

Enhancing trappable antiproton populations through an induction unit followed by frictional cooling

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COOL11, Pansionat "Dubna", Alushta
(Crimea, Ukraine) 12 - 16 September, 2011.

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Motivation

Particular Motivation

To prepare a paper so Max and I could attend this conference. In the course of thinking about what we should work on we came up with this subject, which – actually -- turns out to be in our opinion rather interesting, and you shall see if you agree..

- **Basic thoughts**

Increase the anti-proton flux prior to the construction of ELENA.

- That is, a fast, and rather inexpensive way, to proceed during the ELENA construction period.
- Currently, starting with more trapped antiprotons does not necessarily yield more trapped anti-hydrogen atoms, but as theoretical understanding and experimental control of trapping and mixing improve, the option to start with an order of magnitude more antiprotons might be quite useful.

- **Caveats**

The ideas to be presented here are very preliminary.

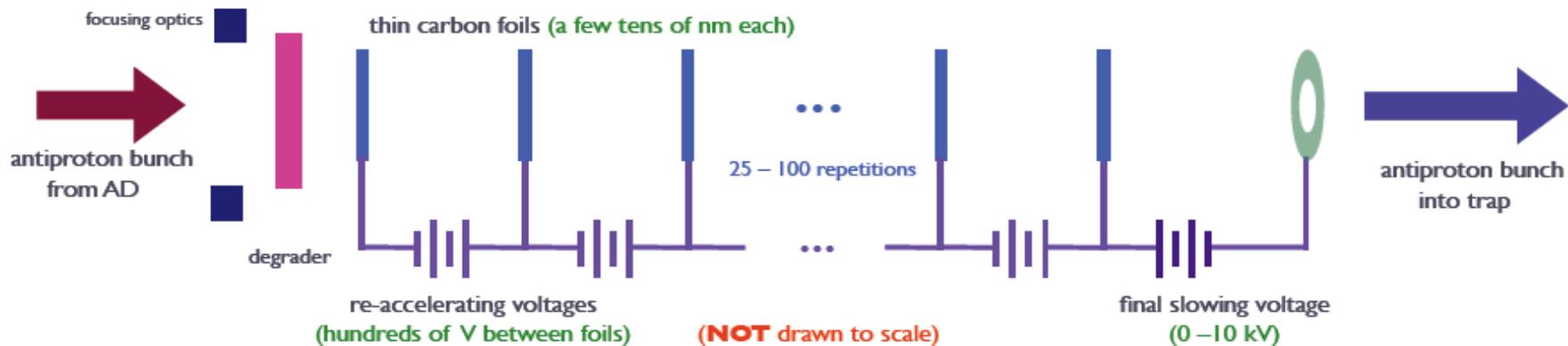
Back-of-the-envelope *estimates* and some Monte Carlo *simulations* suggest that a reasonably simple and compact design would result in an increased antiproton trapping by about one order-of-magnitude.

Multiple-scattering may be significant, and strong solenoidal fields are probably required. More detailed simulations will be needed to verify preliminary results and optimize performance.



Overview of the Concept

- a large mismatch exists between the average kinetic energy of antiprotons exiting the AD and the kinetic energy of antiprotons that can be trapped
 - ~ 5.3 MeV vs. ~ 3 keV
- to enhance yields, one must first increase the number of antiprotons in the region of phase space that subsequently can be trapped:
 - particles must have sufficiently low kinetic energy
 - and sufficiently low divergence angle (outside solenoid) and/or gyro-radius (inside solenoid)
- at sufficiently low velocities inside matter, antiprotons experience a stopping power due to collisions that is approximately linear in the particle velocity
- this “viscous drag”-like damping force can be exploited to try to produce:
 - longitudinal slowing: bringing down the average kinetic energy/longitudinal momentum
 - longitudinal cooling: decreasing the RMS spread in kinetic energy/longitudinal momentum
 - possibly transverse cooling: controlling the RMS divergence angle/gyro-radius
- we propose using a series of thin carbon foils separated by re-accelerating electrostatic gradients in order to slow and cool a portion of the antiproton beam after it exits the degrader...



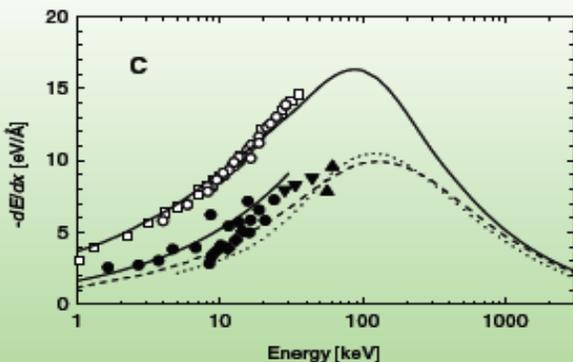
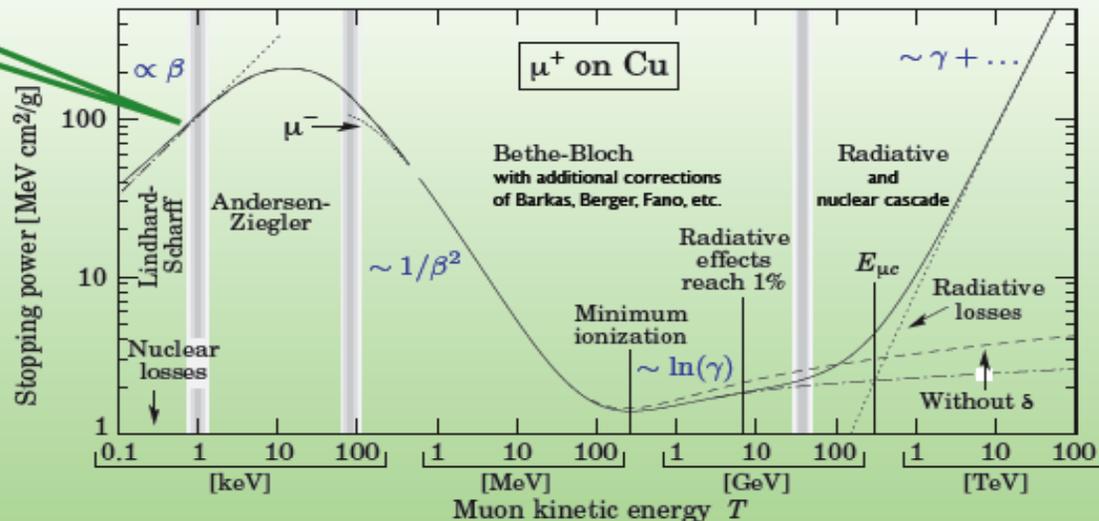
Stopping Power & Frictional Cooling

