



Pushing the Limits of Small Period Permanent Magnet Undulators

Johannes Bahrtdt, HZB, FEL Conference, Shanghai, China, August, 2011

Pushing the limits in various directions...

- Undulators for X-FELs
- Magnet material development
- Cryogenic undulators
- Sub-cm period devices
- Measurement equipment

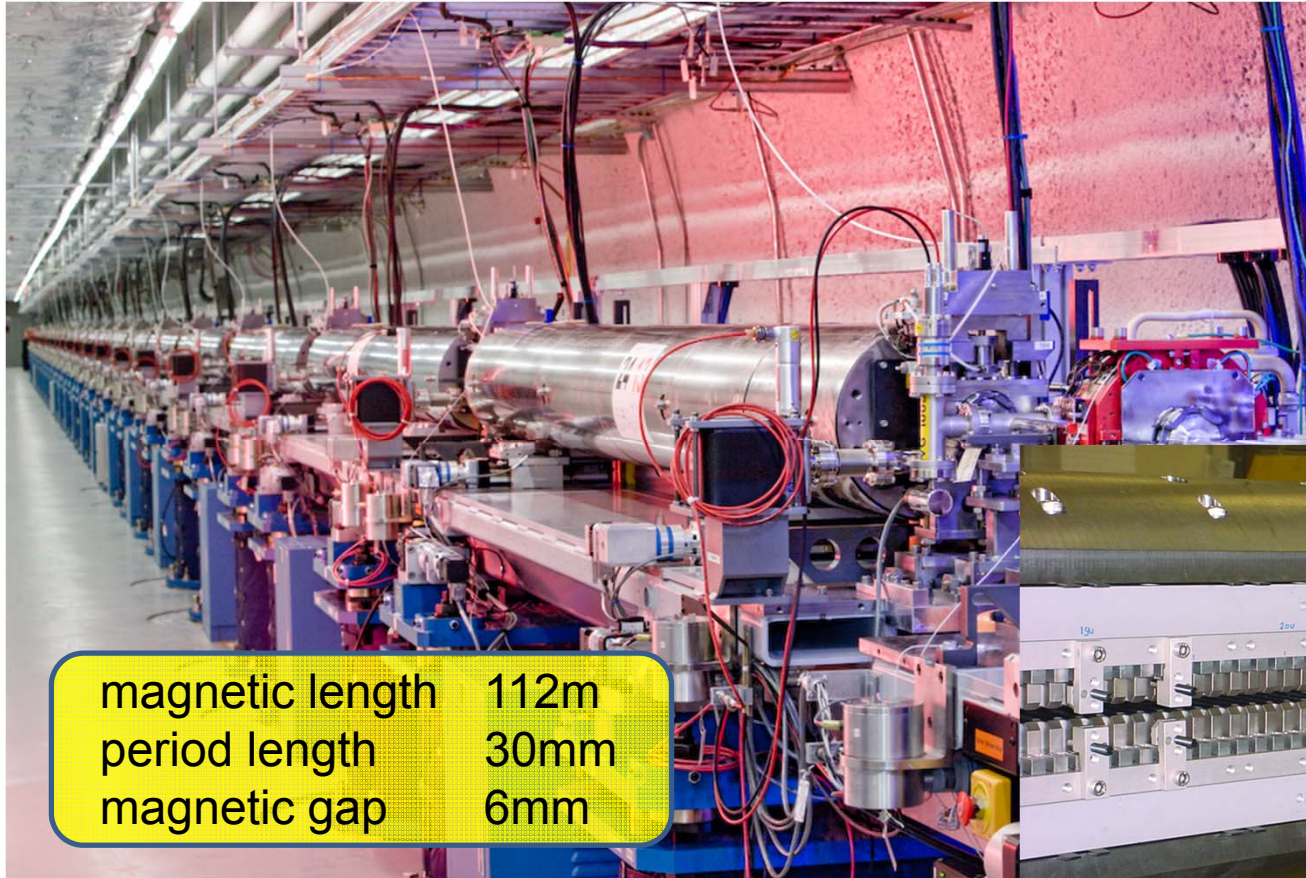
Out of vacuum:

- Trajectory shimming over 5m length, conventional granite bench
- Tight mechanic tolerances (magnetic gap, Taper) over 5m
- Motion control, synchronisation of several dozens of modules
- More complexity added with APPLE II undulators

In-Vacuum

- Trajectory shimming over 5m length with new measurement systems
- Mechanic rigidity under vacuum
- In-vacuum gap measurement system
- Motion control, synchronisation of several dozens of devices

LCLS Fixed Gap Undulators



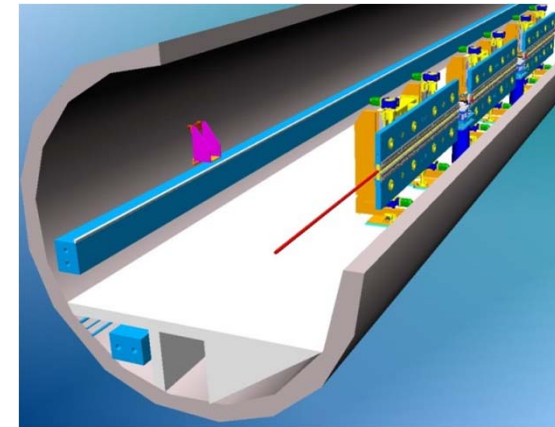
- Next steps:
- self seeding
 - 2nd harmonic afterburner
 - variable gap
 - variable polarization

