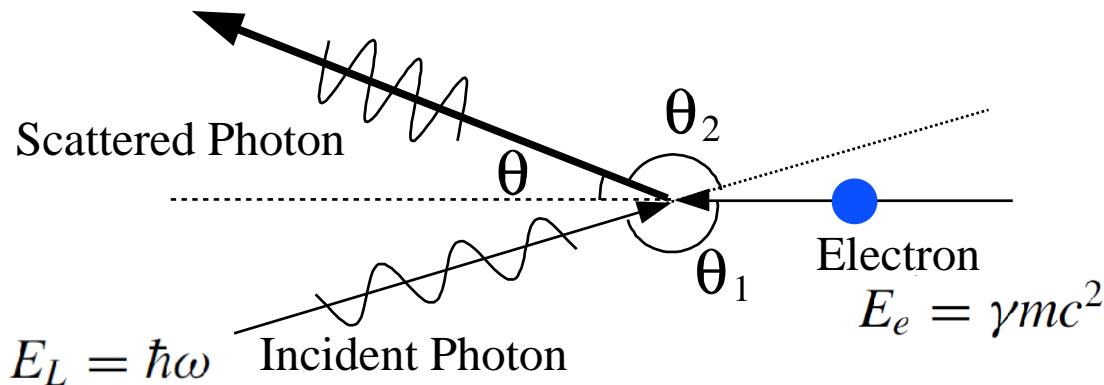


# Beam Physics Issues in ERL $\gamma$ -ray Sources

R. Hajima  
Quantum Beam Science Directorate,  
Japan Atomic Energy Agency

October 19, 2011

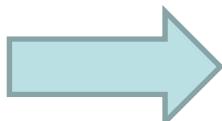
# Laser Compton scattering $\gamma$ -ray sources



$$E_\gamma = \frac{E_L(1 - \beta \cos \theta_1)}{1 - \beta \cos \theta + (E_L/E_e)(1 - \cos \theta_2)}$$

$$E_\gamma \simeq \frac{4\gamma^2 E_L}{1 + (\gamma\theta)^2 + 4\gamma E_L/(mc^2)} \quad \text{for head-on collision}$$

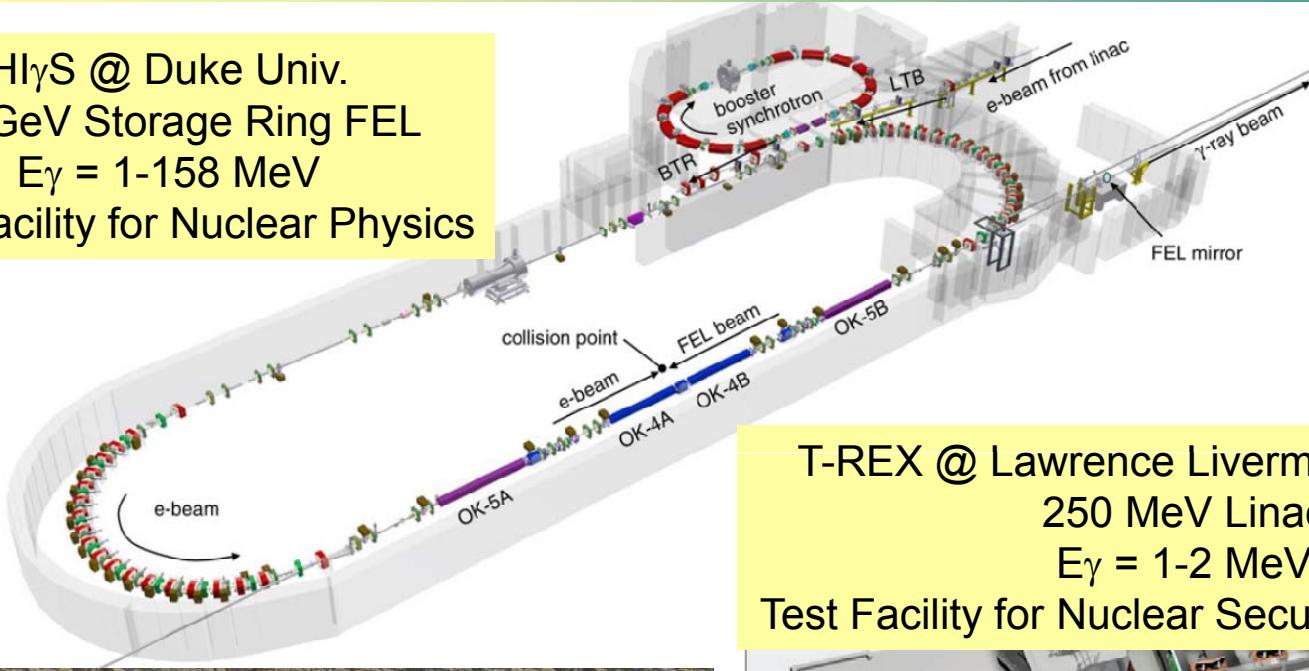
- Narrow bandwidth with a collimator
- Energy Turntable
- Sharp energy edge
- Pencil like beam



New generation of  $\gamma$ -ray source  
(c.f. Bremsstrahlung, isotopes)

# LCS $\gamma$ -ray Sources in the World

H $\gamma$ S @ Duke Univ.  
1.2 GeV Storage Ring FEL  
 $E_\gamma = 1\text{-}158 \text{ MeV}$   
User Facility for Nuclear Physics



T-REX @ Lawrence Livermore Natl. Lab.  
250 MeV Linac  
 $E_\gamma = 1\text{-}2 \text{ MeV}$   
Test Facility for Nuclear Security Applications

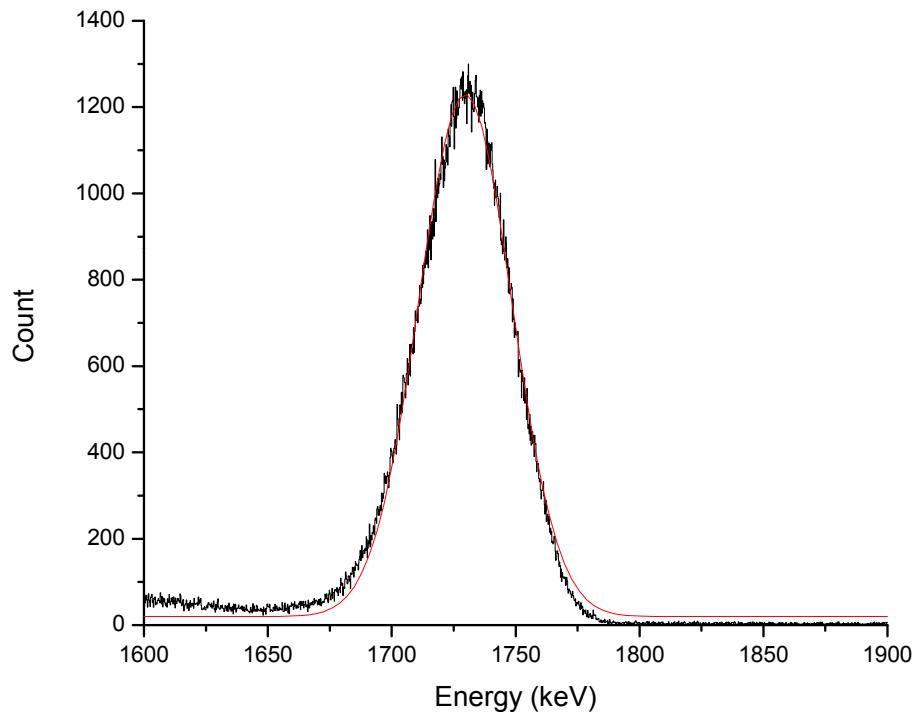


Total cost including facility modifications for 250 MeV system,  
R&D, controls and additional test stand ~ \$30M

