  an integer of  $\lambda_j$ :

$$s = N\lambda_j = (N+1)\lambda_{j-1}$$

$$\Rightarrow \omega_j = \frac{2\pi c N}{s}; \omega_{j-1} = \frac{2\pi c(N+1)}{s} \Rightarrow \Delta\omega_s = \omega_{j-1} - \omega_j = \frac{2\pi c}{s}$$

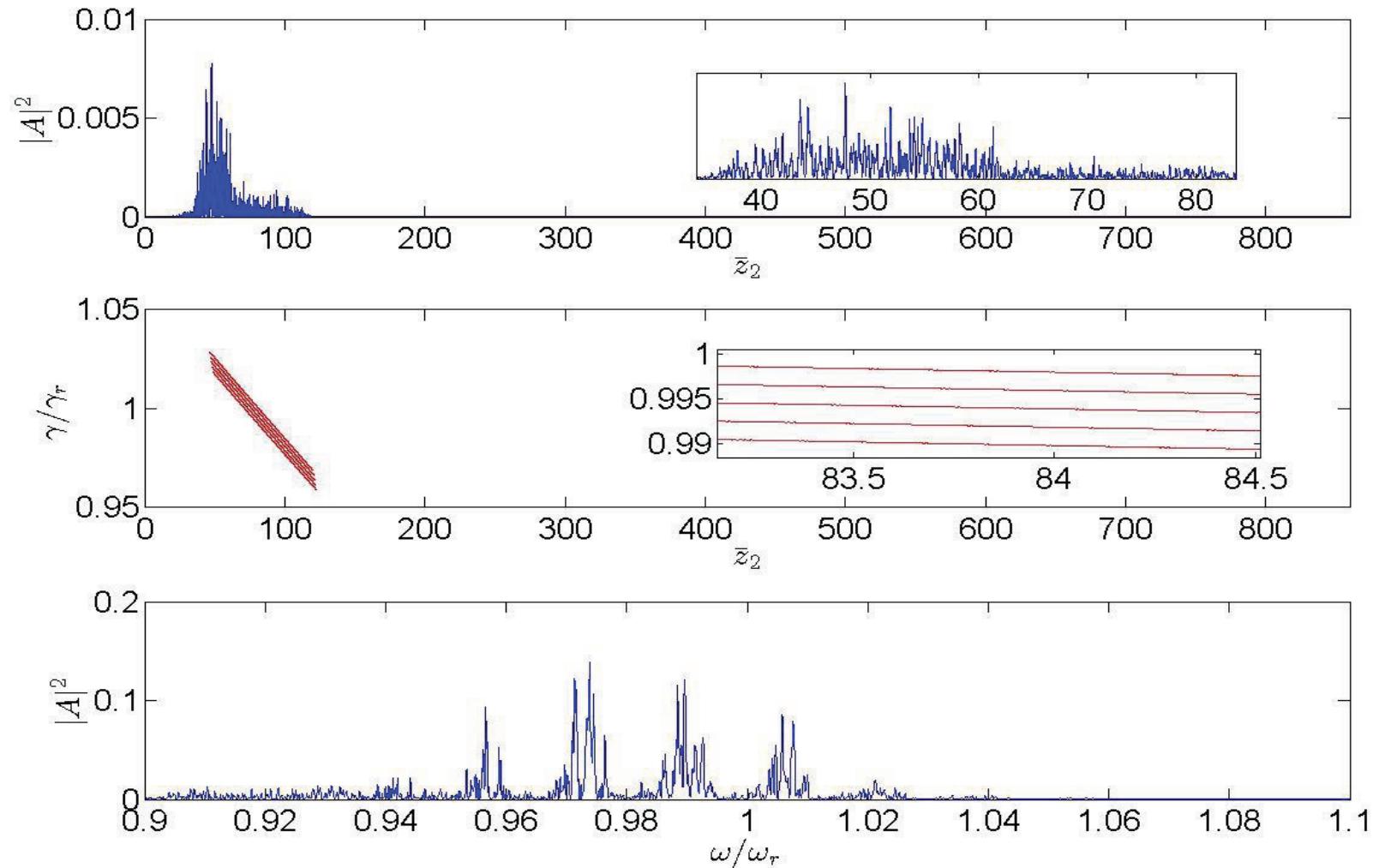
# •Mode generation

— Wavefront



*The spectrum is the same as a ring cavity of length  $s$ . A ring cavity of length equal to the total slippage in each undulator/chicane module has been synthesized*

