

the full-scale prototype are shown in Figs. 4 and 5 respectively. The highest gradient reached by this series is 252 T/m.

The MQXB has been developed by Fermilab[3], who is also responsible for assembling both the KEK quadrupoles and their own into the inner triplet cryostats, as shown in Fig. 6. The quench training for five 1.8 m models and the full-scale prototype are shown in Fig. 7; the maximum gradient reached was 249 T/m.

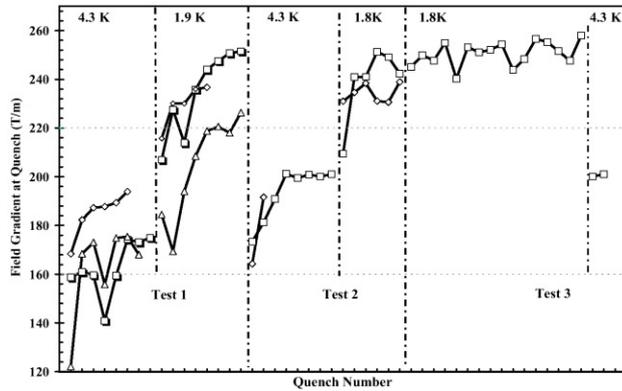


Figure 2: Quench training history of three single aperture MQY model quadrupoles.



Figure 3: Full-scale prototype MQXB quadrupole at KEK.

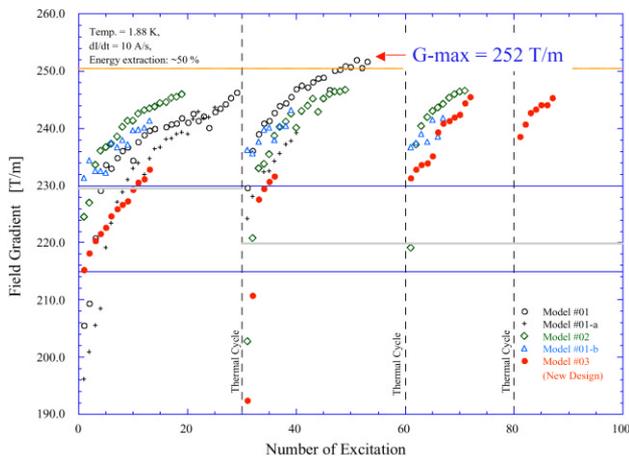


Figure 4: Quench training history of three MQXB model quadrupoles

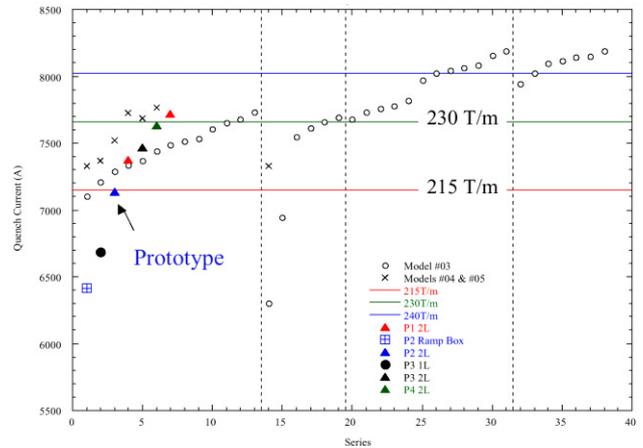


Figure 5: Quench training history of the full scale prototype MQXB quadrupole compared with the model magnets.



Figure 6: Full-scale prototype MQXB quadrupole being prepared for cryostat insertion at Fermilab.

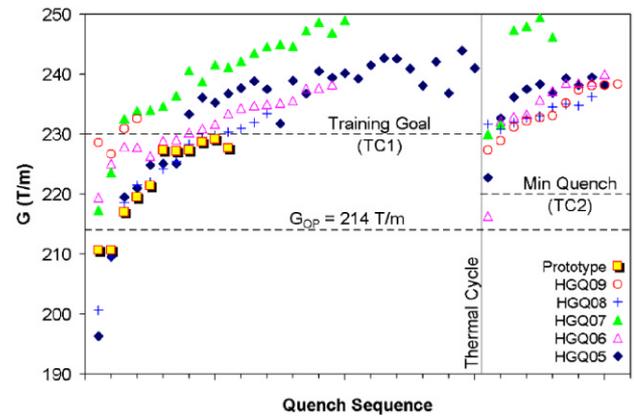


Figure 7: Quench training history of five model and one full-scale prototype MQXB quadrupole.

The LHC arc quadrupole, Fig. 8, is a 56 mm aperture two-in-one magnet, optimized for large-scale production, which has been developed by Saclay and CERN[4]. It is mated to correction coils, beam instrumentation and a cryogenic service module to form a short straight section (SSS). See Fig. 9. The three pre-production quadrupoles easily reach the nominal gradient of 223 T/m, and, as shown in Fig. 10, reach 240-250 T/m after only a few training quenches.

