

SASE FEL PULSE DURATION ANALYSIS FROM SPECTRAL CORRELATION FUNCTION

FEL Conference 2011
Shanghai, 24. August. 2011

Alberto Lutman
Jacek Krzywinski, Yuantao Ding, Yiping Feng,
Juhao Wu, Zhirong Huang, Marc Messerschmidt

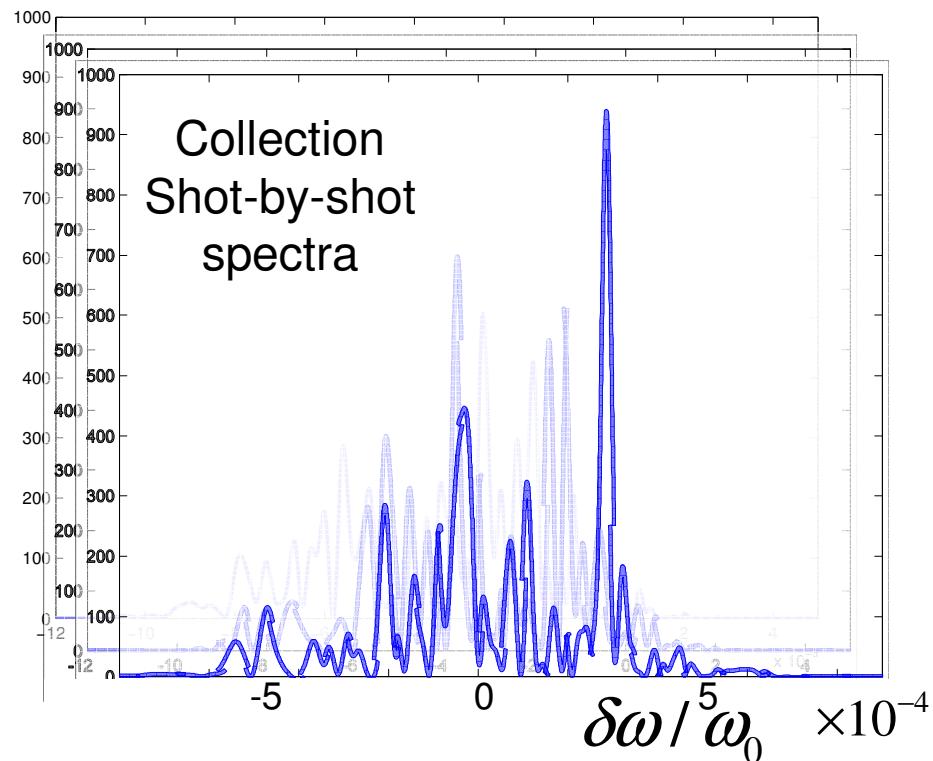


U.S. DEPARTMENT OF
ENERGY

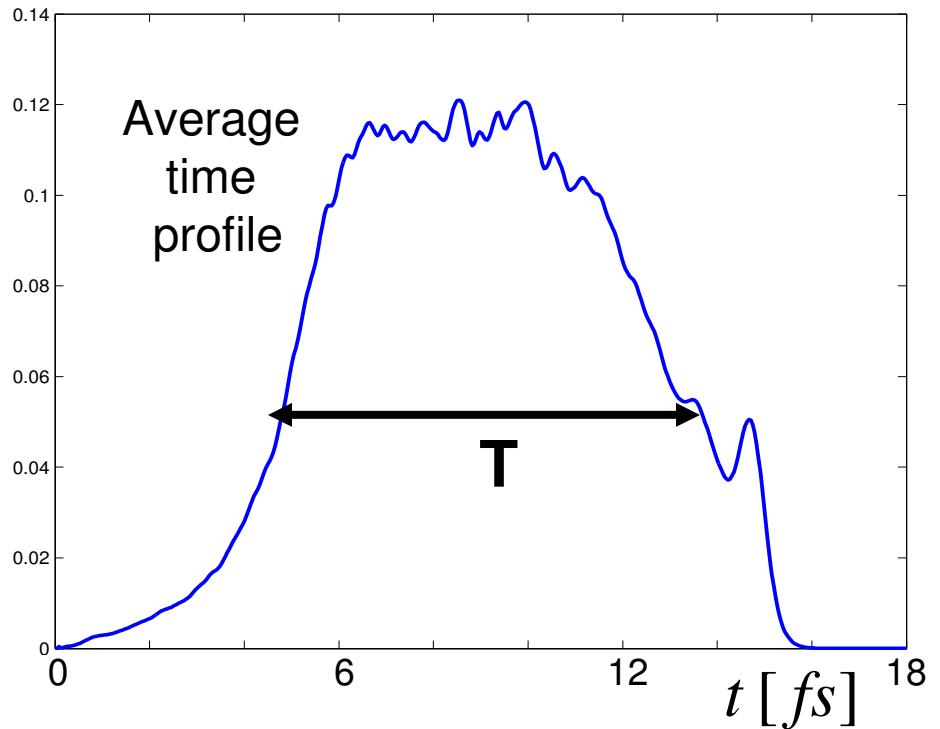
Office of
Science



X-ray pulse duration measurement



Collection
Shot-by-shot
spectra



Average
time
profile

We recover the average x-ray
time profile length T
from a collection of single shot spectra

Model description

Electron beam current:

$$I(t) = (-e) \sum_{k=1}^N \delta(t - t_k)$$

t_k are independent random variables with probability density $f(t)$

SASE FEL Amplifier in linear regime:

$$E(t) = \int_{-\infty}^{+\infty} h(t, \tau) I(\tau) d\tau = \int_{-\infty}^{+\infty} h_{ti}(t - \tau) h_{td}(\tau) I(\tau) d\tau$$
$$h_{ti}(t - \tau) = A_0(z) e^{i(k_0 z - \omega_0(t - \tau))} e^{-\frac{(t - \tau - z/v_g)^2}{4\sigma_t^2}} \left(1 + \frac{i}{\sqrt{3}}\right) \quad \text{time independent}$$

$h_{td}(\tau)$ time dependent

Model description

Electric field spectrum:

$$\tilde{E}(\omega) = (-e)\tilde{H}_{ti}(\omega) \sum_{k=1}^N h_{td}(t_k) e^{i\omega t_k}$$