Courtesy of S. Sasaki

5m Prototype module at measurement bench



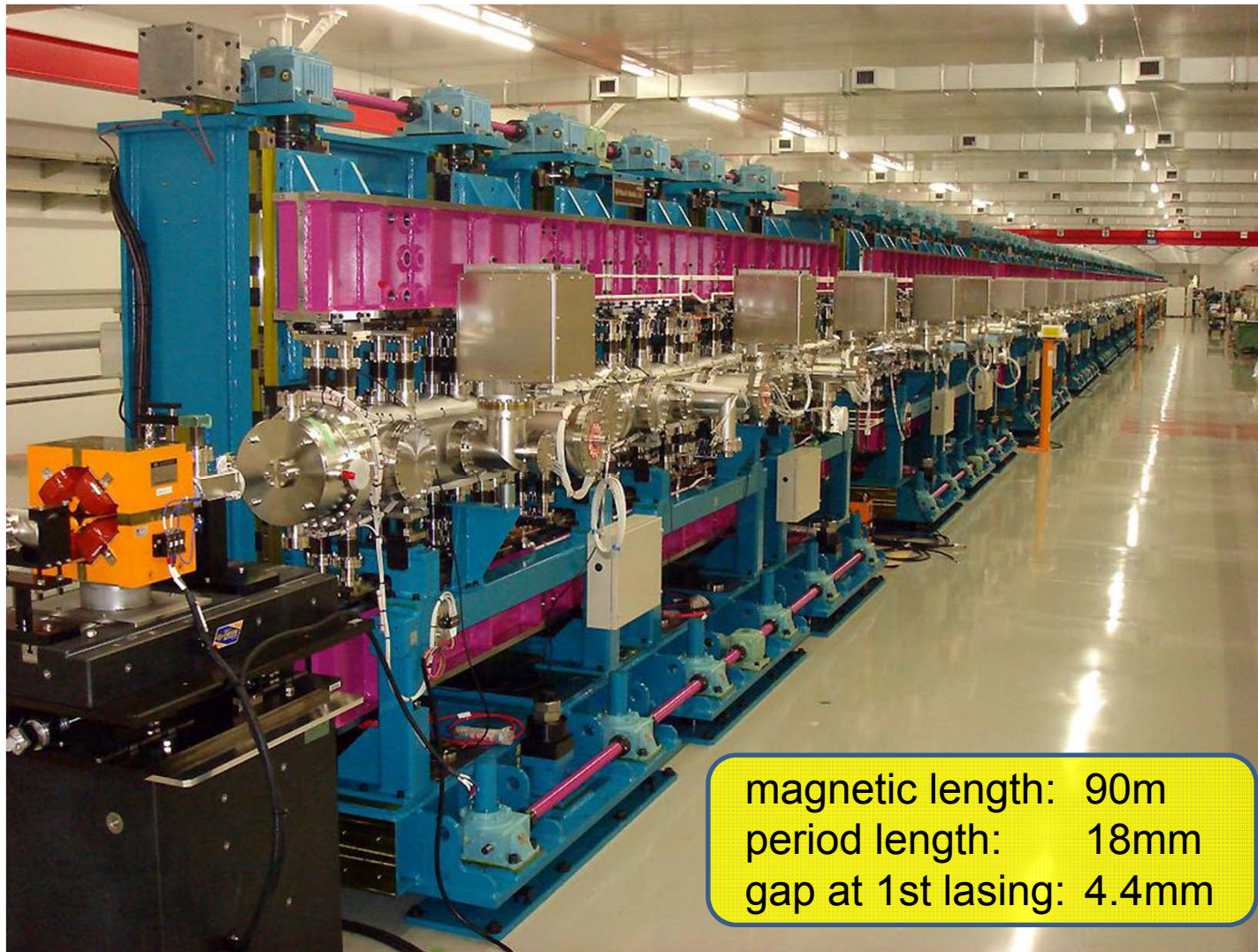
Courtesy of J. Pflüger



magnetic length	175m
period length	40mm
minimum gap	10mm

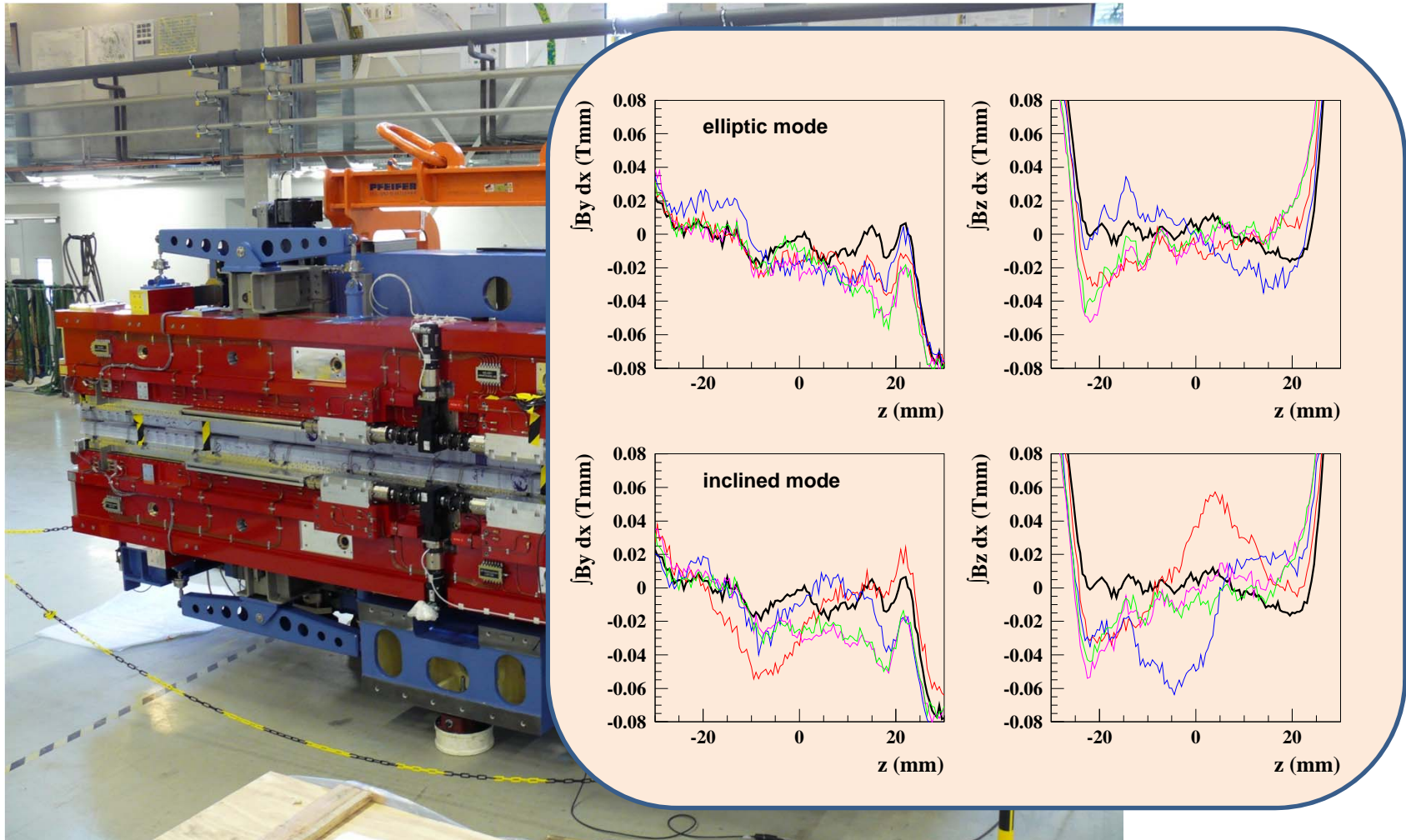


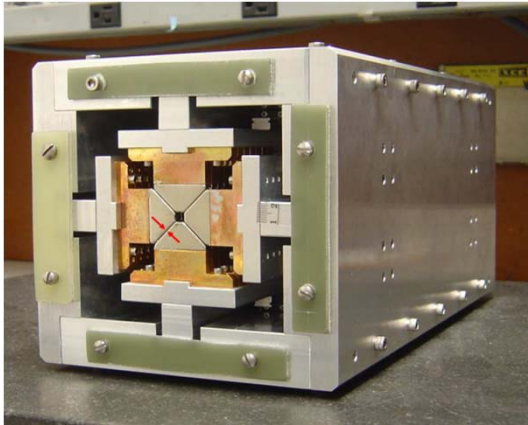
SPRING-8 XFEL In-Vacuum Undulators



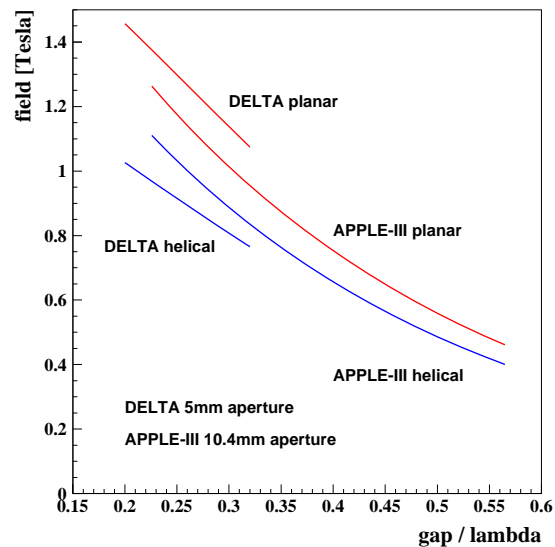
Courtesy of T. Tanaka

PETRA III APPLE II Undulator as built at HZB

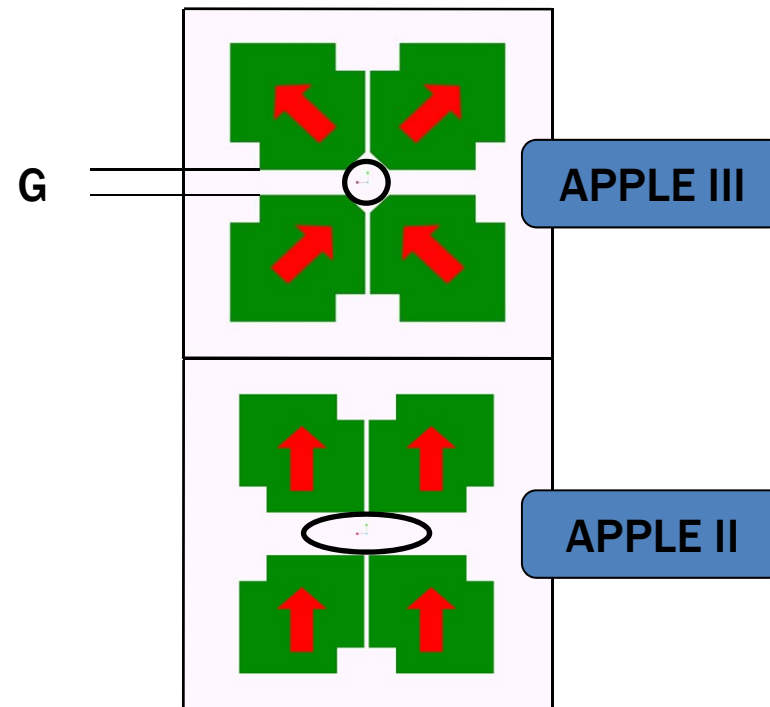




A. Temnykh, *PRST-AB*, 11, 120702 (2008)



Currently under discussion
2nd harmonic afterburner

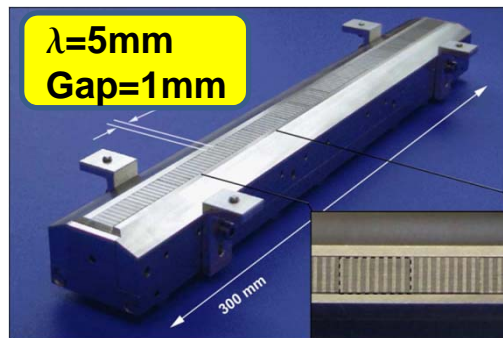


J. Bahrdt et al., *Proc. of FEL Conf., Trieste, Italy (2004)* 610-613

SACLA: first lasing at 4.4mm gap
minimum gap = 3.5mm
 $\lambda=18\text{mm}$

SWISS FEL: minimum gap = 3.2mm
 $\lambda=15\text{mm}$

Prototyping for Laser Plasma Accelerator, Group of F. Gruener



T. Eichner et al., Phys. Rev. Spec. Topic. AB, 10, 082401 (2007) 1-9.

Collaboration between LMU-Munich and HZB

Further prototypes employing:

- cryogenic design
- new materials
- challenging magnet design

Remanence

$$B_r(20^\circ\text{C}) = B_{r-sat}(20^\circ\text{C}) \cdot \frac{\rho}{\rho_0} \cdot (1 - V_{nonmagnetic}) \cdot f_\varphi$$

$$f_\varphi = \cos(\varphi)$$

$$\varphi = \arctan\left(2 \frac{B_{r-perp}}{B_{r-par}}\right)$$

Current status:

- $\rho/\rho_0 > 0.995$, is achieved with the so-called liquid phase sintering
- $f_\varphi > 0.98$, high alignment factor, sophisticated timing during pressing
- Non-magnet particles typical $< 4\%$