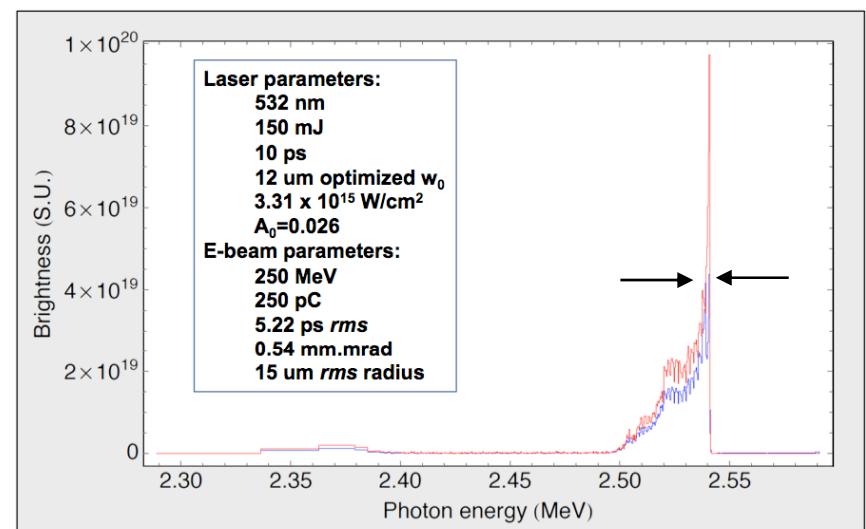


# Example of LCS $\gamma$ -ray spectrum

Storage Ring-based LCS



Linac-based LCS

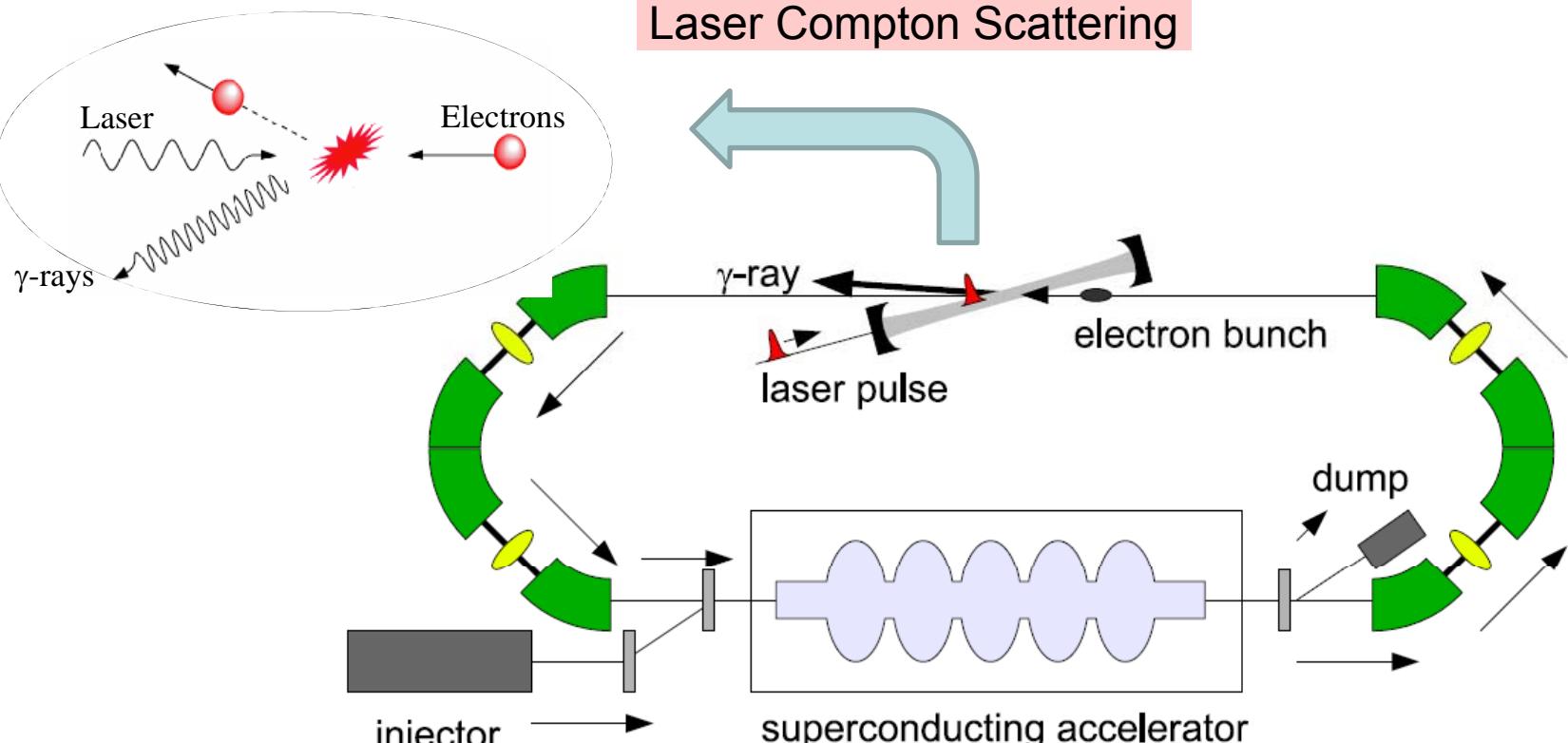


ELI-NP Whitebook (2010)

H $\gamma$ S @ Duke Univ.  
measured spectrum  
 $\Delta E/E \sim 5\%$  (FWHM)

