



MAX PLANCK INSTITUTE

for the science of light

Friedrich-Alexander-Universität
Erlangen-Nürnberg



Pulse-Splitting in Seeded Free Electron Lasers

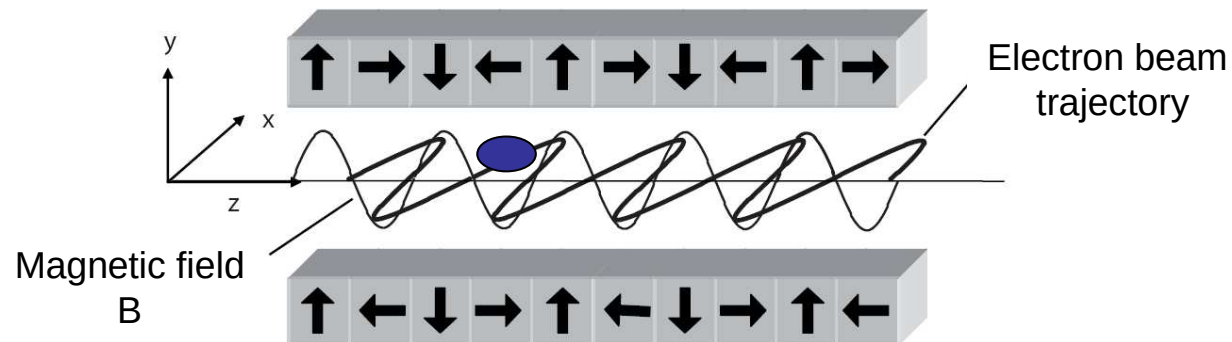
M. Labat, N. Joly, S. Bielawski,
C. Sz waj, C. Bruni, M.E. Couprie

Introduction to seeded FELs

Principle of an FEL

- **Amplifier medium:**

- Relativistic e- @ $E \approx \text{MeV-GeV}$
from accelerator (storage ring or LINAC)
- Periodic magnetic field
generated in an undulator

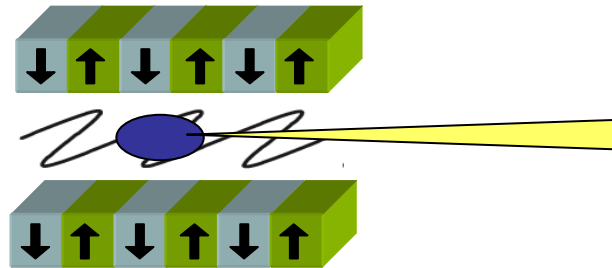


Introduction to seeded FELs

Principle of an FEL

- **Optical radiation:**

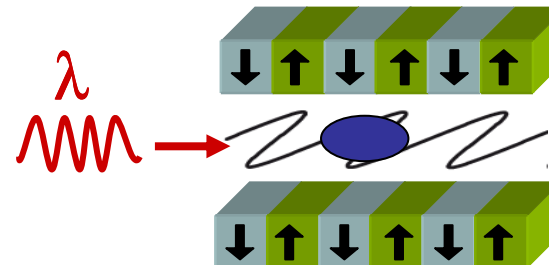
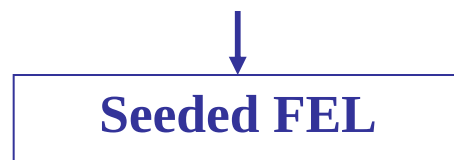
- Spontaneous emission:
Synchrotron radiation



Resonant wavelength:

$$\lambda = \frac{\lambda_u}{2n\gamma^2} \left(1 + \frac{K^2}{2}\right)$$

- External source or **seed**:
Conventional laser, etc..



Introduction to seeded FELs

Advantages of seeded FELs

- **Injection of a coherent seed** enables:

L.H. Yu, Phys. Rev. A 44 (1991).

- Reduction of the saturation length
- Higher temporal coherence
- Higher shot to shot stability
- Shorter and controlled pulse duration
- Stronger nonlinear harmonic generation
- More efficient cascading configurations

- **Seed sources:**

- Conventional lasers: $\lambda > 250 \text{ nm}$
- High order Harmonics Generated in gas (HHG): $270 > \lambda > 1 \text{ nm}$

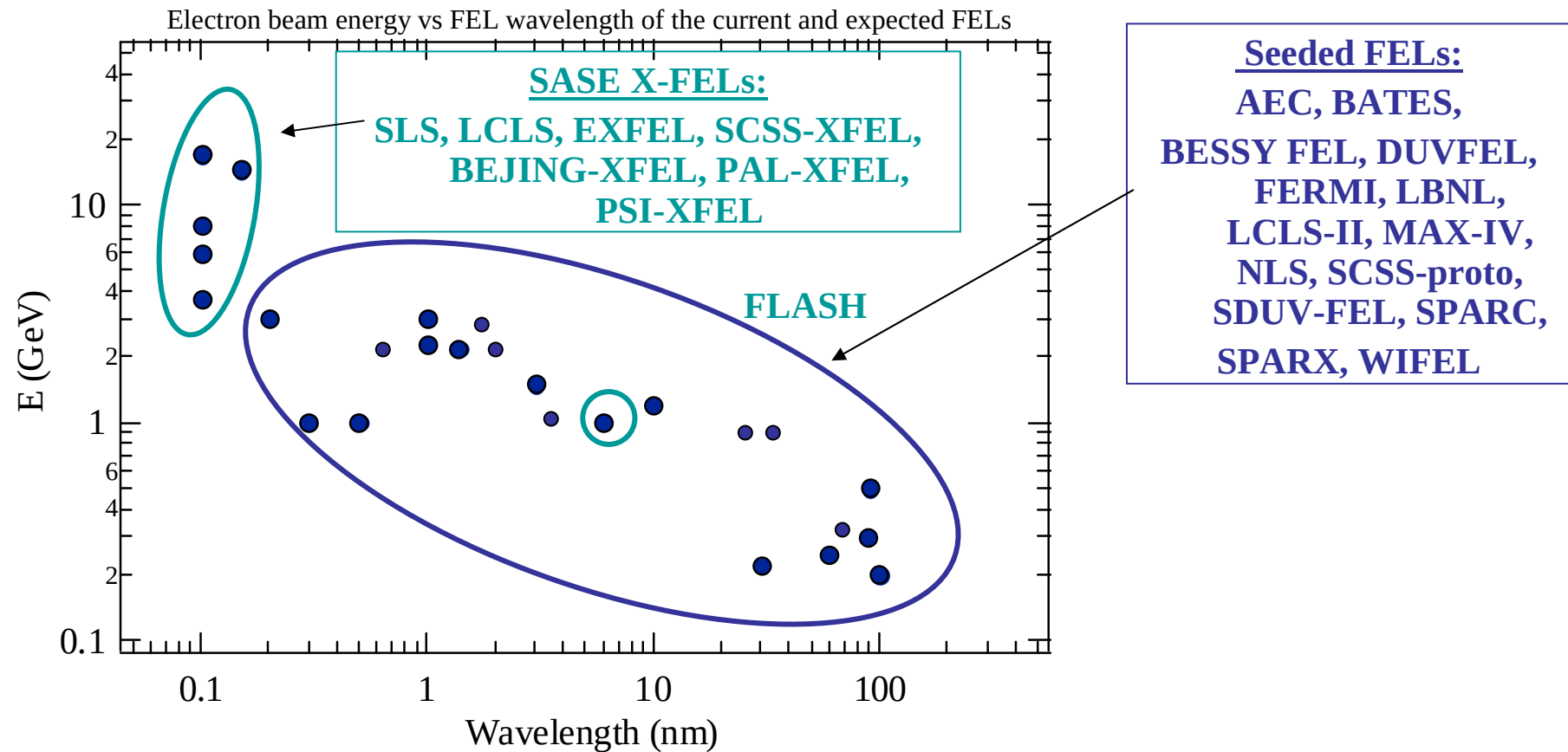
D. Garzella et al., NIMA 528, 502 (2004).

M.E. Couprie et al., NIMA 528, 507 (2004).

→ **HHG + HGHG : Coherent compact FELs at very short wavelengths**

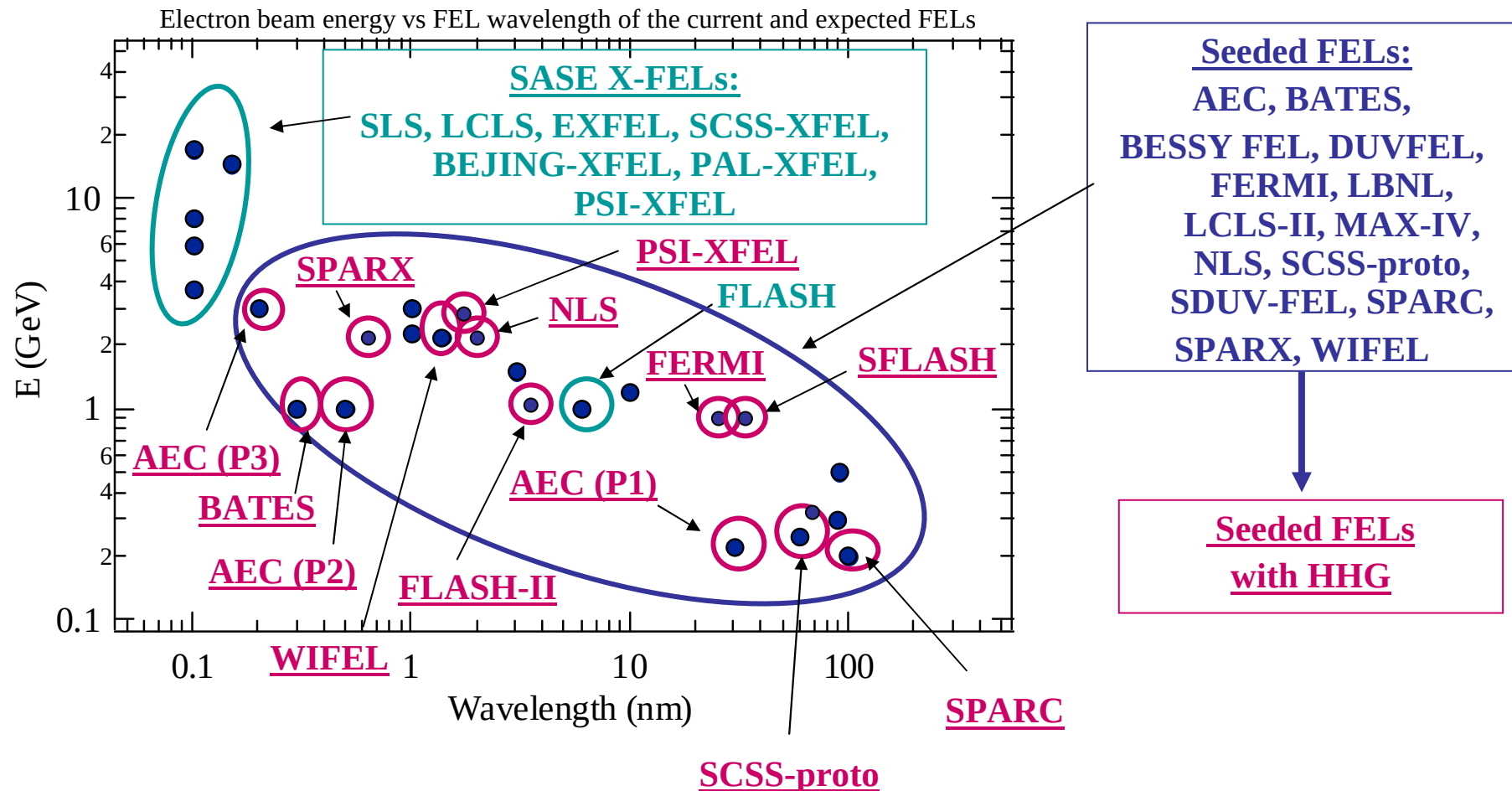
Introduction to seeded FELs

Seeded FELs around the world



Introduction to seeded FELs

Seeded FELs around the world



The ARC-EN-CIEL project



Introduction

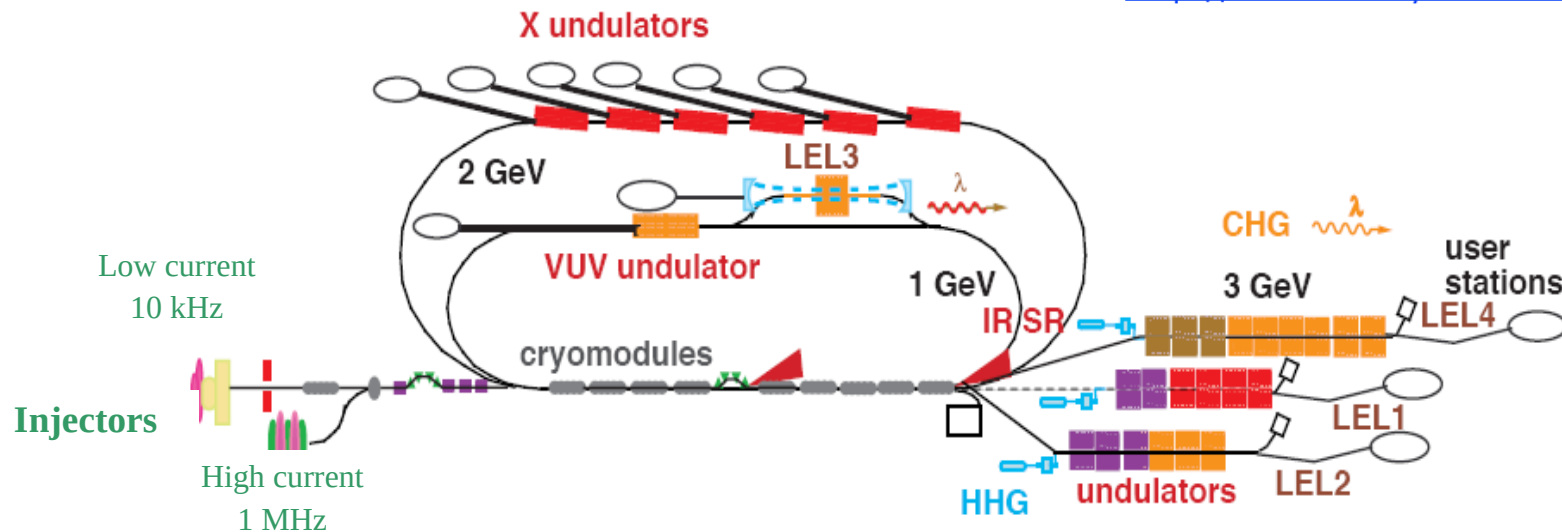
<http://arcenciel.synchrotron.fr/ArcEnCiel>

