

SIMULATION STUDY ON TRANSVERSE LASER COOLING AND CRYSTALLIZATION OF HEAVY-ION BEAMS AT THE COOLER STORAGE RING S-LSR

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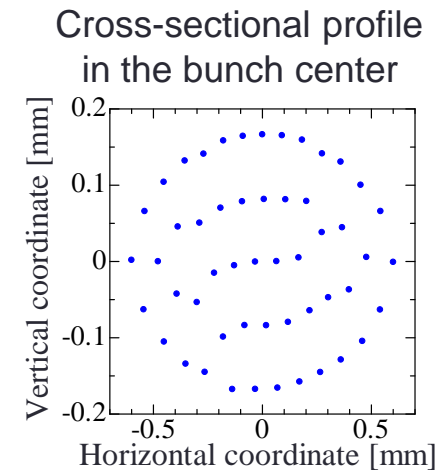
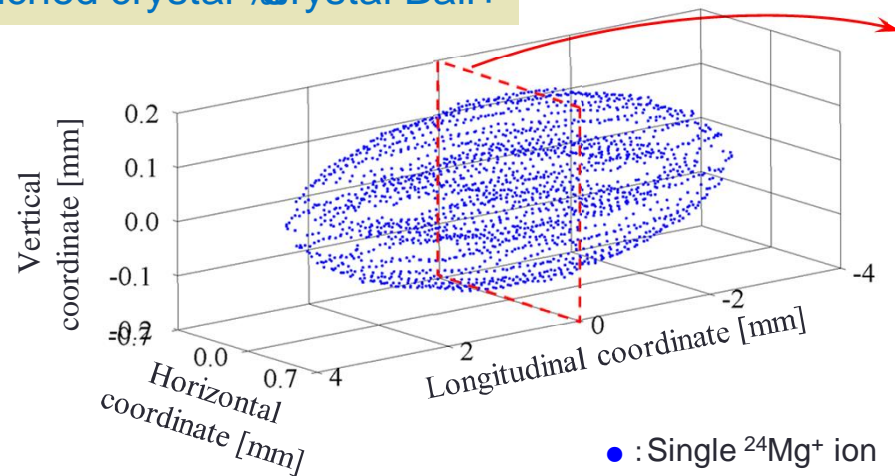
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Outline

- ” Intro: Purpose of the present study
- ” Molecular dynamics (MD) simulation
- ” Simulation conditions
- ” MD results
 - ” Three-dimensional laser cooling
 - ” Crystallization
- ” Summary

Beam Crystallization

3D bunched crystal %Crystal Ball+



- ” Coulomb crystalline state of an ion beam strongly cooled in a storage ring
- ” Characteristics:
 - ” Ultralow emittance
 - ” Coulomb coupling constant > 170
 - ” Periodic oscillation with the external focusing force
 - ” Stable after removing the cooling force

Purpose of the Present Study

- “ Feasibility of beam crystallization was already predicted if the ring and laser conditions were sufficient. (PRL2004, PRSTAB2005)
- “ However, laser-cooling conditions have been limited in the recent experiments at S-LSR.
 - “ Single laser beam, low power, and fixed detuning.
- “ To show numerically how to attain a low-emittance beam using **Resonant Coupling** and **Laser Cooling** by assuming actual parameters at S-LSR.
 - “ Optimization of a cooling laser for high cooling efficiency
(To be presented at NA-PAC13)
 - “ **Fast 3D cooling of low-current beams**
 - “ **Feasibility of beam crystallization**
- “ Numerical study using a Molecular Dynamics (MD) simulation technique.

Molecular Dynamics (MD) Simulation

“ The most reliable simulation technique for the study of beam cooling and crystallization.

“ Hamiltonian

$$H = \frac{p_x^2 + p_y^2 + p_z^2}{2} - \frac{\gamma}{\rho} x p_z + \frac{x^2}{2\rho^2} - \frac{K(s)}{2} (x^2 - y^2) + \frac{r_p}{\beta^2 \gamma^2} \phi.$$

“ Motion of real particles is integrated in a symplectic manner.