adiabatic “frictional” regime

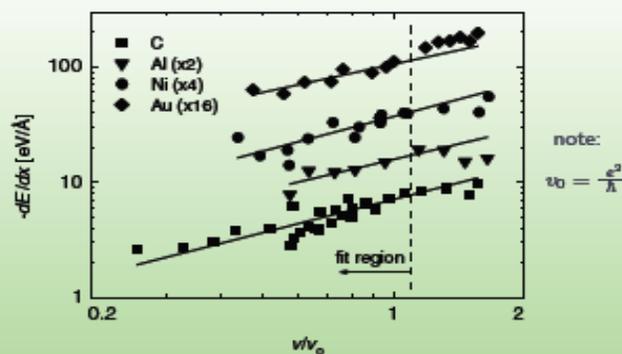
at particle velocities below bound electronic velocities in atoms and above a low threshold where nuclear recoil losses are important, the **average force** is approximately *proportional to the speed and directed opposite the velocity*, leading to an approximately **linear damping**

frictional cooling was first discussed in the context of muons, specifically to cool antimuons for various high-precision measurements including anomalous magnetic moments, or the hyperfine structure of muonium

typical $\left\langle \frac{1}{\rho} \frac{dE}{dx} \right\rangle$ for charge particles in solid matter (from reference [1])



measured low-velocity stopping power in solid carbon for protons (open symbols) and antiprotons (filled symbols)



measured antiproton stopping power normalized to Bohr velocity on log-log plot (both plots from reference [12])

although the general features of the stopping power curves are similar, the magnitude of the damping forces in a given medium differ for a particle and its anti-particle, due primarily to material polarization (as well as exchange effects)

This **Barkas Effect** was first observed for pions, but is clearly evident for protons and antiprotons...

Comments

- anywhere **stopping power is positive**, collisional energy losses lead to **slowing** on average:
 - decrease in **average** kinetic energy
 - decrease in **average** longitudinal momentum (the only component with non-zero average)
 - if **damping foils are alternated with longitudinal re-accelerating gradients**, **average** longitudinal momentum of (a suitably low-energy portion) of beam can approach an equilibrium
- equilibrium energy/momentum determined by balance of energy gains between foils and losses in foils
- note only particles in some sufficiently low initial energy range will have time to reach equilibrium over an energy range where **stopping power increases monotonically** with particle kinetic energy, particles can also be cooled longitudinally
- because **faster particles experience more slowing** so **variance** in energy or longitudinal momentum may be reduced *for particles within coolable energy range*
- but stochastic nature of collisions leads to unavoidable fluctuations in the number and extent of the individual energy/momentum transfers



Comments (Concl.)

- diffusion necessarily accompanies damping (fluctuation-dissipation theorem)
- leads to **straggling**, or non-zero **variance** in energy changes
- interplay of momentum diffusion, damping, and deceleration determines achievable longitudinal cooling
- since damping forces point on average in direction **opposite total momentum**, but only longitudinal momentum is restored, particles may be **cooled transversely** as well
- in absence of fluctuations, this decreases the RMS divergence angle or (in a solenoidal field) the gyro-radius
- **but** *fluctuations* contribute a diffusive **heating** term, described by *multiple-scattering rates*
- **multiple scattering** between re-accelerations tends to increase heating term without improving cooling term
- some particles may be stopped or back-scattered, and hence **lost** from the beam
- **lower-Z** materials lead to lower multiple scattering rates — we are focusing on **carbon**
- coupling between the angular, spatial, and energy drift/diffusion may complicate the Fokker-Planck dynamics
- but it appears that here the energy/angular coupling can be approximated simply
- various other effects may also occur, but are expected to be less important:
- space-charge emittance growth, intra-beam scattering, annihilation....