ELI-NP  
calculated spectrum  
 $\Delta E/E < 1\%$

# Energy-Recovery Linac as a $\gamma$ -ray Source

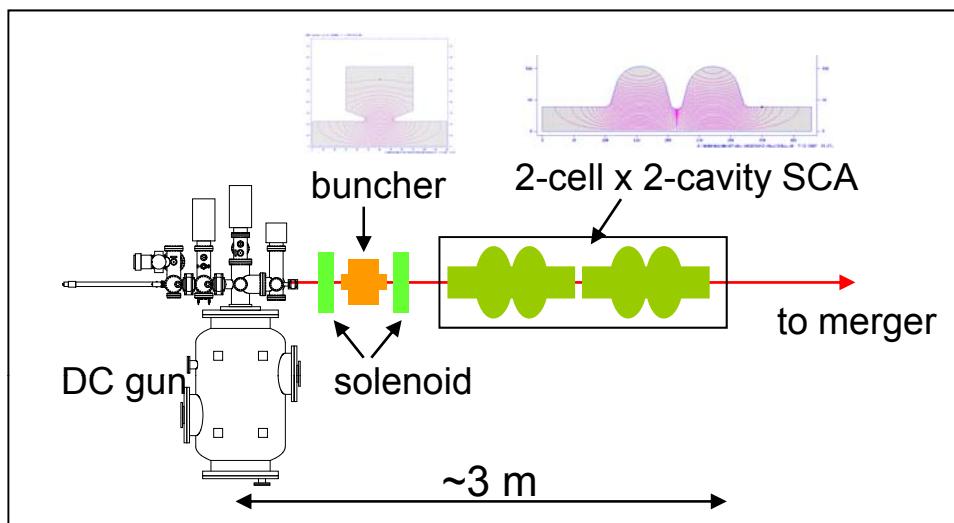


R. Hajima et al., AccApp'07 (2007).

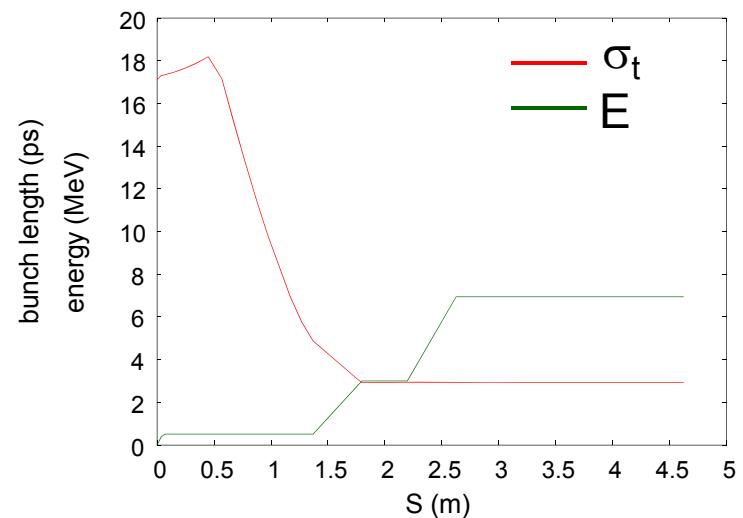
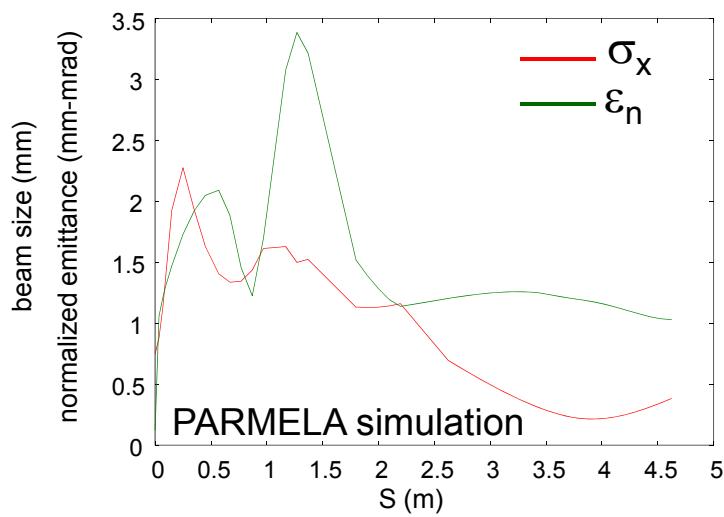
- ✓ energy-recovery for high-average current beams
- ✓ always fresh electron beam

$\gamma$ -ray beam with high-flux and narrow bandwidth

# Design of an Injector for an ERL $\gamma$ -ray source (7 MeV, 13 mA injector)

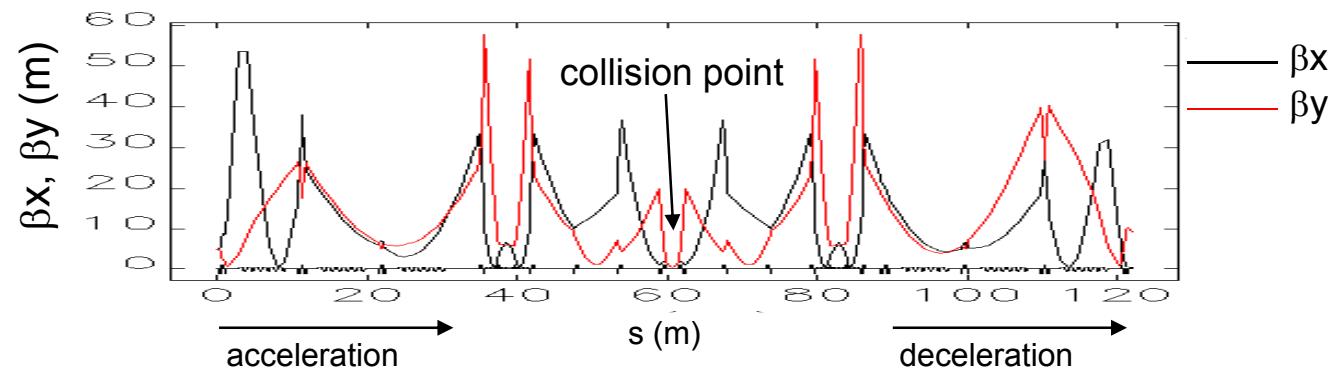
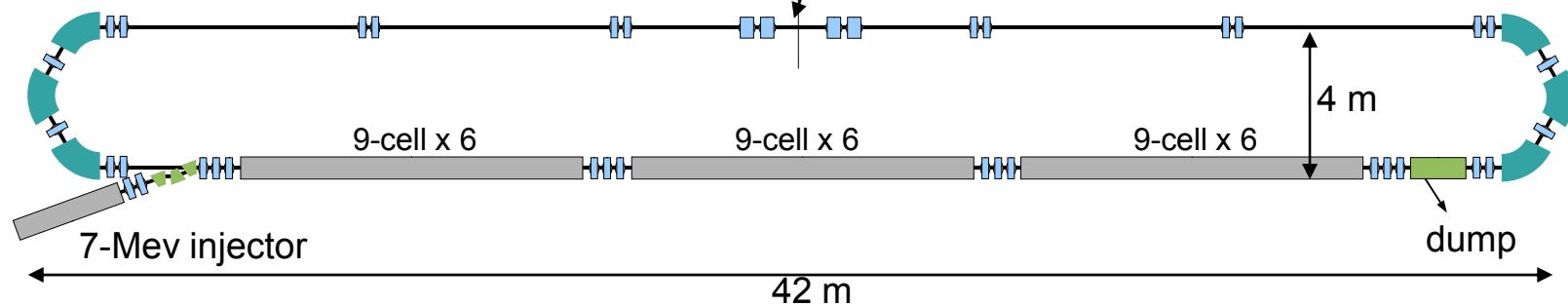


gun voltage	500 kV
buncher field	1.25 MV/m
1 <sup>st</sup> cavity field	10 MV/m
2 <sup>nd</sup> cavity field	15 MV/m
bunch charge	100 pC
norm. emittance	1.0 mm-mrad
bunch length	3 ps (rms)
energy	7 MeV
energy spread	0.1 %
repetition	130 MHz
average current	13 mA