Average x-ray profile:

$$X(t) = |h_{td}(t)|^2 f(t)$$

First and second order correlations

$$\langle \tilde{E}(\omega') \tilde{E}^*(\omega'') \rangle = e^2 N \tilde{H}_{ti}(\omega') \tilde{H}_{ti}(\omega'') \tilde{X}(\omega' - \omega'')$$

$$\left\langle |\tilde{E}(\omega')|^2 |\tilde{E}(\omega'')|^2 \right\rangle = e^4 N^2 \left| \tilde{H}_{ti}(\omega') \right|^2 \left| \tilde{H}_{ti}(\omega'') \right|^2 \left(\tilde{X}(0) + |\tilde{X}(\omega' - \omega'')|^2 \right)$$

Model description

First and second order correlations functions

$$g_1(\omega', \omega'') = \frac{\langle \tilde{E}(\omega') \tilde{E}^*(\omega'') \rangle}{\sqrt{\langle |\tilde{E}(\omega')|^2 \rangle \langle |\tilde{E}(\omega'')|^2 \rangle}}$$

$$g_2(\omega', \omega'') = \frac{\langle |\tilde{E}(\omega')|^2 |\tilde{E}(\omega'')|^2 \rangle}{\langle |\tilde{E}(\omega')|^2 \rangle \langle |\tilde{E}(\omega'')|^2 \rangle}$$

$$|g_1(\omega', \omega'')| = \frac{|\tilde{X}(\omega' - \omega'')|}{\tilde{X}(0)}$$

Valid in the linear regime

$$g_2(\omega', \omega'') = 1 + |g_1(\omega', \omega'')|^2$$

For non-linear regime, we run numerical simulations

Model description

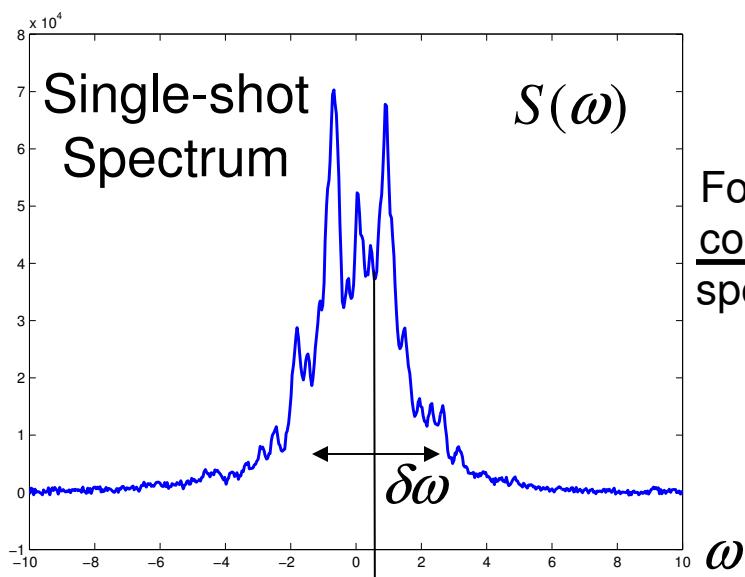
Correlation of the intensity at the exit of the spectrometer

$$G_2(\delta\omega) = \frac{\langle S(\omega_0 + \delta\omega/2)S(\omega_0 - \delta\omega/2) \rangle}{\langle S(\omega_0 + \delta\omega/2) \rangle \langle S(\omega_0 - \delta\omega/2) \rangle} - 1$$

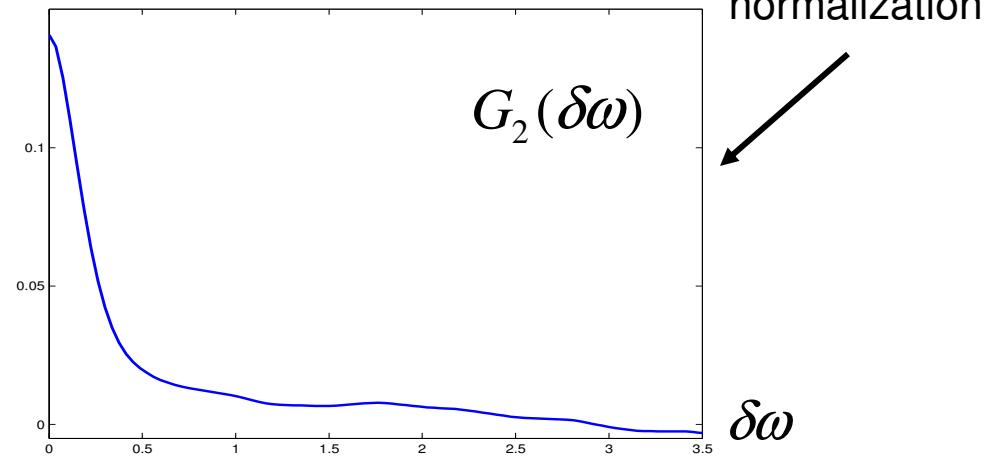
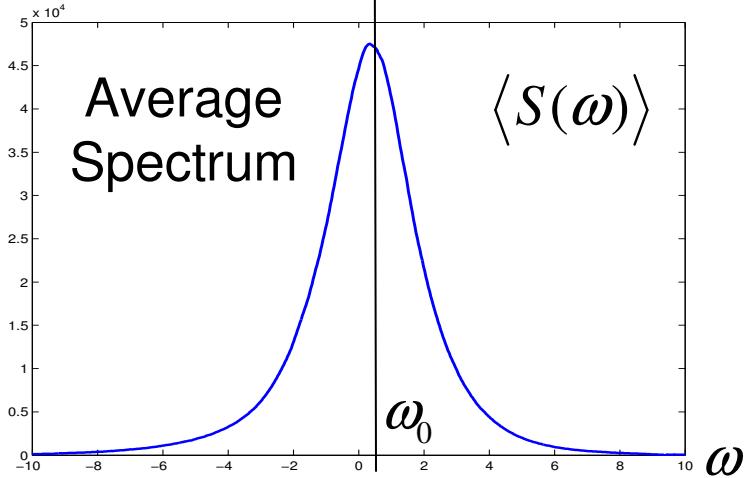
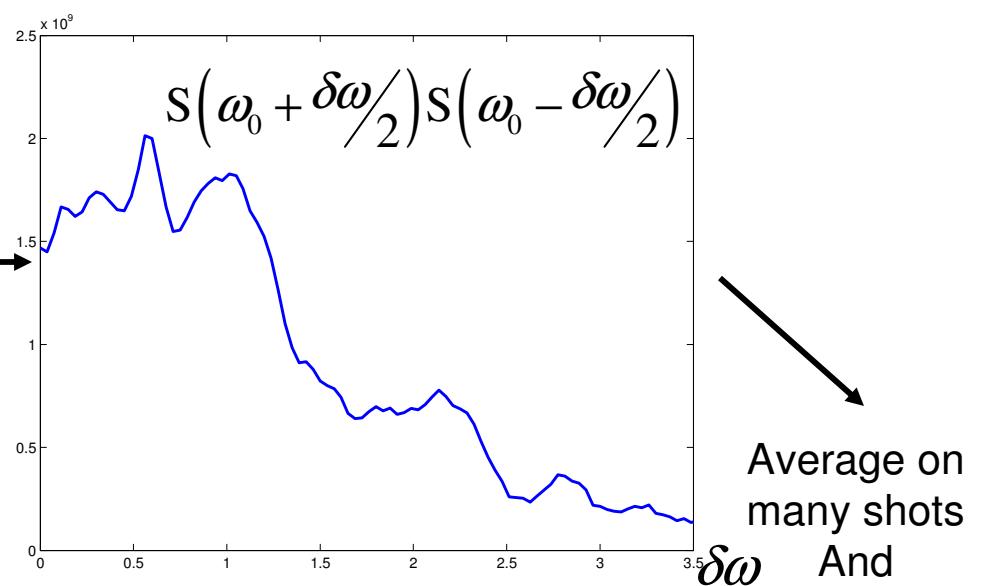
$S(\omega)$ Intensity spectrum at frequency ω

