

Towards Diffraction Limited Storage Ring Based Light Sources

Liu Lin

Brazilian Synchrotron Light Laboratory - LNLS

A wide-angle photograph of a long, multi-span bridge stretching across a body of water under a hazy, golden sky. The bridge has numerous vertical supports and a few tall, thin towers in the distance.

8th International Particle Accelerator Conference

COPENHAGEN, DENMARK, 2017 MAY 14—19

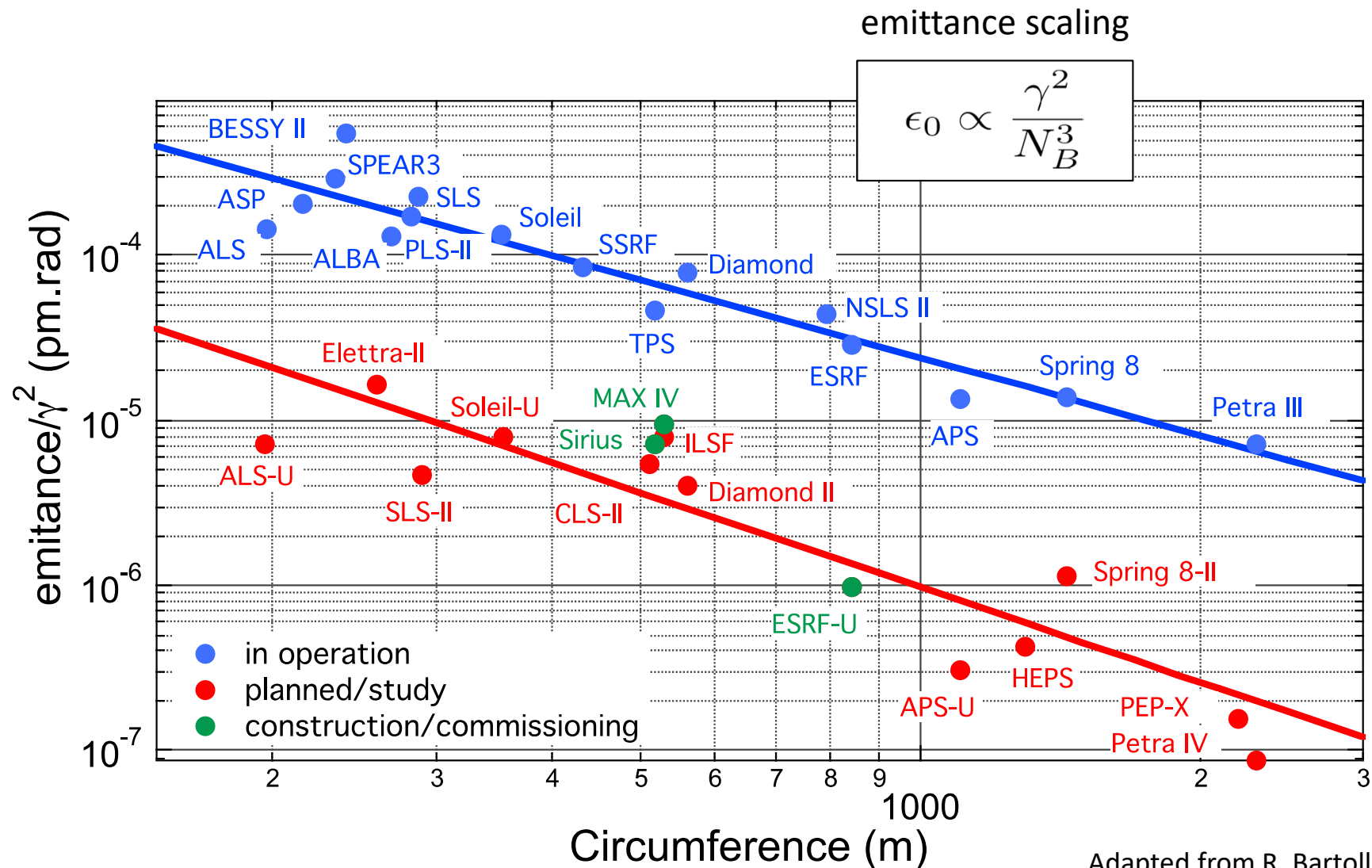
Towards Diffraction Limited Storage Ring Based Light Sources