The ARC-EN-CIEL project



Light sources

<http://arcenciel.synchrotron.fr/ArcEnCiel>



- **4th generation light source:**

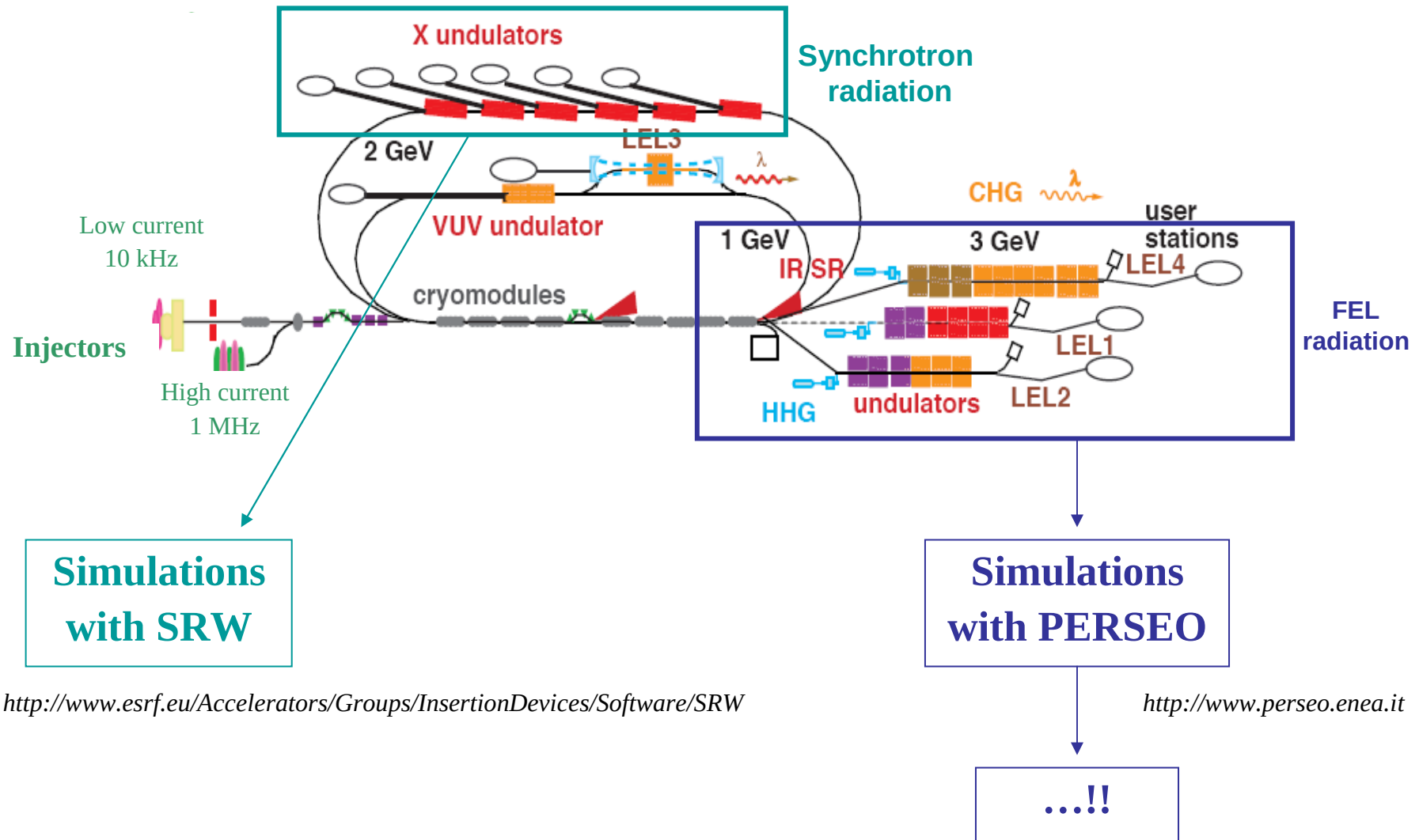
- Synchrotron radiation from undulators (IR to X-rays):
U20 and U30 (planar + helical)
- 4 Free Electron Lasers:

3 seeded HGHG ($\lambda \rightarrow 0.3$ nm) + 1 oscillator (VUV)

The ARC-EN-CIEL project



Writing of the CDR



Simulation of a VUV seeded FEL

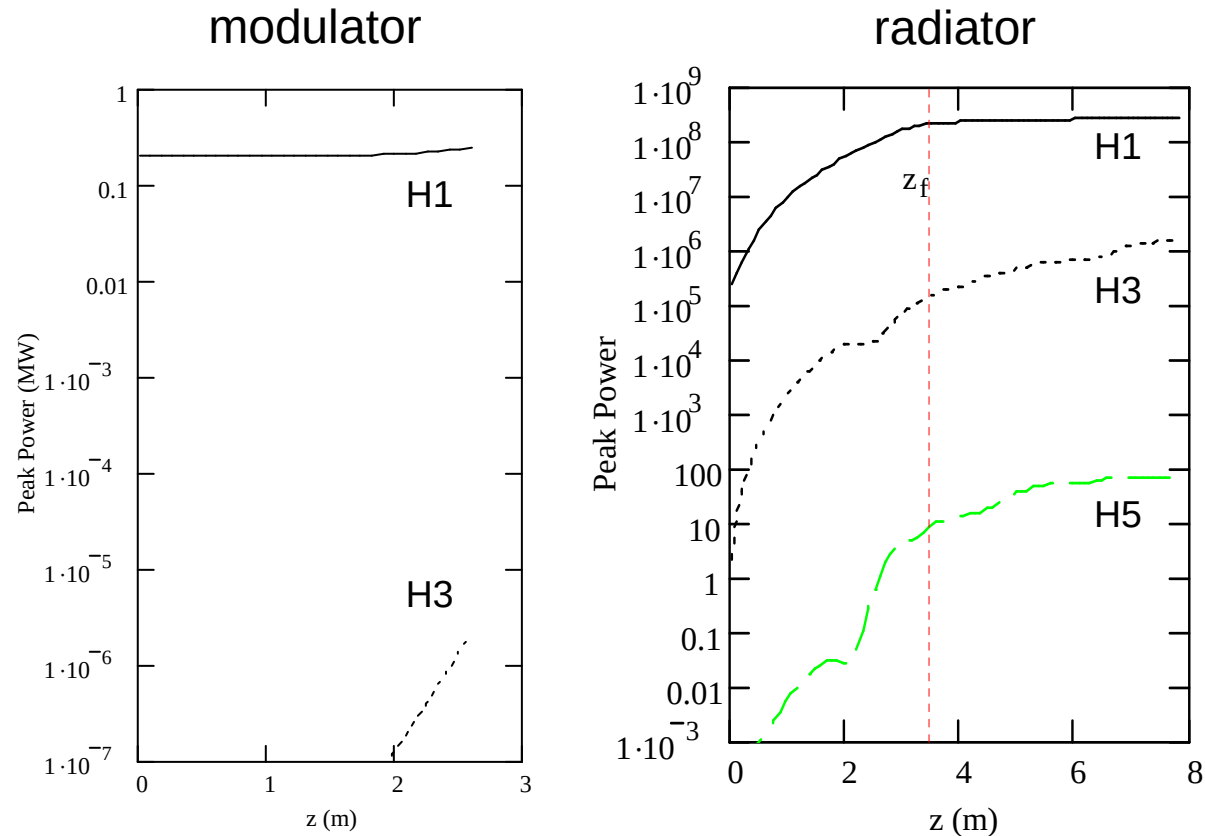
Parameters for AEC – FEL1

- Electron beam:
 - $E=220$ MeV, $I=300$ A, $\sigma_z=450$ fs-rms, $\varepsilon=1.2.\pi$.mm.mrad
- Seed:
 - 114 nm
 - 200 kW in 150 fs-fwhm
- Undulators:
 - Modulator: 100 periods of 26 mm
 - Dispersive section: $R_{56}=1.5$ μ m
 - Radiator: 260 periods of 30 mm
- **Final wavelength = 114 nm**

Simulation of a VUV seeded FEL

Performances of AEC – FEL1

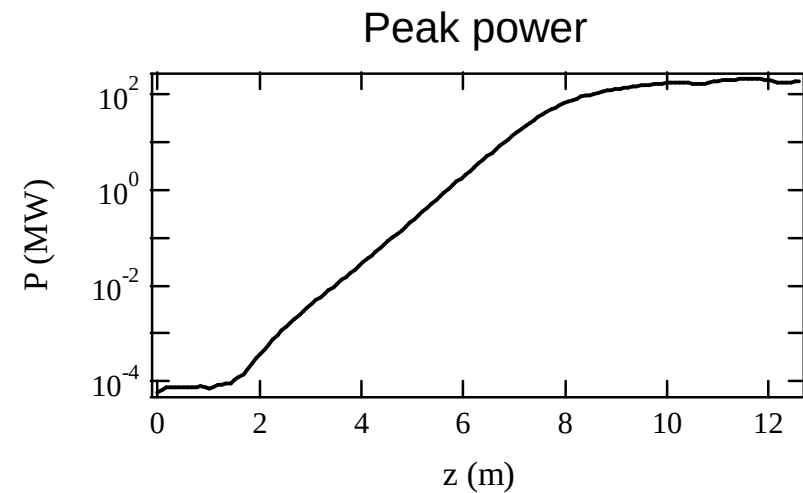
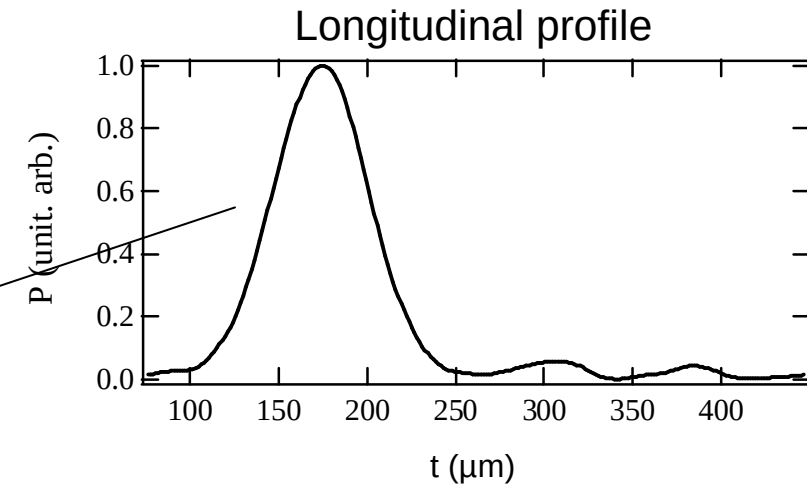
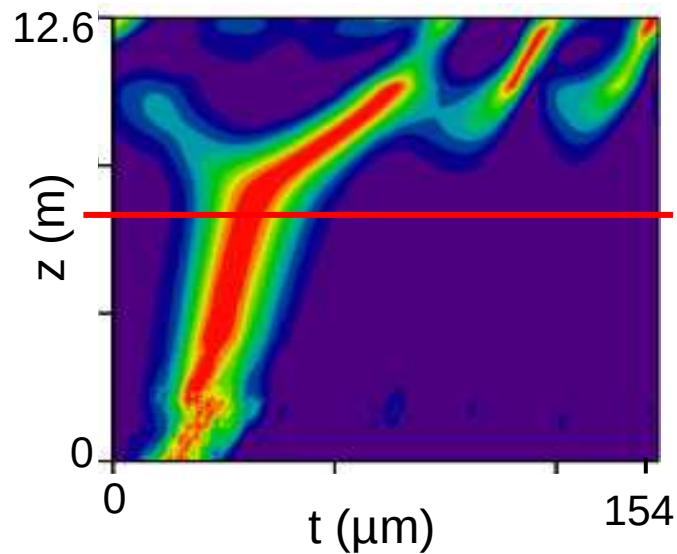
Peak power evolution in the:



Simulation of a VUV seeded FEL

« Performances » of AEC – FEL1

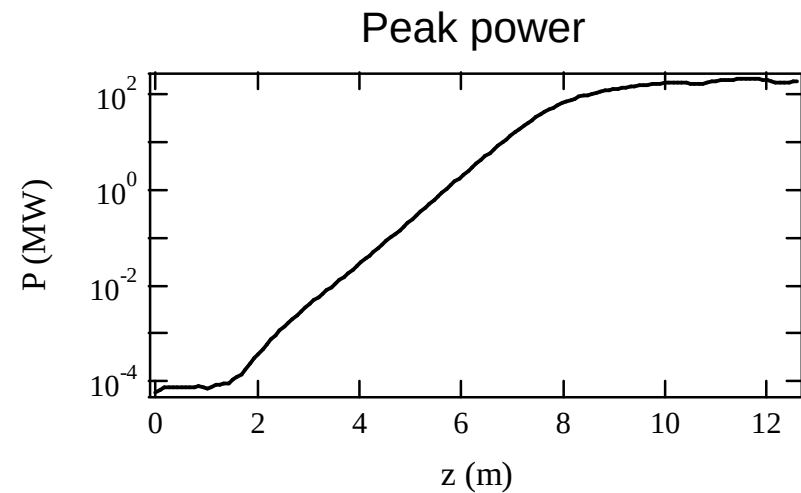
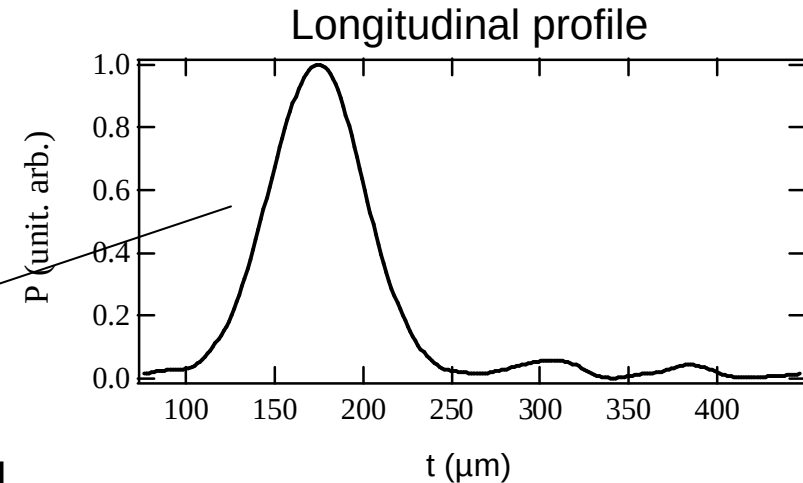
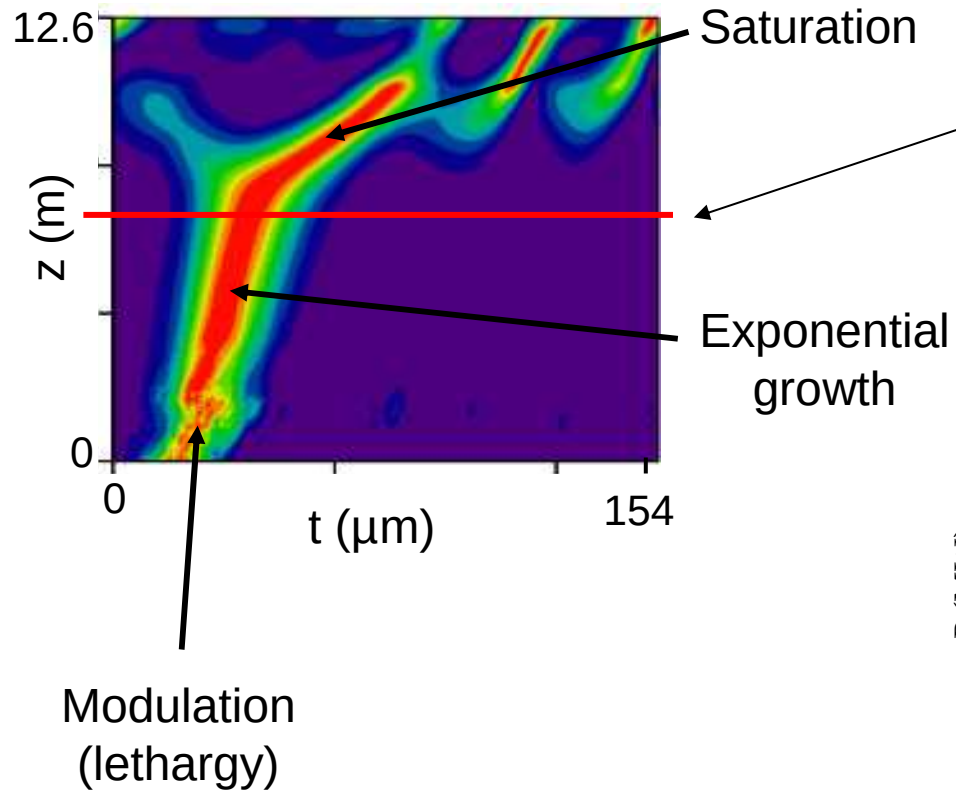
Looking at spatio-temporal diagrams