ω_0 Central frequency of amplification

Procedure to calculate G_2



For each
collected
spectrum



G_2 function for different $X(t)$ profiles

$$|\tilde{H}_{ti}(\omega)|^2 \propto e^{-\frac{(\omega-\omega_0)^2}{2\sigma_a^2\omega_0^2}}$$

σ_a relative FEL bandwidth

We assume that the spectrometer has Gaussian resolution function

$$\sigma_m \text{ relative rms spectrometer resolution}$$

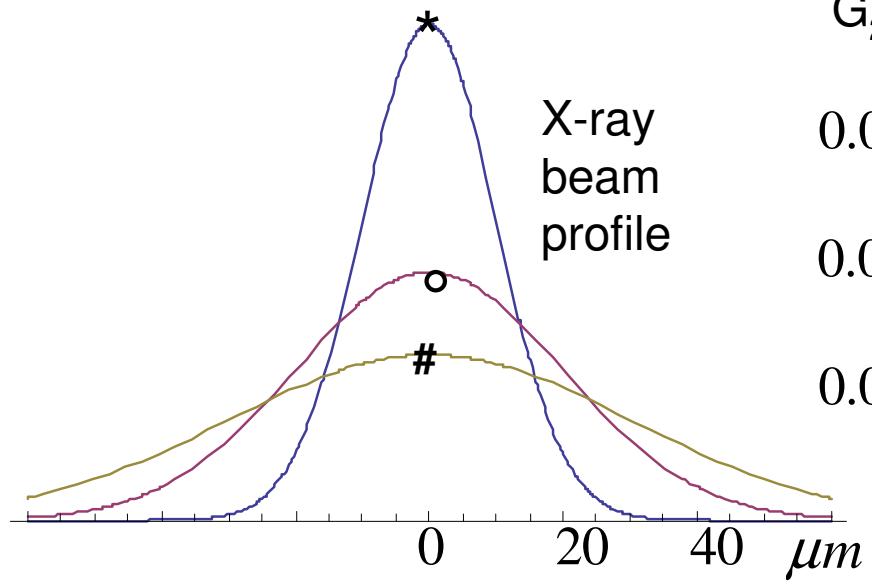
$$G_2(\delta\omega) = \int_{-\infty}^{+\infty} \frac{e^{-\frac{(\zeta-\delta\omega\xi_0)^2}{2\sigma^2}} |\tilde{X}(\zeta, T)|^2}{\sqrt{2\pi}\sigma} d\zeta$$

$$\xi_0 = \frac{\sigma_a^2}{\sigma_m^2 + \sigma_a^2} \omega_0$$
$$\sigma = \sqrt{2} \frac{\sigma_a \sigma_m}{\sqrt{\sigma_m^2 + \sigma_a^2}} \omega_0$$

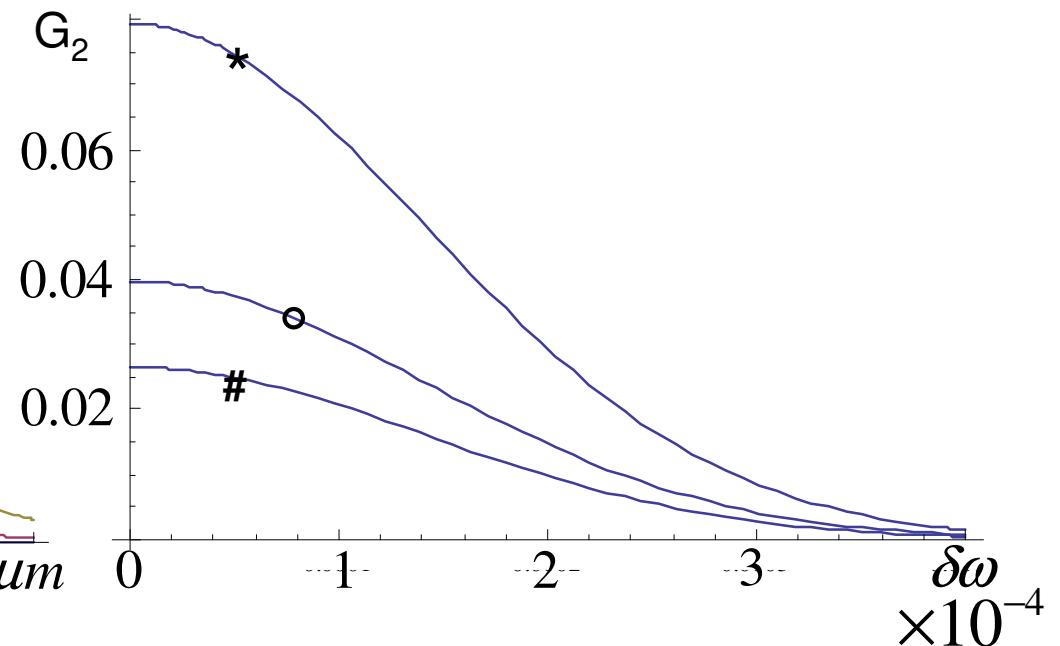
To find an analytical expression for G_2 we need just to plug in the $X(t)$ average profile