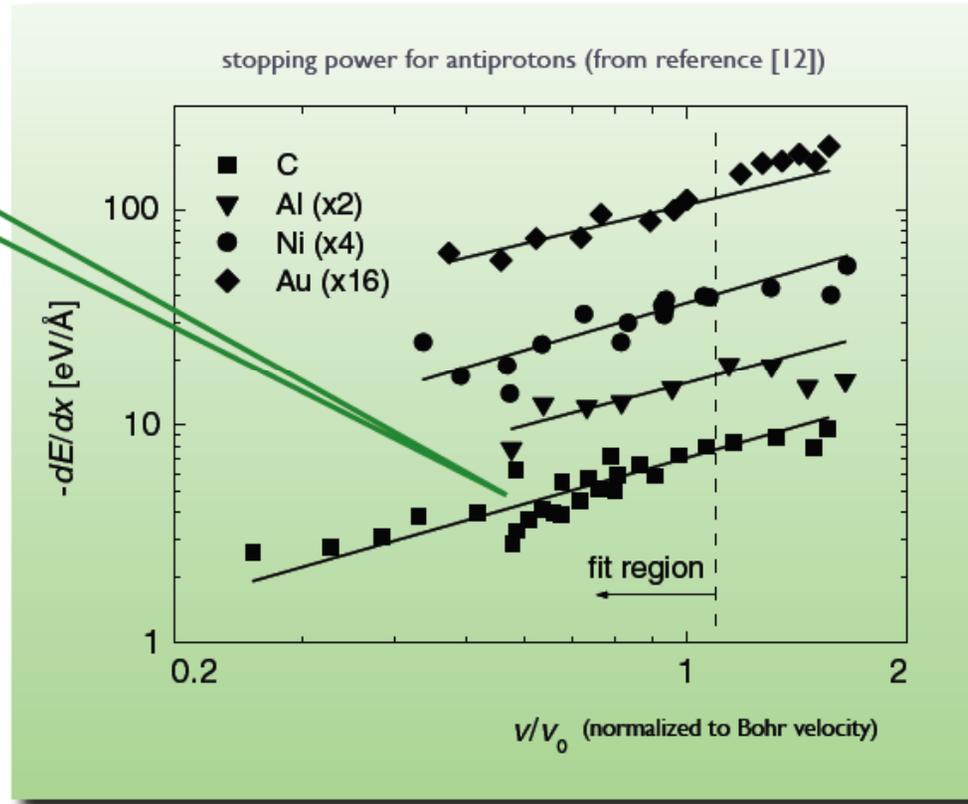
Antiproton Stopping Power

average stopping power estimated from data for antiprotons in carbon

we use these data to estimate the average stopping power as a function of kinetic energy:

$$\left\langle -\frac{dE}{dz} \right\rangle \approx 0.45\sqrt{E} \quad \left[\frac{\text{eV}}{\text{nm}} \right]$$

where energy E is in eV
for kinetic energies less than about 60 keV



Energy Straggling: Fluctuations in Energy Loss

dashed blue line shows predictions based on a semiclassical binary collision model

solid black line includes fitted effects of foil inhomogeneities

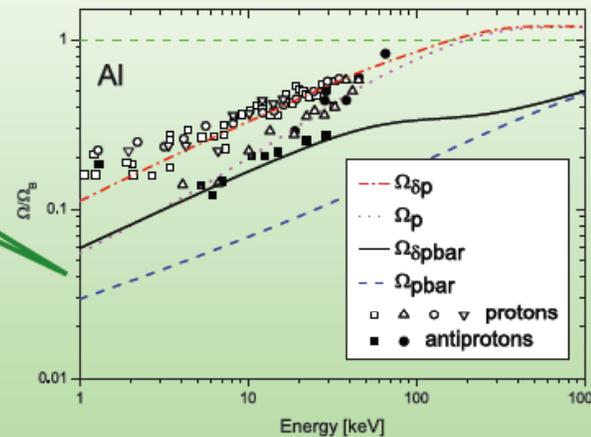
we use these data to make an estimate for the energy straggling term — on the order of **10% of the Bohr value** at low energies, and scaling linearly with particle speed:

$$\frac{d\Omega^2}{dz} \approx .05 E \left[\frac{\text{eV}^2}{\text{nm}} \right]$$

(E is in eV)

Ω is the rms energy spread

straggling for protons/antiprotons in aluminum foils (from reference [11])



Measured antiproton (■, 260 Å foil, ● 350 Å) and proton (Δ 260 Å, ▽ 260 Å, □ 320 Å, ○ 320 Å) fluctuation in aluminum, normalized to the Bohr straggling Ω_B .

data for antiprotons *in carbon* have not been reported, but similar data for Al and Au (not shown) suggest Z-dependence is weak...

$\Omega_B^2 = 4\pi Z e^4 n t$, where Z is the foil atomic number and n is the foil density and t is the foil thickness



Multiple Scattering and Angular Diffusion

data suggest that at kinetic energies of a few keV to a few tens of keV, the multiple scattering rates for *protons* are about **one order-of-magnitude lower** than the level suggested by the commonly-used Molière analytic estimate

Barkas Effect suggests that *antiprotons* might suffer **even less** angular diffusion from multiple scattering, but the magnitude of the Barkas effect is uncertain....