Example: energy product of Nd₂Fe₁₇B

Theory: 509 kJ/m³  Achieved: 469 kJ/m³

Not much space for improvement for known materials

Coercivity H_{cj} still one order of magnitude below theory

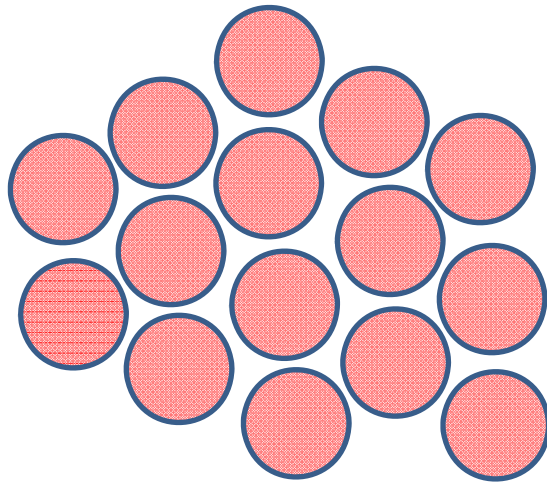
„Brown's paradoxon“

Many reasons (Kronmüller)

- non-magnetic extended grains between magnetic grains
- misoriented grains
- missing grain boundary surfaces resulting in large grains
- extended transition regions of tilted crystal structures between grains

R & D ongoing:

- Optimizing the microcrystalline structure, grain size, alignment angle
- Dy enriched powder
- Two alloy method (Shin-Etsu)
- Grain Boundary Diffusion (GPD)

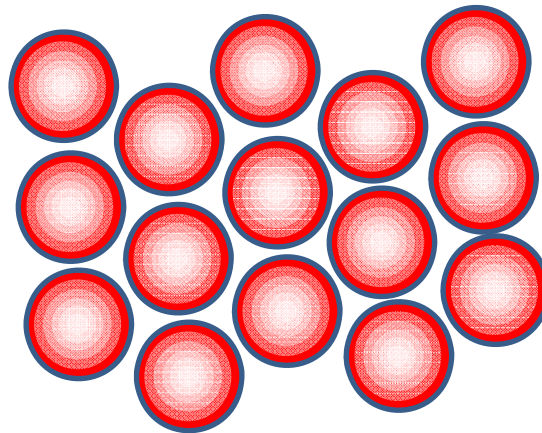


grades differ in amount of Dy in grain bulk

Procedure:

add Dy to powder before pressing

- gain in coercivity
- loss in remanence

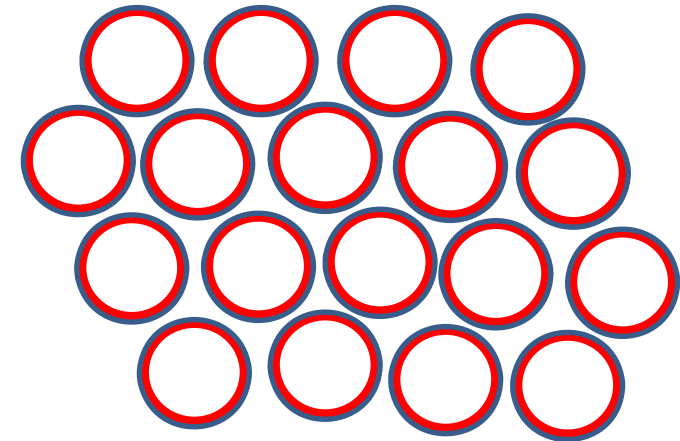


slight Dy gradient to grain boundaries

Procedure:

Two alloy method (Shin-Etsu)

- gain in coercivity
- smaller loss in remanence



sharp concentration peak of Dy at grain boundary

Procedure:

Grain Boundary Diffusion (GPD)

- higher coercivity
- no loss in remanence

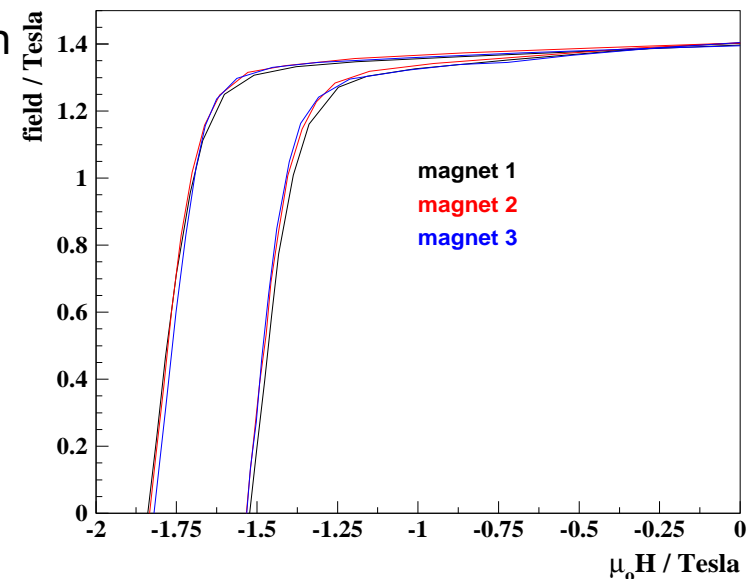
Grain Boundary Diffusion (GBD)

Known already for 10 years, now commercially available

- Coating of magnets with Dy or Tb (expensive)
- Heating the magnets to 800-900 degrees
- Fast diffusion of Dy in liquid phase (boundary)
- Slow diffusion from boundary into the bulk
- Thin highly enriched Dy shells around grain
- Depth profile: $\Delta H_{cj} = Ae^{-x^2/d^2}$
- Penetration depth of Dy about $d = 2.2\text{mm}$

*H. Nakamura et al., J. Phys. D:
Appl. Phys. 44 (2011) 064004-1-5.*

→ Extremely useful for
high end small period devices



PrFeB samples before / after GBD
Courtesy of Vacuumschmelze

Advantage of Cryogenic Devices:

- gain in remanence, larger tuning range for small period IDs
- stability recovered at low temperatures
- less sensitive on thermal dependence of magnetic properties

Advantage of PrFeB Undulators

- no spin reorientation at 135K, lower temperatures possible
- cooling with liquid N₂, no heating necessary
- gain in stability (coercivity), though, gain declines at low temperatures due to decreasing thermal capacity (zero at 0K)

Availability of PrFeB

- ingredients comparable prices as for NdFeB
- however, more expensive due to small number of customers
- melting batch 1-2 to, powder batch 500kg
- laboratory batch: 30kg, reproducibility, performance incr. with batch size

Existing devices: **ESRF, PSI, DIAMOND, SOLEIL (PrFeB)**

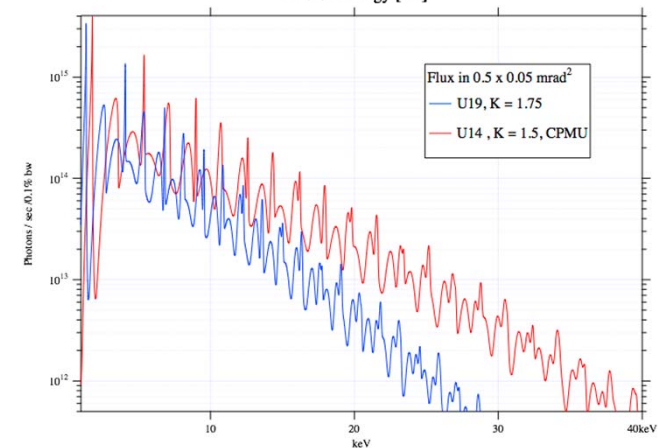
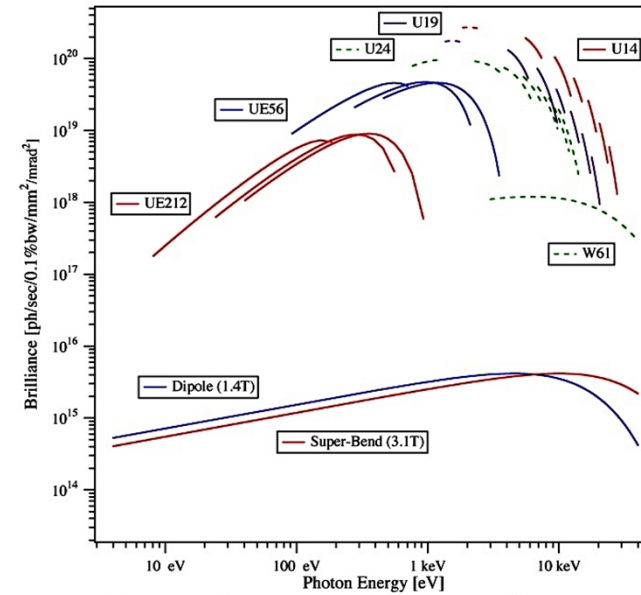
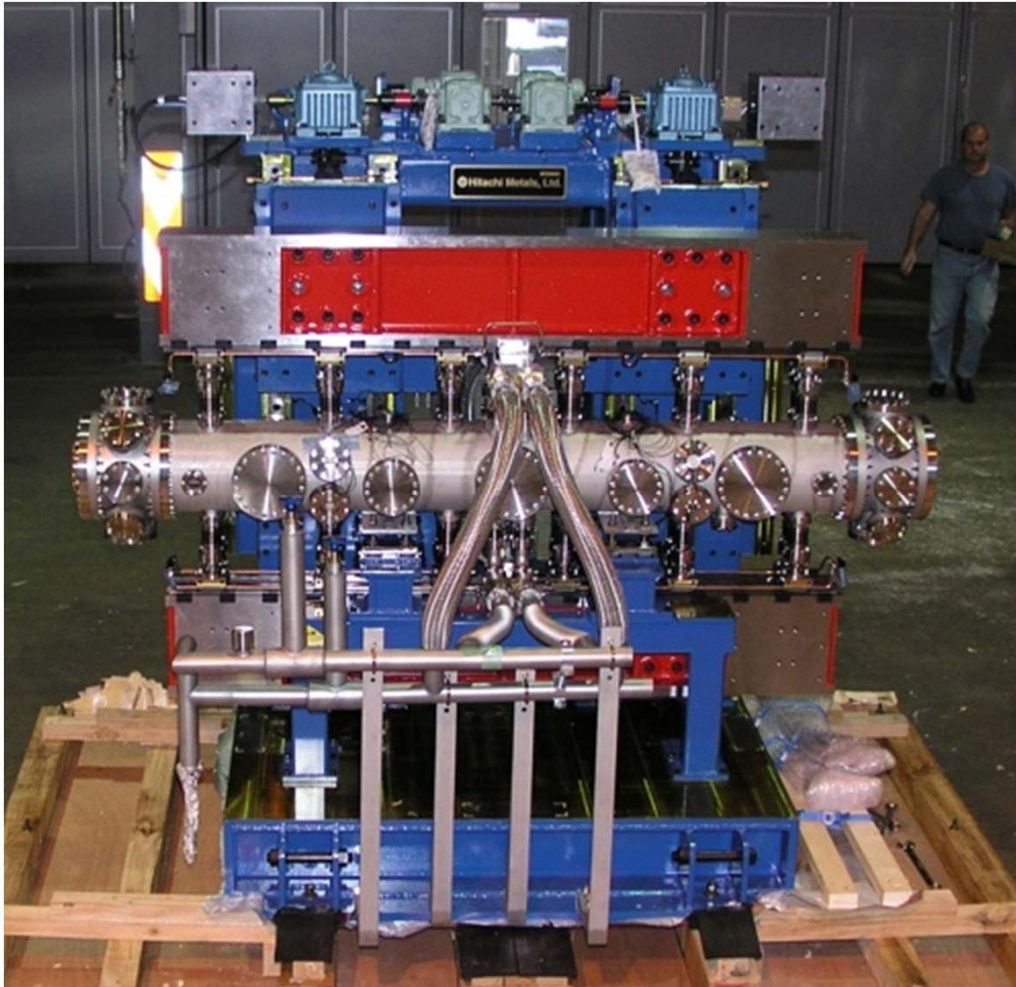
Installed: **ESRF, PSI, (soon SOLEIL)**