G_2 with Gaussian profile

$$X(t, \sigma_T) = \frac{e^{-\frac{t^2}{2\sigma_T^2}}}{\sqrt{2\pi}\sigma_T}$$



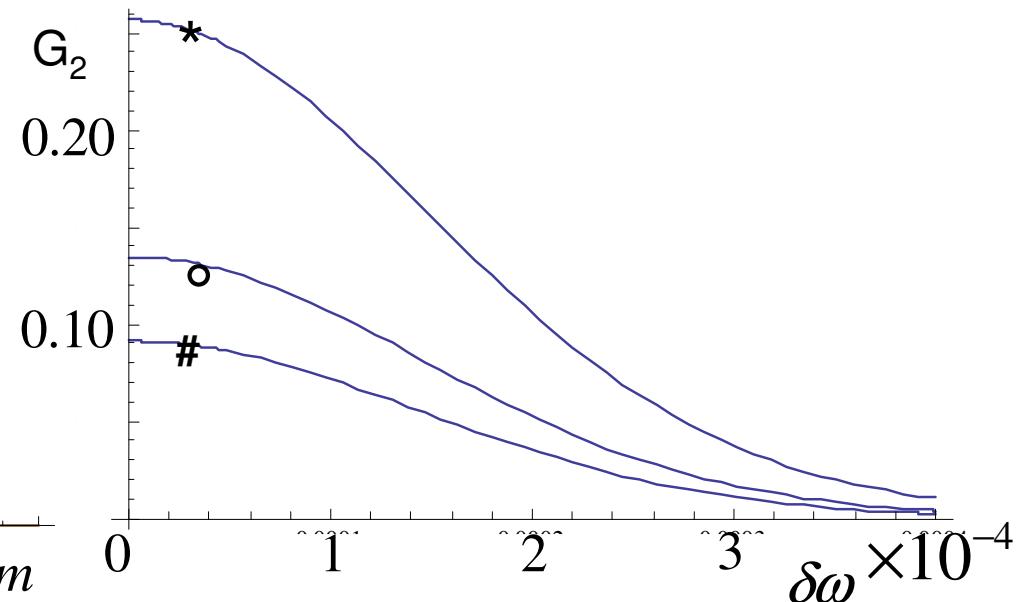
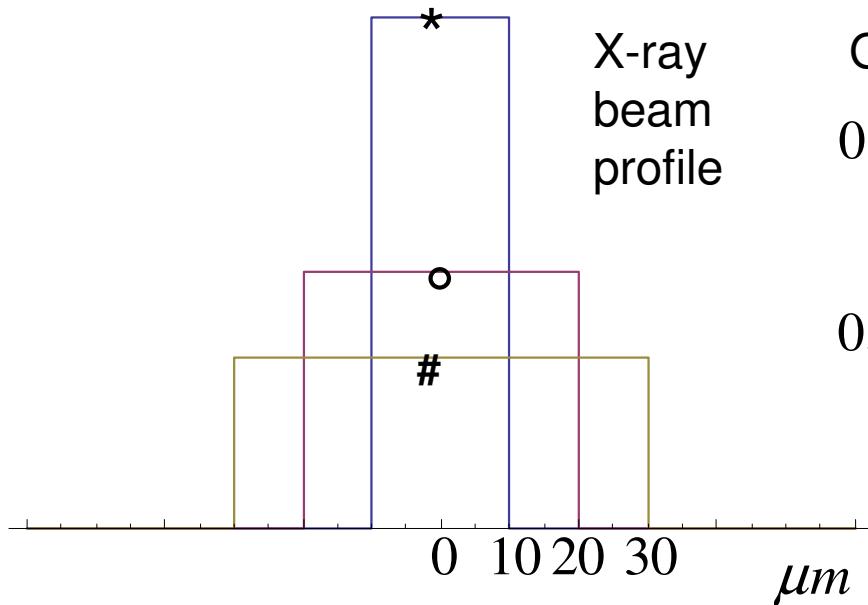
$$G_2(\delta\omega) = \frac{e^{-\frac{\delta\omega^2 \xi_0^2 \sigma_t^2}{1+2\sigma^2\sigma_t^2}}}{\sqrt{1+2\sigma^2\sigma_t^2}}$$



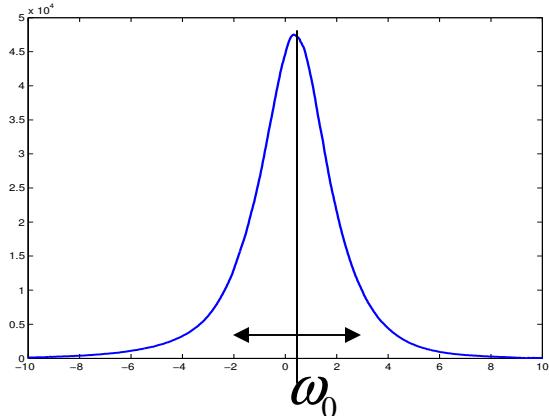
G_2 with flat top profile

$$X(t, T) = \begin{cases} \frac{1}{T} & |t| \leq \frac{T}{2} \\ 0 & |t| > \frac{T}{2} \end{cases}$$

$$G_2(\delta\omega) = 2 \int_0^1 e^{-\frac{\zeta^2 \sigma^2 T^2}{2}} (1 - \zeta) \cos(\delta\omega \xi_0 T \zeta) d\zeta$$