Simulation of a VUV seeded FEL

« Performances » of AEC – FEL1

Looking at spatio-temporal diagrams

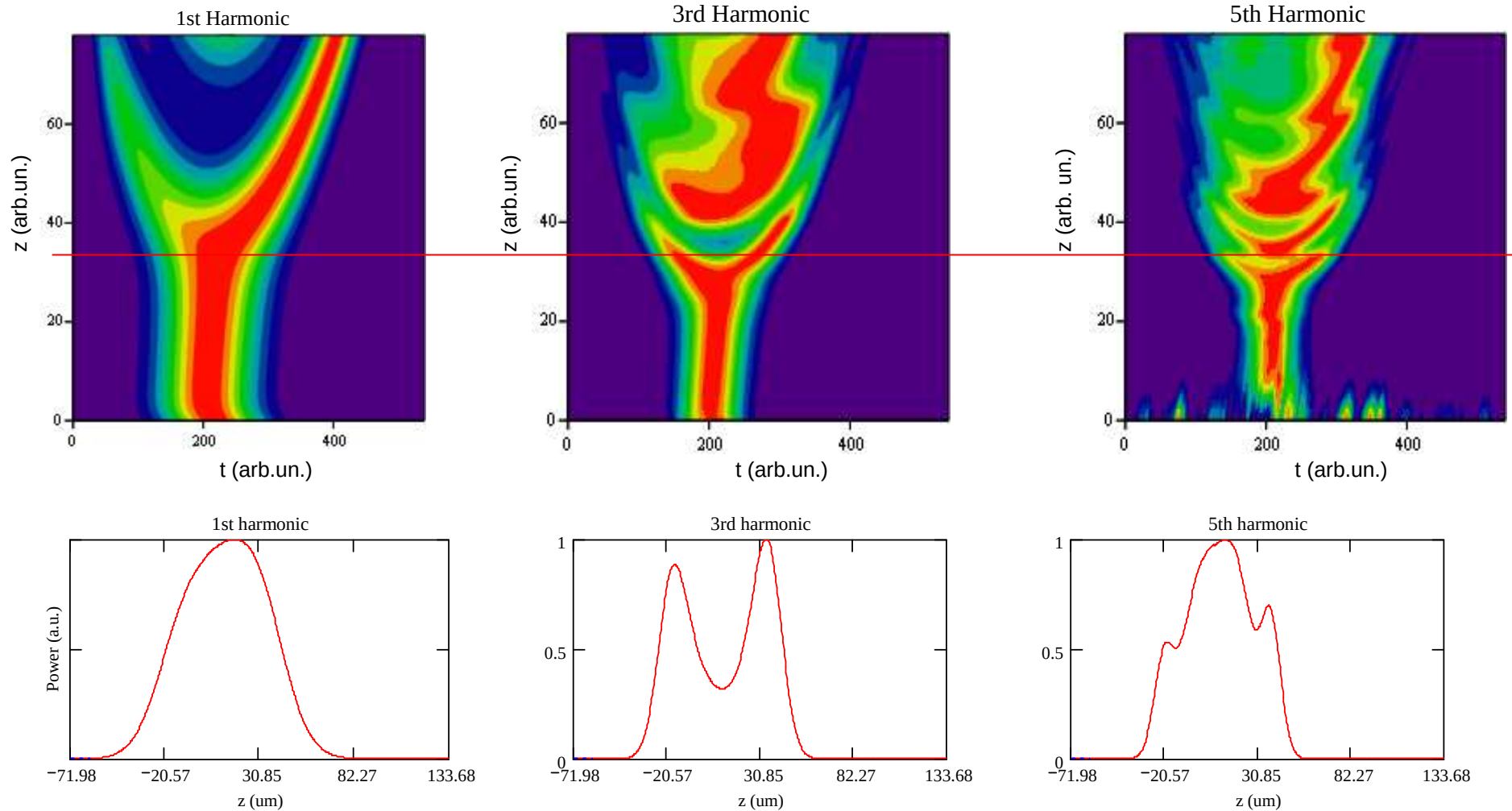


Simulation of a VUV seeded FEL

« Performances » of AEC – FEL1

Simulation of a VUV seeded FEL

« Performances » of AEC – FEL1



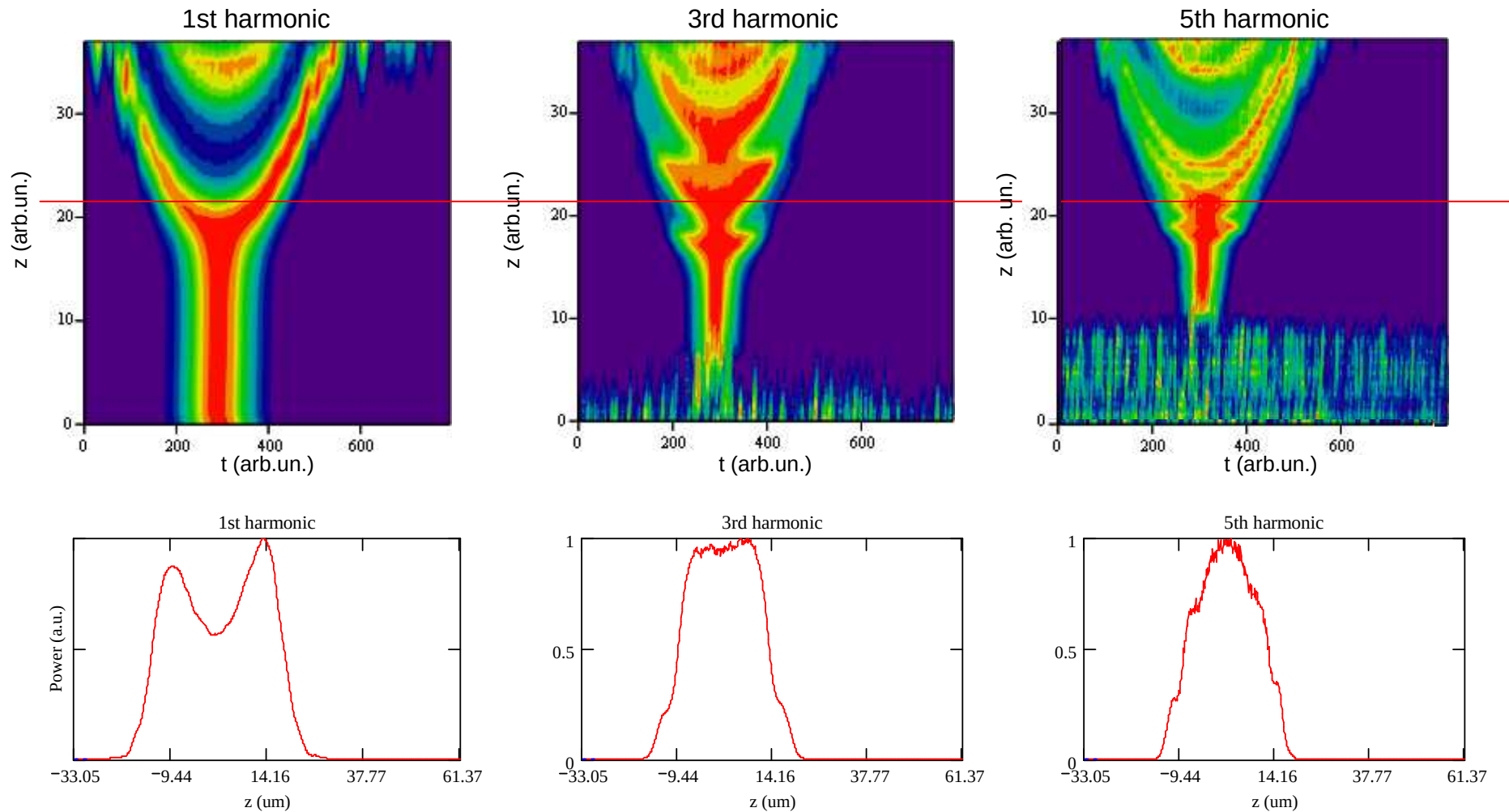
Simulation of a XUV seeded FEL

Parameters of AEC – FEL2

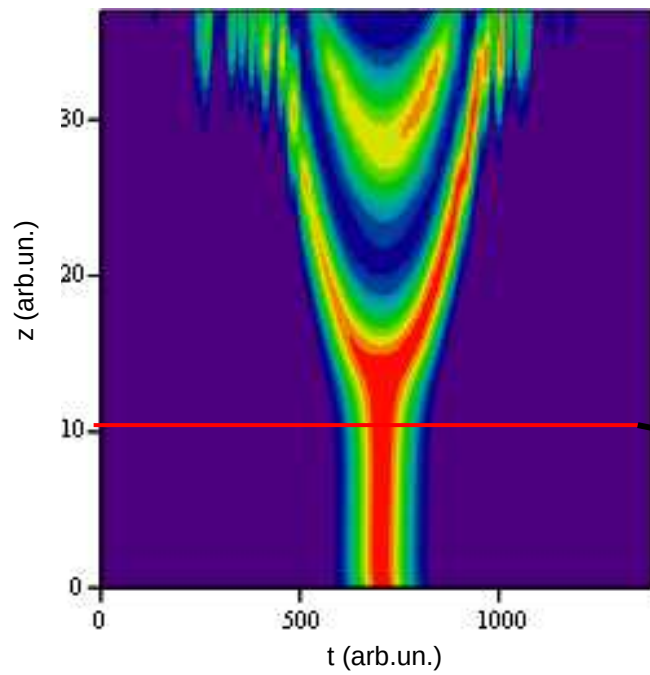
- Electron beam:
 - $E=1$ GeV, $I=1500$ A, $\sigma_z=200$ fs-rms, $\varepsilon=1.2.\pi$.mm.mrad
- Seed:
 - 8.9 nm
 - 30 kW in 50 fs-fwhm
- Undulators:
 - Modulator: 200 periods of 26 mm
 - Dispersive section: $R_{56}=1.5$ μ m
 - Radiator: 500 periods of 15 mm
- **Final wavelength = 8.9 nm**

Simulation of a **XUV** seeded FEL

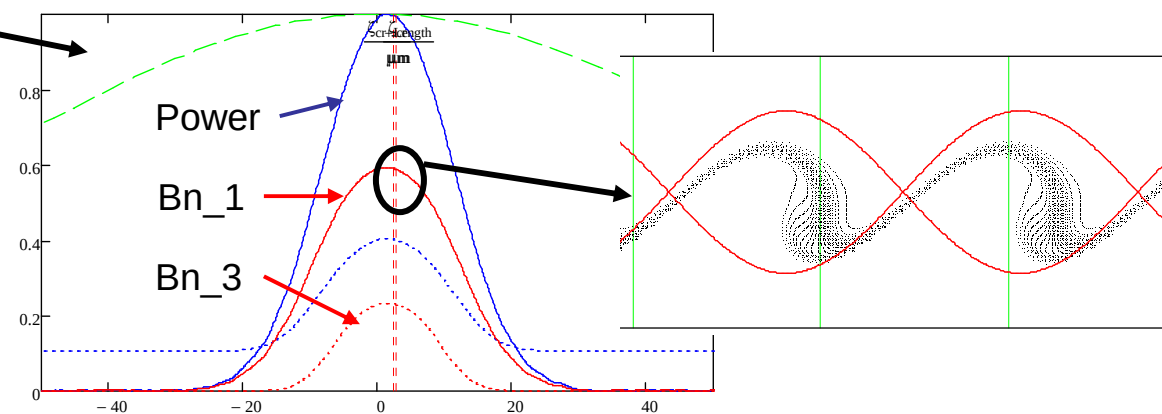
« Performances » of AEC – FEL2



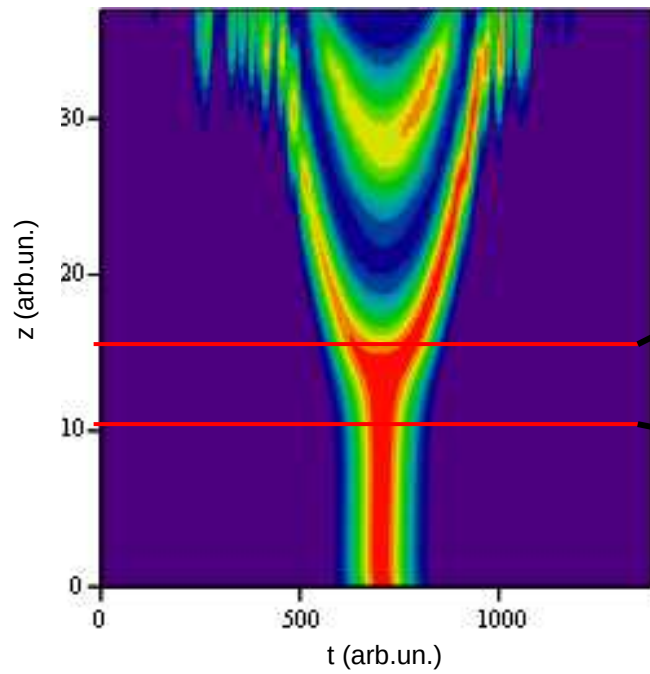
Pulse splitting and overbunching



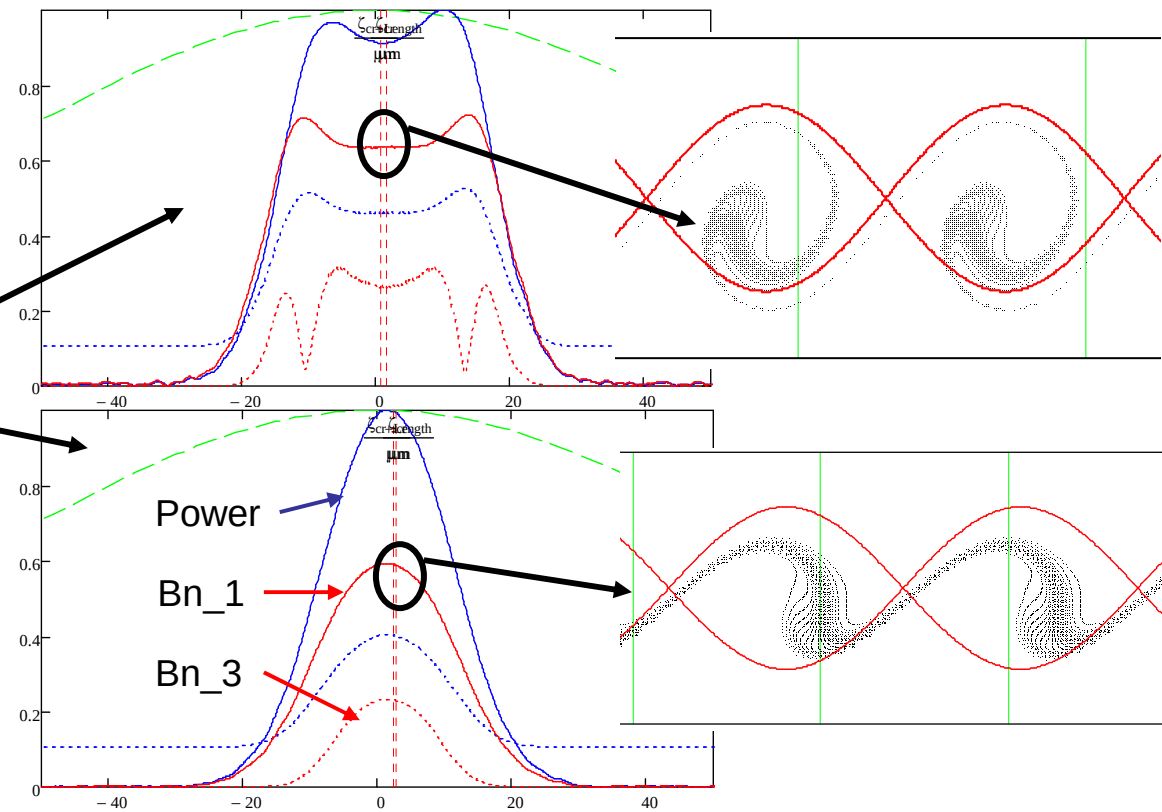
1 GeV e- beam.
Seed: 300 kW @ 8.9 nm.
200xU26 – R₅₆ – 500xU26.