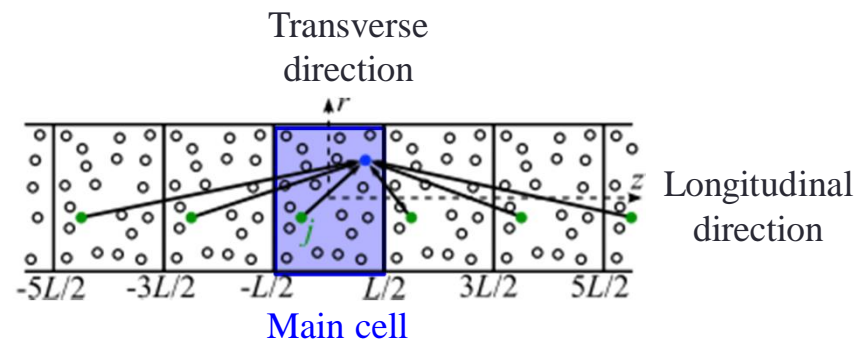
“ Coulomb potential --- Periodic boundary condition imposed

$$\phi = \sum_j (\phi_{short}^{(j)} + \phi_{long}^{(j)}).$$

$$\phi_{short}^{(j)} = \frac{1}{\sqrt{(x-x_j)^2 + (y-y_j)^2 + (z-z_j)^2}}$$

$$\phi_{long}^{(j)} = \frac{2}{L} \int_0^\infty \frac{\cosh(kz^{(j)}/L) J_0(kr^{(j)}/L) - 1}{e^k - 1} dk$$

where $z^{(j)} = |z - z_j|$ and $r^{(j)} = \sqrt{(x-x_j)^2 + (y-y_j)^2}$.



For a bunched beam, L can be set as a bucket length (C/h).

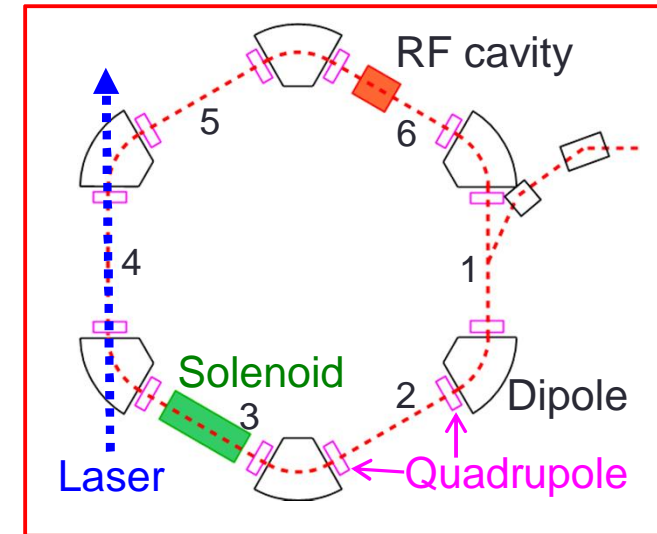
Resonant Coupling for 3D Cooling

- “ A possible scheme for efficient transverse cooling
 - “ H. Okamoto, D. Mohl, and A. M. Sessler, (PRL1993, PRE1994)
- “ First, introduce a coupling source in the ring.
 - “ RF cavity placed where the dispersion is finite for X-Z coupling
 - “ Solenoid magnet for X-Y coupling
- “ Then, operate the ring at a difference resonant condition;
 - $\nu_x - \nu_z \approx \text{integer}$ for X-Z coupling
 - $\nu_x - \nu_y \approx \text{integer}$ for X-Y coupling

MD Simulation Conditions (1)

- “ Machine (S-LSR at Kyoto Univ.)
 - “ Circumference 22.56 m
 - “ Superperiodicity 6
- “ Lattice
 - “ Tunes **Case-I** (v_x, v_y, v_z)~(2.07, 1.12, 0.07)
Case-II (v_x, v_y, v_z)~(2.07, 1.07, 0.07)
 - “ RF bunching voltage ~40 V
 - “ Harmonic number 100
 - “ Adiabatic capture 5,000 turns (0.2sec)
- “ Beam
 - “ Ion species 40-keV $^{24}\text{Mg}^+$
 - “ Lorentz factors $\beta=1.89 \times 10^{-3}$, $\gamma=1.00000179$
 - “ Revolution frequency (period) 25 kHz (40 μsec)
 - “ Initial RMS emittance ($\varepsilon_x=\varepsilon_y$)
 - $1 \times 10^{-9} \pi \text{ m.rad}$ (Normalized)
 - $5 \times 10^{-7} \pi \text{ m.rad}$ (Un-normalized)
 - “ Initial dp/p (rms) 3×10^{-4}

Schematic view of S-LSR



From the
measurement
result

MD Simulation Conditions (2)

“ Laser (1 co-propagating laser)