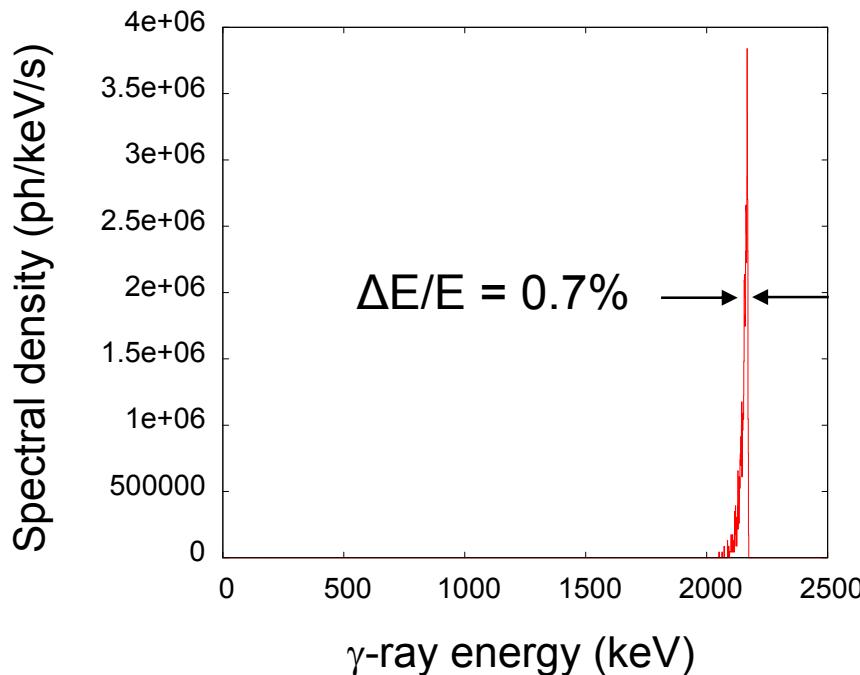
# Design of a 350 MeV ERL for $\gamma$ -ray source

full energy	350 MeV
injection energy	7 MeV
bunch charge	100 pC
repetition	130 MHz
average current	13 mA
$\varepsilon_n (x/y)$	2.5/1 mm-mrad
bunch length	3 ps (rms)
energy spread	$3 \times 10^{-4}$



# Design Performance of 350-MeV ERL LCS $\gamma$ -ray Source

electron beam		laser	
Energy	350 MeV	wavelength	1064 nm
Bunch Charge	100 pC	pulse energy	1.8 $\mu$ J
Bunch Length (rms)	3 ps	rms pulse length	2 ps
Bunch Repetition	130 MHz	pulse repetition	130 MHz
Norm. emittance (x/y)	2.5/1 mm-mrad	enhancement cavity	3000
rms collision spot (x/y)	38/24 $\mu$ m	rms collision spot	29 $\mu$ m
rms energy spread	0.03%	collision angle	3.5 deg.
$\gamma$ -ray flux	$1.0 \times 10^{13} /s$		



total flux  
 $= 1 \times 10^{13} \text{ ph/s}$

Exceeds existing facilities  
by order of  $10^5$ - $10^8$

with a collimator of 0.1mrad  
 $\Delta E/E = 0.7\%$  (FWHM)

# Spent beam

Electron beam degradation due to LCS is summarized in [1]

[1] Z. Huang and R. Ruth, PRL 80, 976 (1998).

Quantum excitation by laser Compton scattering

$$\Delta(\varepsilon_n) = \frac{3}{10} \frac{\lambda_c}{\lambda_L} \frac{(\Delta E)_\gamma}{E} \beta^* \quad \Delta(\sigma_E) = \sqrt{\frac{7}{10}} \hbar \omega \Delta E_\gamma$$

For the designed 350-MeV ERL, 2-MeV  $\gamma$ -ray source,

$$\Delta(\varepsilon_n) = 5 \times 10^{-13} \text{ m} - \text{rad} \quad \Delta(\sigma_E) = 40 \text{ keV} \quad \sim 10^{-4}$$