Non stationary ω_0 and FEL gain jitter



Central frequency jitters with
Gaussian law with rms $\sigma_\omega \omega_0$

$$G_2(\delta\omega) = K(\delta\omega) \int_{-\infty}^{+\infty} \frac{e^{-\frac{(\zeta - \delta\omega\xi_0)^2}{2\sigma^2}} |\tilde{X}(\zeta, T)|^2}{\sqrt{2\pi}\sigma} d\zeta$$
$$K(\delta\omega) = \frac{\left(\sigma_a^2 + \sigma_m^2 + \sigma_\omega^2\right) e^{-\frac{\delta\omega^2 \sigma_\omega^2}{4(\sigma_a^2 + \sigma_m^2)(\sigma_a^2 + \sigma_m^2 + \sigma_\omega^2)}}}{\sqrt{(\sigma_a^2 + \sigma_m^2)(\sigma_a^2 + \sigma_m^2 + \sigma_\omega^2)}}$$

Non stationary ω_0 and FEL gain jitter

Shot to shot gain as random variable with average \bar{G}

correlate spectral intensities $I' = (\bar{G} + \Delta G)(\bar{S}' + \Delta S')$ at ω'
 $I'' = (\bar{G} + \Delta G)(\bar{S}'' + \Delta S'')$ at ω''

$$\frac{\langle I'I'' \rangle}{\langle I' \rangle \langle I'' \rangle} = \left(1 + \frac{\langle \Delta S' \Delta S'' \rangle}{\bar{S}' \bar{S}''} \right) \left(1 + \frac{\langle \Delta G^2 \rangle}{\bar{G}^2} \right)$$

$$G_2(\omega', \omega'') = \frac{\frac{\langle I'I'' \rangle}{\langle I' \rangle \langle I'' \rangle}}{1 + \frac{\langle \Delta G^2 \rangle}{\bar{G}^2}} - 1$$

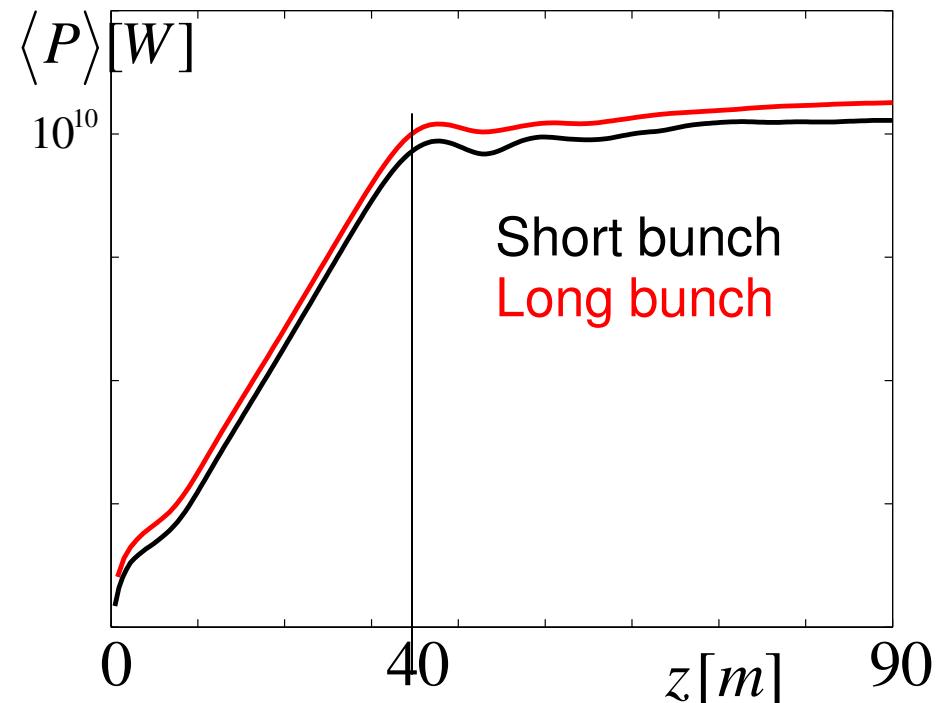
Numerical simulations

- 1) Verify that relations hold well enough in saturation
- 2) Recover bunch length and spectrometer resolution