“ Power	8mW
“ Spot radius w (2sigma)	0.66 mm (Peak Saturation Power~4.6)
“ Detuning Δ (fixed)	-200 MHz
“ Cooling time	3 sec

} From the
experiment

These parameters are rather limited as compared to past experiments in TSR & ASTRID.

$$\text{Cooling force: } F = \frac{1}{2} \hbar k_L \Gamma \frac{S_L}{1 + S_L + (2\Delta / \Gamma)^2}$$

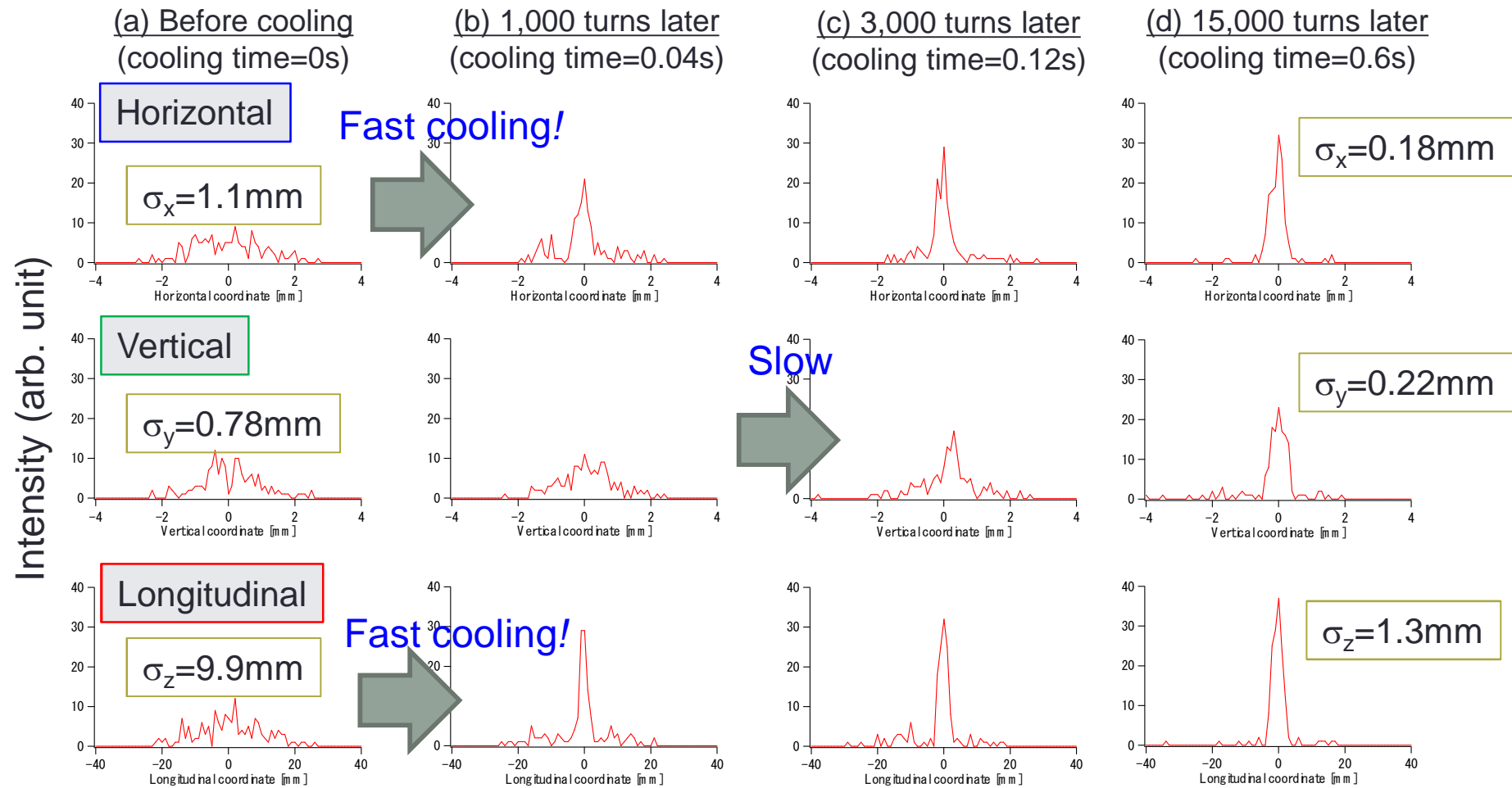
$$\text{Saturation parameter : } S_L = S_0 \exp\left[-\frac{2(x^2 + y^2)}{w^2}\right]$$

$$\text{Laser detuning : } \Delta \approx \omega \gamma \left[1 - \beta \left(1 + \frac{\delta p}{p}\right)\right] - \omega_0$$