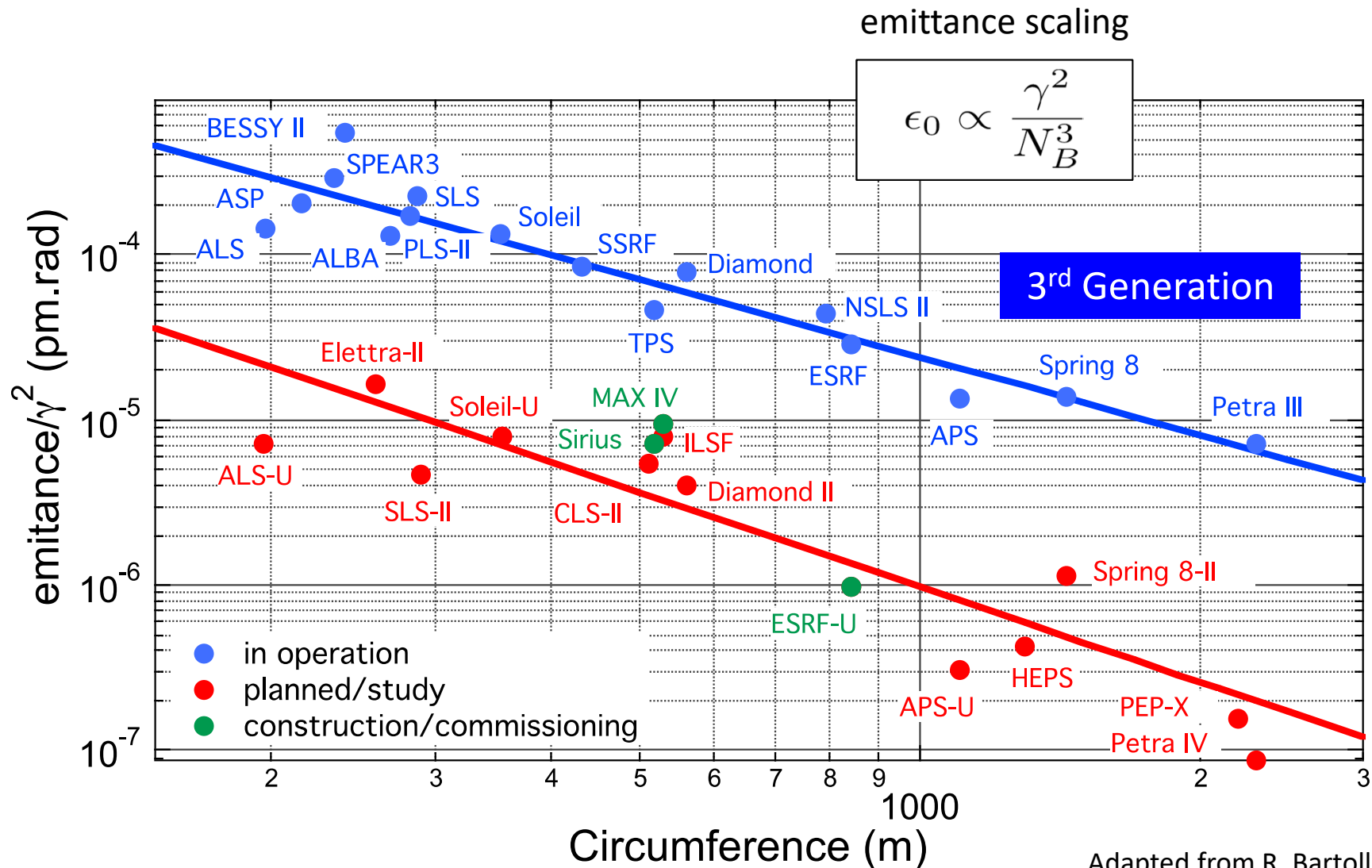
Highlights from Sirius, the Brazilian Light Source Project