# The beamlet method



U-CD-U-CS -special undulator-chicane modules

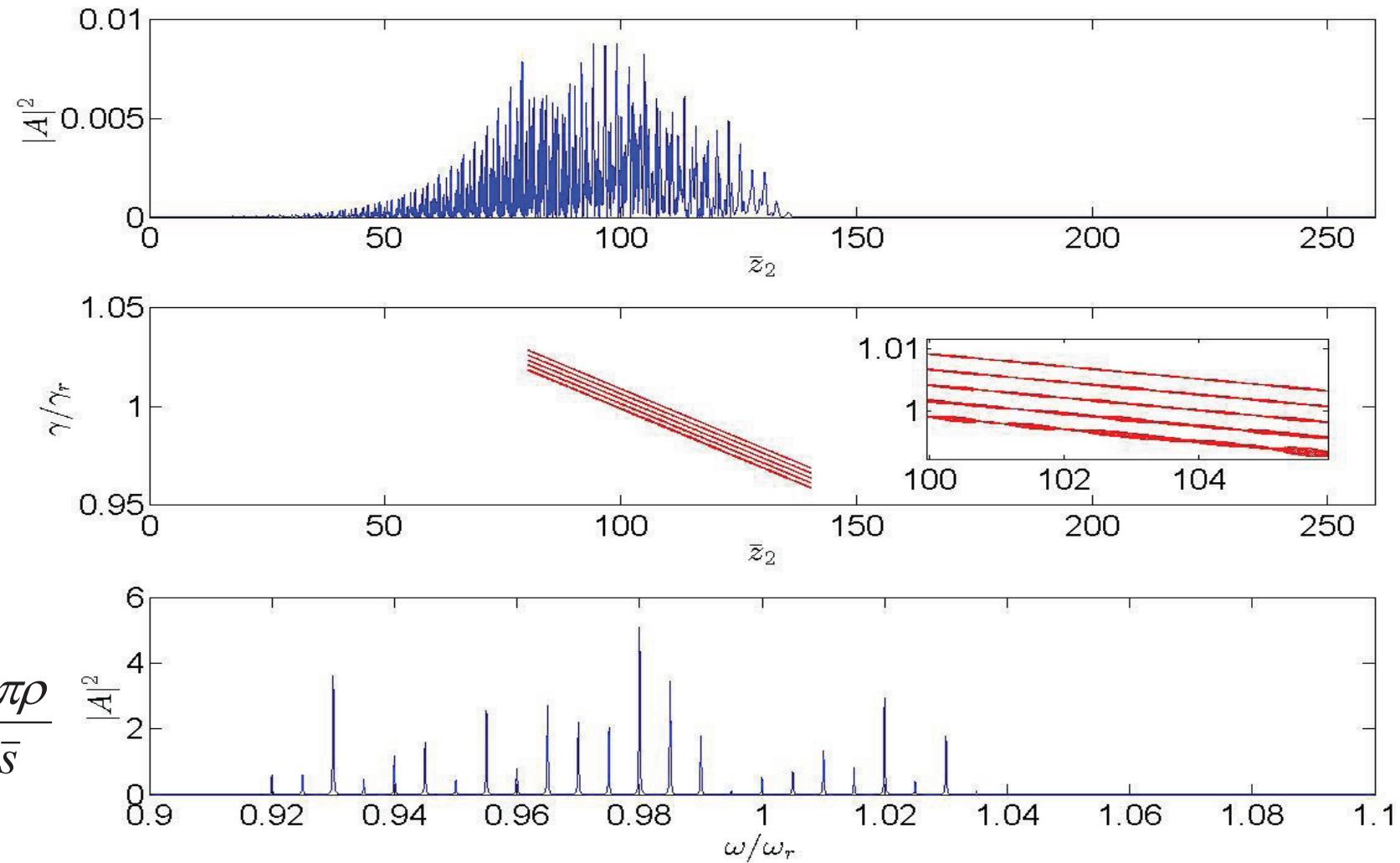
U – undulator

CD – dispersion only chicane

CS – slippage only chicane

$$\frac{\Delta\omega}{\omega_r} = \frac{4\pi\rho}{\bar{\delta}}$$

# The beamlet method



$$\bar{s} = 2.51, \bar{l} = 1.88, \bar{\delta} = 0.63, D = -1.88$$

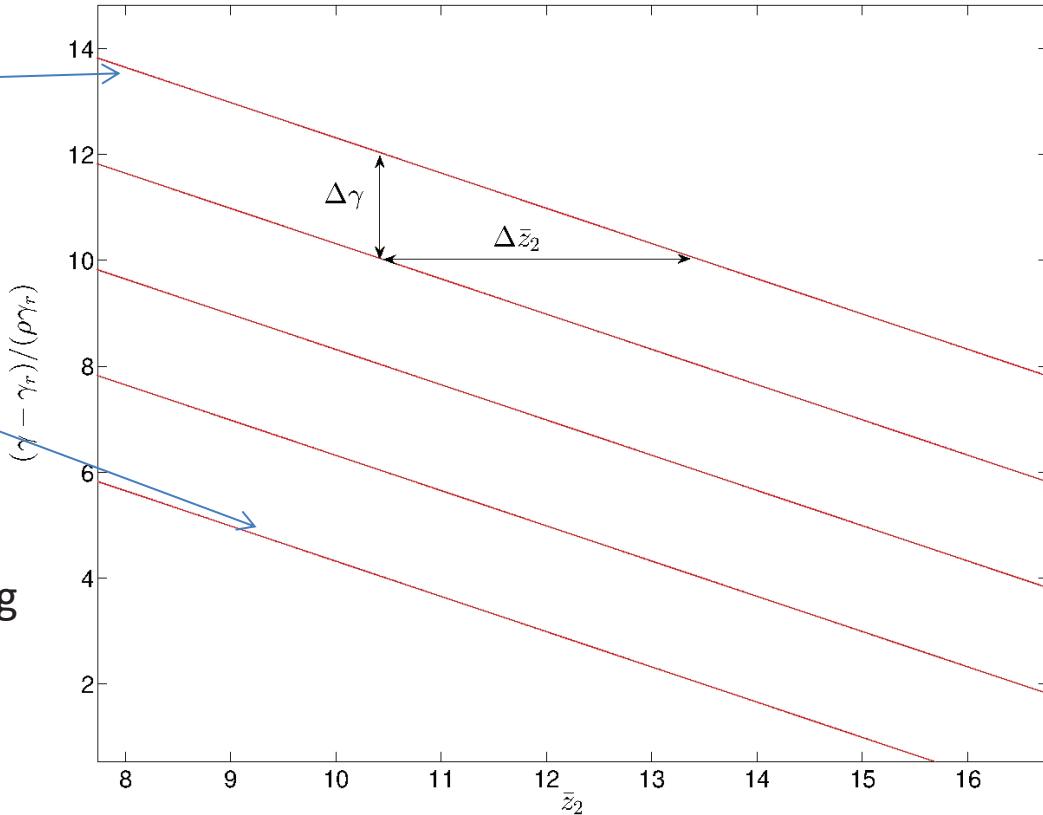
$$\bar{s}_\gamma = 2 \left( \frac{\gamma_r - \gamma_j}{\gamma_r} \right) (\bar{l} + D) + \bar{s}$$

# The beamlet method

- Electron pulses with a large energy chirp exhibit an energy dependent slippage

$$\bar{s}_\gamma = 2 \left( \frac{\gamma_r - \gamma_j}{\gamma_r} \right) (\bar{l} + D) + \bar{s}$$

- These electrons take a shorter path through the undulator than lower energy electrons



- Can lead to mismatching when passing radiation from beamlet to beamlet

# Potential Solution