- $30\mu\text{m}$ electron bunch
- $3\mu\text{m}$ electron bunch

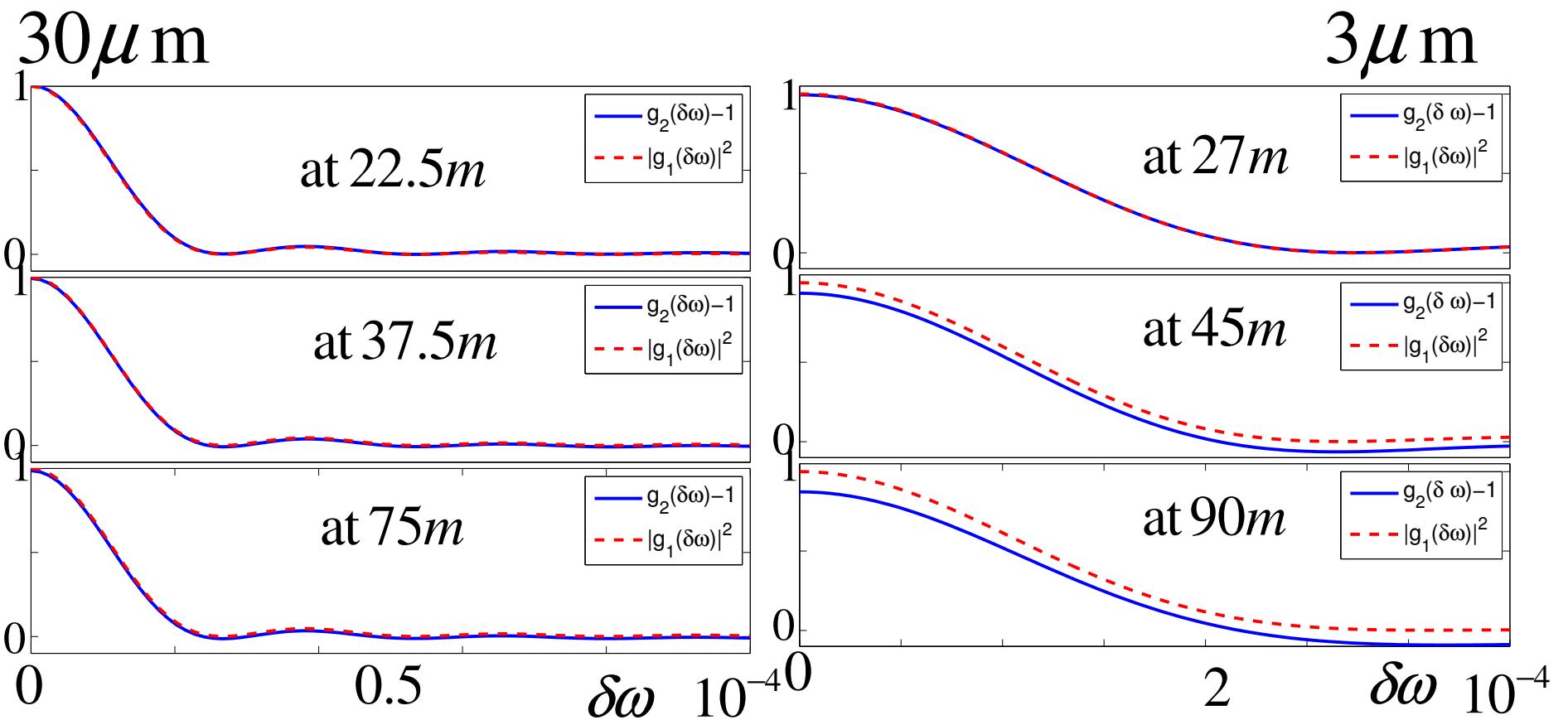
Wavelength 0.8 nm
Undulator period 3 cm

- $\sigma_m = 1 \times 10^{-4}$
- $\sigma_m = 2 \times 10^{-4}$

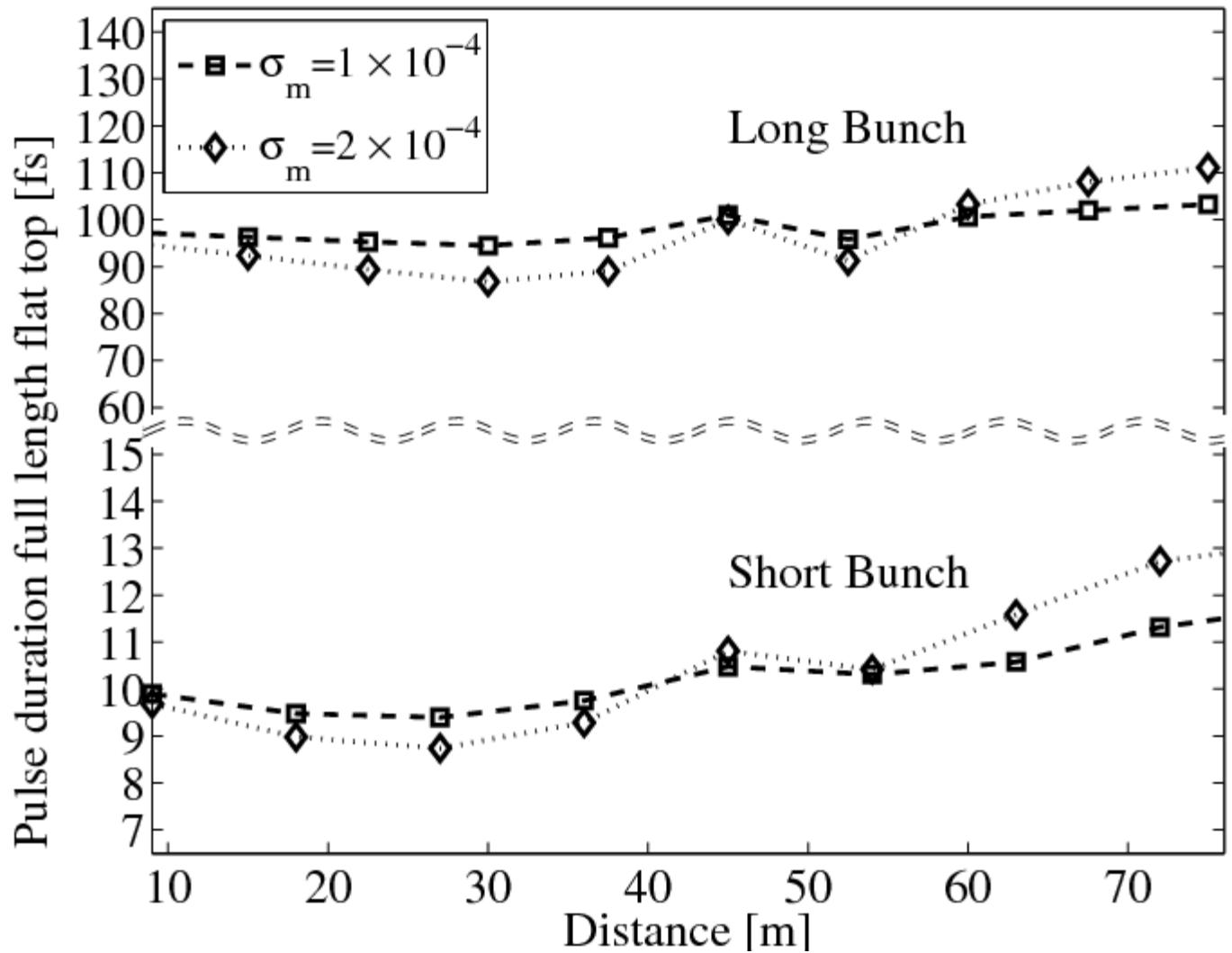


Verify relation

$$g_2(\omega', \omega'') - 1 = |g_1(\omega', \omega'')|^2$$



Simulated measurements results