$$\frac{\Delta E_\gamma}{E} = 0.6\% \quad \text{maximum energy loss}$$

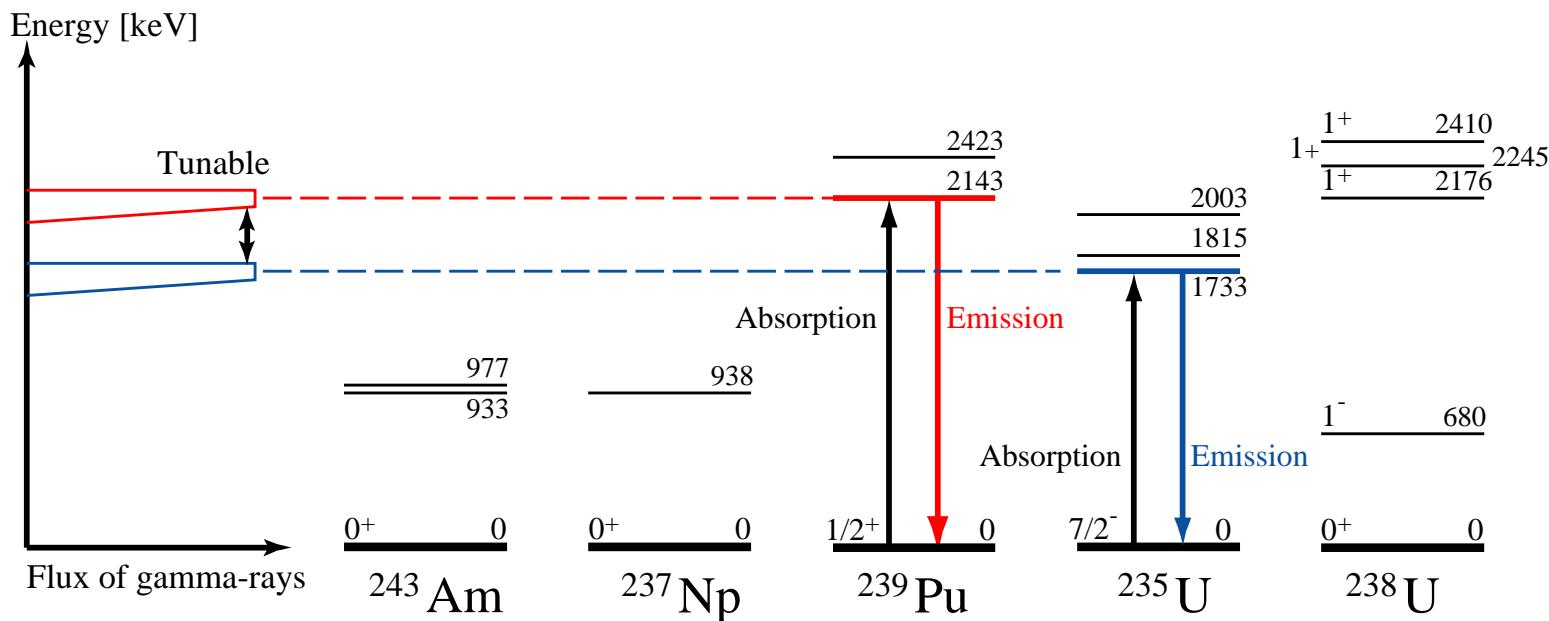


Multiple LCS in the ERL loop is possible with keeping the  $\gamma$ -ray quality

# Requirement from nuclear physics experiments

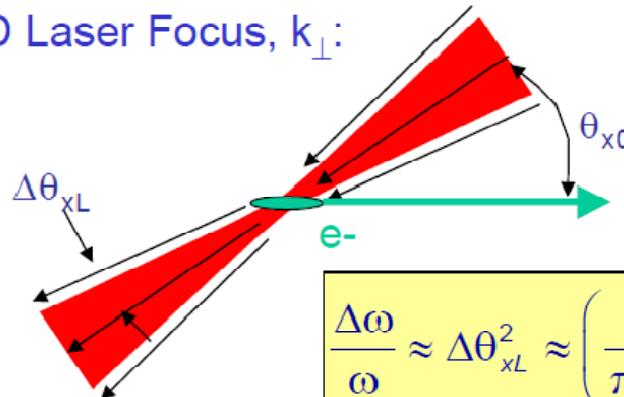
In Nuclear Resonance Fluorescence experiment

- energy width of resonance level is very small  
typically  $\Delta E_{\text{NRF}} \sim 30\text{meV}$  for 2MeV excitation level
- incident  $\gamma$ -ray out of  $\Delta E_{\text{NRF}}$  results in background noise

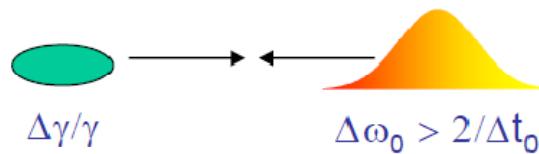


## Bandwidth and Divergence Effects

3D Laser Focus,  $k_{\perp}$ :

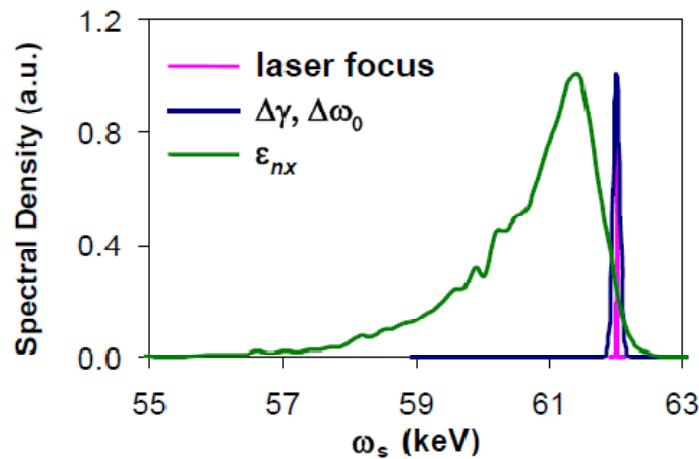
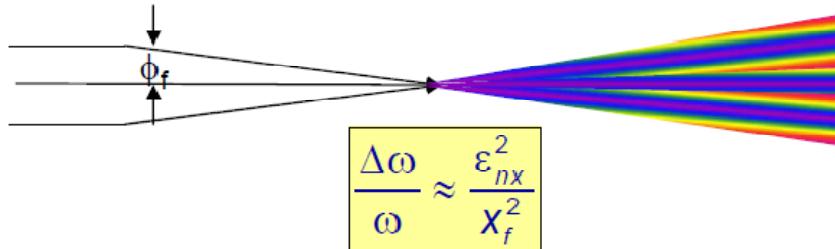


Laser Bandwidth/Electron Energy Spread:



$$\frac{\Delta\omega}{\omega} \approx 2 \sqrt{\frac{\Delta\omega_0^2}{\omega_0^2} + \frac{\Delta\gamma^2}{\gamma^2}}$$

Electron Emittance:



# emittance for the ultimate LCS $\gamma$ -ray source

## diffraction limited electron beam

(energy spread from emittance becomes comparable to that from the laser diffraction)

$$\varepsilon_n \approx \frac{\lambda}{4\pi}$$

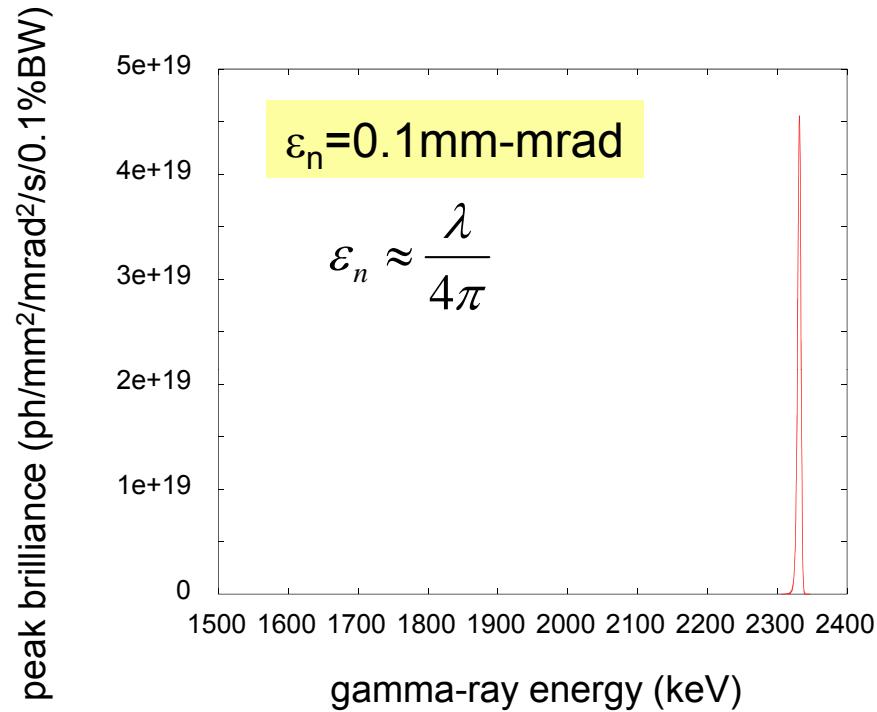
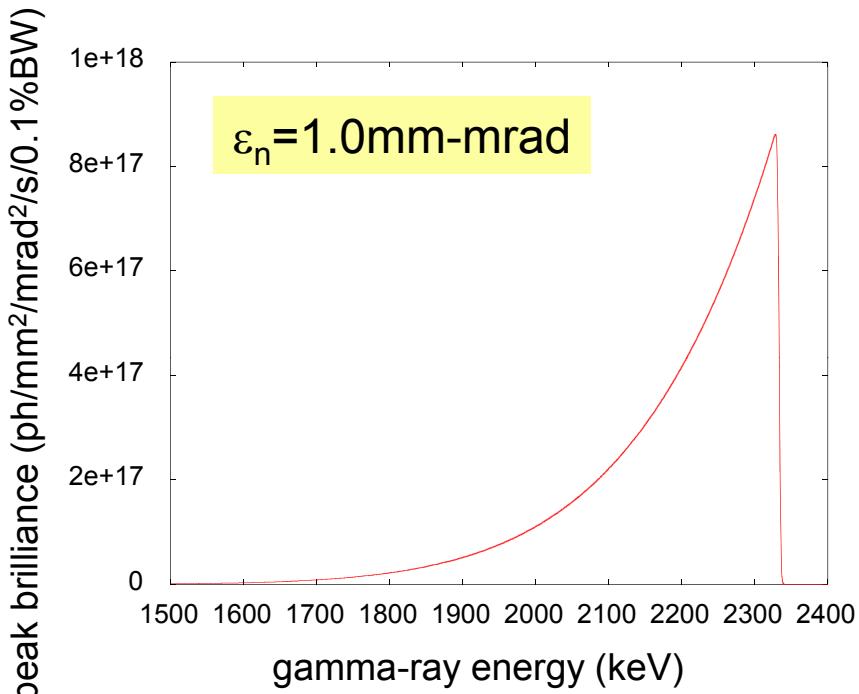
“normalized emittance”