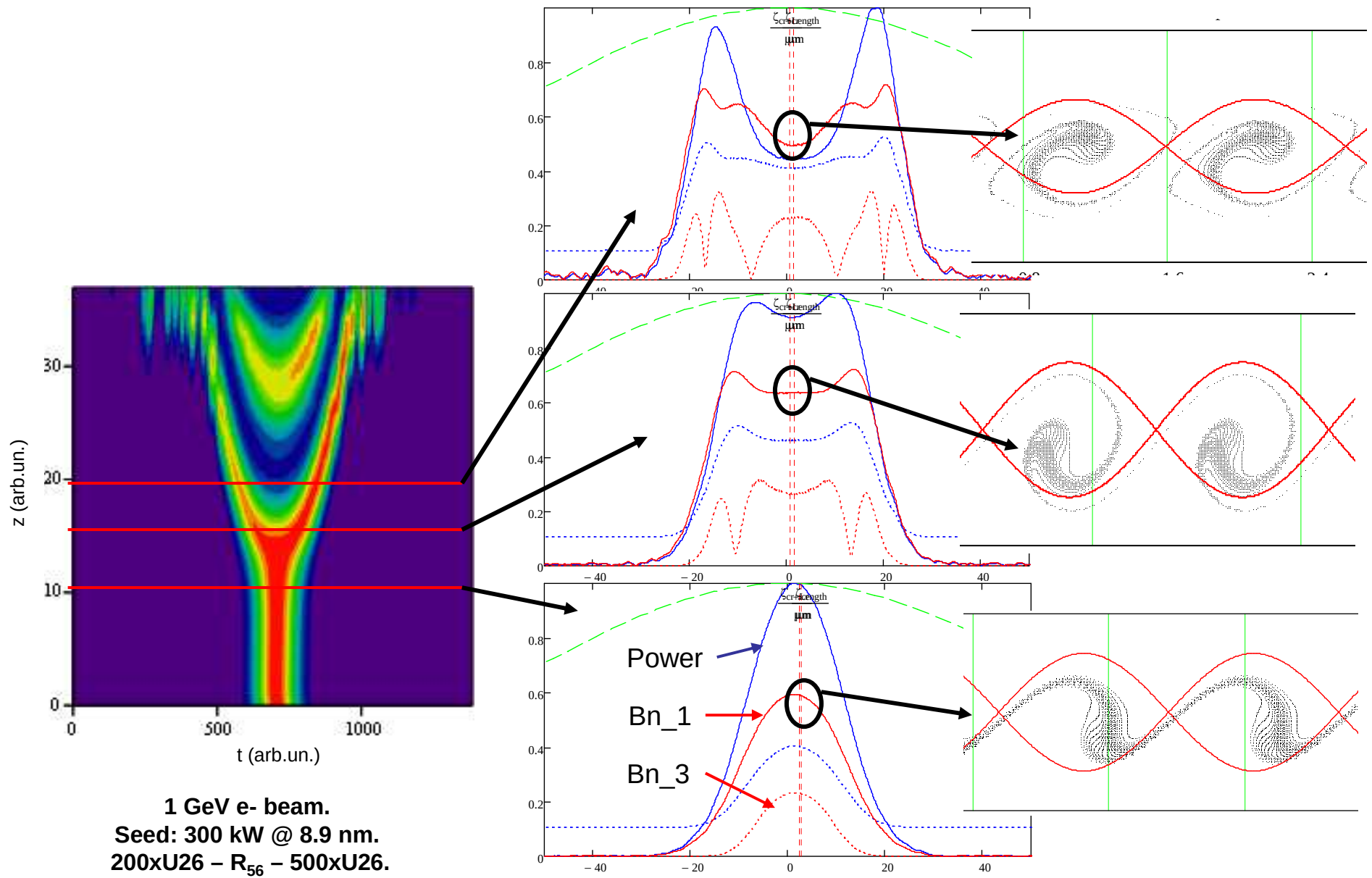
Pulse splitting and overbunching



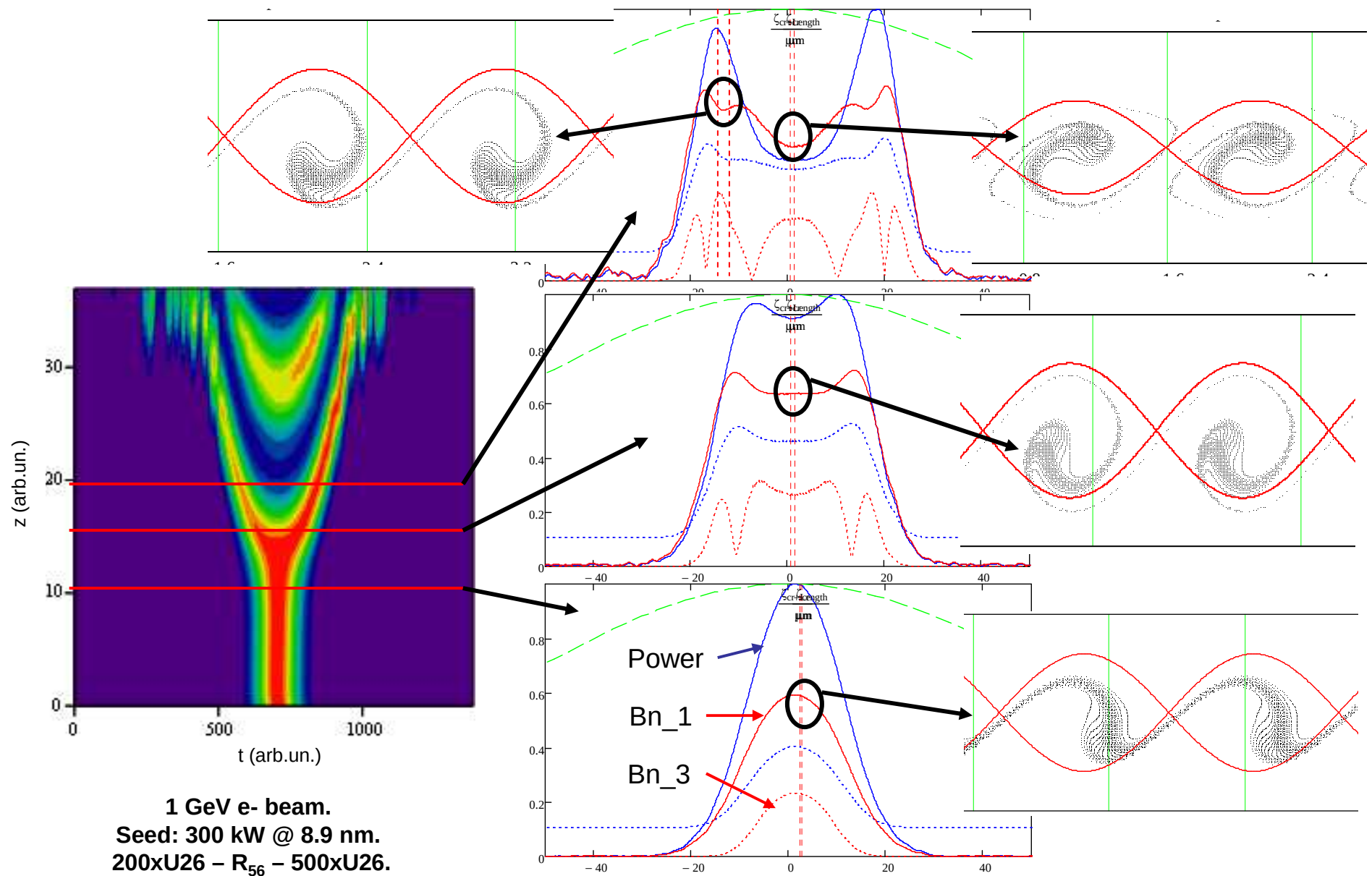
1 GeV e- beam.
Seed: 300 kW @ 8.9 nm.
200xU26 – R₅₆ – 500xU26.



Pulse splitting and overbunching

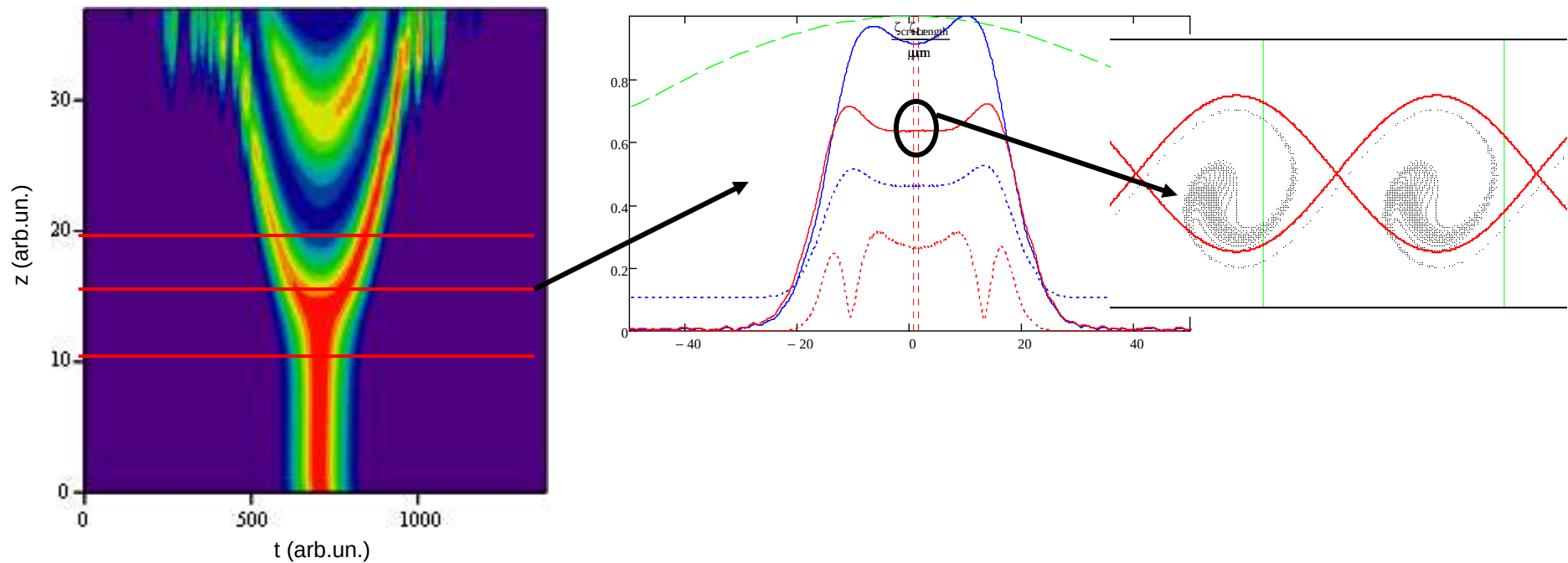


Pulse splitting and overbunching



Simulation of a **XUV** seeded FEL

Pulse splitting and overbunching

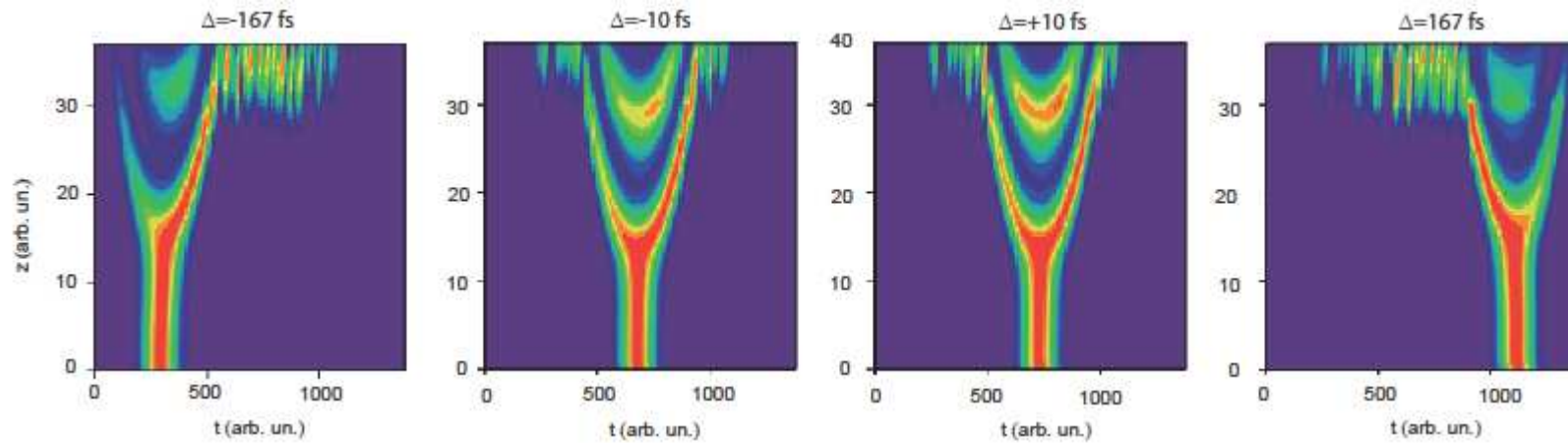


Pulse splitting is related to local overbunching

Pulse splitting dependencies

Pulse splitting dependencies

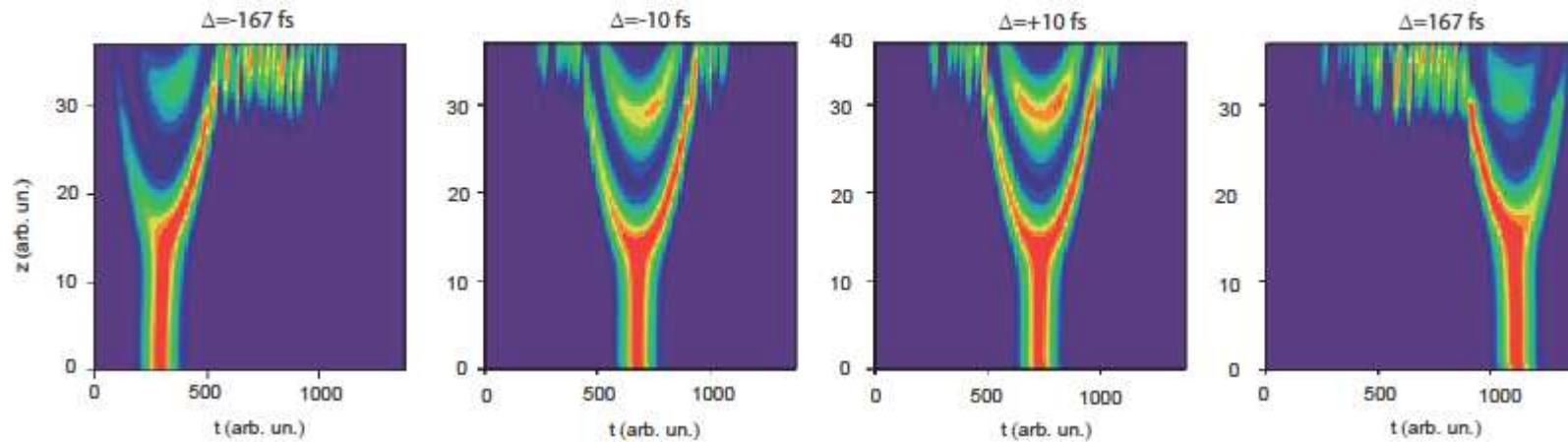
Vs synchronization



1 GeV e- beam. Seed: 300 kW @ 8.9 nm. 200xU26 – R₅₆ – 500xU26.

Pulse splitting dependencies

Vs synchronization



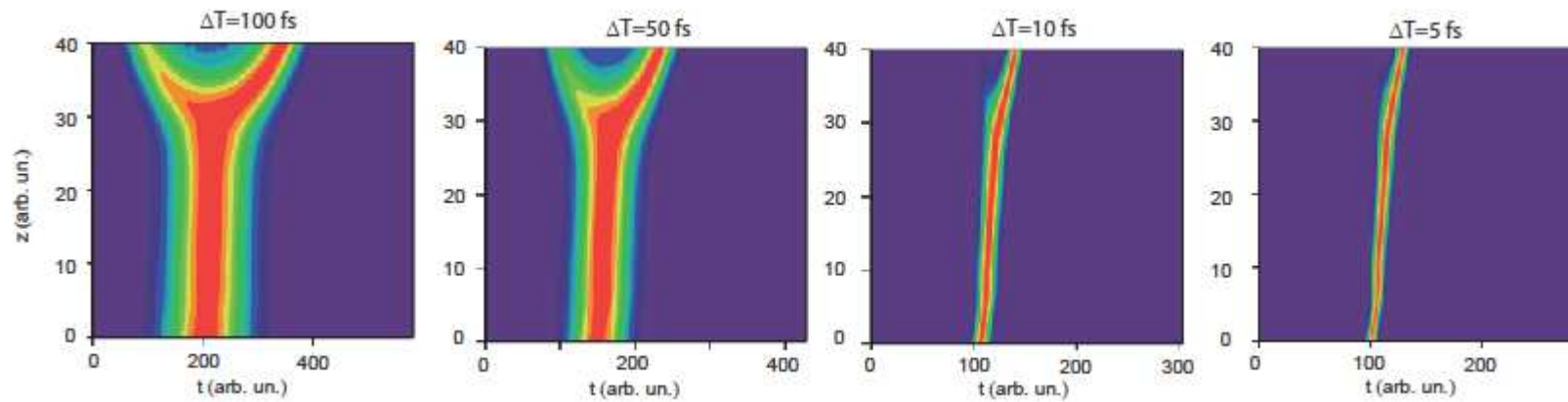
1 GeV e- beam. Seed: 300 kW @ 8.9 nm. 200xU26 – R₅₆ – 500xU26.