MD Results (1: Time evolution)

Case-I

(v_x, v_y, v_z)
 $\sim (2.07, 1.12, 0.07)$



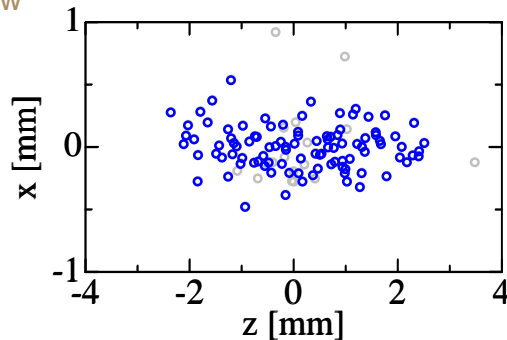
The vertical direction is cooled through the Coulomb interaction between ions, although no artificial cooling force is introduced.

MD Results (2: Equilibrium state)

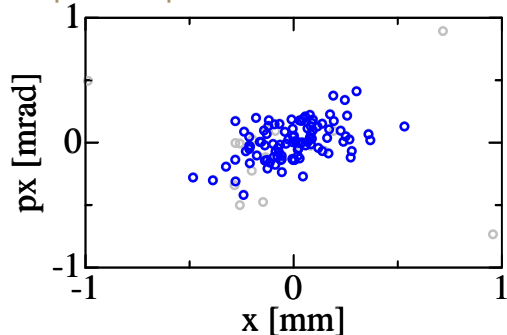
Case-I

(v_x, v_y, v_z)
 $\sim (2.07, 1.12, 0.07)$

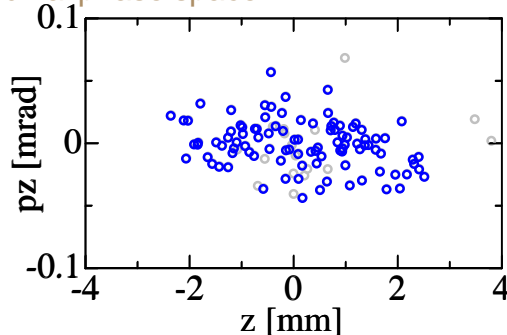
Top-view



Horizontal phase space



Longitudinal phase space



“ The ion number of the cooled part (**blue ions** in the picture) is about 100. Namely, the cooling efficiency is about 70%.

“ Horizontal

“ Norm. rms $\varepsilon=4.6 \times 10^{-11}$ [$\mu\text{m}\cdot\text{rad}$]

“ $T_x=18[\text{K}]$

“ **Radius $\sigma=0.18\text{mm}$**

“ Vertical

“ Norm. rms $\varepsilon=2.6 \times 10^{-11}$ [$\mu\text{m}\cdot\text{rad}$]