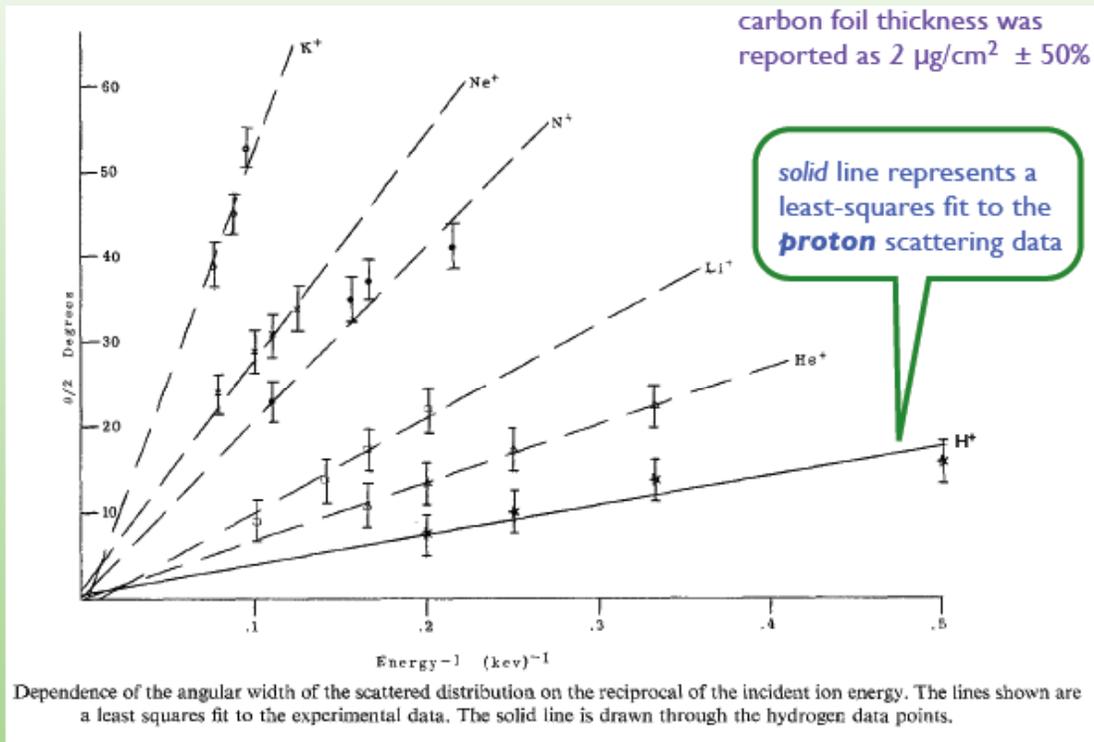
multiple scatter estimated as:

$$\frac{d\sigma_{\theta}^2}{dz} \approx \kappa \frac{2.3 \cdot 10^5}{E^2} \left[\frac{\text{rad}^2}{\text{nm}} \right]$$

for some $\kappa \leq 1$

from the graph, we estimate $\kappa \approx .04$ for protons in carbon

angular divergence vs. inverse incident energy for ions in thin carbon foils (reference [3])



data for *antiprotons* in carbon have not been reported to our knowledge....

Although scattering is expected to be *weaker* for *antiprotons* than *protons*, in absence of theoretical model or experimental data, we **conservatively** use values in the range $0.10 \leq \kappa \leq 1.0$

Simple model

Homogeneous media with homogeneous electric field to compensate average energy losses, in the low-energy limit where the energy loss dE_{fr}/dx is proportional to $E^{1/2}$ (i.e., velocity).

$$\frac{d\sigma^2}{dx} = -\frac{dE_{fr}}{E_{eq} dx} \sigma^2 + \frac{d\Omega_s^2}{E_{eq}^2 dx} + \frac{1}{4} \left(\frac{dE_{fr}}{E_{eq} dx} \right)^{-1} \left(\frac{d\sigma_\theta^2}{dx} \right)^2$$

Cooling

Straggling
heating

Scattering
heating

Here, σ is the relative energy spread, E_{eq} is the energy where losses equal the energy gain in the applied electric field, and σ_θ is the multiple scattering angle. In the region of interest one can neglect the “straggling” term Ω_s .

In equilibrium

$$\sigma = \frac{1}{2} \left(\frac{dE_{fr}}{E_{eq} dx} \right)^{-1} \frac{d\sigma_\theta^2}{dx}$$



From the experimental data:

Energy in eV , distance in nm and angles in radian

$$q\mathfrak{R} = \frac{dE_{fr}}{dx} + E_{eq} \frac{d\sigma_{\theta}^2}{dx}$$

$$\frac{dE_{fr}}{dx} = .45 \sqrt{E_{eq}} \left[\frac{eV}{nm} \right]$$

$$\frac{d\Omega_s^2}{dx} = .05 E_{eq} \left[\frac{eV^2}{nm} \right]$$

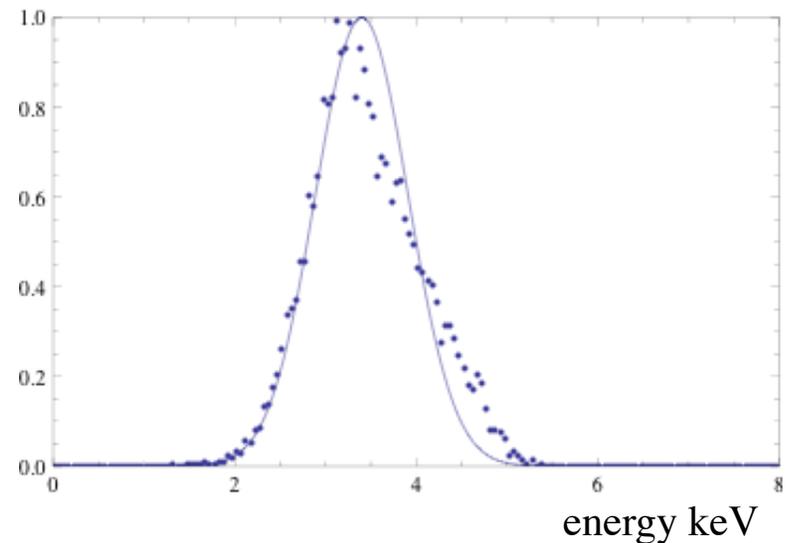
$$\frac{d\sigma_{\theta}^2}{dx} = \frac{\kappa 2.3 10^5}{E_{eq}^2} \left[\frac{1}{nm} \right]$$

\mathfrak{R} is the applied electric field

q is the charge of pbar

Results of simple model and Monte Carlo simulations

mean energy $E_{eq} = 3.5$ keV
RMS energy spread = 600 eV
RMS divergence angle = 0.4 radians
RMS spot size = 6 mm
losses = 18%
RMS time-spreading = 17 ns
pop. enhancement in 3 keV window
centered at 3.5 keV = 12X
70 carbon foils (20 nm each)
 $V = 540$ V
total $V \approx 38$ kV



Solid line is a Gaussian distribution with parameters as listed above.