Pulse splitting follows the gain medium

Pulse splitting dependencies

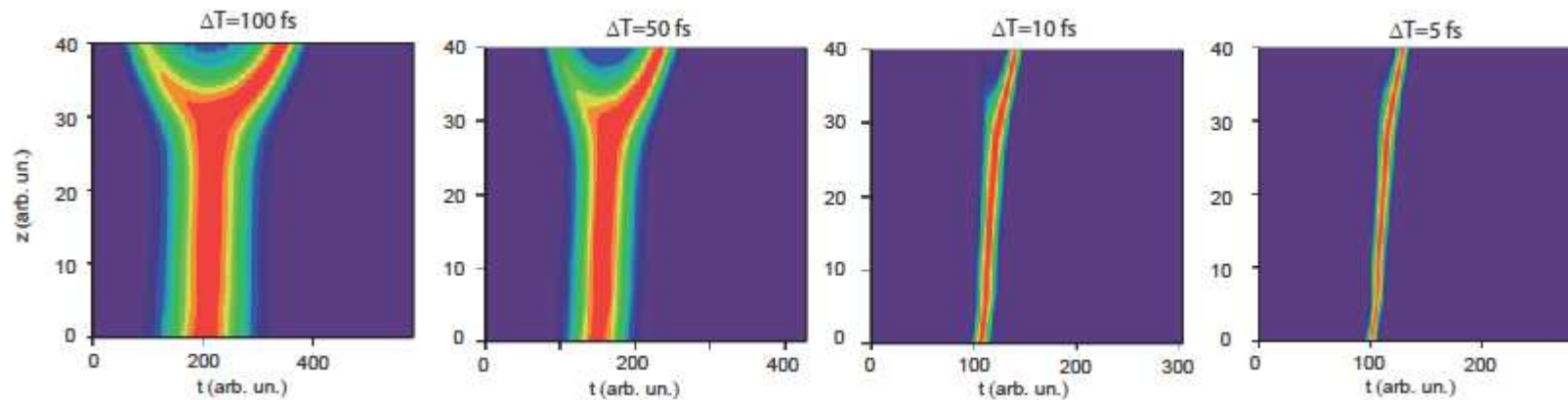
Vs pulse duration



1 GeV e- beam. Seed: 50 kW @ 14 nm. 700xU26.

Pulse splitting dependencies

Vs pulse duration



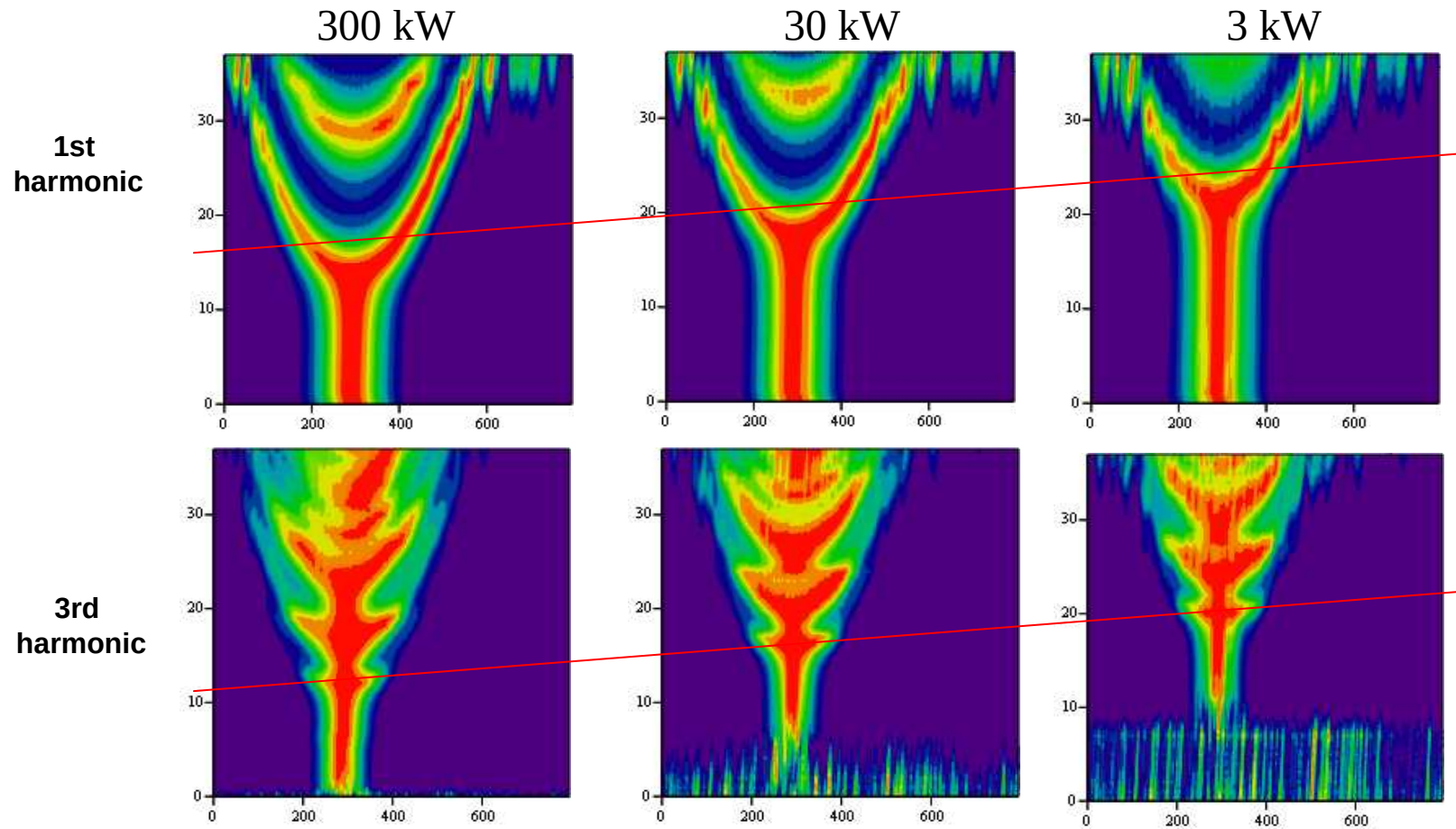
1 GeV e- beam. Seed: 50 kW @ 14 nm. 700xU26.



Pulse splitting occurrence depends in seed duration

Pulse splitting dependencies

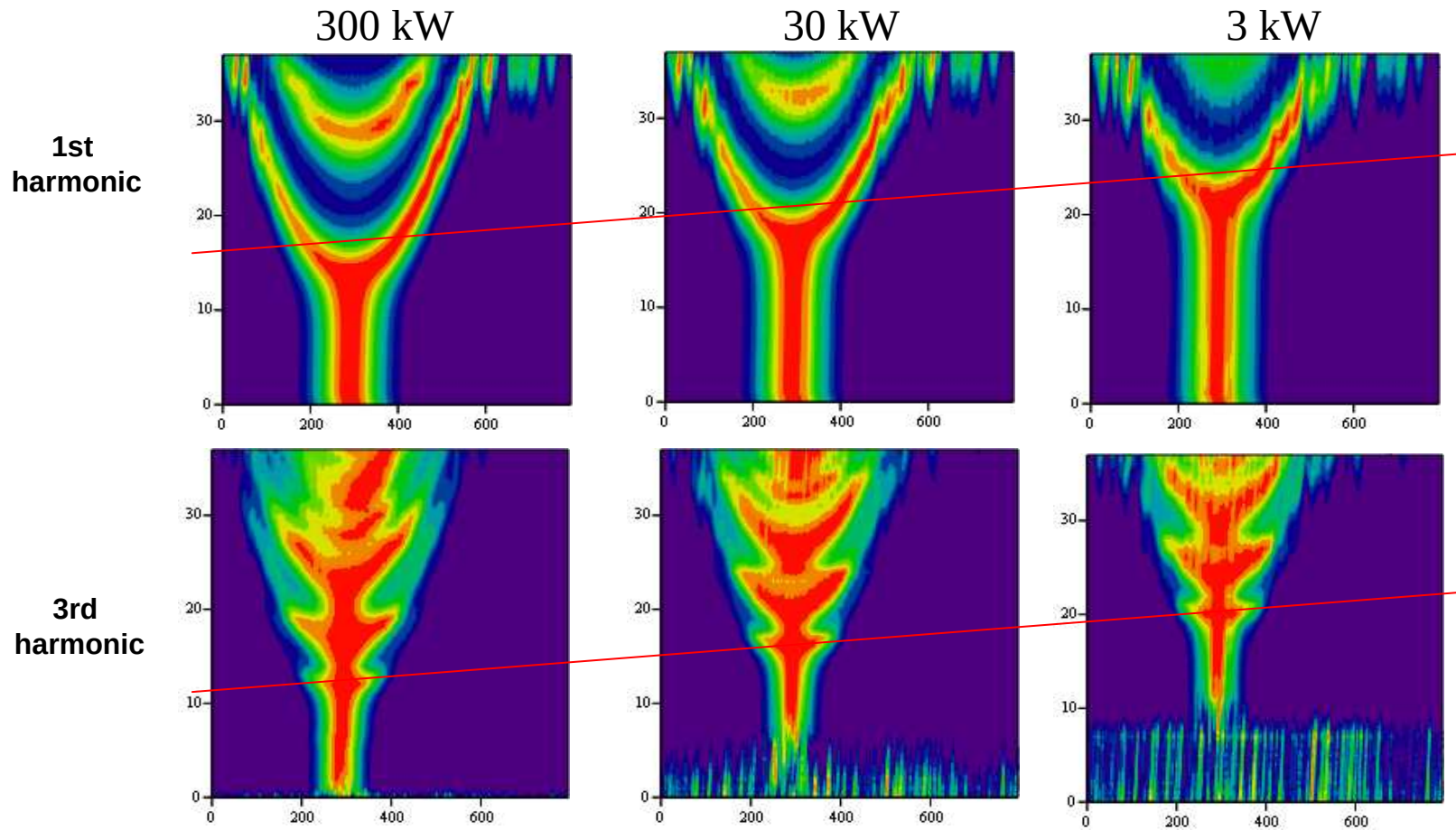
Vs power



1 GeV e- beam. Seed @ 8.9 nm. 200xU26 – R₅₆ – 500xU26

Pulse splitting dependencies

Vs power



Pulse splitting occurrence does not depend in power

Pulse splitting in seeded FELs

Intermediate conclusion

- Simulation of VUV and XUV seeded FELS (ARC-EN-CIEL):
→ Observation of a pulse splitting
 - Study of the double pulse behavior:
 - Related to local overbunching / follows gain medium
 - Related to seed pulse duration
 - Not due to excessive seeding power
- Requires a better understanding



**Start over from
FEL fundamental equations**

Modeling of a seeded FEL

The 1-D Colson-Bonifacio model

W.B. Colson, Phys. Lett. A 59 (1976)
B. Bonifacio et al., PRA 40 (1989)

Modeling of a seeded FEL

The 1-D Colson-Bonifacio model

- **Modeling of the e- beam:**
 - N_e particles in the phase space (Φ, p) :
 - Φ_j : relative phase
 - p_j : relative normalized energy
 - Particles distribution:
 - Φ_j : uniform distribution within $[-\pi; \pi]$
 - p_j : normal distribution centred at 0 with standard deviation σ_y .

Modeling of a seeded FEL

The 1-D Colson-Bonifacio model

- **Modeling of the e- beam:**
 - N_e particles in the phase space (Φ, p)
 - Particles distribution
- **Modeling of the radiation:**
 - Radiation potential vector $A(z, t)$
- **Dimensionless coordinates:**
 - z : along the undulator in slippage length units
 - t : along the electron bunch in cooperation length units

Modeling of a seeded FEL

The 1-D Colson-Bonifacio model

W.B. Colson, Phys. Lett. A 59 (1976)
B. Bonifacio et al., PRA 40 (1989)

- FEL equations:**

$$\frac{\partial \phi_j}{\partial z} = p_j$$

$$\frac{\partial p_j}{\partial z} = - [A(z, t) e^{i\phi_j} + c.c.]$$

Evolution of the particles
in the phase space
(Φ_j, p_j)

$$\left(\frac{\partial}{\partial z} + \frac{\partial}{\partial t} \right) A(z, t) = \chi(t) b(z, t)$$

Propagation of the field A
in the electronic medium

Modeling of a seeded FEL

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W.B. Colson, Phys. Lett. A 59 (1976)
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Evolution of the particles
in the phase space
(Φ_j, p_j)

$$\left(\frac{\partial}{\partial z} + \frac{\partial}{\partial t} \right) A(z, t) = \chi(t) b(z, t)$$