First Spectra from SLS / PSI Device

PAUL SCHERRER INSTITUT

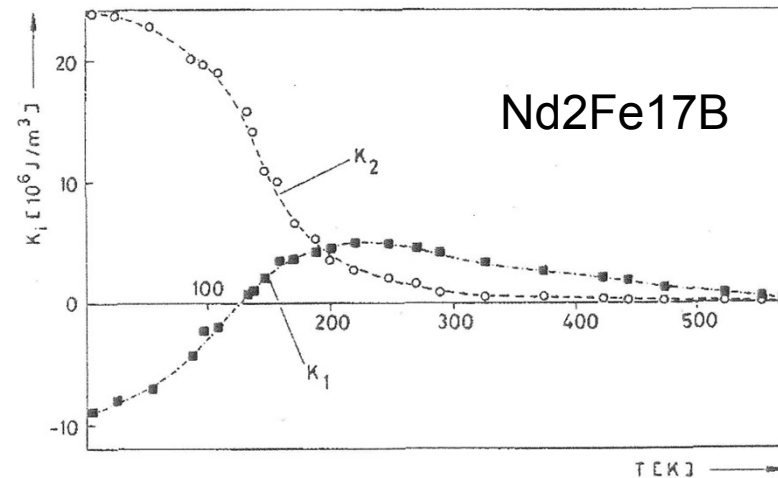
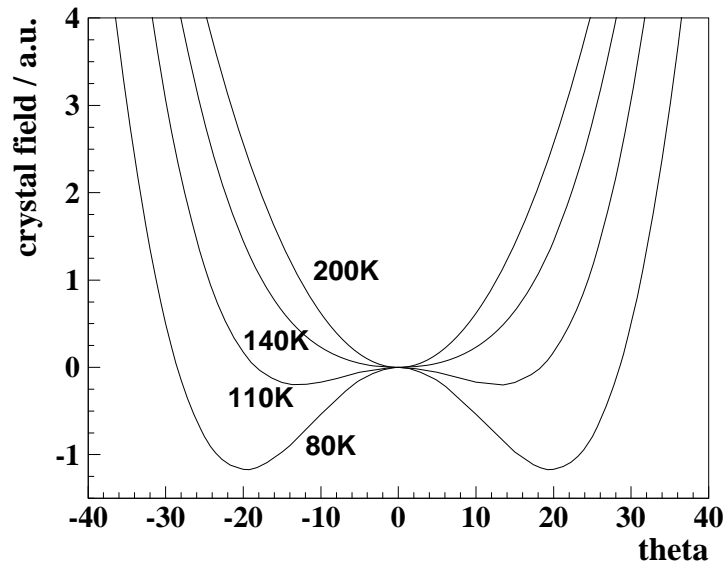


Marco Calvi, Thomas Schmidt, Antonio Cervellino, Phil Willmott



Anisotropy energy

$$E = K_1 \sin^2(\vartheta) + K_2 \sin^4(\vartheta)$$

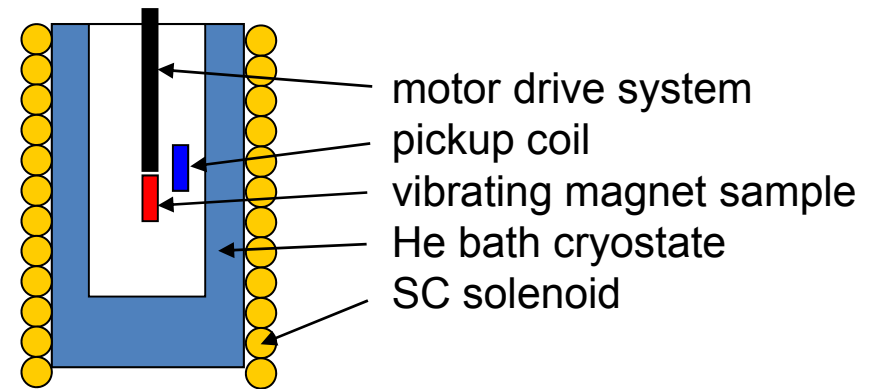
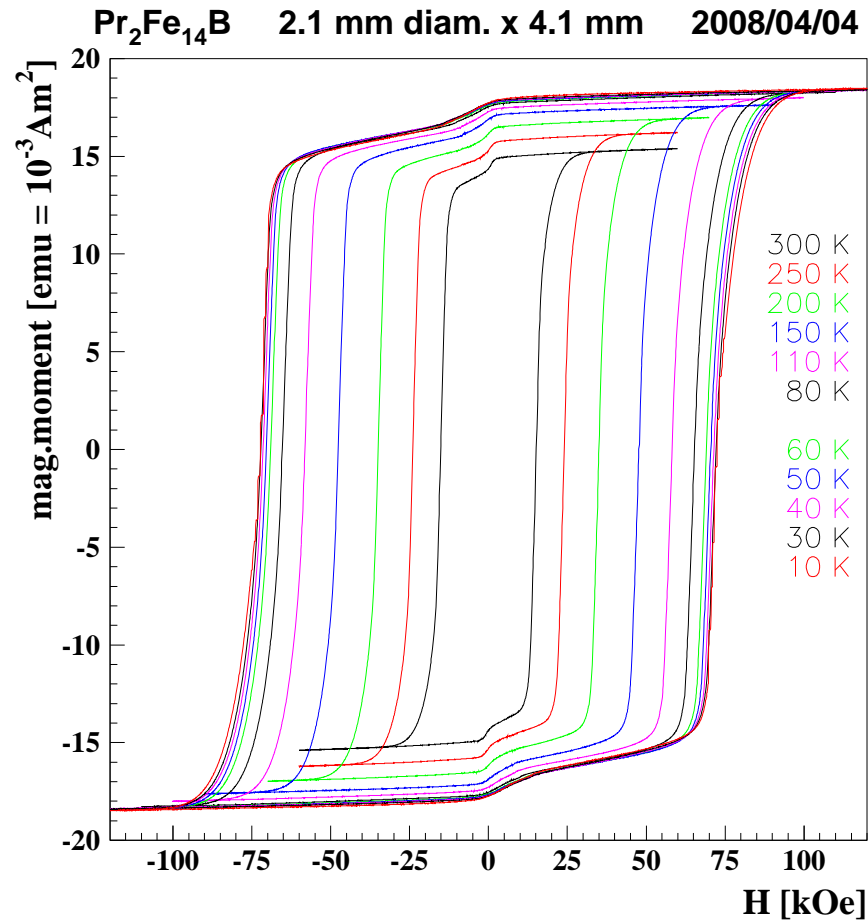


*K.-D. Durst, H. Kronmueller,
J. Mag. And Mat. 59 (1986) 86-89*

Implementation of these equations in RADIA:

*G. Le Bec, J. Chavanne, C. Benabderrahmane,
Proc. of PAC, Vancouver, BC, Canada (2009) 327-329*

New PrFeB grade as developed by VAC, raw data



Analysis requires:

- demagnetization factors
- surface model

fitting model

$$M = J_{bulk} \cdot bulk_volume + J_{side} \cdot side_volume + J_{pole} \cdot pole_volume$$

$$J_{bulk}(H_{bulk}) = \sum_{i=1}^3 J_i \cdot \tanh\left(\frac{\xi_i \cdot \mu_0}{J_i} (H_{bulk} - H_{cj})\right) + F$$

$$H_{bulk} = H_{ext} - (D \cdot J_{bulk}(H_{bulk})) / \mu_0$$

with

$$F = \mu_0 \cdot \chi_0 \cdot (H_{bulk} - H_{cj}) \quad \text{for } H_{bulk} \leq H_{cj1}$$

$$F = \mu_0 \cdot \chi_0 \cdot (H_{cj1} - H_{cj}) + \mu_0 \cdot \chi_{01} \cdot (H_{bulk} - H_{cj1}) \quad \text{for } H_{bulk} > H_{cj1}$$

$$J_{pole} = B_{bulk} = J_{bulk} + \mu_0 H_{bulk}$$

$$J_{pole} = |J_{bulk}(0)| \frac{B_{bulk}}{|B_{bulk}|} \quad \text{for } |B_{bulk}| \leq |J_{bulk}(0)|$$

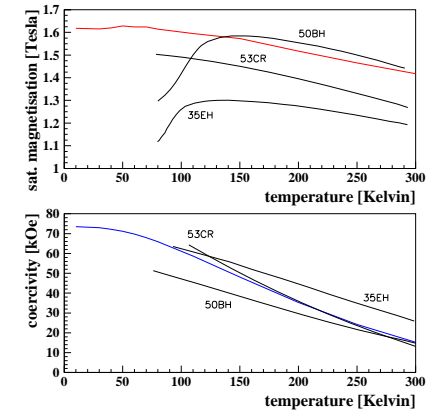
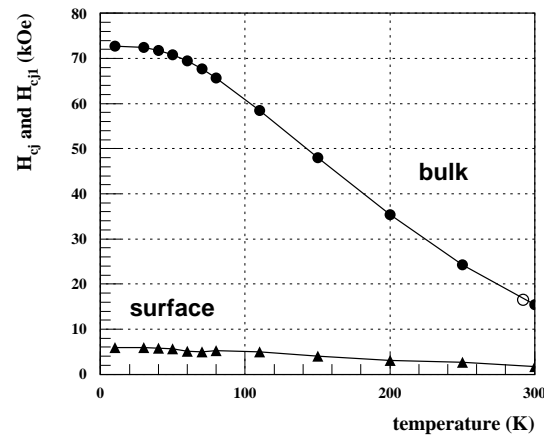
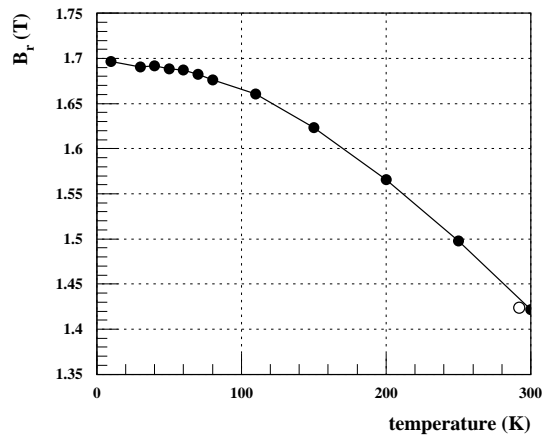
$$J_{pole} = |J_{bulk}(0)| \frac{B_{bulk}}{|B_{bulk}|} \quad \text{for } |B_{bulk}| > |J_{bulk}(0)|$$

$$J_{side} = J_{bulk}(0) \frac{H_{bulk} - H_{cj1}}{|H_{bulk} - H_{cj1}|}$$

12 Fitting parameters $J_1, J_2, J_3, \xi_1, \xi_2, \xi_3, \chi_0, \chi_1, H_{cj}, H_{cj1}, d_s, d_p$

Results from fit:

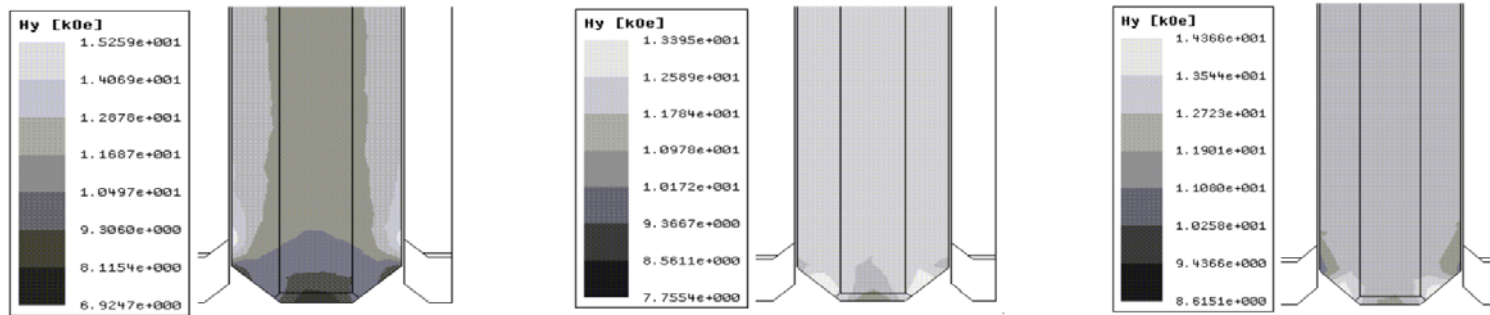
- energy product (BH)max @ 85 K = 520 kJ/m³
- $B_r(10K)=1.7T$
- $H_{cj}=73kOe$



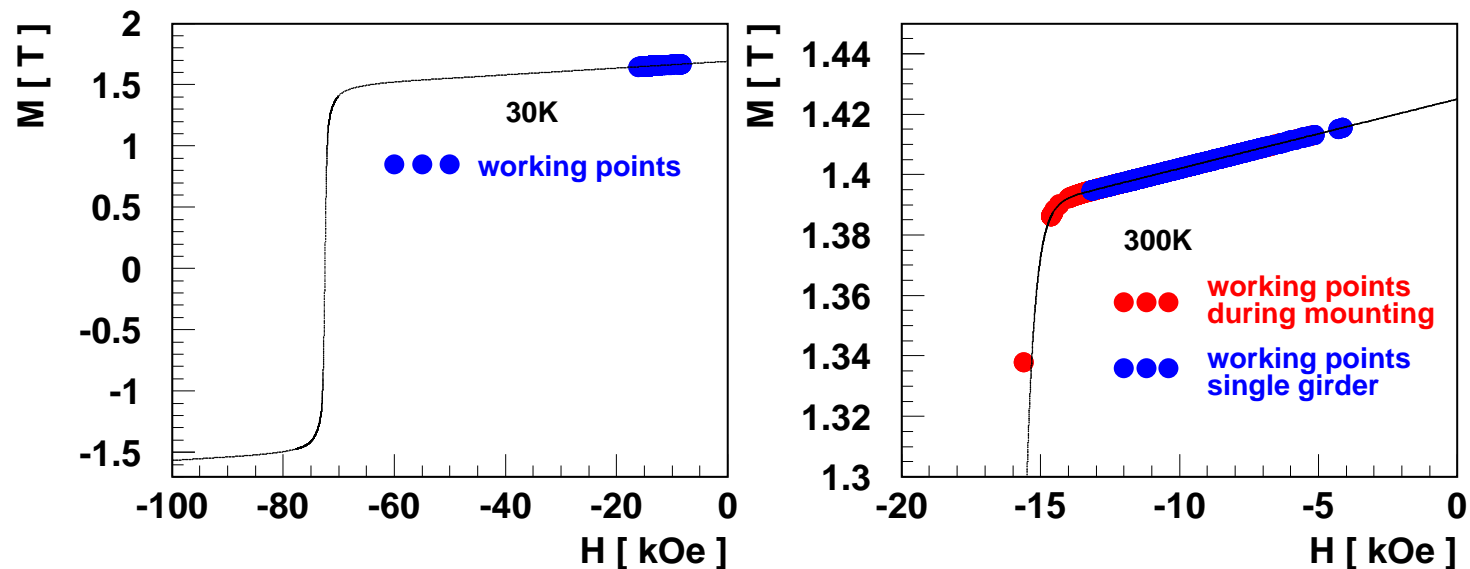
K. Üstüner et al., 20th Conference on Rare Earth Permanent Magnets, Crete 2008