“ $T_y=3.8[\text{K}]$

“ **Radius $\sigma=0.22\text{mm}$**

“ Longitudinal

“ Rms $dp/p = 2.2 \times 10^{-5}$

“ $T_z=0.45[\text{K}]$

“ Radius $\sigma=1.3\text{mm}$

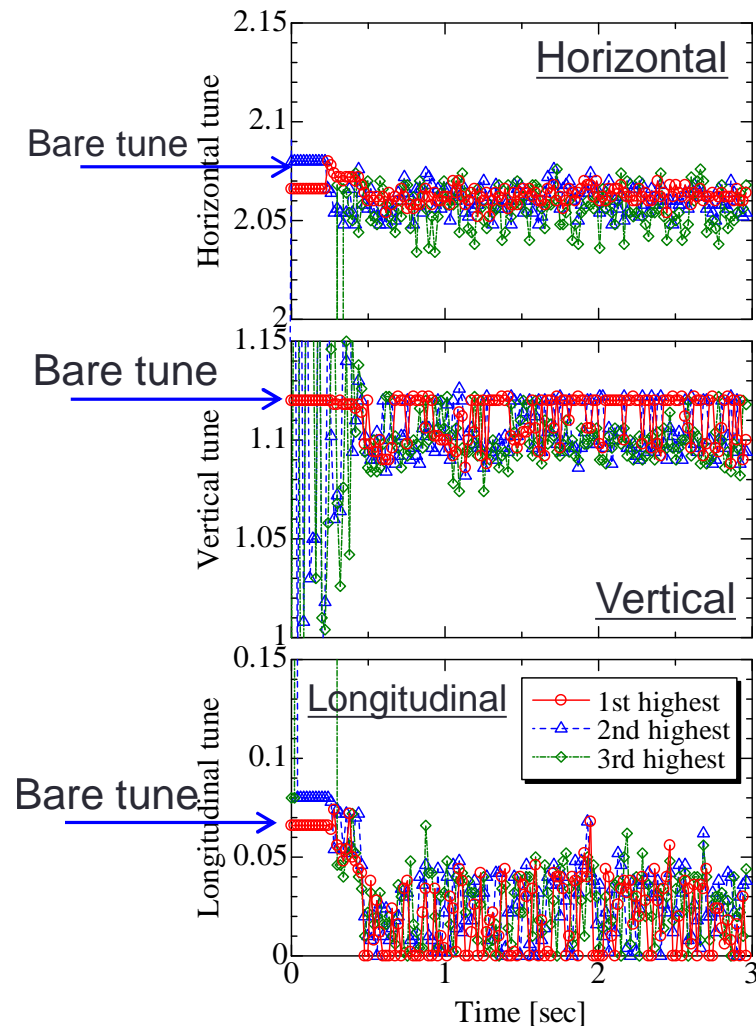
These values agree well with the observation result in S-LSR!!

The beam is three-dimensionally cooled, but the ordered configuration cannot be seen.

Case-I

 (v_x, v_y, v_z) $\sim(2.07, 1.12, 0.07)$

MD Results (3: Tune shift)



“ The orbits of several ions are Fourier-transformed to see the time evolution of tunes in all three directions.

“ The three highest peaks in the power spectrum (right pictures) are plotted.

“ Result: tune shift

“ Horizontal 2.07 --> 2.05~2.06

“ Vertical 1.12 --> 1.09~1.10

“ Longitudinal 0.07 --> 0.00~0.04

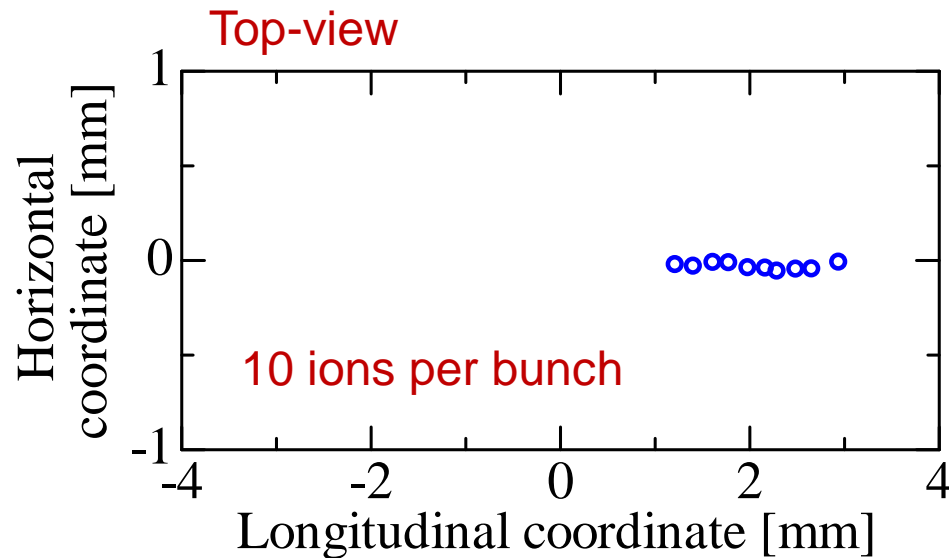
“ The synchrotron tune is almost damped by laser cooling.

“ The beam is still oscillating in the transverse direction.