Retrieved
spectrometer
resolutions

Long bunch
 $(1.01 \pm 0.01) \times 10^{-4}$
 $(2.05 \pm 0.05) \times 10^{-4}$

Short bunch
 $(0.98 \pm 0.03) \times 10^{-4}$
 $(1.98 \pm 0.09) \times 10^{-4}$

Experimental results

Experimental demonstration performed at LCLS

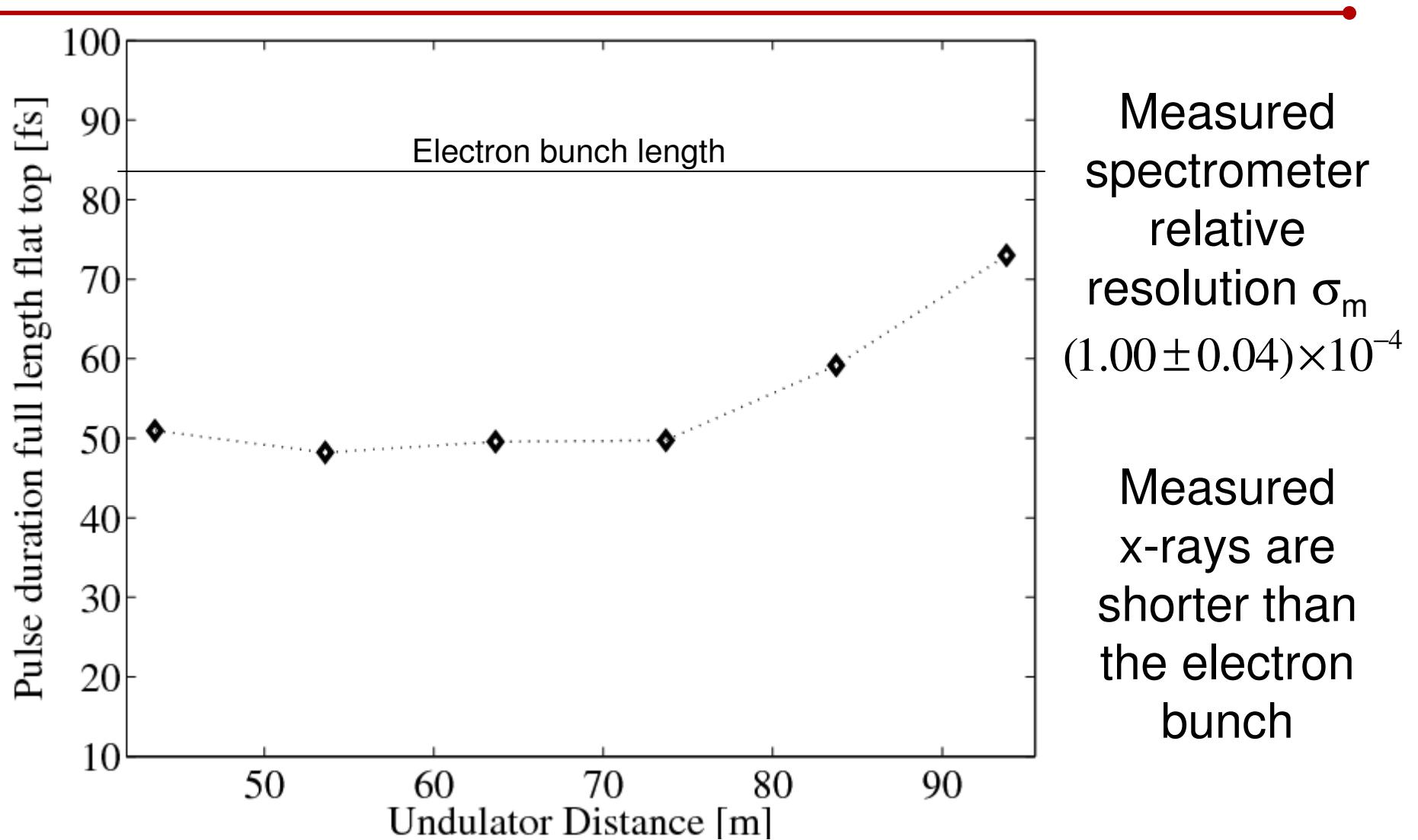
Photon energy 1.5 keV
Electron charge 250 pC