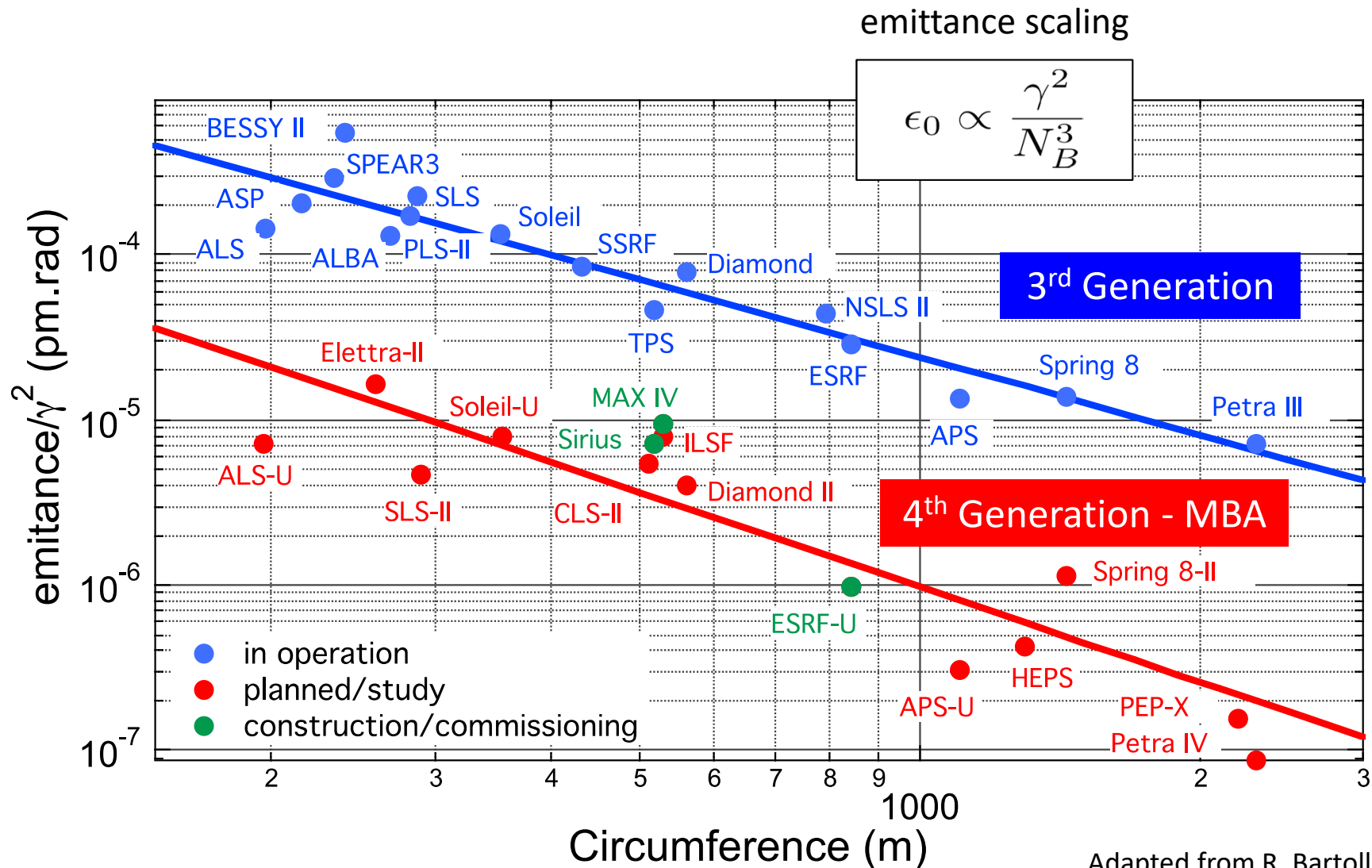
Liu Lin

Brazilian Synchrotron Light Laboratory - LNLS

- The new generation of Storage Ring Light Sources
- High brightness and coherence
 - Low emittance
 - Phase-space matching
- Sirius, the Brazilian Light Source Project
 - Lattice design highlights: low beta sections
 - Light source & beamline integration: improved solution for CARNAÚBA beamline
- Conclusion