Propagation of the field A
in the electronic medium

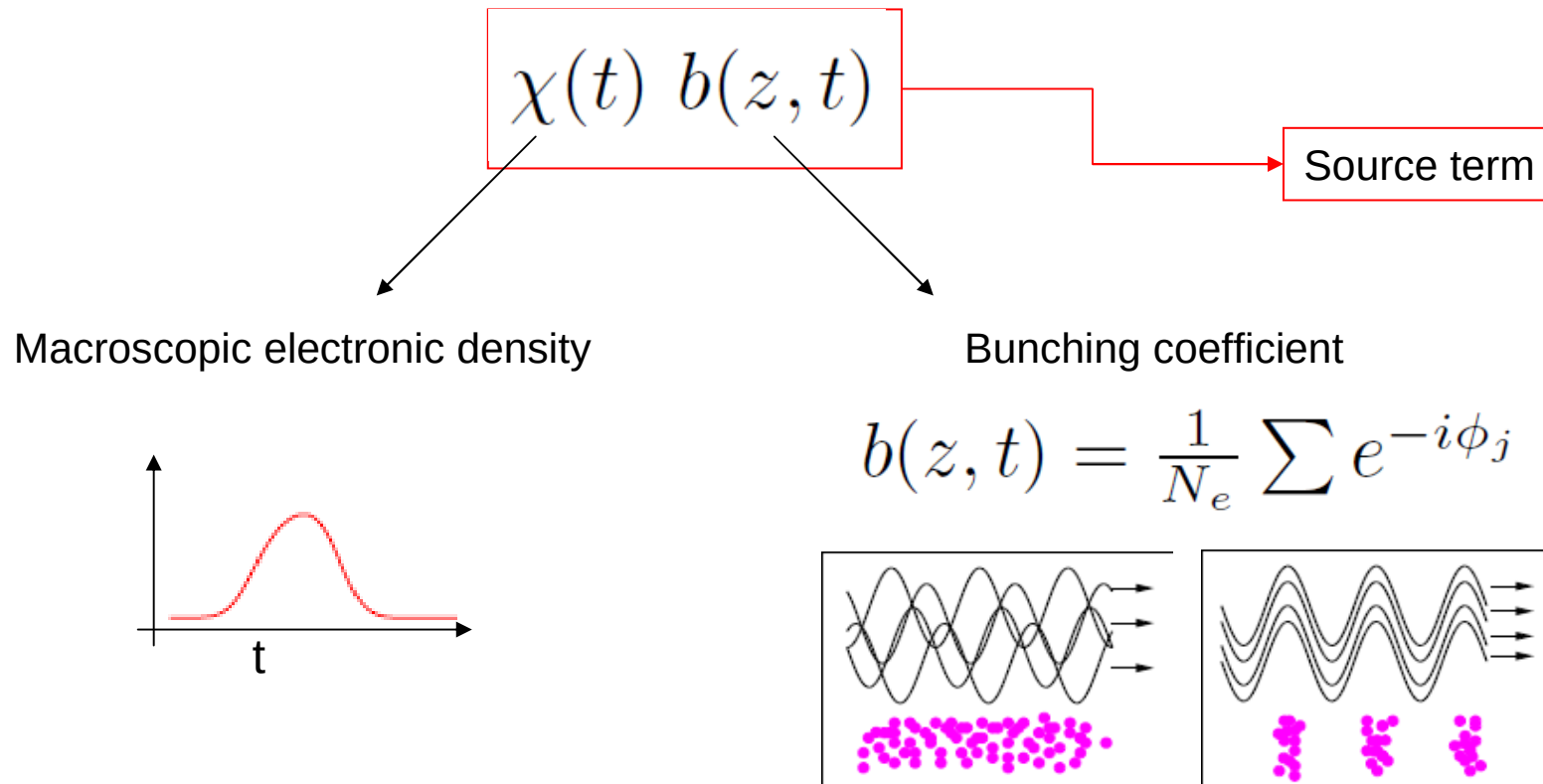
Source term

Modeling of a seeded FEL

The 1-D Colson-Bonifacio model

W.B. Colson, Phys. Lett. A 59 (1976)
B. Bonifacio et al., PRA 40 (1989)

- FEL equations:**



Modeling of a seeded FEL

Numerical tool

N. Joly

- **Integration of the FEL equations:**
 - Description of the particles dynamics
 - Description of the evolution of the radiation pulse
 - Analysis of the **spatio-temporal regimes** of the seeded FEL with dimensionless parameters

Analysis of the spatiotemporal regimes

A few dimensionless parameters

B. Bonifacio et al., PRA 40 (1989)

- « Slippage parameter »:

$$S_e = \frac{\text{slippage length}}{e- \text{ bunch length}}$$

- « Superradiant parameter »:

$$K = \frac{\text{cooperation length}}{e- \text{ bunch length}}$$

Analysis of the spatiotemporal regimes

New tools for seeded FELs

- **Radiation potential vector:**

$$A(z = 0, t) = A_0 \Rightarrow A(z = 0, t) = A_0 e^{-\frac{1}{2}(\frac{t}{\sigma_t})^2}$$

The diagram illustrates the decomposition of the radiation potential vector $A(z = 0, t)$ into two components. The initial expression $A(z = 0, t) = A_0$ is shown with A_0 enclosed in a red box. A red arrow points from this box to a red box labeled "shot noise". The subsequent expression $A(z = 0, t) = A_0 e^{-\frac{1}{2}(\frac{t}{\sigma_t})^2}$ is shown with the entire right-hand side enclosed in a red box. A red arrow points from this box to a red box labeled "seed".

- Additional dimensionless « **seed** » parameter:

$$S_{seed} = \frac{\text{slippage length}}{\text{seed pulse length}}$$

Analysis of the spatiotemporal regimes

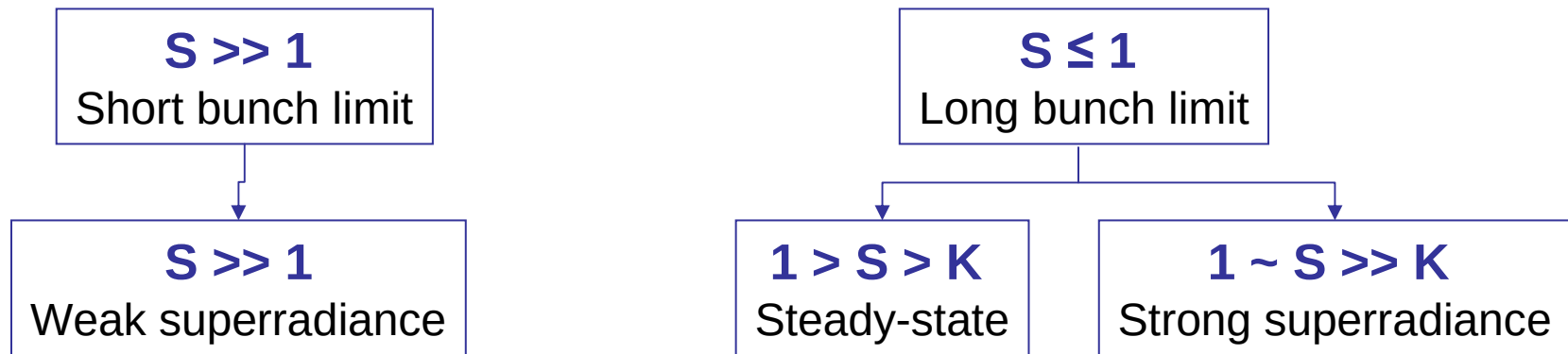
« Initial » regimes

B. Bonifacio et al., PRA 40 (1989)

Analysis of the spatiotemporal regimes

« Initial » regimes

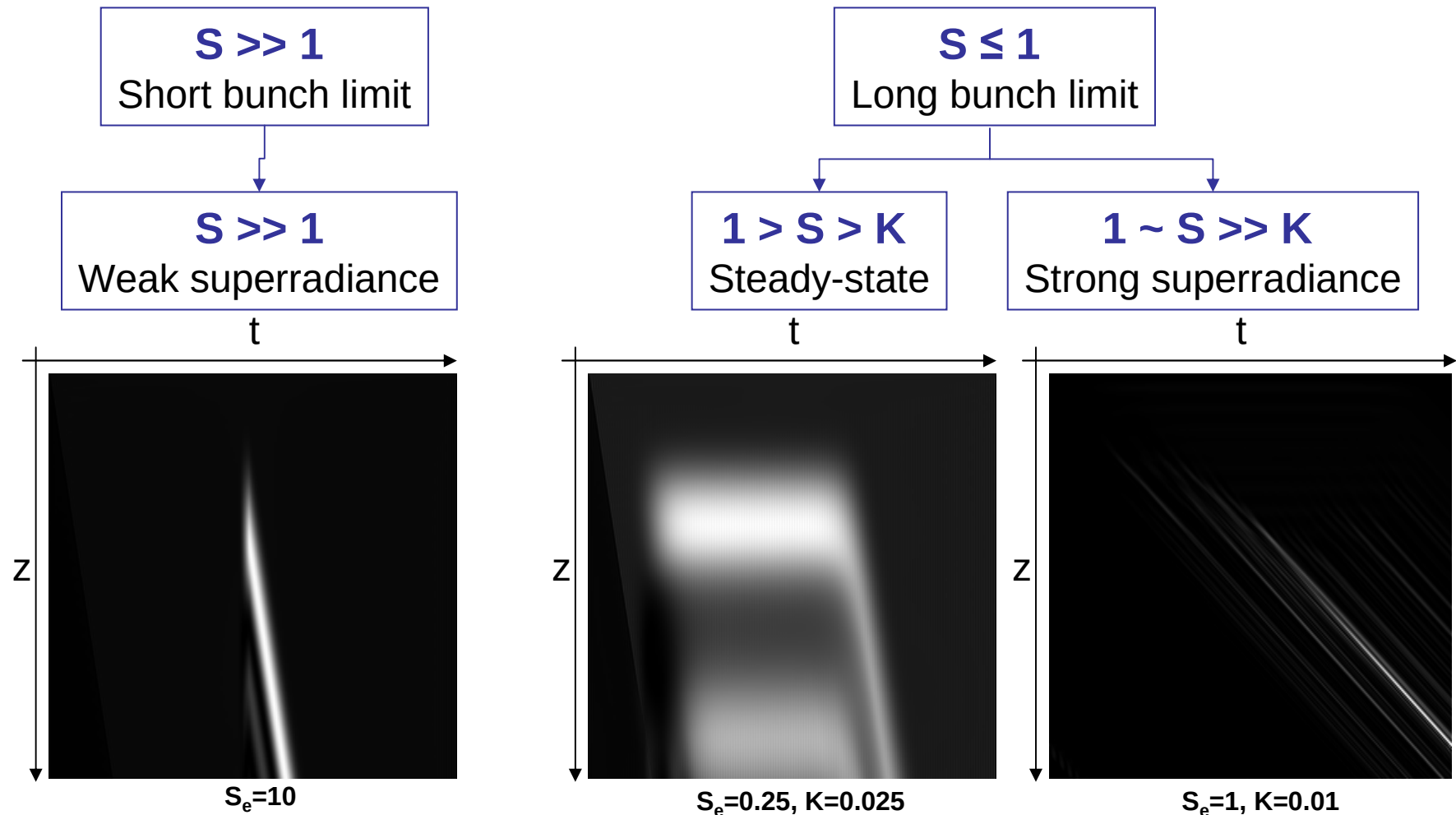
B. Bonifacio et al., PRA 40 (1989)



Analysis of the spatiotemporal regimes

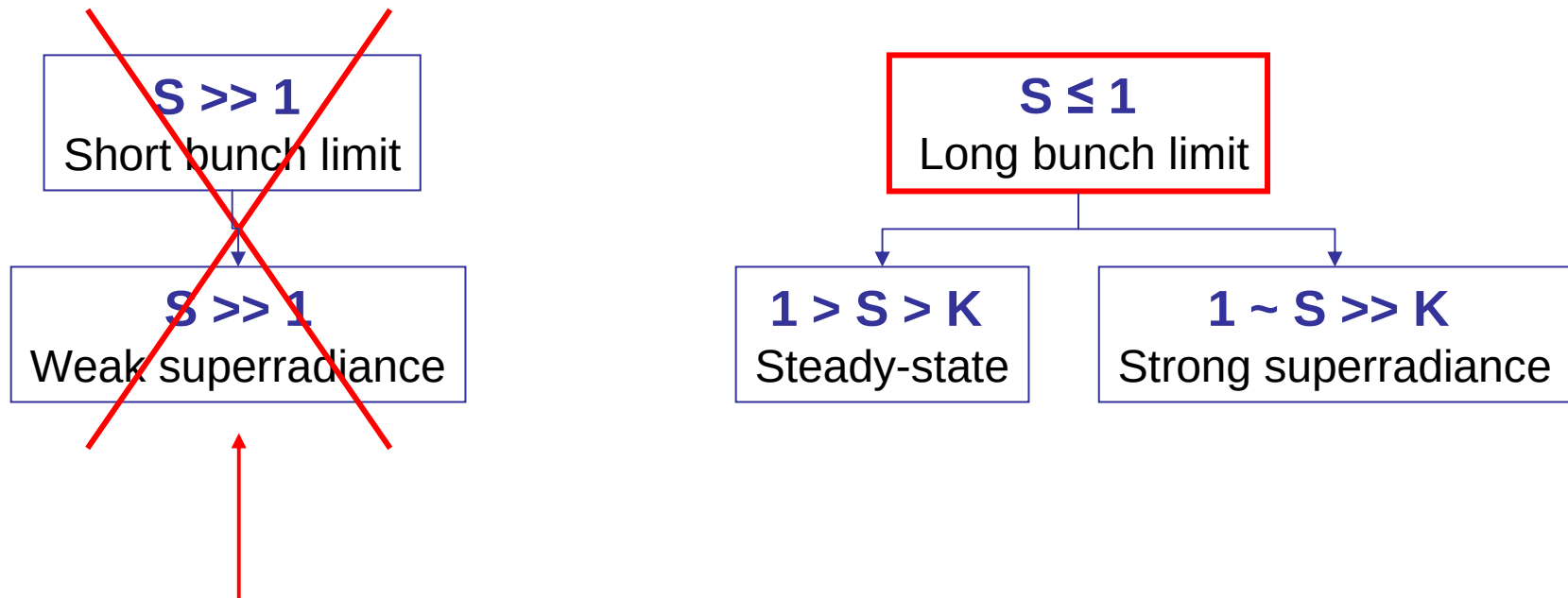
« Initial » regimes

B. Bonifacio et al., PRA 40 (1989)



Analysis of the spatiotemporal regimes

« Initial » regimes



Today designs for $S \leq 1$

LCLS: $S \sim 0.01$

EXFEL: $S \sim 0.02$

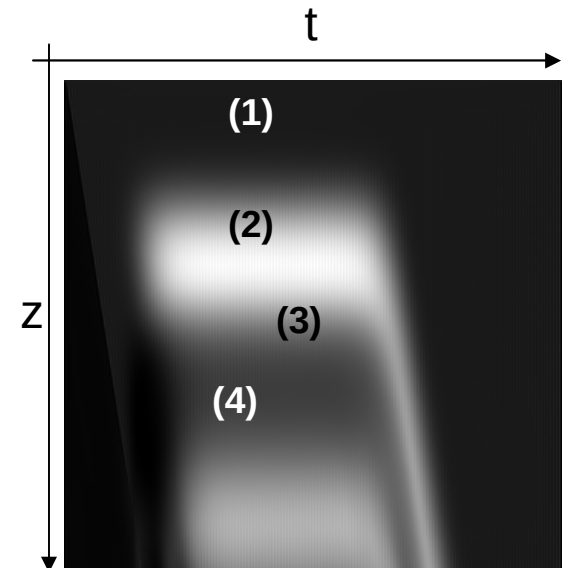
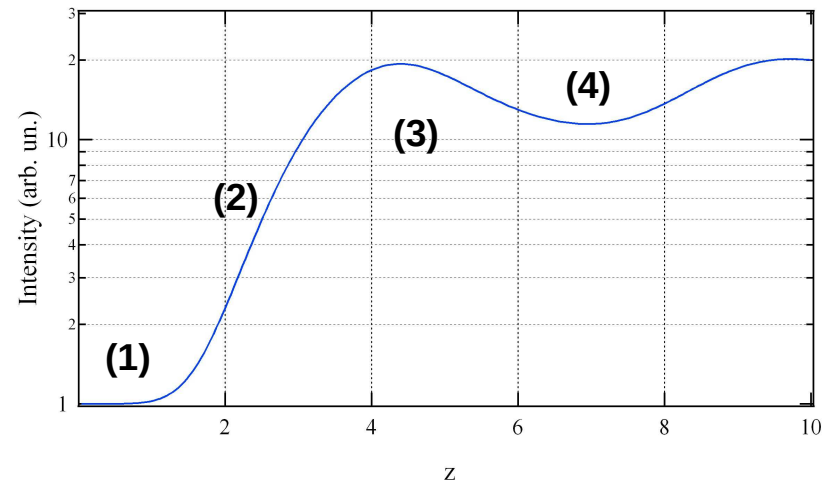
FERMI: $S \sim 0.1$

SPARC: $S \sim 0.07$

Analysis of the spatiotemporal regimes

The steady-state FEL

- Condition for occurrence:
 - $S \leq 1$ and $1 > S > K$
(small slippage)
- Properties:
 - (1) Lethargy
 - (2) Exponential growth
 - (3) Saturation
 - (4) Power oscillations