$$E_\gamma \simeq \frac{4\gamma^2 E_L}{1 + (\gamma\theta)^2 + 4\gamma E_L/(mc^2)}$$

$$\lambda=1 \text{ } \mu\text{m} \text{ laser} \rightarrow \varepsilon_n \sim 0.08 \text{ mm-mrad}$$

The target value is same as ERL X-ray sources and XFEL

# On-axis Spectral Brightness (analytical estimation)



$$\hat{B}_x = \frac{4 \times 10^{-15}}{\pi^2} \frac{\gamma_0^2}{\varepsilon^2} \frac{N_e N_\lambda}{\Delta \tau} \frac{r_0^2}{w_0^2} \exp\left[\frac{\chi - 1}{2\chi \Delta u_\perp^2} \left[ 2 + \frac{\delta\omega^2 + \delta\gamma^2\chi^2}{2\chi(\chi - 1)\Delta u_\perp^2} \right]\right] \left[ 1 - \Phi\left\{\frac{\chi - 1}{\sqrt{\delta\omega^2 + \delta\gamma^2\chi^2}} \left[ 1 + \frac{\delta\omega^2 + \delta\gamma^2\chi^2}{2\chi(\chi - 1)\Delta u_\perp^2} \right]\right\} \right]$$

$$\times \frac{\eta e^{1/\mu^2} [\Phi(1/\eta) - 1] - \mu e^{1/\mu^2} [\Phi(1/\mu) - 1]}{\mu^2 - \eta^2}, \quad (50)$$

calculation by using a formula in [1].

[1] F.V. Hartemann et al. Phys. Rev. ST AB 8, 100702 (2005).

# e-beam energy spread for the ultimate LCS $\gamma$ -ray source

contribution of emittance

$$\left( \frac{\Delta E_\gamma}{E_\gamma} \right)_{rms} = 2 \left( \frac{\varepsilon_n}{\sigma_x} \right)^2$$

normalized emittance  $\varepsilon_n = 0.1\text{mm-mrad}$ , collision spot  $\sigma_x = 10\mu\text{m}$   $\rightarrow \Delta E_\gamma/E_\gamma = 0.02\%$

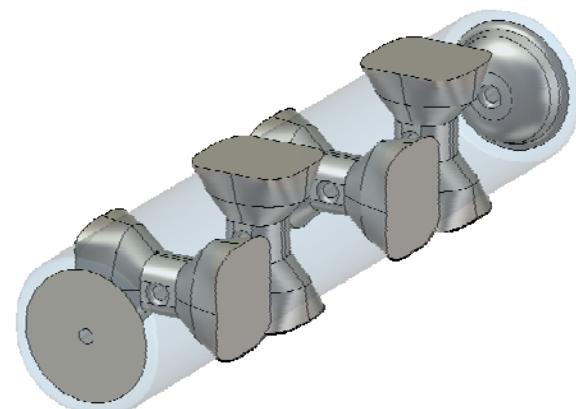
contribution of  
electron energy spread

$$\left( \frac{\Delta E_\gamma}{E_\gamma} \right)_{rms} = \left( \frac{2\sigma_E}{E} \right)$$