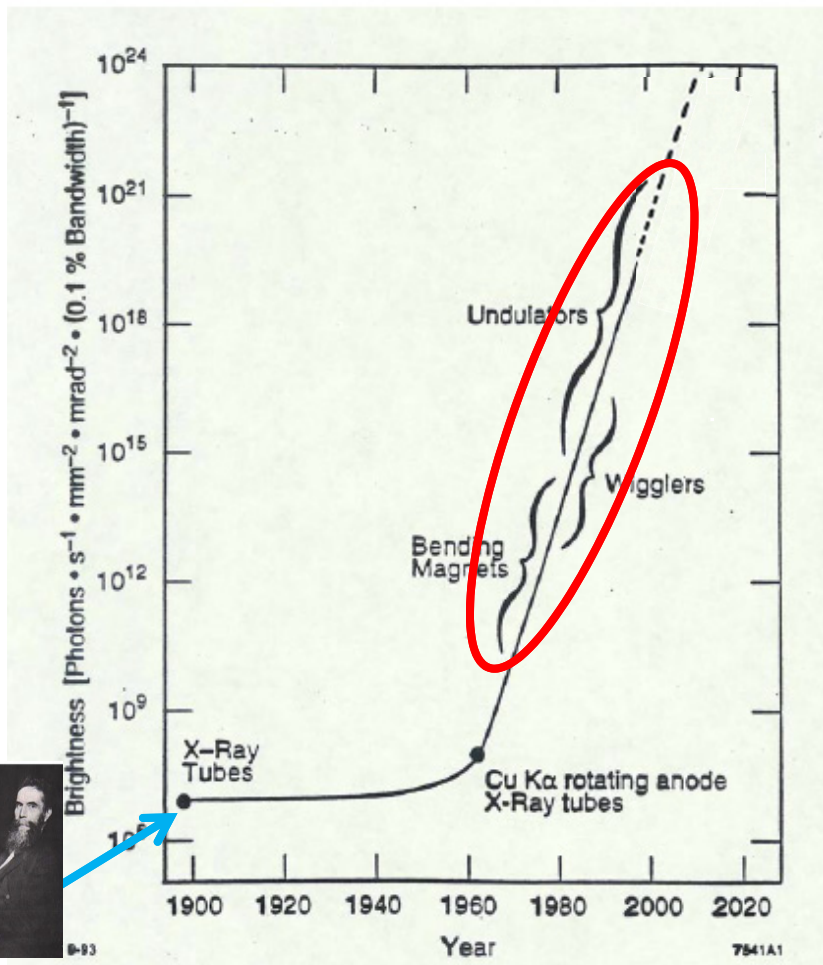






Machine	Energy [GeV]	Circum. [m]	N_B	Emit. [μm]	status
MAX-IV	3	528	140	330	operation, new
Sirius	3	518	100	250	construction, new
ESRF-U	6	844	224	135	construction, upgrade
ALS-U	2	196	108	109	planned, upgrade
APS-U	6	1104	280	42	planned, upgrade
CLS-II	3	510	147	186	planned, new
Diamond-II	3	561	144	140	planned, upgrade
Elettra-II	2	259	72	250	planned, upgrade
HEPS	6	1296	336	59	planned, new
ILSF	3	528	100	275	planned, new
PEP-X	4.5	2199		12	planned, upgrade
PETRA-IV	6	2304	504	12	planned, upgrade
SLS-II	2.4	290	84	103	planned, upgrade
Soleil-II	2.75	354	104	230	planned, upgrade
Spring8-II	6	1436	200	157	planned, upgrade

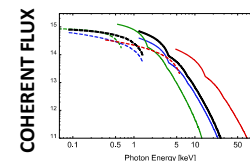
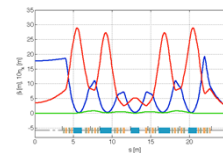
Planned machines are at different planning stages.