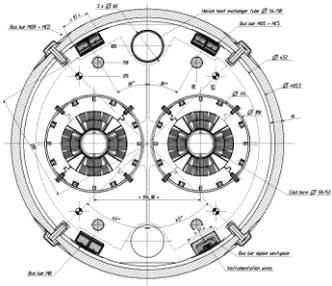


Figure 8: LHC arc quadrupole magnet (MQ).



Figure 9: Pre-production LHC short straight section.

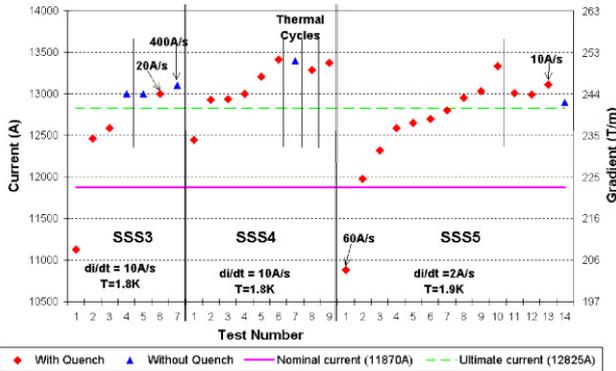


Figure 10: Quench training history of three pre-production LHC arc quadrupoles.

3 THE NEXT GENERATION

The quadrupoles described above have essentially reached the limit of what is possible with NbTi conductor. The peak field in the conductor under operating conditions is 8 T or more, and the short sample limits approach 10 T. They require superfluid helium cooling to achieve the high gradients demanded by LHC. Higher gradients will require the use of higher performance conductor, such as Nb₃Sn. Figure 11[5] shows the substantial progress in the development of this material over the past 17 years. Current densities at 12 T can now be achieved that are comparable to those reached at half that field in NbTi.

An example of a Nb₃Sn based quadrupole, which is proposed for the final focus system for TESLA, is shown in Fig. 12[6]. It uses the same coil package design as for the LHC, but replaces the NbTi cable with Nb₃Sn. Even without an iron yoke, and when imbedded in the 4 T field of the experimental solenoid, it is designed to

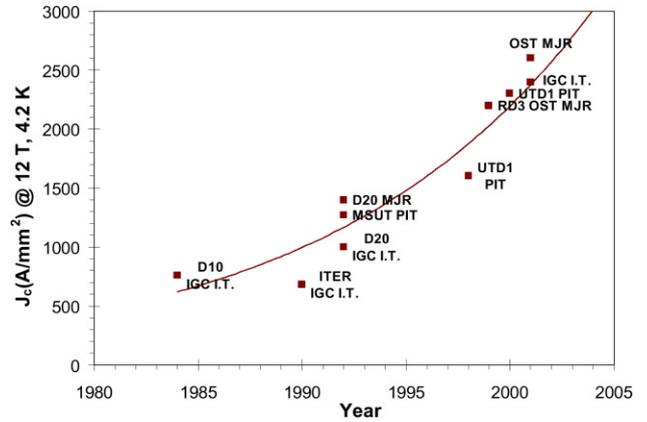


Figure 11: Progress in Nb₃Sn critical current density.

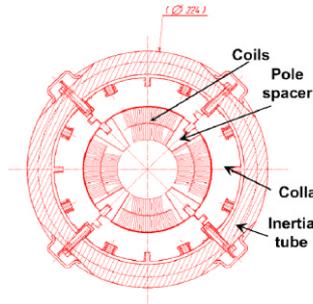
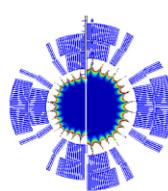


Figure 12: Candidate design for TESLA final focus quadrupole.

operate at 250 T/m. It requires $J_c > 1800 \text{ A/mm}^2$, well within the range available today.

Another proposed application for Nb₃Sn quadrupoles is in second generation LHC interaction regions. It is expected that the inner triplets currently under construction at Fermilab and KEK will eventually limit the luminosity of the LHC and will need to be replaced as part of the natural evolution of the machine. The simplest upgrade would be to rebuild the existing inner triplets either with higher gradient or larger aperture quadrupoles. A study presented at this conference concludes[7] that the latter is preferred, as it allows a more substantial decrease in β^* . Figure 13 shows the two candidate quadrupole designs considered. These are two-layer coils similar to that of the MQXB, but with a smaller coil to yoke spacing. Figure 14 shows how the short sample gradient depends on conductor performance. For example, if



	MQXB	IRQ70	IRQ90
Superconductor	NbTi	Nb ₃ Sn	Nb ₃ Sn
Bore diameter, mm	70	70	90
Yoke ID, mm	185.12	160.00	180.00
Number of turns	120	120	148
L, mH/m	3.5	3.17	5.07
G/I @ I _{nom} , T/m/kA	18.09	17.43	15.02
G _{nom} , T/m	205	275	205
I _{nom} , kA	11.33	15.78	13.65

Figure 13: Candidate designs for 2nd generation LHC IR quadrupoles using Nb₃Sn cable, compared with the MQXB design.