$$\sigma_E/E = 0.01 \% \rightarrow \Delta E_\gamma/E_\gamma = 0.02\%$$

for a fixed bunch length  
energy spread due to the RF curvature

$\sigma_E \propto f^2$   possible use of a low-frequency spoke cavity

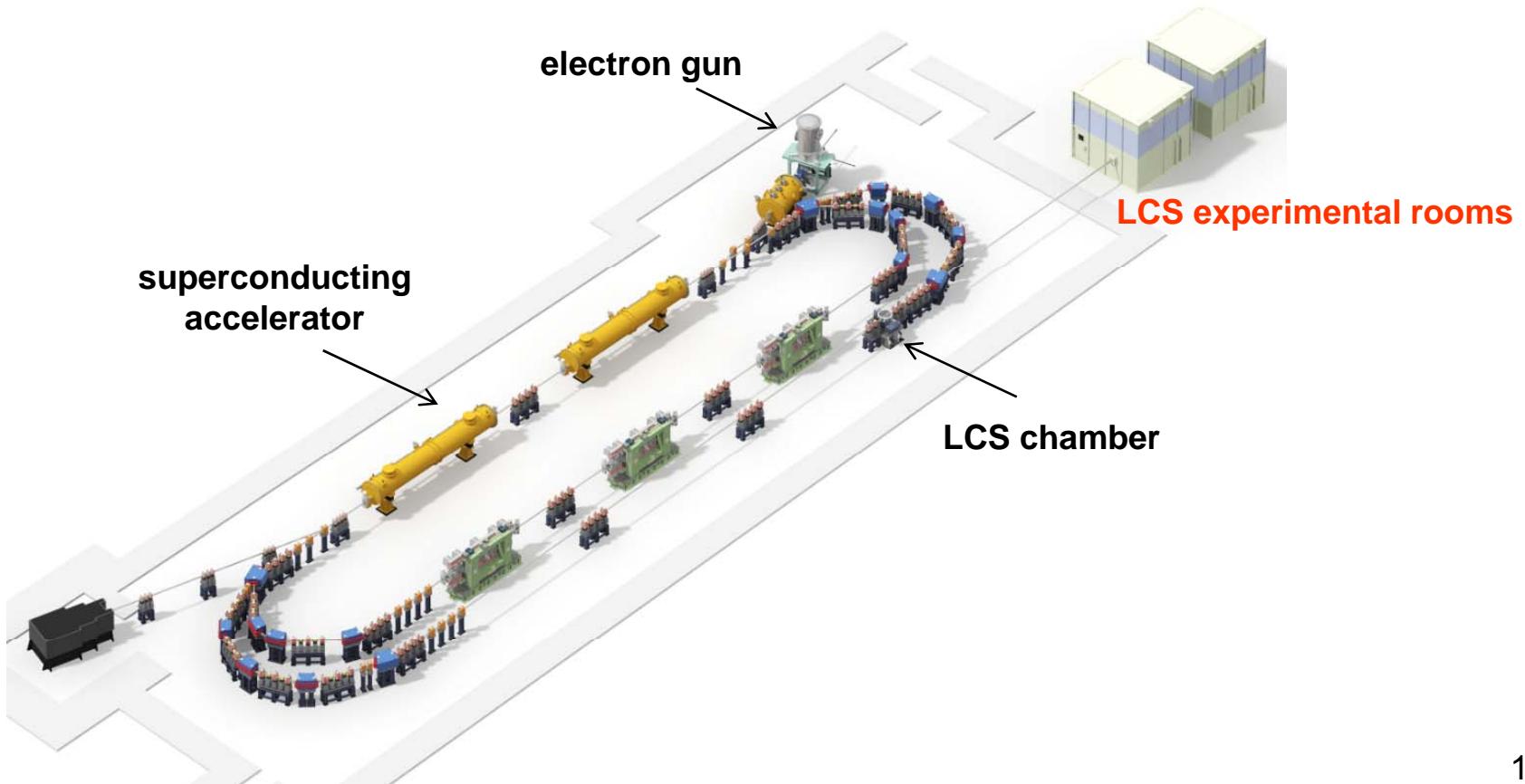


M. Sawamura et al., Proc. SRF-2011 and ERL-2011 Poster

# LCS- $\gamma$ experiment at the Compact ERL

3-year R&D program funded from MEXT

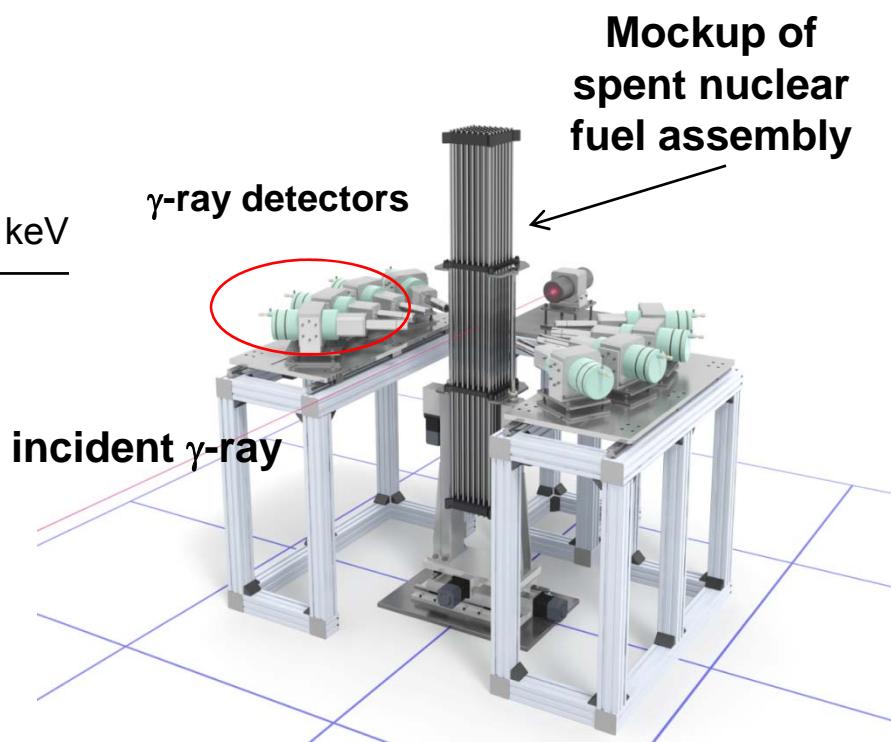
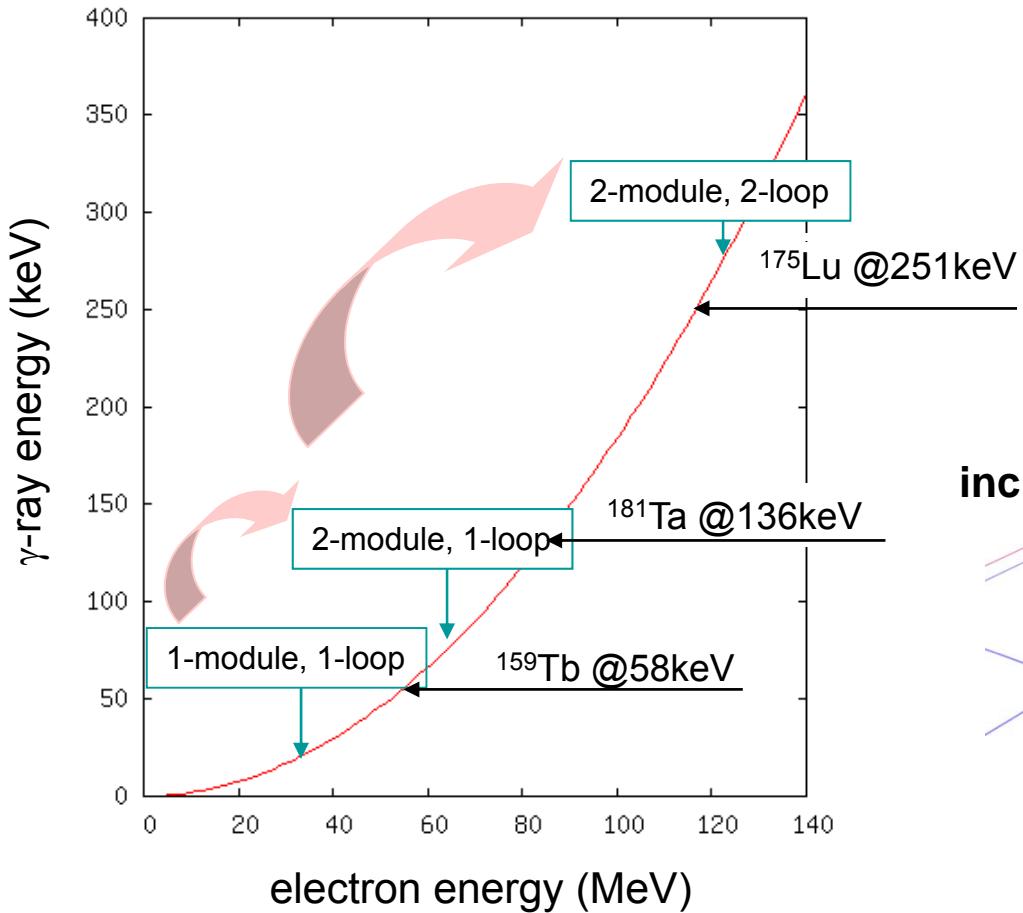
- Installation of a LCS chamber
- Generation of LCS  $\gamma$ -rays
- Demo-Experiment of NRF measurement