Herman Winick March 26, 2007

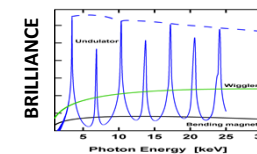
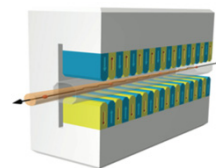
4th Generation

emittance reduction with MBA lattices, high performance IDs, **high coherent flux**



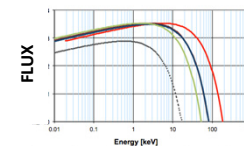
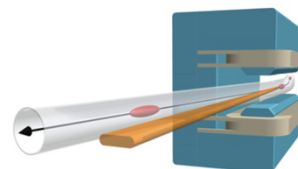
3rd Generation

DBA, TBA lattices with straight sections for wigglers and undulators, **high brilliance**



2nd Generation

dedicated sources from bending magnets, **high flux**



1st Generation

parasitic operation in colliders, bending magnets

- Spectral brilliance: Flux density in phase space

$$B(\lambda) \propto \frac{F(\lambda)}{(\epsilon_{x,e^-} \otimes \epsilon_r(\lambda)) (\epsilon_{y,e^-} \otimes \epsilon_r(\lambda))}$$

- Spectral brilliance: Flux density in phase space

$$B(\lambda) \propto \frac{F(\lambda)}{(\epsilon_{x,e^-} \otimes \epsilon_r(\lambda)) (\epsilon_{y,e^-} \otimes \epsilon_r(\lambda))}$$

← Photon flux [photons/s/0.1% bw]

- Spectral brilliance: Flux density in phase space

$$B(\lambda) \propto \frac{F(\lambda)}{(\epsilon_{x,e^-} \otimes \epsilon_r(\lambda)) (\epsilon_{y,e^-} \otimes \epsilon_r(\lambda))}$$

Photon flux [photons/s/0.1% bw]

electron beam emittance

- Spectral brilliance: Flux density in phase space

$$B(\lambda) \propto \frac{F(\lambda)}{(\epsilon_{x,e^-} \otimes \epsilon_r(\lambda)) (\epsilon_{y,e^-} \otimes \epsilon_r(\lambda))}$$

Photon flux [photons/s/0.1% bw]

electron beam emittance

photon limiting emittance

- Spectral brilliance: Flux density in phase space

$$B(\lambda) \propto \frac{F(\lambda)}{(\epsilon_{x,e^-} \otimes \epsilon_r(\lambda)) (\epsilon_{y,e^-} \otimes \epsilon_r(\lambda))}$$

← Photon flux [photons/s/0.1% bw]

electron beam
emittance

photon limiting
emittance

$\left\{ \begin{array}{l} \epsilon_r = \sigma_r \sigma_{r'} = \frac{\lambda}{4\pi} \text{ for Gaussian beam} \\ \epsilon_r = \sigma_r \sigma_{r'} \approx \frac{\lambda}{2\pi} \text{ for undulator beam} \end{array} \right.$

- Spectral brilliance: Flux density in phase space

$$B(\lambda) \propto \frac{F(\lambda)}{(\epsilon_{x,e^-} \otimes \epsilon_r(\lambda)) (\epsilon_{y,e^-} \otimes \epsilon_r(\lambda))}$$

← Photon flux [photons/s/0.1% bw]

electron beam
emittance

photon limiting
emittance

$\left\{ \begin{array}{l} \epsilon_r = \sigma_r \sigma_{r'} = \frac{\lambda}{4\pi} \text{ for Gaussian beam} \\ \epsilon_r = \sigma_r \sigma_{r'} \approx \frac{\lambda}{2\pi} \text{ for undulator beam} \end{array} \right.$

- Coherent fraction for undulator radiation

$$f_{coh} = \frac{(\lambda/2\pi)^2}{(\epsilon_{x,e^-} \otimes \epsilon_r(\lambda)) (\epsilon_{y,e^-} \otimes \epsilon_r(\lambda))}$$

- Spectral brilliance: Flux density in phase space

$$B(\lambda) \propto \frac{F(\lambda)}{(\epsilon_{x,e^-} \otimes \epsilon_r(\lambda)) (\epsilon_{y,e^-} \otimes \epsilon_r(\lambda))}$$