Dots are from simulations using the experimental data from the previous slide.

Numerical Simulations

We consider four cases:

1. The Anti-proton De-accelerator (AD) giving 5 MeV anti-protons followed by a degrader foil (which is the present situation) and then followed by a frictional cooling section.
2. The AD with an induction accelerator operating from 5 MeV to 50 keV followed by a frictional cooling section.
3. The AD with an RFQ to 50 keV followed by a frictional cooling section.
4. The performance of ELENA. ELENA followed with a small degrader foil. No frictional cooling needed.



Monte Carlo Model

•initial conditions:

- energy distribution based on assumed output beam from degrader:
- degrading foil comparable in thickness to typical *range* of antiprotons at mean energy of AD output
- produces wide energy distribution
- but approximately **uniform kinetic energy spectrum** between 0 keV and ~300 keV
- an estimated ~4% of the original antiproton population in bunch will lie below 50 keV
- simulations considered sub-population of antiprotons with kinetic energy below 50 keV
- assumed a uniform energy distribution between 0 keV and 50 keV— at higher energies particles cannot be cooled
- focused on *relative enhancement* of population in a window around a few keV— absolute populations not needed
- transversely, used a gaussian beam with 2 mm spot size
- and 0.03 radian RMS divergence (likely too small, but largely irrelevant as it blows up after first foil anyway)

•other physical assumptions:

- non-relativistic kinematics
- 20 nm thick carbon foils
- equal DC voltage drop between successive foils
- ambient longitudinal magnetic field of magnitude 3.5 T everywhere
- annihilation was ignored
- tracked individual sample particle trajectories and collisions

•transport parameters:

- adopted stopping powers and straggling as in equilibrium theory
- tried various values of κ , e.g., 0.11, 0.25, 1.0

•performance:

- custom simulations were performed, because *ICOOOL*'s results were not reliable at low energies
- runs performed with about 105 sample particles
- targeted various final mean kinetic energies: ~3 keV, 5 keV, 10 keV
- chose number of foils and voltage drop to match target energy and to reach equilibrium



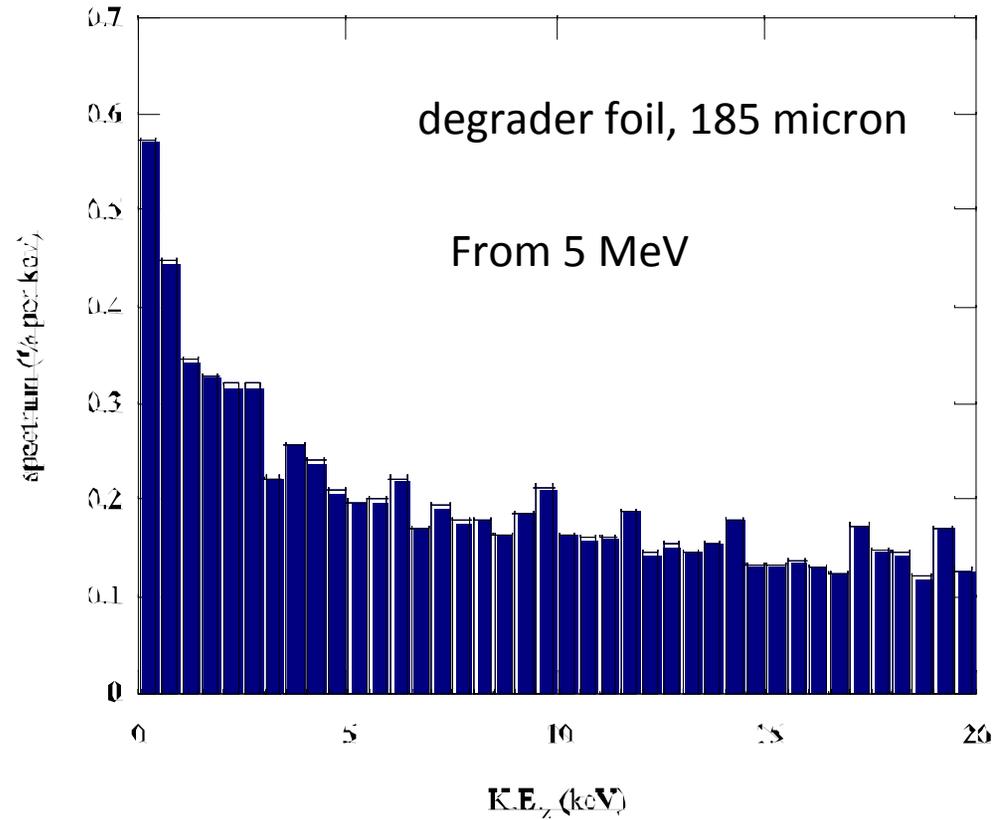
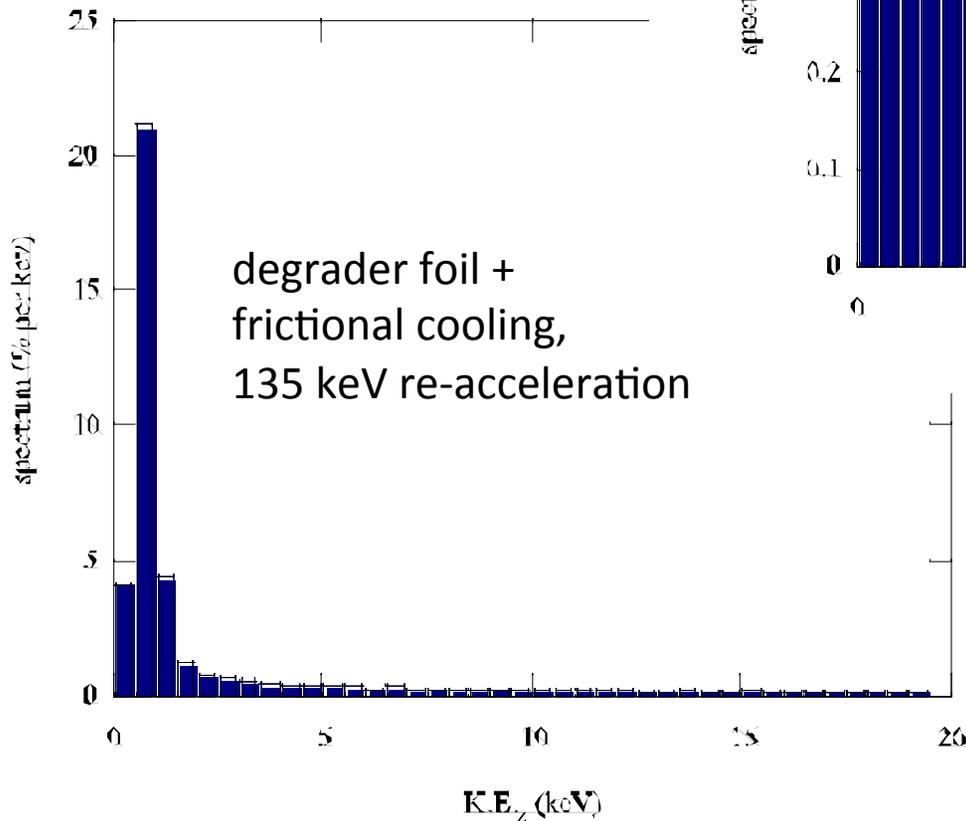
AD parameters