Comparison with other magnet grades

Magnetic Design of HZB-LMU Prototype



high reverse fields at pole tips
even higher reverse fields during assembly



Assembly and First Test with Beam

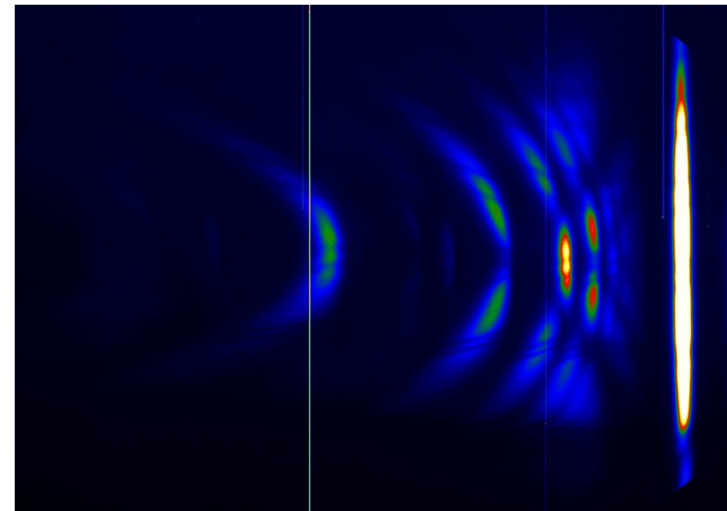


Many thanks to Tasso

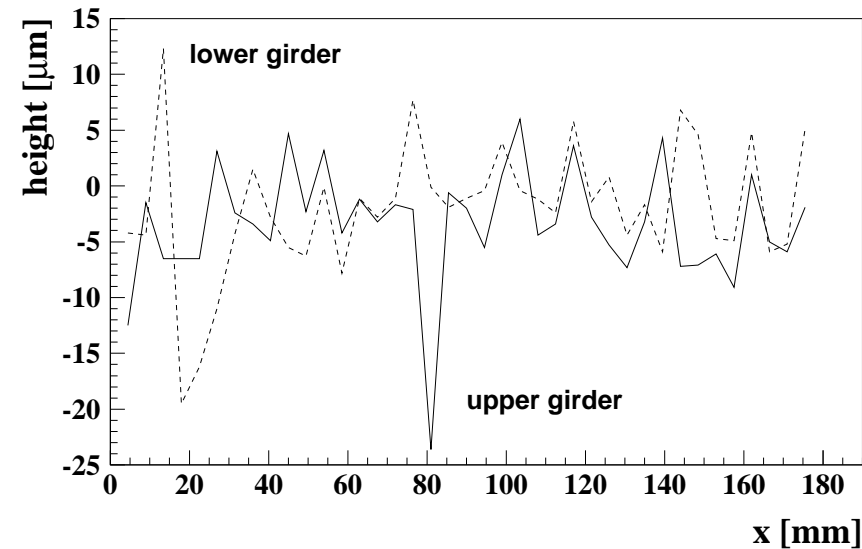
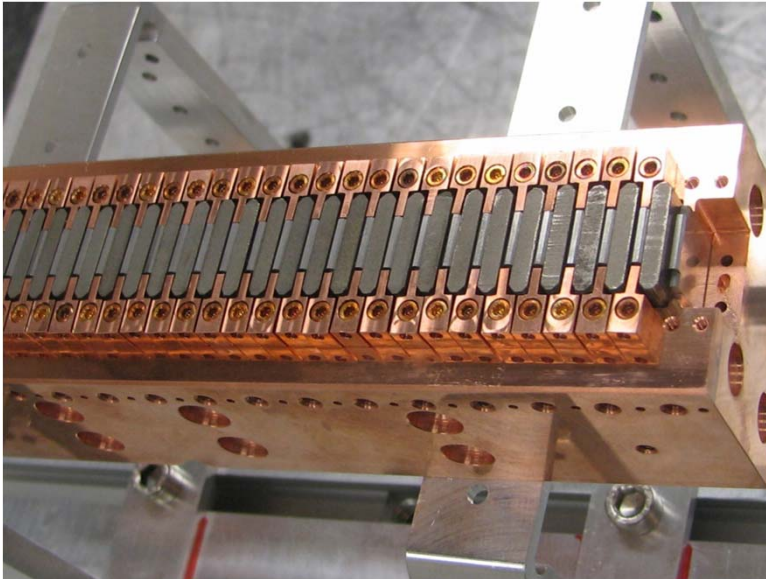
LYKOS

-22°C
Huge Windchill-Factor!
(when the fans are blowing)

Fixed gap device
Magnet material: PrFeB
Pole material: CoFe
Design: Hybrid
Period length: 9mm
Period number: 20
Field gain at low temperatures: 15%



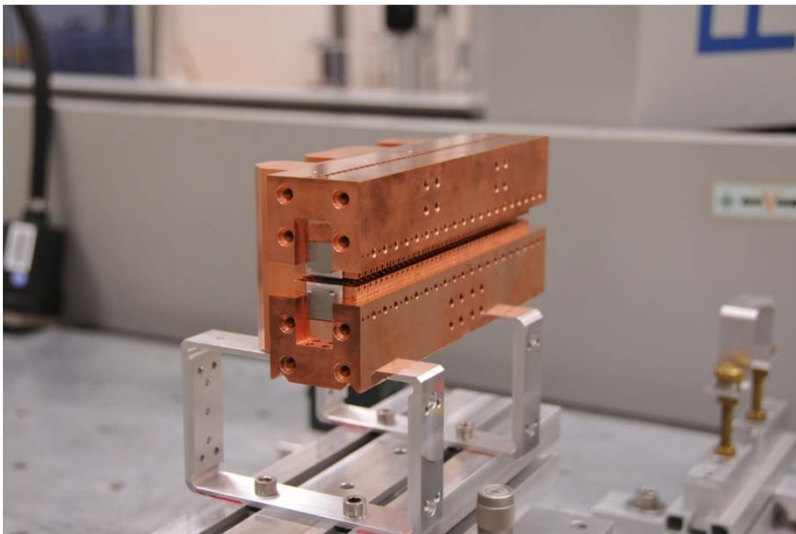
First spectrum taken at MAMI, 850 MeV
Transmission grating, F. Holy et al.

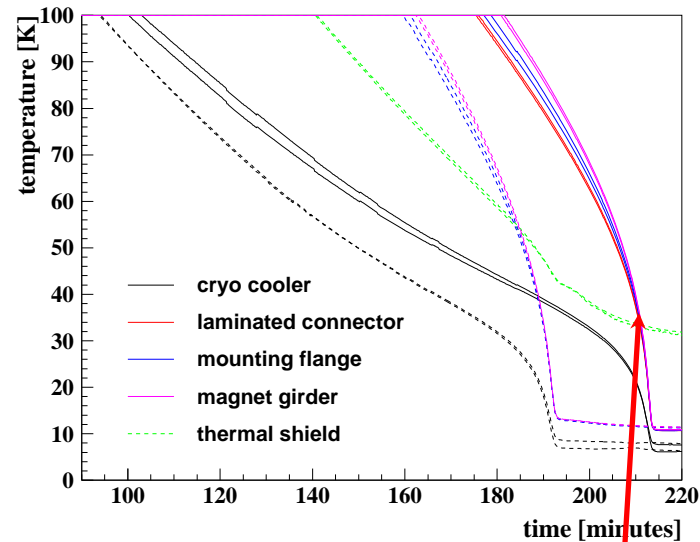
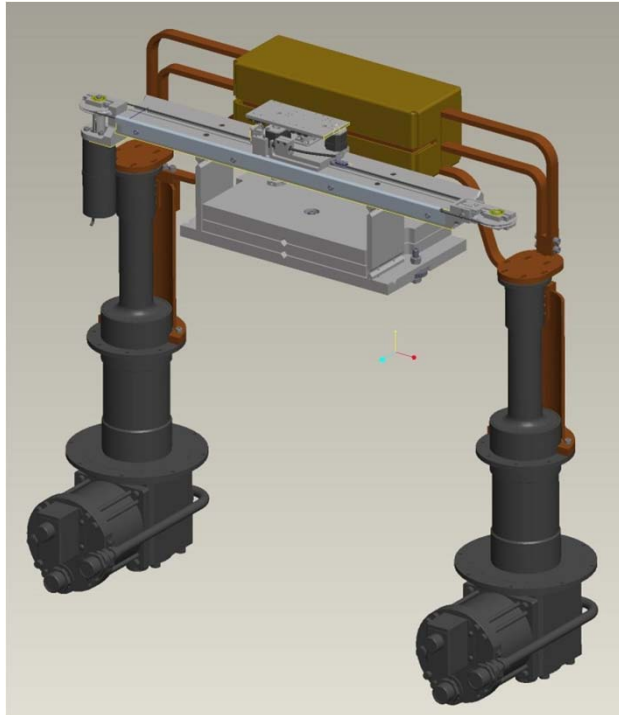


Pole heights after initial assembly

Field quality is determined
by geometric tolerances
high precision machining required

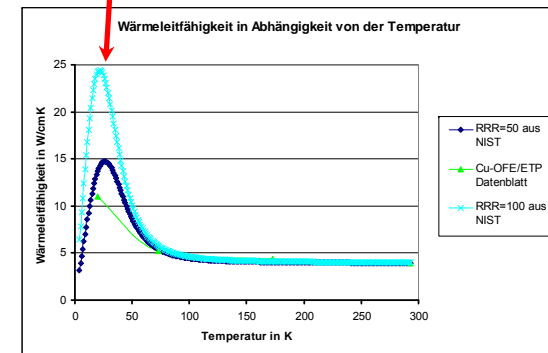
field tuning by pole height adjustment
use of pole clamps with various heights





Cooling time:
3h 10min without shield
2h 50min with shield
final temperature is
the same in both cases

Calculated end temperature of the undulator: 17K
Measured end temperature of undulator: 11K



Reason for discrepancy is a pessimistic assumption for the Cu heat conductivity

Mechanism of reversible demagnetization:

- Local heating (at grain boundaries) above critical temperature
- Demagnetization of single domain and subsequent flip of the whole grain (nucleation type magnets)

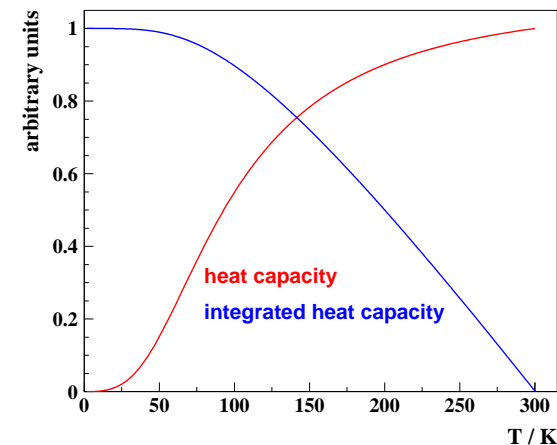
Cryogenic devices enhance allowed heat deposition below critical T the stability increases at lower temperatures but:

- the heat capacity decreases at low temperatures
- **high coercivity alone does not help**

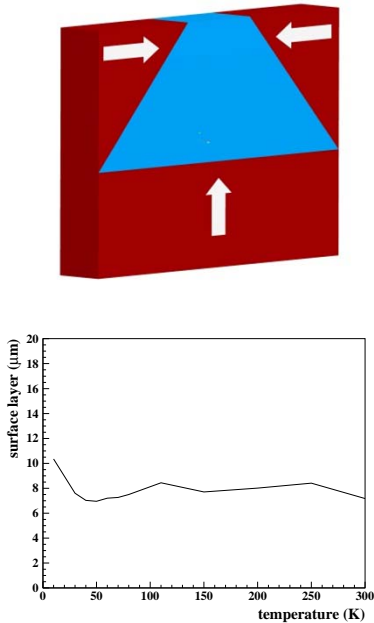
$$C_p(T) \sim \left(\frac{T}{T_D}\right)^3 \int_0^{\frac{\theta_D}{T}} \frac{x^4 e^x}{(e^x - 1)^2} dx$$

$$h \theta_D \approx 400.$$

H. Fujii et al., J. of Magnetism and Magn. Materials 70 (1987) 331-333



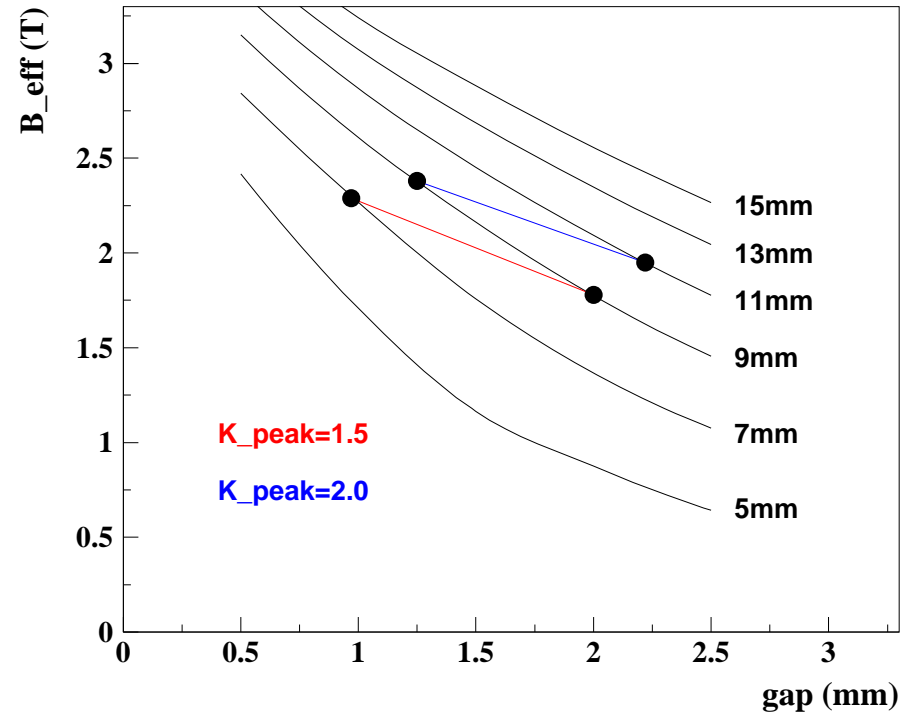
Pushing the Parameters with Shaped Poles



surface layer thickness: $8.5 \pm 1.5 \mu\text{m}$
Slightly more than grain size

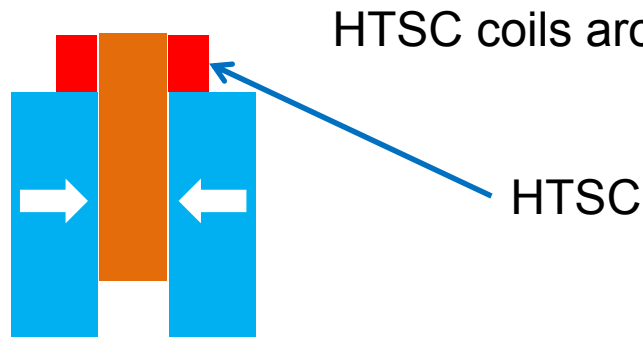
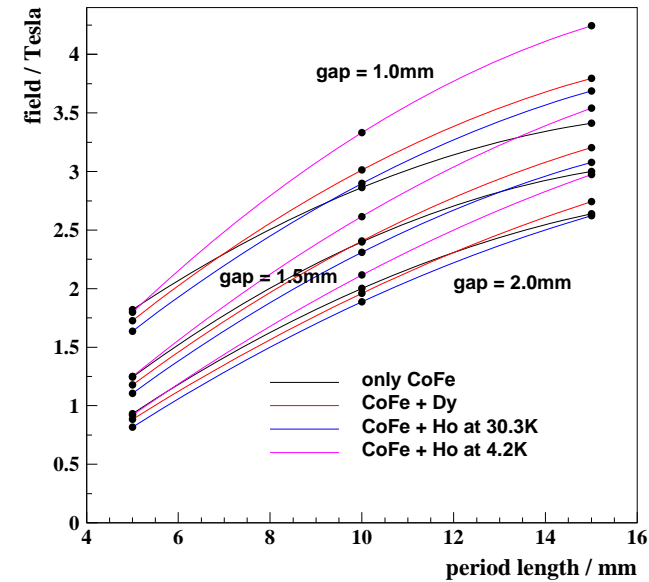
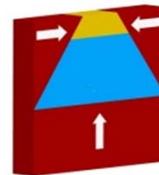
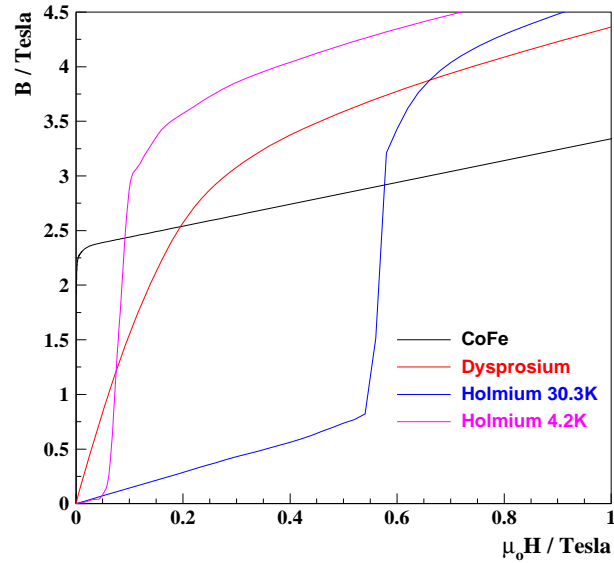
Simulations:

- minimum dead layer ($9 \mu\text{m}$)
- TiN coating ($3\text{-}4 \mu\text{m}$)
- tight shape tolerances ($<10 \mu\text{m}$)
- small air gaps ($10\text{-}20 \mu\text{m}$)



	<i>1st prototype</i>	<i>2nd prototype</i>
Poles	straight	transverse wedged
Period	9mm	9mm
Gap	2.5mm	1.5mm
Beff	1.15 T	1.45T @ 2.5mm 2.15T @ 1.5mm

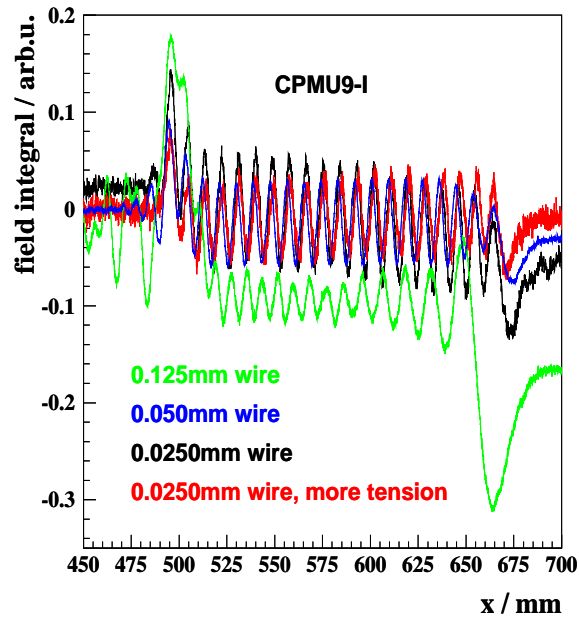
Substitution of part of CoFe poles with Dy or Ho



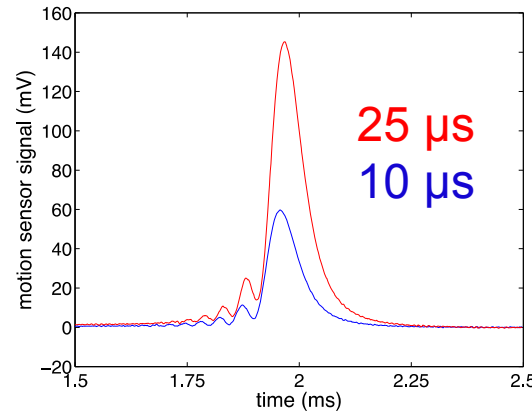
HTSC coils around pole pieces

Coil support:
impregnation with resin epoxy
enhanced stability
further R & D necessary