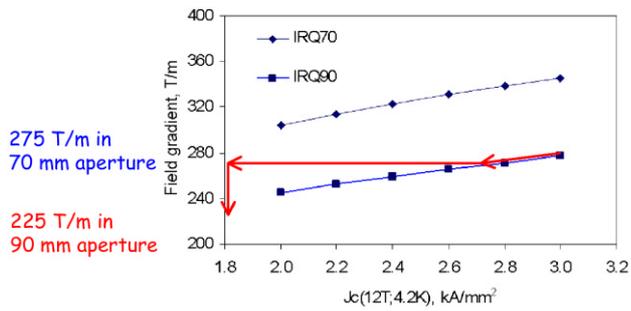


Figure 14: Operating gradients for proposed 70 mm and 90 mm aperture 2nd generation LHC IR quadrupoles.

3000 A/mm² is achieved in the wire, then allowing for 10% cabling degradation (diagonal red arrow), and 20% operating margin (vertical arrow), an operating gradient of 225 T/m can be achieved in a 90 mm aperture.

Arc quadrupoles for a very large hadron collider (VLHC) with energy above 100 TeV[5] will also require the use of Nb₃Sn. Examples of two designs for a $\sqrt{s} = 200$ TeV VLHC are shown in Fig. 15[8]. Both operate at 400 T/m with a 43.5 mm aperture. One design, with horizontal beam separation, is focusing for one beam and defocusing for the other. As with the LHC arc quadrupoles, the two apertures are magnetically coupled. The other, with vertical separation, is designed to be focusing or defocusing for both beams, as would be natural to use with flat-beam optics (see below). In this case the two apertures are magnetically decoupled, requiring a larger yoke.

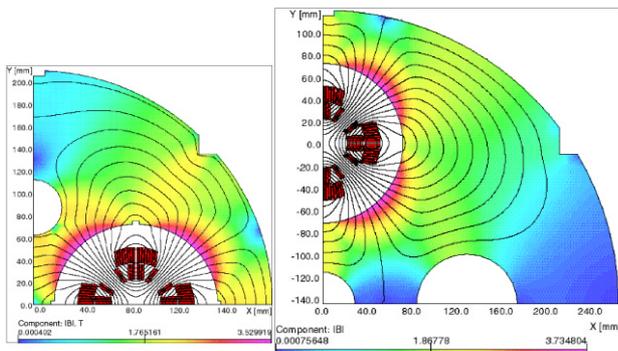


Figure 15: Two candidate designs for a $\sqrt{s} = 200$ TeV VLHC arc quadrupoles for horizontal (left) or vertical (right) beam separation.

4 BEYOND THE STATE-OF-THE-ART

The quadrupoles described in section 2 are relatively straightforward extrapolations from familiar designs used in many accelerators, with the use of Nb₃Sn being the main innovation. The gradients specified can be achieved with conductor that is currently available, or will be in the foreseeable future. The main challenge is to master the difficult technology of working with the brittle A15 material. More speculative designs have been proposed to address specific requirements of hadron collider interaction regions.

An alternate upgrade path to the simple replacement of the existing LHC inner triplets discussed above would be to add a doublet of very high gradient quadrupoles (“Q0”) between the existing triplets and the interaction point (IP). An example of the optics for such a case is shown in Fig. 16[9]. Here $\beta^* = 18$ cm is reached, with β_{\max} comparable to that in the existing LHC optics. The inner-most of the pair of Q0 magnets operates at 540 T/m with a 50 mm aperture, which yields a peak field in the conductor of over 16 T. In addition, this magnet would absorb more than 1 kW of power from the collision debris, since it would displace absorbers that exists in the current IR layout. A design concept using HTS material is sketched in Fig. 17. It is made with conductor-friendly racetrack coils in which the outer blocks are the return path of the conductors near the bore. This coil design makes use of HTS material conceivable. In addition, the active flux return results in a smaller yoke than with a conventional $\cos 2\theta$ coil design. However, at 500 mm OD and several meters length, this is not a small magnet and it would be very difficult to integrate this within the forward detectors of the LHC experiments.