fraction of particles < 3 keV:

1.1% degrader only

16% frictional cooling

factor 14 improvement



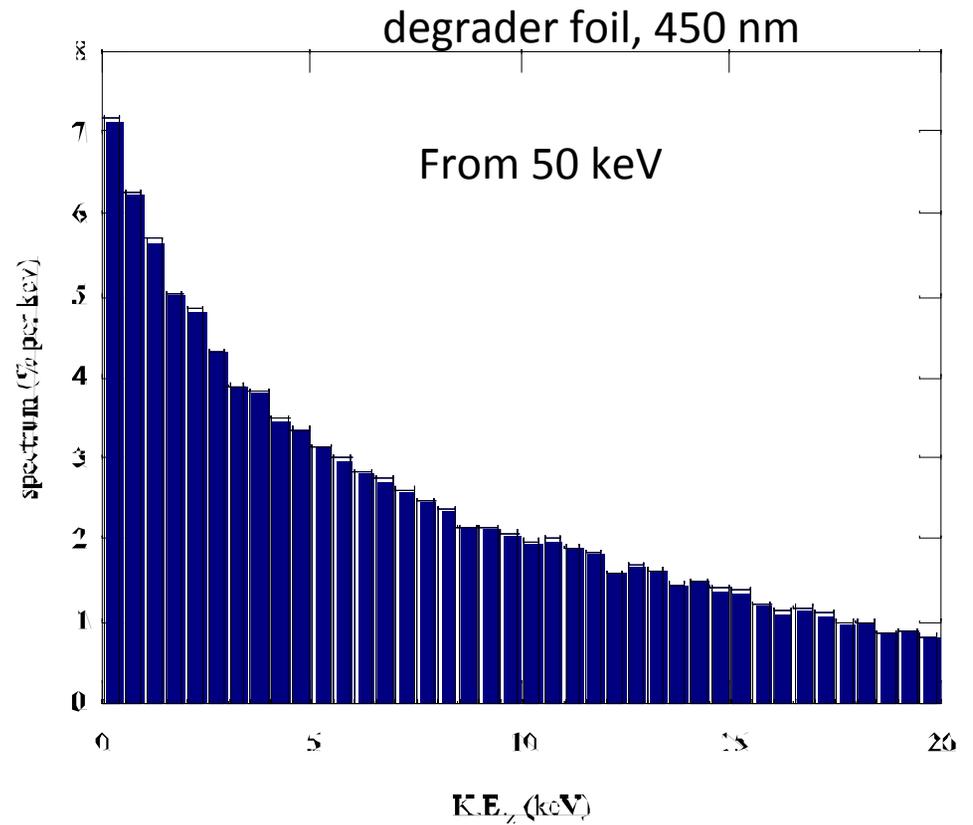
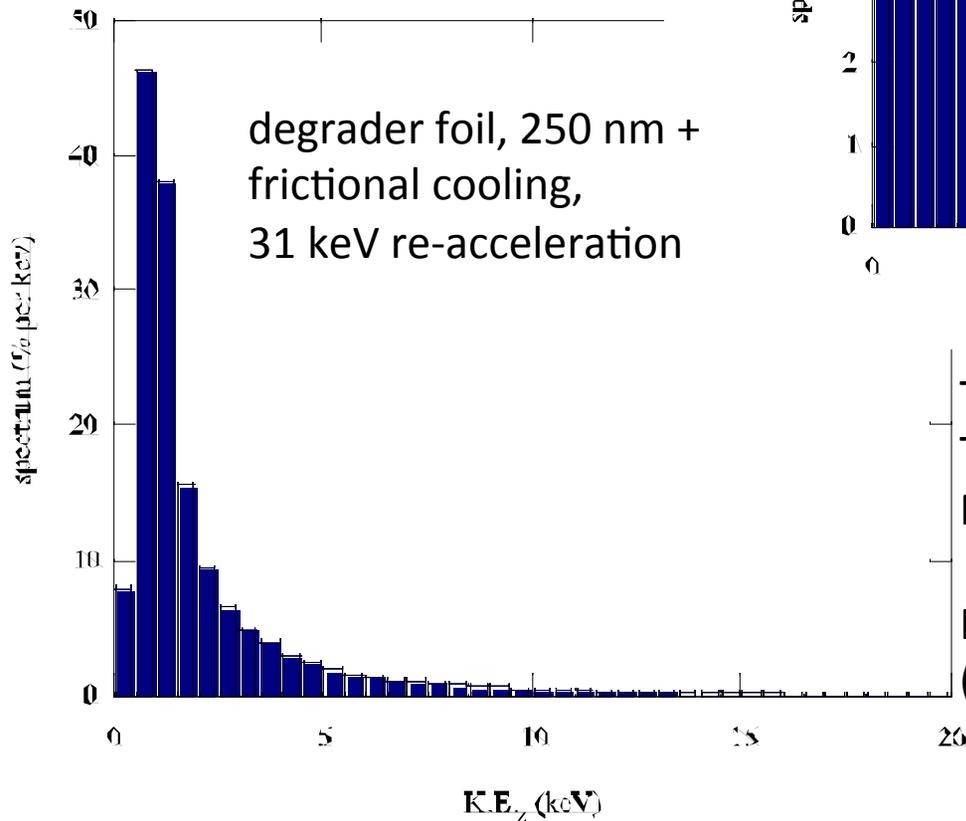
induction linac:

fraction of particles < 3 keV:

16% degrader only

62% frictional cooling

factor 4 improvement



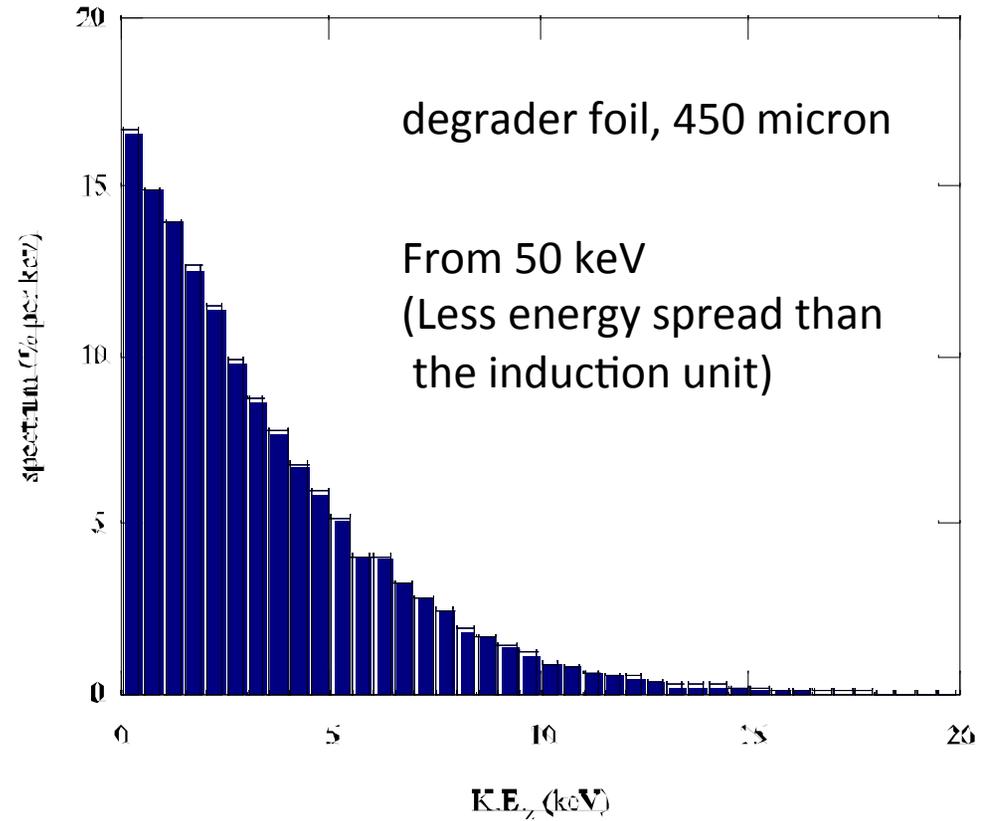
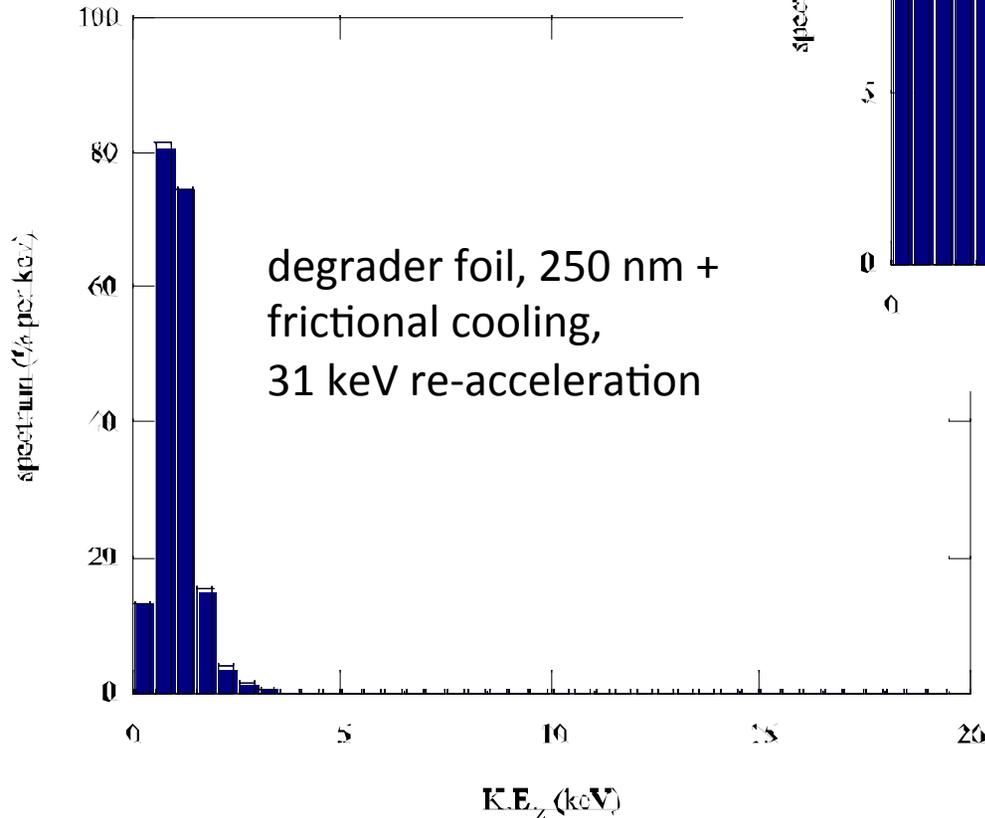
The induction unit runs at 1 MeV/m for
The full AD pulse of (about) 200 ns.
Need about 5 m