# Demo-Experiment at the Compact ERL

Electron Energy Upgrade  
35 MeV → 65 MeV → 125 MeV

Non-destructive measurement of  
“substitute” isotope ( $^{159}\text{Tb}$  for example)



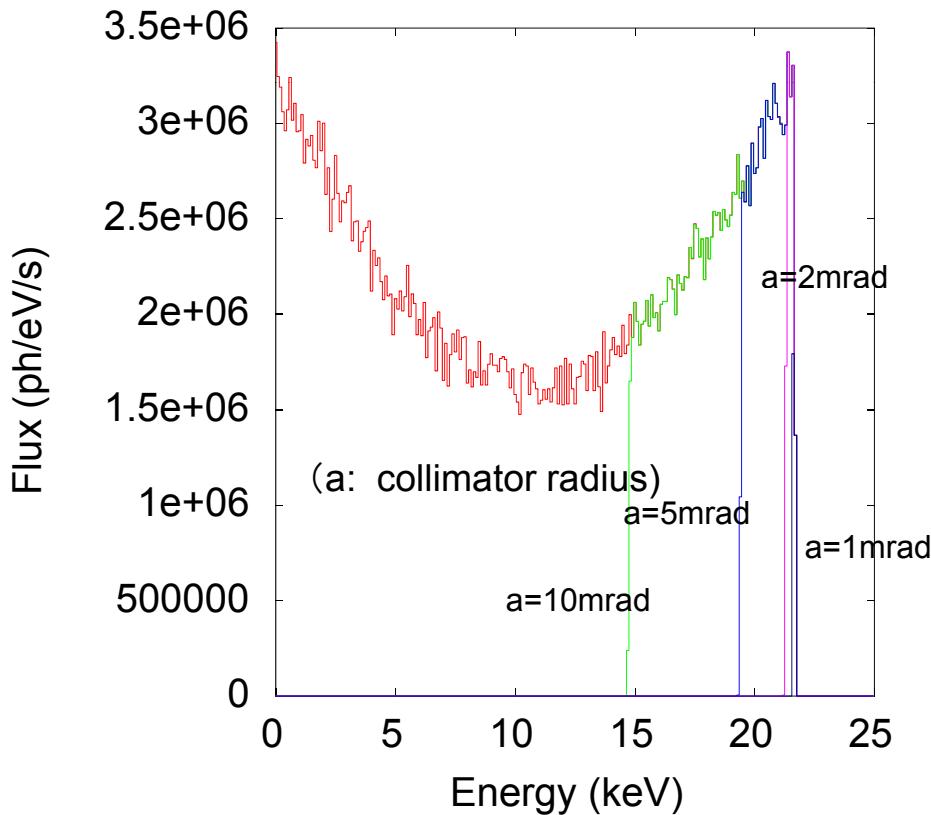
# LCS photon flux expected in 35-MeV cERL

rep. rate 130 MHz, laser intracavity power = 100kW, collision spot = 40 $\mu$ m

8pC, 0.1mm-mrad

total flux = 4.8E10/s

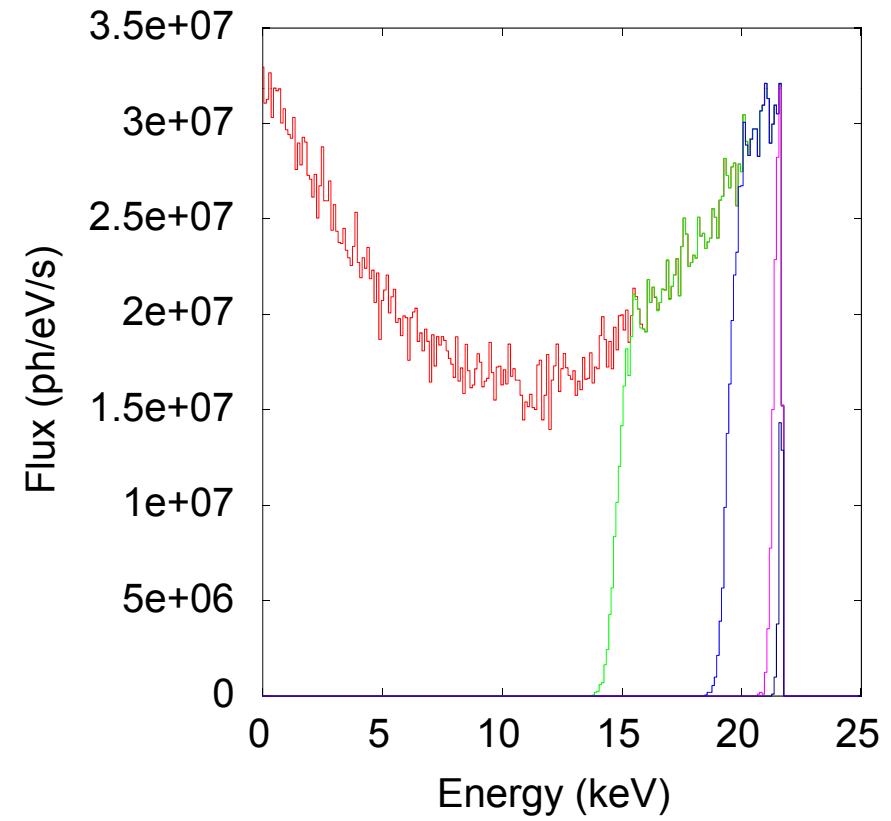
$\Delta E/E$  (rms) = 100eV for  $a=1.8$  mrad



80pC, 1mm-mrad

total flux = 4.8E11/s

$\Delta E/E$  (rms) = 100eV for  $a=1.4$  mrad



# Summary

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- ERL is an ideal driver for laser Compton  $\gamma$ -ray sources.
  - no accumulation of quantum excitation by hard photon emission
- High-flux and narrow-bandwidth  $\gamma$ -rays are available.
- There are many scientific and industrial applications.  
(See Hayakawa's talk)
- Acc. technologies for “diffraction limited” electron beams can be shared with ERL X-ray sources and XFELO.
- Beam dynamics issues such as emittance growth, BBU, multi-loop design are also common to other ERLs.
- For the further reduction of  $\gamma$ -ray bandwidth, we need some developments: spoke cavity, precise LLRF, energy-width compensation ...