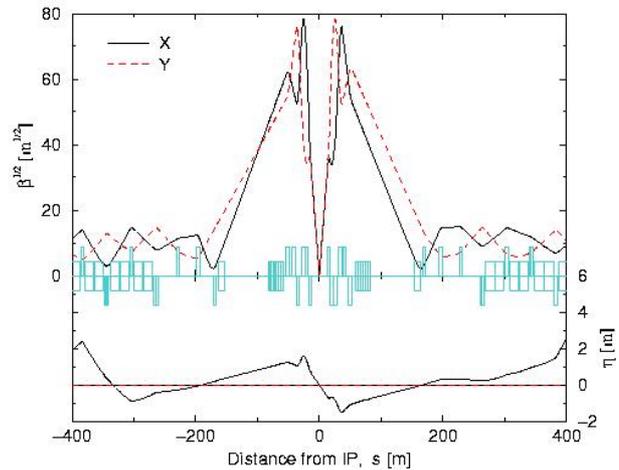


Figure 16: LHC IR optics with “Q0” doublet inserted in between the IP and the existing triplet.

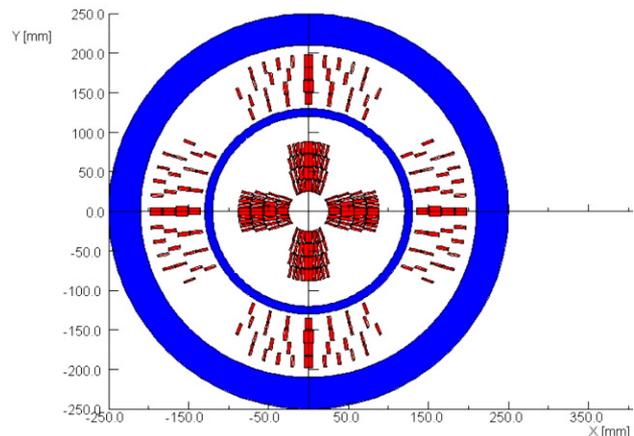


Figure 17: Design concept for an LHC Q0 quadrupole using high-temperature superconductor.

A VLHC with beam energy above about 50 TeV would be the first synchrotron radiation damped hadron machine, which opens the way to consider flat beams, as is routine in electron machines. To exploit the flat beams at the IP, the quadrupoles must be vertically focusing for both beams on both sides of the IP. This, in turn, requires the use of twin-aperture final focus quadrupoles. To minimize β_{\max} , the quadrupoles must be as close as possible to the IP and therefore have as small as possible an aperture spacing. Optics for such a case is shown in Fig. 18[5]. The final focus quadrupole doublet is placed between the two beam separation dipoles, and the spacing between apertures is only about 70 mm. Here Nb₃Sn or HTS material will be required not only due to the high field, but also due to the very high heating from the collision debris (10's of kW). A design concept for such a quadrupole, based on Nb₃Sn racetrack coils and operating at 400 T/m, is shown in Fig 19[5]. This "design" is only an idea at this stage, and it is clear that it will be very difficult to achieve good field quality and the high gradient required, while containing the enormous forces on the conductor and providing adequate cooling. This illustrates the great challenges that remain if accelerator based particle physics is to advance into the multi-TeV mass range.

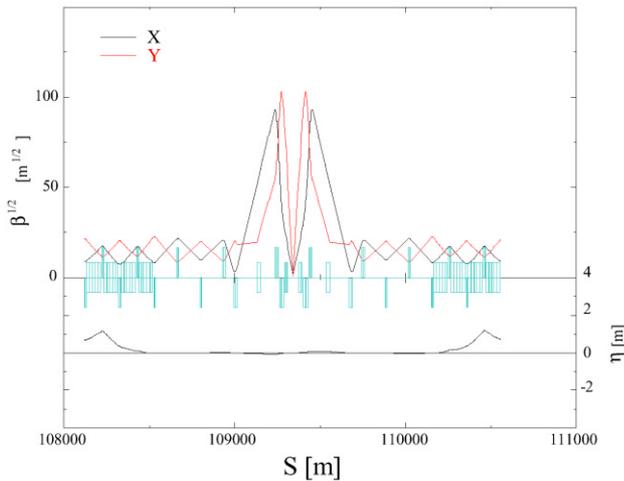


Figure 18: Flat-beam IR optics for a $\sqrt{s}=200$ TeV VLHC.

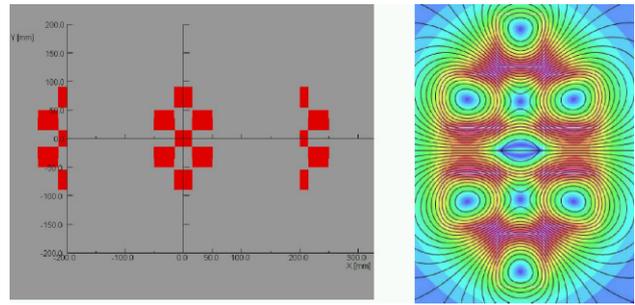


Figure 19: Possible twin-aperture "F-F" IR quadrupole for a VLHC with flat-beam optics.

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