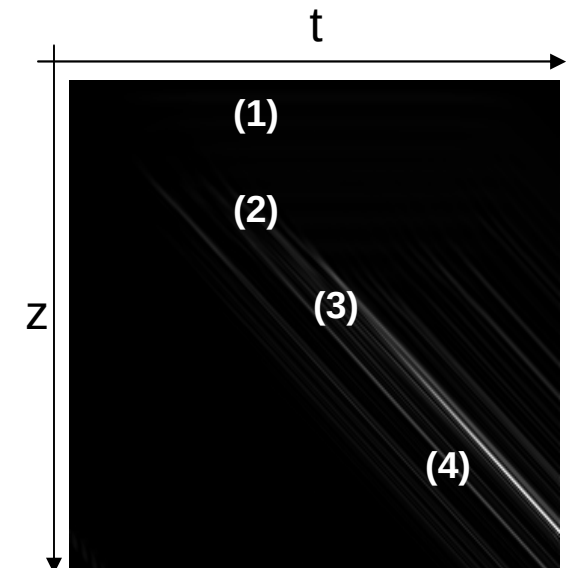
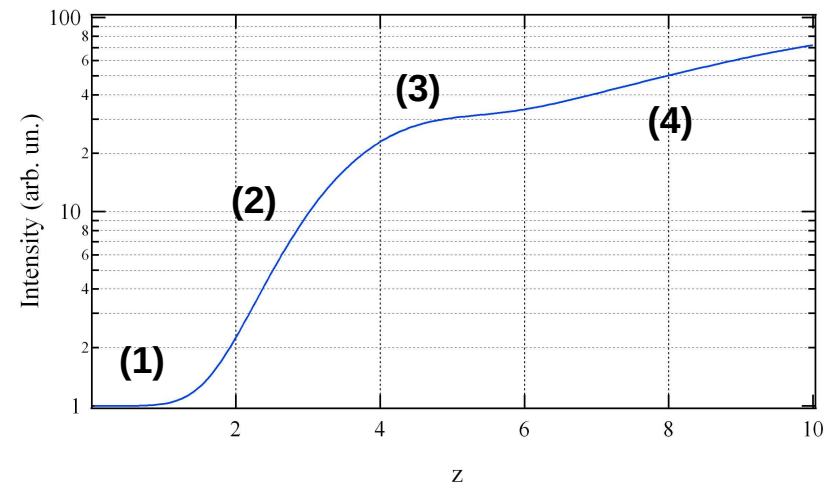


Analysis of the spatiotemporal regimes

The strong superradiance

R. Bonifacio et al., PRA 40 (1989).

- Condition for occurrence:
 - $1 \sim S \gg K$
(strong slippage)
- Properties:
 - (1) Lethargy
 - (2) Exponential growth
 - (3) No saturation
 - (4) ...



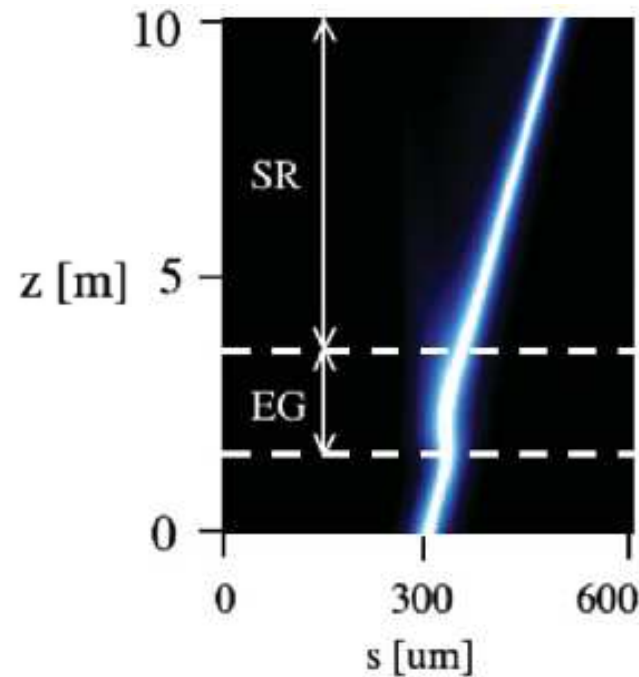
$$P \propto z^2$$
$$\sigma_z \propto z^{-1/2}$$

Analysis of the spatiotemporal regimes

The strong superradiance

T. Watanabe, PRL 98 (2007).

- Demonstration @ NSLS in 2007:

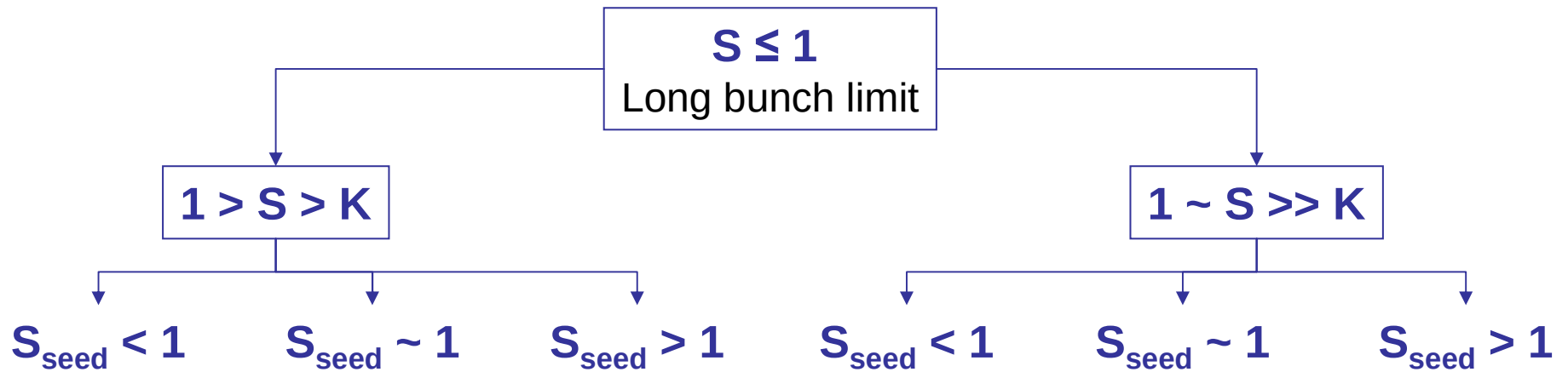


Analysis of the spatiotemporal regimes

The new regime

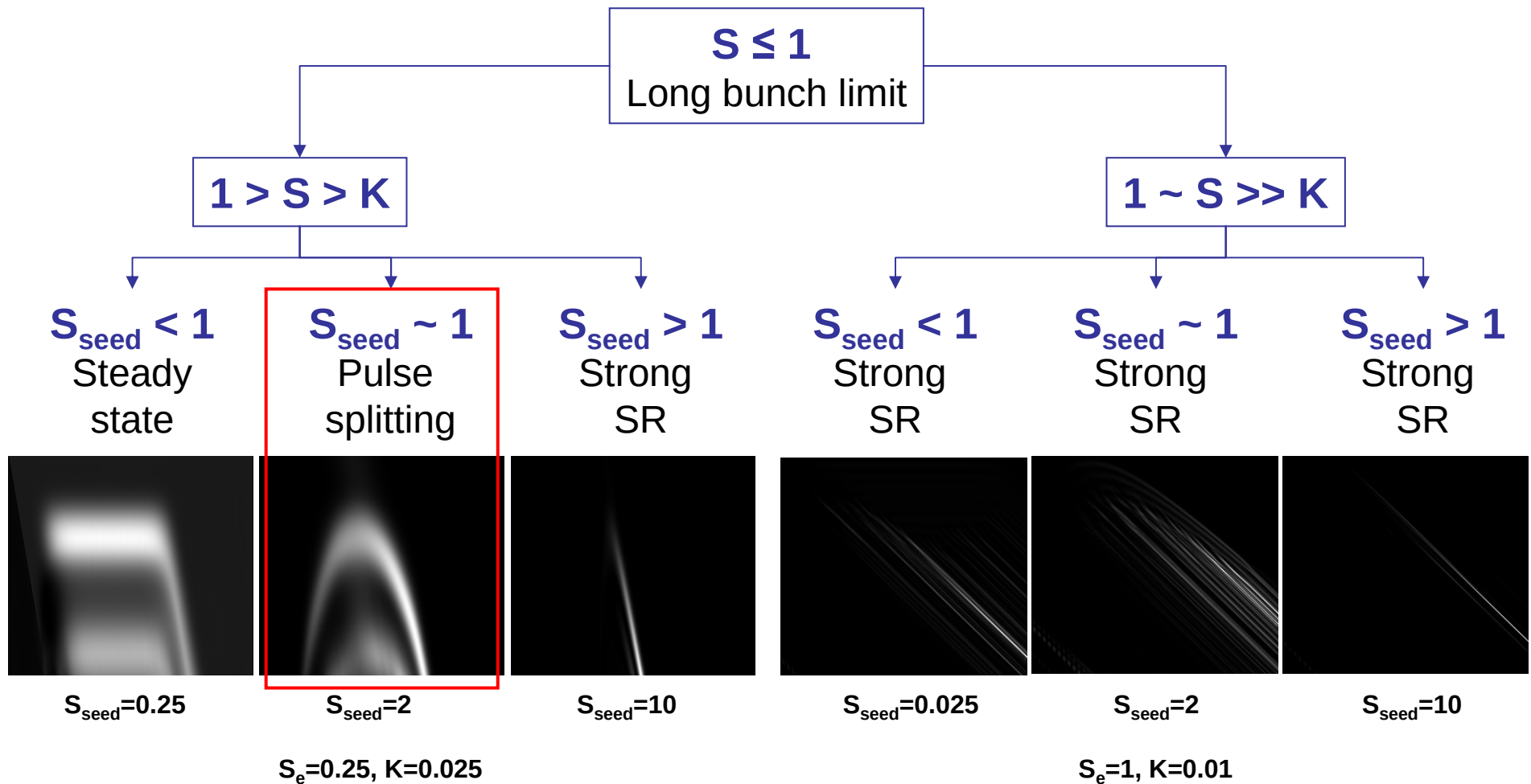
Analysis of the spatiotemporal regimes

The new regime



Analysis of the spatiotemporal regimes

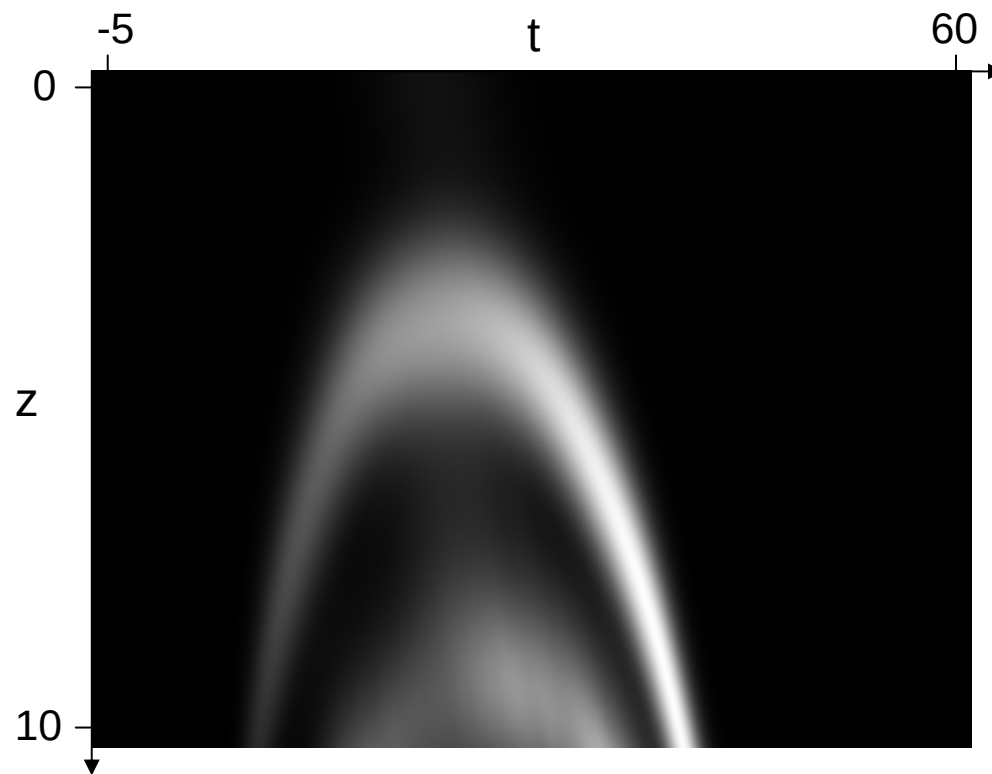
The new regime



Analysis of the spatiotemporal regimes

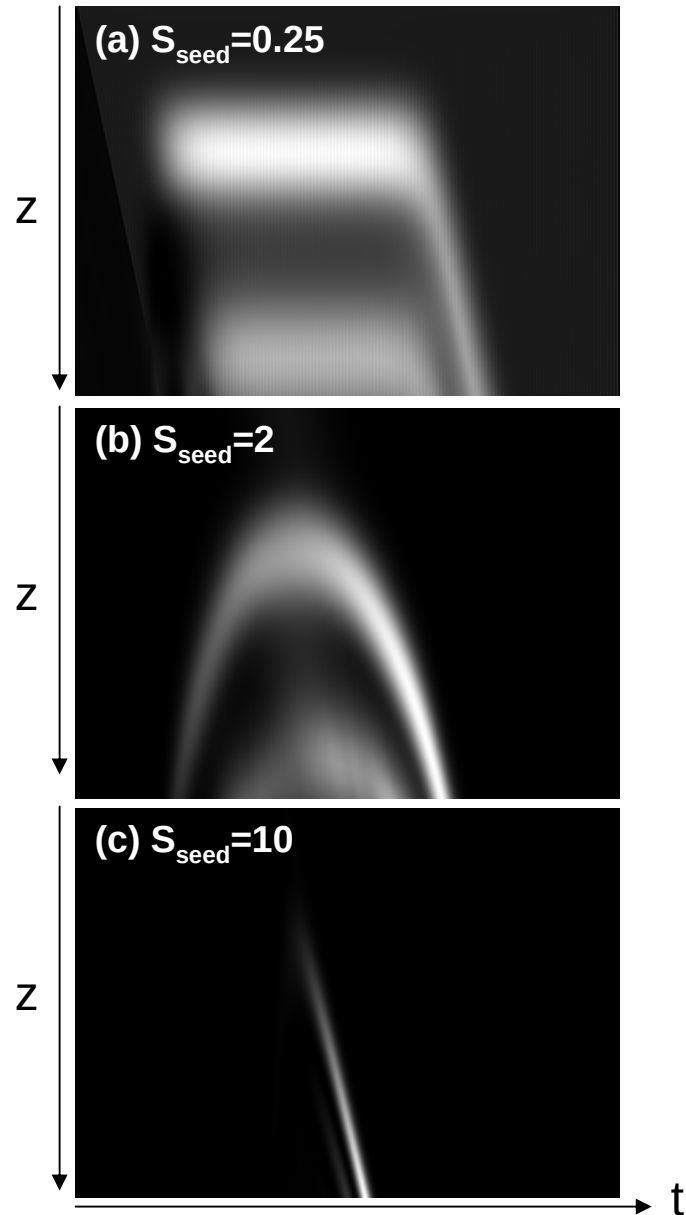
The pulse splitting

M. Labat et al., PRL 103 (2009)

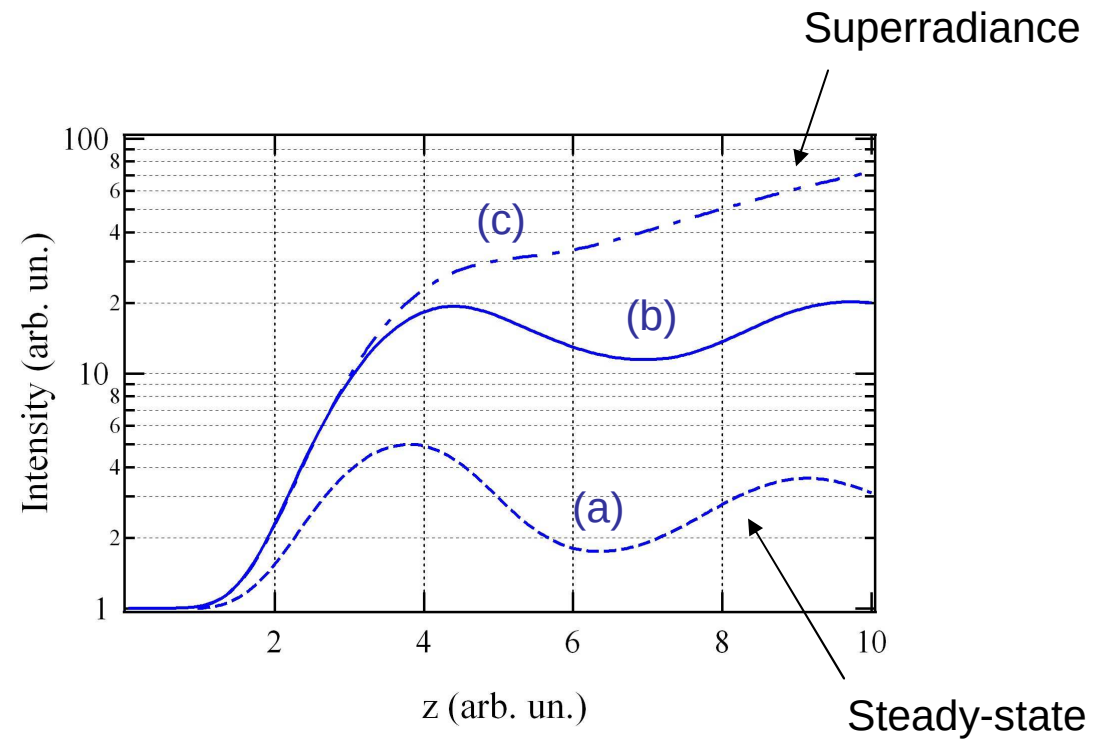


$S_e=0.25, K=0.025, S_{\text{seed}}=2$

The pulse splitting regime



Output power

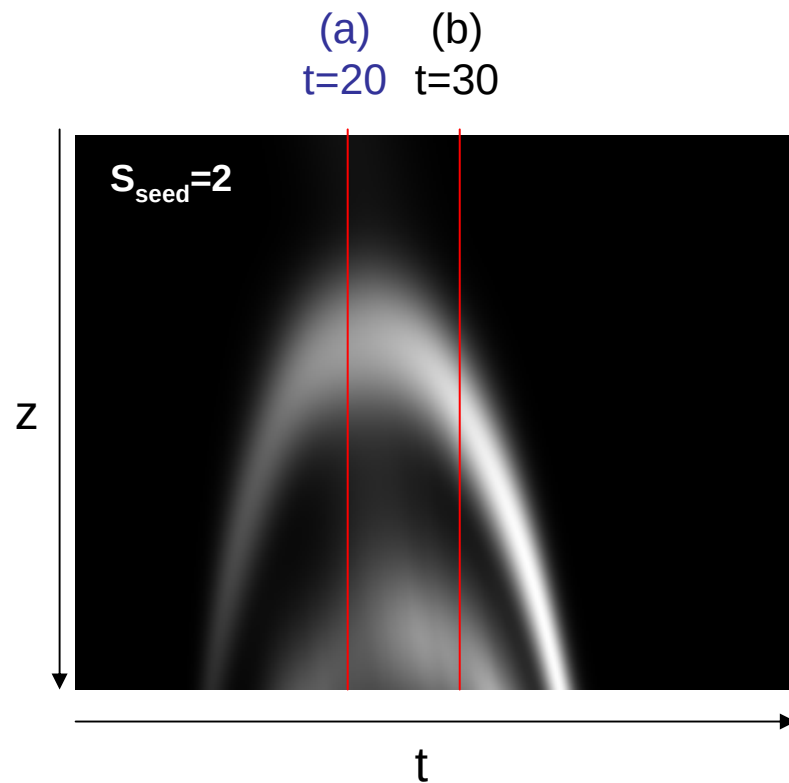


The pulse splitting regime

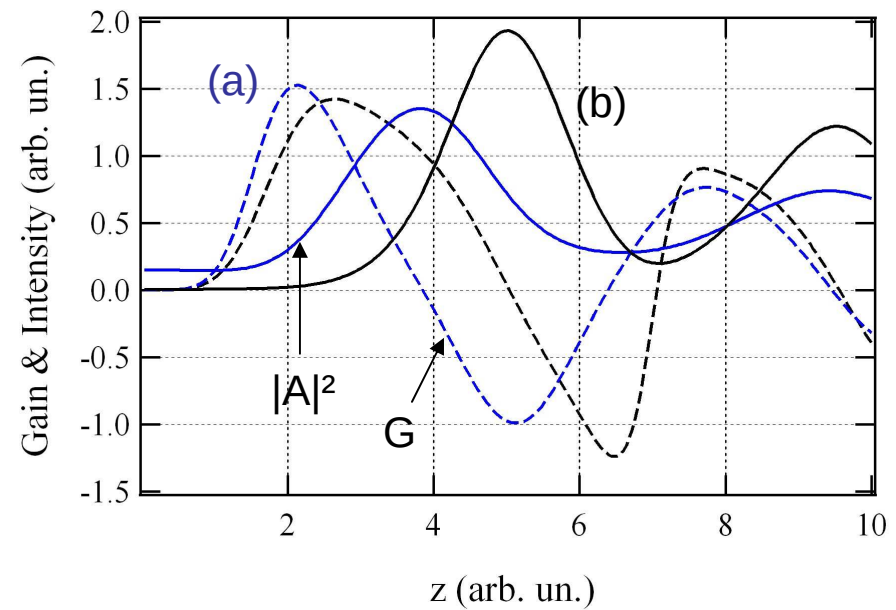
Gain

The pulse splitting regime

Gain

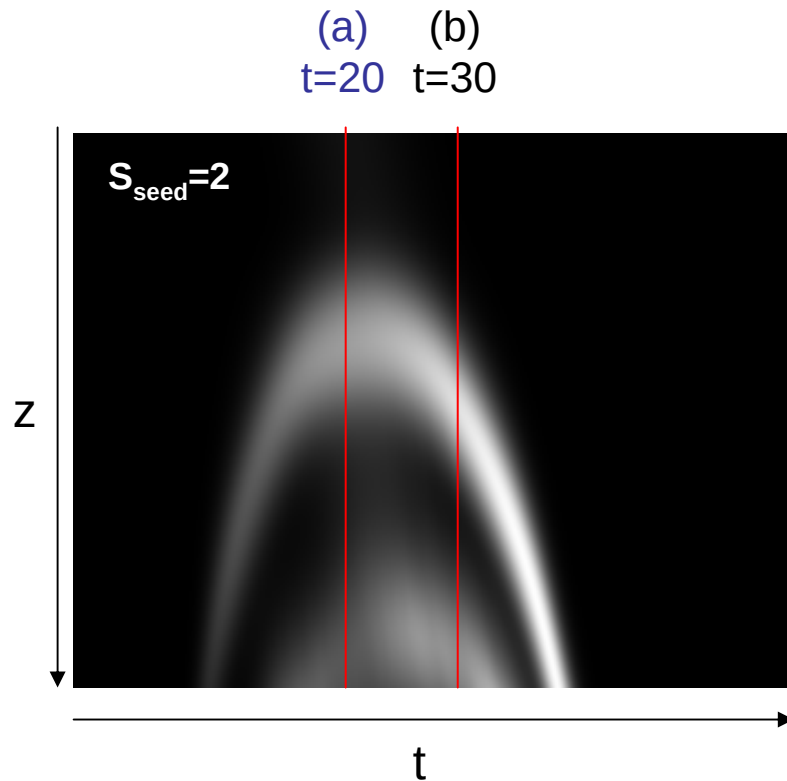


$$G(z, t) = \text{Re}\left[\frac{\chi(t)b(z, t)}{A(z, t)}\right]$$

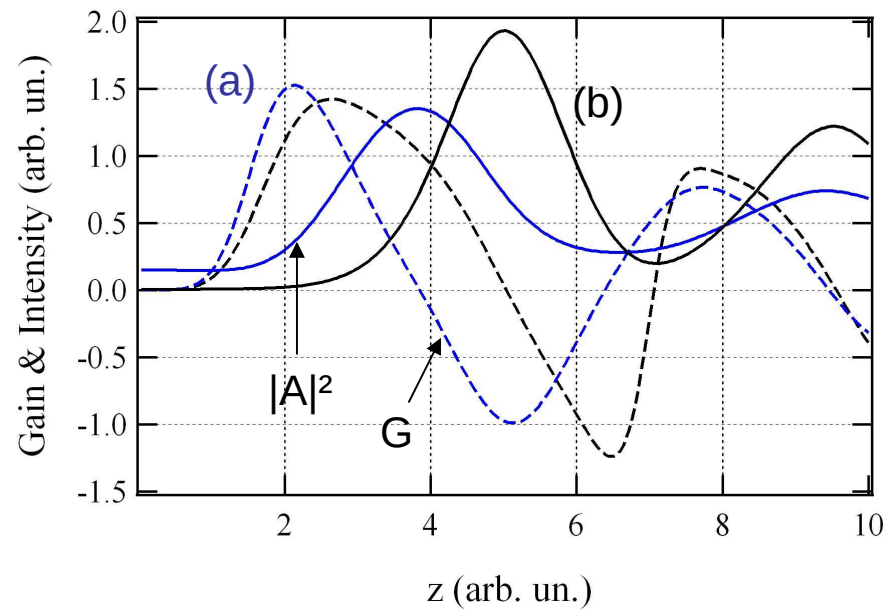


The pulse splitting regime

Gain



$$G(z, t) = \text{Re}\left[\frac{\chi(t)b(z, t)}{A(z, t)}\right]$$



Local gain $\rightarrow$ **Local saturation** $\rightarrow$ Progressive saturation along the pulse

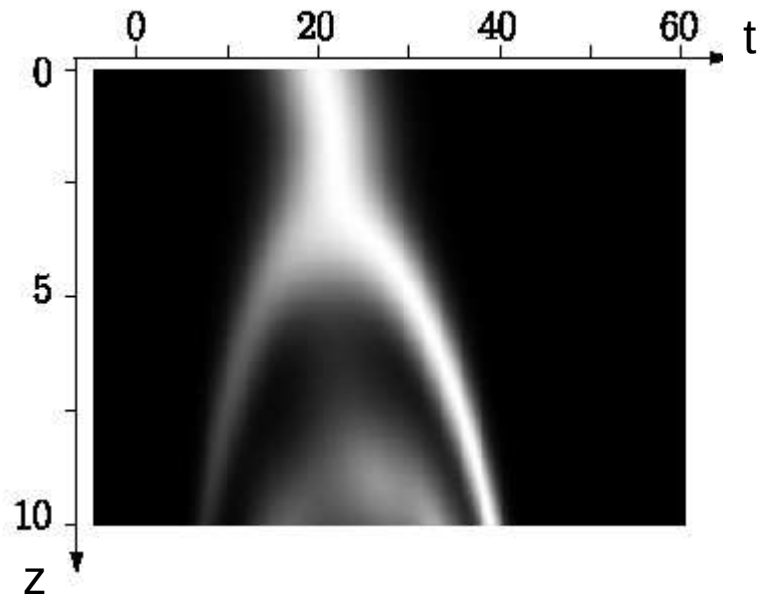
$\rightarrow$ Pulse splitting

The pulse splitting regime

Approximative splitting equation

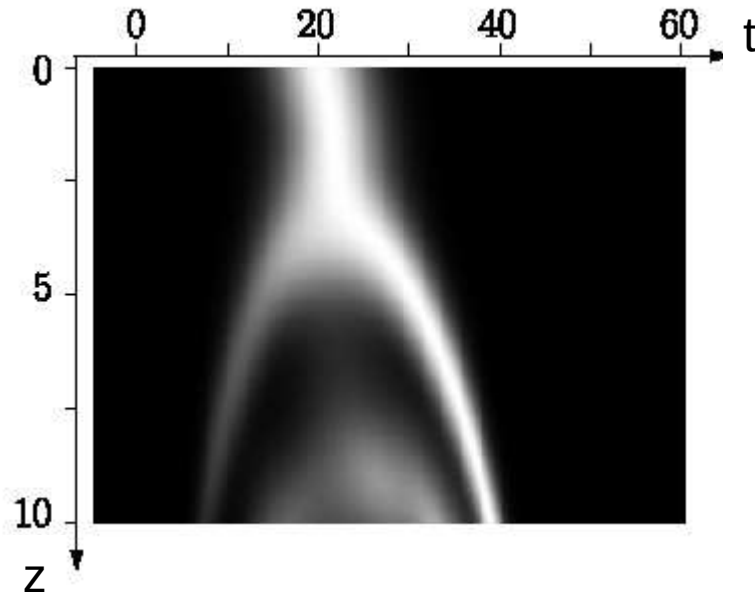
The pulse splitting regime

Approximative splitting equation



The pulse splitting regime

Approximative splitting equation



- For $S_e < 1$ and $S_{seed} \sim 1$:

$$\left(\frac{\partial}{\partial z} + \frac{\partial}{\partial t}\right) A(z, t) \approx \frac{\partial}{\partial z} A(z, t)$$

- Solution of FEL equations:

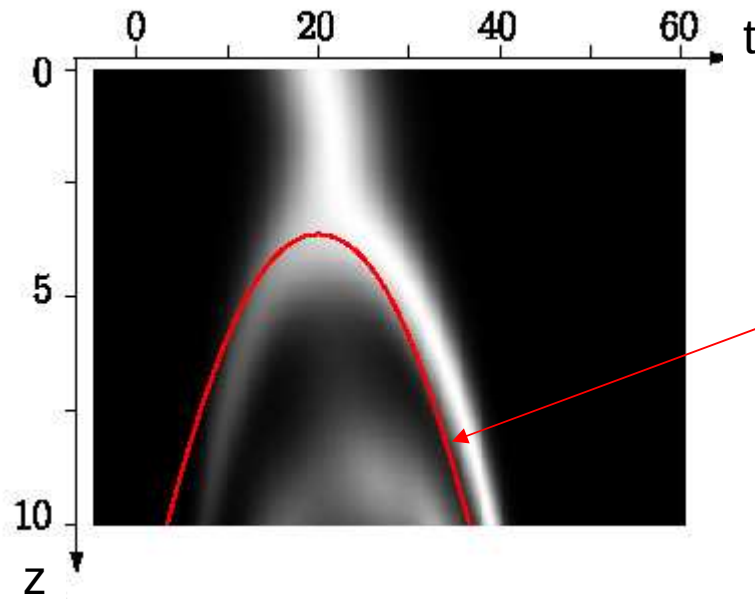
$$|A|^2(z, t) \approx \frac{1}{9} \exp[\sqrt{3}z] |A|^2(z=0, t)$$

- At saturation:

$$|A|^2(z_{sat}, t) \cong O(1) \cong 1.4$$

The pulse splitting regime

Approximative splitting equation



$$z_{sat}(t) = \frac{1}{\sqrt{3}} \ln \left[\frac{9 \times 1.4}{|A|^2(z=0, t)} \right]$$

CONFIRMS

Local gain → **Local saturation** → Progressive saturation along the pulse



Pulse splitting

The pulse splitting regime

Next step

Can we observe a pulse splitting
on an existing seeded FEL ??

Pulse splitting in seeded FELs

Conclusion

- Simulation of VUV and XUV seeded FELS (ARC-EN-CIEL):
 - Observation of a pulse splitting
- Studies under PERSEO:
 - Rough understanding
- Integration of FEL equations:
 - Observation of a pulse splitted regime
- Analysis of dimensionless parameters:
 - Definition of the regime conditions of occurrence
 - Better understanding of the phenomenon



Let's make experiment....