It can run at 5 MeV/m for (about) 50 ns
(or 1/4 of the AD pulse)



RFQ:

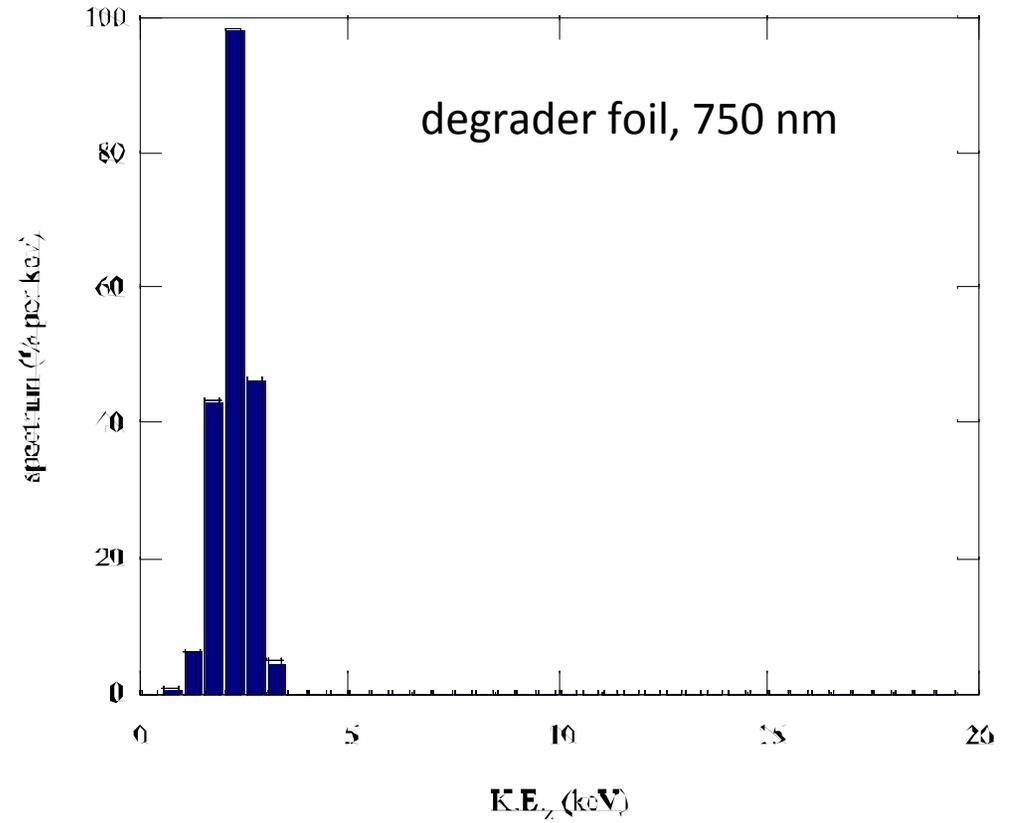
fraction of particles < 3 keV:
40% degrader only
95% frictional cooling
factor 2.5 improvement



Elena parameters

fraction of particles < 3 keV:
98% degrader only

No frictional cooling needed



Conclusions

A frictional cooling section can be easily made (fast and inexpensive) and it would significantly increase the flux to the anti-hydrogen experiments in the interim while ELENA is under construction and commissioning.

A frictional cooling section, in its own right, would be interesting; that is, it brings in new physics which might be of future importance and, furthermore, is a real-world application of the frictional cooling concept (which has been shown in-principle, but not yet in a practical device).



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Thank you for your attention!!

Any questions?

(They will be answered by Max. Your choice as to whether in Ukrainian or Russian)