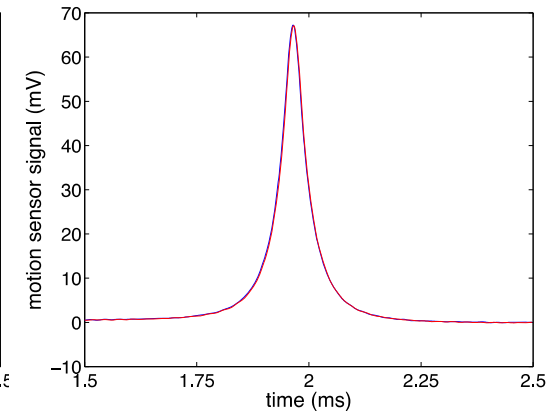
T. Tanaka et al., New Journal of Physics, 8 (2006) 287



before correction



after correction



Courtesy of D. Arbelaez, LBNL

$$\rho A \frac{\partial^2 y}{\partial t^2} - T \frac{\partial^2 y}{\partial z^2} + EC \frac{\partial^4 y}{\partial z^4} = B_x I$$

$$C = \frac{\pi d^4}{64}$$

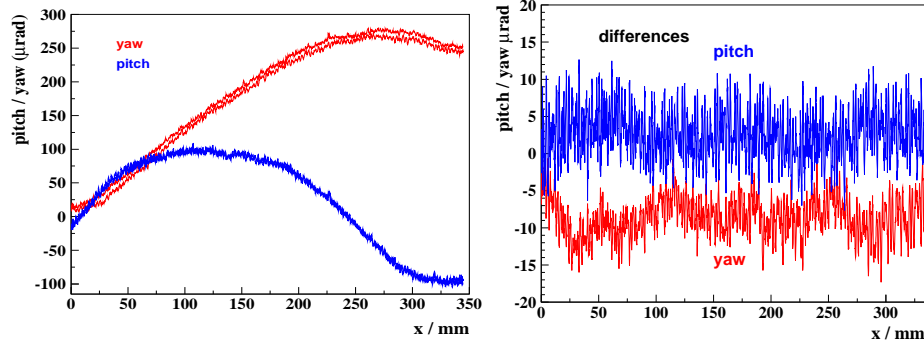
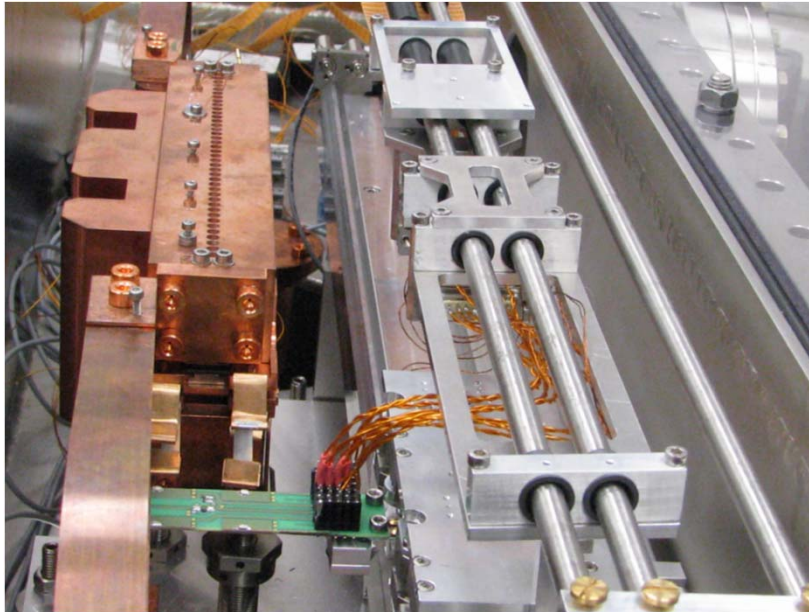
Free space propagation: $c = c_0 \sqrt{1 + \frac{EC}{T} k^2}$

$$y = e^{ik(z-ct)}$$

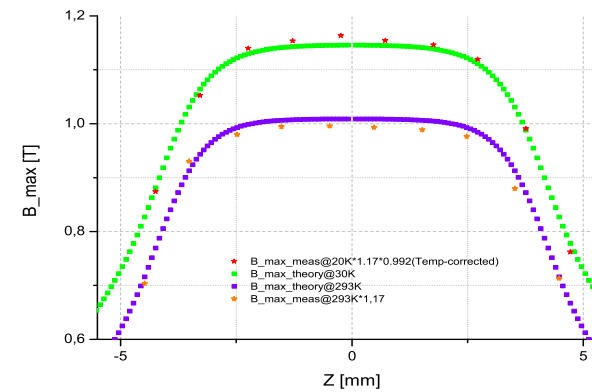
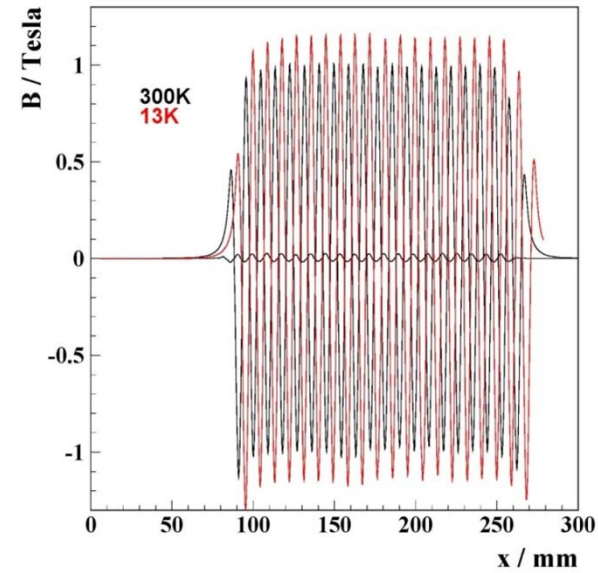
with $c_0 = \sqrt{\frac{T}{\rho A}}$

Numerical solution of inhomogenous
time dependent differential equation

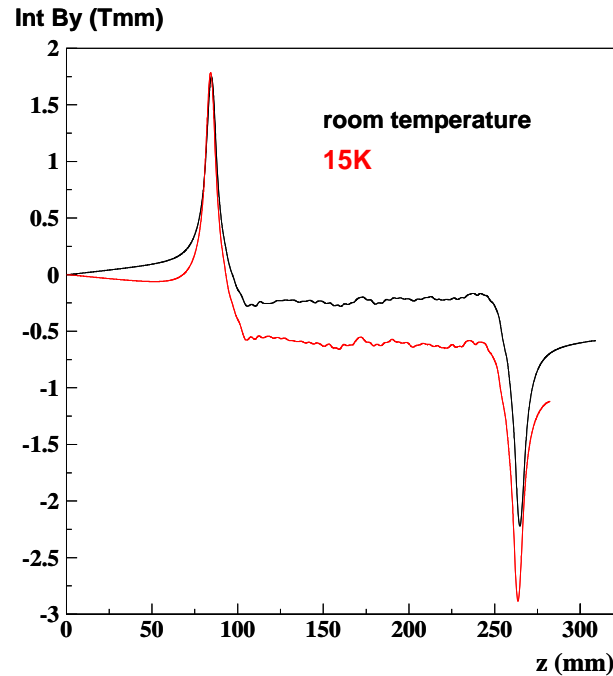
In-Vacuum Hall probe bench for local measurements



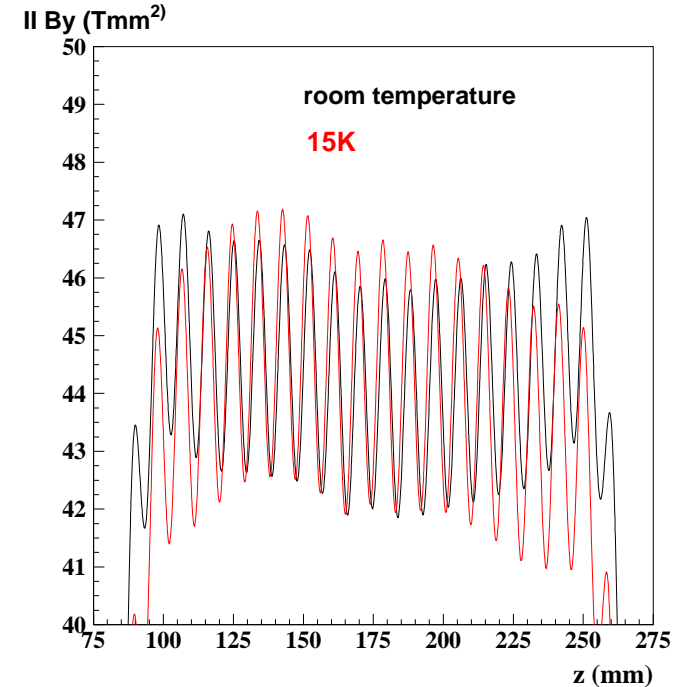
Field data of 9mm prototype



Differences Between Room Temperature and 15K



Field integral (filtered) jumps by 0.4Tmm due to temperature dependent partial saturation of endpoles



2nd field integral considerable deviations observed only at the ends

Precise analysis requires HP-calibration (T-dependent)
- HP-temperature at 15K: room temperature -15K

Higher fields possible with

- Cryogenic temperatures
- New materials, PrFeB, GBD
- Ambitious magnet designs, narrow good field region

Magnetic quality assurance with improved systems

- pulsed wire, correcting for dispersion
- In-vacuum Hall probe bench

Mechanic challenges

- Tighter fabrication tolerances
- Advanced gap measurement systems
- Careful cryogenic design