Different undulator length

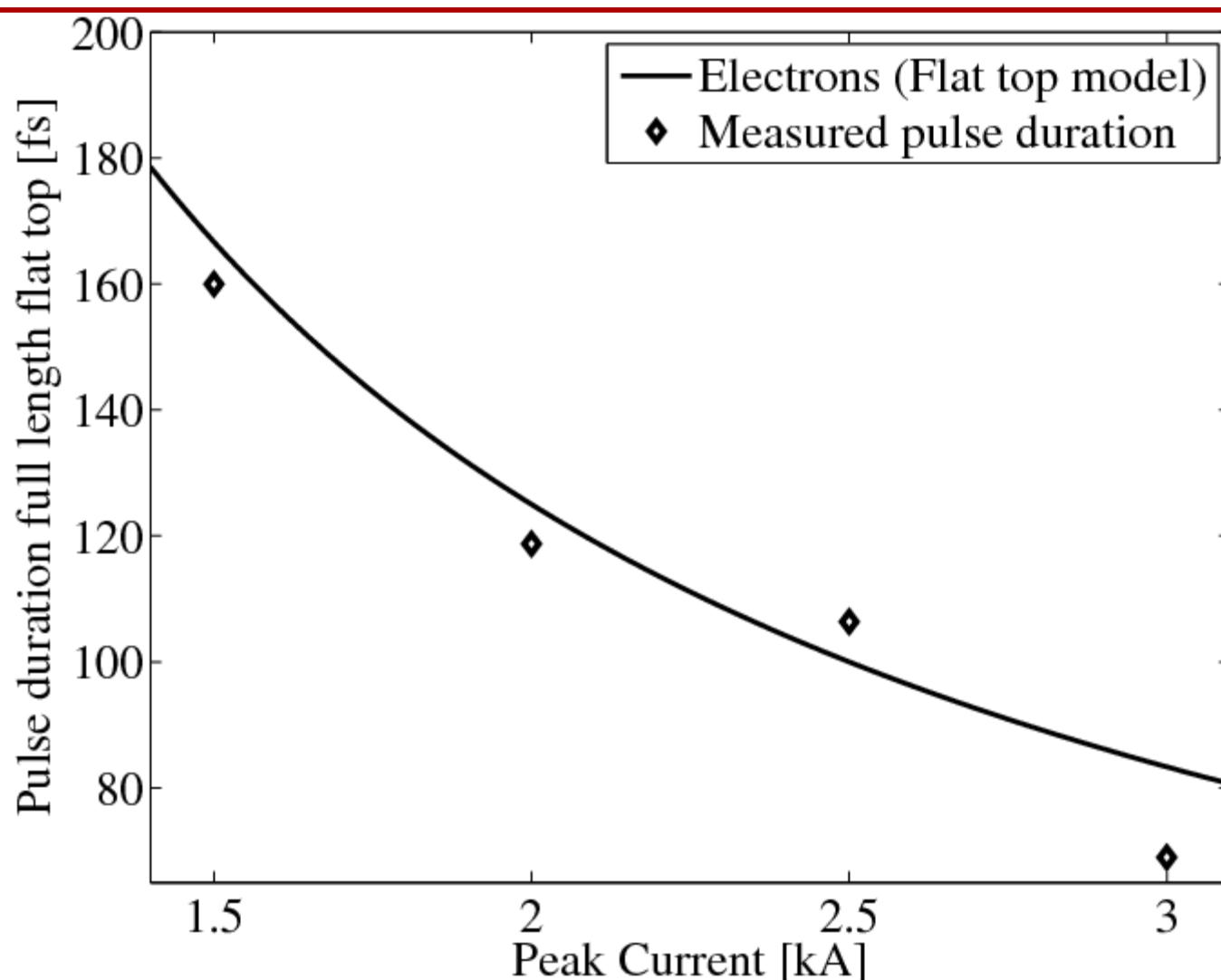
Different peak current

Controlling lasing part of electron bunch with slotted foil

Different undulators distance

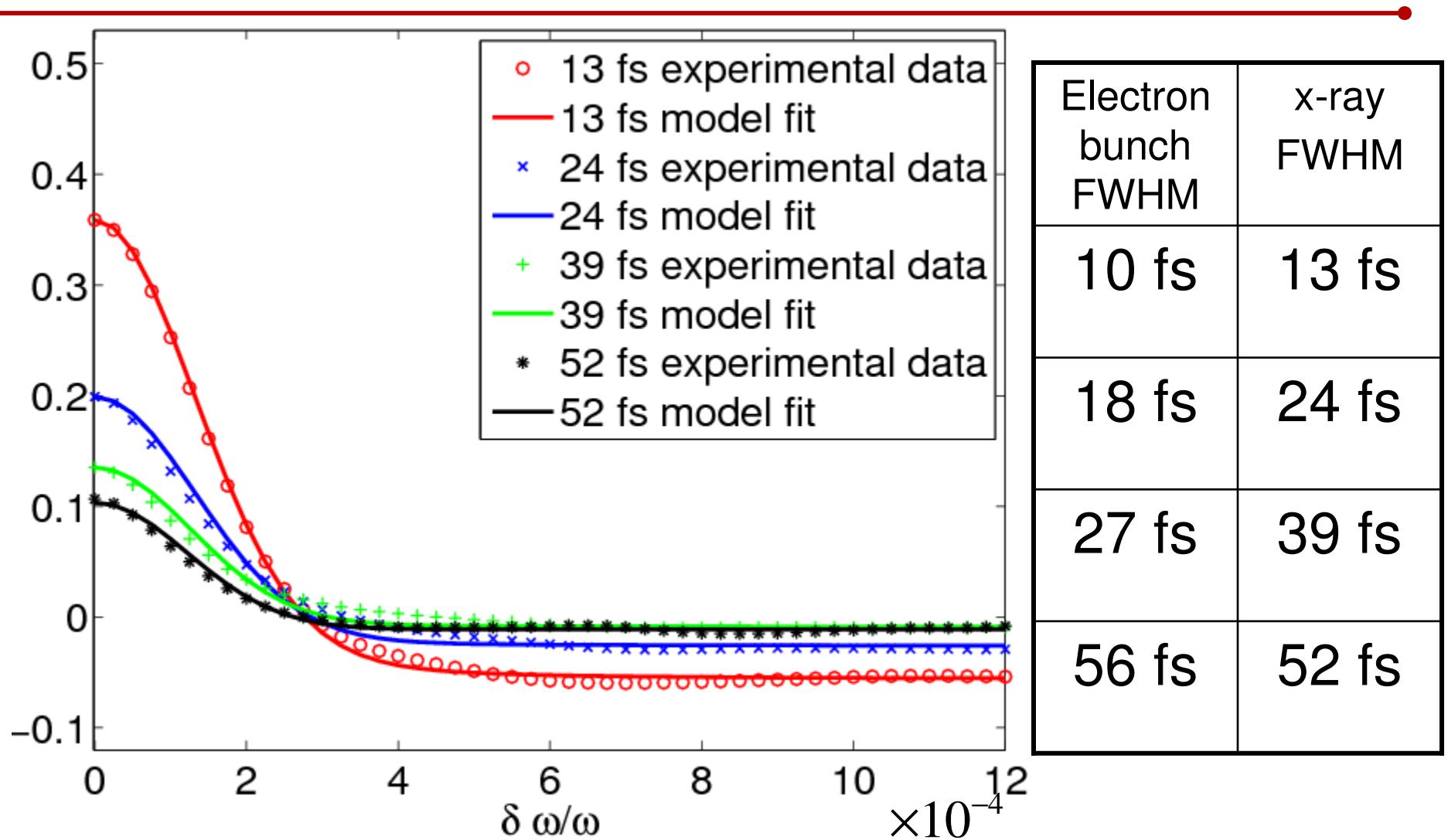


Different peak currents



Measured x-rays pulse durations are consistent with electron bunch length change.

Slotted foil measurements





Thank you for your attention



20

FEL Conference 2011
Shanghai, 24.August.2011