

Figure 2: Quasimonoenergetic electron beams produced in 4 mm-long gas jet at various laser powers (27-50 TW) and plasma densities $6.5 \times 10^{18} - 7.5 \times 10^{18} \text{ cm}^{-3}$.

Fig. 2(b). Again, by further increasing the laser power to 50 TW while the plasma density was $\sim 6.6 \times 10^{18} \text{ cm}^{-3}$, the central energy of the quasimonoenergetic electron beam moved to 350 MeV and the maximum energy increased up to 880 MeV (GeV-class), as shown in Fig. 2 (c). The total charge of the electron beam in this case was measured to be about 500 pC and its divergence angle was about 5 mrad (FWHM). The laser-plasma channel lengths were observed to be longer at higher laser powers due to the self focusing, channels of lengths ($\pm 100 \mu\text{m}$) of 2.5 mm, 3.5 mm and 4 mm were produced at the laser powers of 27 TW, 37 TW and 50 TW, respectively. Those channel lengths were approximately equal to the dephasing lengths for the plasma densities mentioned above, and this was important in obtaining quasimonoenergetic spectrum. Experimental results of Fig. 2 clearly shows the doubling of a quasimonoenergetic electron beam energy by doubling the laser power, i.e., the electron beam energy scales as $\sim a_0^2$ (for given plasma parameters). This is an important experimental result on this scaling for the multi-hundreds MeV quasimonoenergetic electron beam generation by a LWFA.

Stability of quasimonoenergetic electron beams

3D-PIC simulations have predicted that plasma bubble

regime has a highly stable particle acceleration structure. Additionally, in recent experiments [4] it was noted that long laser focusing geometries (f number 21.4, in our case) are necessary for the stability of the laser pulse propagation through the plasma and the generation of stable wakefields and electron beams. Motivated by those predictions, we have performed a stability test for an electron beams produced by our laser-plasma accelerator with the following parameters: laser intensity of 37 TW, 4-mm long gas jet and plasma density of $7 \times 10^{18} \text{ cm}^{-3}$. Figure 3 shows raw images of the electron beam energy spectra obtained in 10 consecutive shots. The averaged peak energy, and the energy fluctuation were $\sim 237 \text{ MeV}$ and $\sim 5\%$, respectively. The average charge per electron bunch and the charge fluctuation were $\sim 205 \text{ pC}$ and $\sim 32\%$, respectively. The measured fluctuation in the laser energy during those 10 consecutive shots was $\sim 4.5\%$ which is very close to the electron beam energy fluctuations. This leads to the conclusion that ultra stable electron beam energies can be realized by stabilizing the shot to shot laser pulse energy.

Plasma channels observed in these shots were very similar to each other and reproducible. Each channel had a single filament with a length of $\sim 3.5 \text{ mm}$. Based on single laser pulse LWFA experiments; our stability test shows quite stable quasimonoenergetic electron generation in the energy range of 240 MeV.

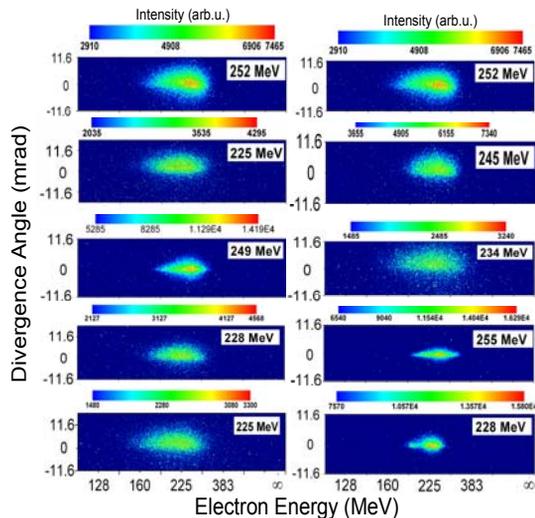


Figure 3: Quasimonoenergetic electron beams of ten consecutive shots

POTENTIAL APPLICATION TO FREE ELECTRON LASERS

Recently Schlenvoigt et al. [5] succeeded to generate synchrotron radiation by injecting an electron beam from a laser-plasma accelerator into 1-m long undulator. The electron beam parameters were: energy- 64 MeV, energy spread 5.3 %, charge- 28 pC and normalized emittance of 1.3π mm.mrad. The synchrotron radiation wavelength was 790 nm in the visible range. In this experiment, the emission wavelength scales with the beam energy just as the theory predicts. This suggests the generation of much shorter wavelengths (EUV-x rays) by this technique is straight forward.

By extending the undulator length to 3 m, and feeding it with our electron beam of energy 350 MeV, we expect to generate synchrotron radiation of wavelength about 23 nm which can be shortened to few nm by using more compact undulator with a shorter period. So, it is clear and has been proven that the existing laser-plasma electron beams can generate synchrotron radiation, we expect this will grow in the near future as there are a lot of groups around the world generating multi-hundred MeV electron beams from laser-plasma interactions. The next important step is to use the laser-plasma electron beams for generating free electron laser. This seems possible in the future because the existing laser-plasma electron beam parameters are only about 1 order of magnitude far from the required parameters for FEL lasing. For example, the beam energy spread required for lasing is about 0.1 % while it is now about 1 % for laser-

plasma electron beams. As for the beam energy and emittance, already the laser-plasma electron beams are good enough: energies up to the GeV and emittances of the order of 1π mm. mrad are already available.

So in conclusion, we believe that the only two parameters which require to be enhanced are the beam charge and energy spread. Currently, we are entering the era of using multi-hundred TW and up to PW lasers in the laser acceleration research and this in addition to using well controlled plasma conditions, will sure lead to enhance the beam parameters to meet the FEL requirements. This will surely be a great achievement because this kind of FEL will be compact and cheap and will be applied in unlimited number of applications in science, medicine and technology.

REFERENCES

- [1] T. Tajima and J. Dawson, Phys. Rev. Lett. 43 (1979) 267.
- [2] Nasr A. M. Hafz et al., Nat. Photonics 2 (2008) 571.
- [3] T. M. Jeong et al., Jap. J. Appl. Phys. 46 (2007) 7724.
- [4] N. Hafz et al., Appl. Phys. Lett. 90(2007) 151501.
- [5] H. -P. Schlenvoigt et al., Nat. Phys. 4 (2008) 130.