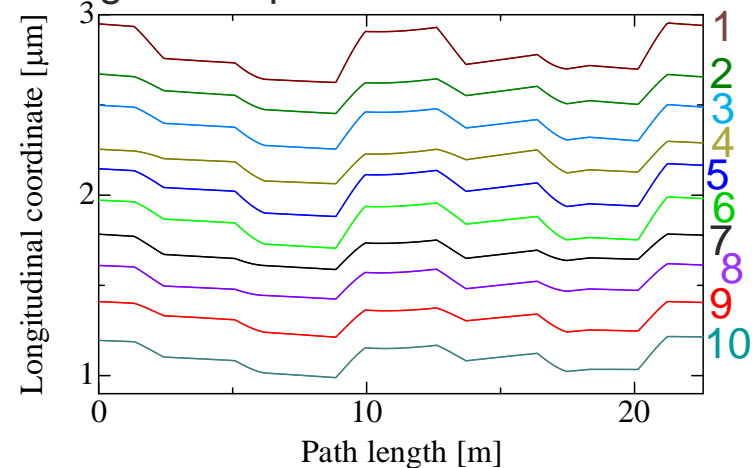
The laser-cooled beam is three-dimensionally space-charge-dominated.

MD Results (4: Crystallization)

Case-I

 (v_x, v_y, v_z) $\sim(-2.07, 1.12, 0.07)$ 

Time evolution (one turn) of the longitudinal positions of the 10 ions



“ Even with the limited laser-cooling condition, **1D string crystal** can be formed when the beam current is sufficiently low and detuning is small.

“ Each ion does not pass by neighboring ions.
 “ The synchrotron oscillation is fully depressed.

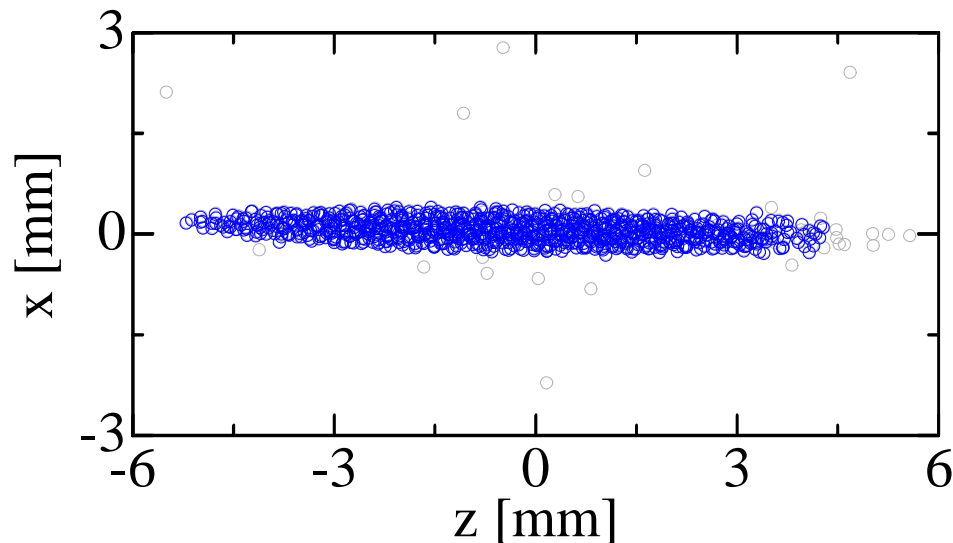
Note that the bunch is positioned forward because the beam is pushed by the co-propagating laser.

Beam crystallization is feasible at S-LSR!!

Case-II

 (ν_x, ν_y, ν_z) $\sim (2.07, 1.07, 0.07)$

MD Results (5: Ideal case)



- " 1000 ions per bunch
- " 3D-resonant tunes:
 - " $(\nu_x, \nu_y, \nu_z) = (2.07, 1.07, 0.07)$
 - " Weak solenoid $B=80\text{G}$
- " Laser conditions:
 - " 2 lasers (co- and counter-propagating)
 - " High power (100mW)
 - " Frequency scanned (-4GHz to -40MHz for 1sec)

- " More than 90% ions are laser-cooled.
- " Transverse norm. rms emittance $\sim 1 \times 10^{-11} \pi \text{m.rad}$ ($T_{x,y} \sim 10\text{K}$)
- " Longitudinal momentum $dp/p \sim 1 \times 10^{-5}$ ($T_z \sim 0.1\text{K}$)

The highest-quality heavy-ion beam can be formed just by improving the laser system in S-LSR!!

Summary

- “ 3D laser cooling of the heavy-ion beam in S-LSR was studied using the MD simulation technique.
- “ The three-dimensionally low-temperature bunched ion beam was generated through resonant coupling.
- “ The MD result agreed well with the observation result in the recent experiment in S-LSR.
- “ Beam crystallization (1D string at low line density) is possible even in a limited cooling situation.
- “ An ultra-low-emittance bunched beam can be formed at a high intensity by a combination of powerful laser cooling and resonant coupling.