← Photon flux [photons/s/0.1% bw]

electron beam
emittance

photon limiting
emittance

$\left\{ \begin{array}{l} \epsilon_r = \sigma_r \sigma_{r'} = \frac{\lambda}{4\pi} \text{ for Gaussian beam} \\ \epsilon_r = \sigma_r \sigma_{r'} \approx \frac{\lambda}{2\pi} \text{ for undulator beam} \end{array} \right.$

- Coherent fraction for undulator radiation

$$f_{coh} = \frac{(\lambda/2\pi)^2}{(\epsilon_{x,e^-} \otimes \epsilon_r(\lambda)) (\epsilon_{y,e^-} \otimes \epsilon_r(\lambda))}$$

- Diffraction limited storage ring

$$\epsilon_{x,y} \approx \epsilon_r(\lambda) = \frac{\lambda}{2\pi}$$

- Spectral brilliance: Flux density in phase space

$$B(\lambda) \propto \frac{F(\lambda)}{(\epsilon_{x,e^-} \otimes \epsilon_r(\lambda)) (\epsilon_{y,e^-} \otimes \epsilon_r(\lambda))}$$

← Photon flux [photons/s/0.1% bw]

electron beam
emittance

photon limiting
emittance

for Gaussian beam

$$\epsilon_r = \sigma_r \sigma_{r'} = \frac{\lambda}{4\pi}$$

for undulator beam

$$\epsilon_r = \sigma_r \sigma_{r'} \approx \frac{\lambda}{2\pi}$$

- Coherent fraction for undulator radiation

$$f_{coh} = \frac{(\lambda/2\pi)^2}{(\epsilon_{x,e^-} \otimes \epsilon_r(\lambda)) (\epsilon_{y,e^-} \otimes \epsilon_r(\lambda))}$$

- Diffraction limited storage ring

$$\epsilon_{x,y} \approx \epsilon_r(\lambda) = \frac{\lambda}{2\pi}$$

$\epsilon_{x,y} \approx 100$ pm.rad
diffraction limit for 2 keV

$\epsilon_{x,y} \approx 20$ pm.rad
diffraction limit for 10 keV

$$\epsilon_x = C_q \frac{\gamma^2}{J_x} \frac{\oint \mathcal{H}(s) h(s)^3 ds}{\oint h(s)^2 ds}$$

$$\epsilon_x = C_q \frac{\gamma^2}{J_x} \frac{\oint \mathcal{H}(s) h(s)^3 ds}{\oint h(s)^2 ds}$$

damping partition

$$\epsilon_x = C_q \frac{\gamma^2}{J_x} \frac{\oint \mathcal{H}(s) h(s)^3 ds}{\oint h(s)^2 ds}$$

damping partition

curvature function

$$h(s) = 1/\rho(s)$$

$$\epsilon_x = C_q \frac{\gamma^2}{J_x} \frac{\oint \mathcal{H}(s) h(s)^3 ds}{\oint h(s)^2 ds}$$

damping partition

curvature function

$$h(s) = 1/\rho(s)$$

dispersion's betatron amplitude

$$\mathcal{H} = \frac{\eta^2 + (\alpha\eta + \beta\eta')^2}{\beta}$$

η : dispersion function
 β, α : Twiss functions

$$\epsilon_x = C_q \frac{\gamma^2}{J_x} \frac{\oint \mathcal{H}(s) h(s)^3 ds}{\oint h(s)^2 ds}$$

damping partition

curvature function

$$h(s) = 1/\rho(s)$$

dispersion's betatron amplitude

$$\mathcal{H} = \frac{\eta^2 + (\alpha\eta + \beta\eta')^2}{\beta}$$

η : dispersion function
 β, α : Twiss functions

Emittance depends on optics at places where radiation is emitted (dipoles).

$$\epsilon_x = C_q \frac{\gamma^2}{J_x} \frac{\oint \mathcal{H}(s) h(s)^3 ds}{\oint h(s)^2 ds}$$

damping partition

curvature function
 $h(s) = 1/\rho(s)$

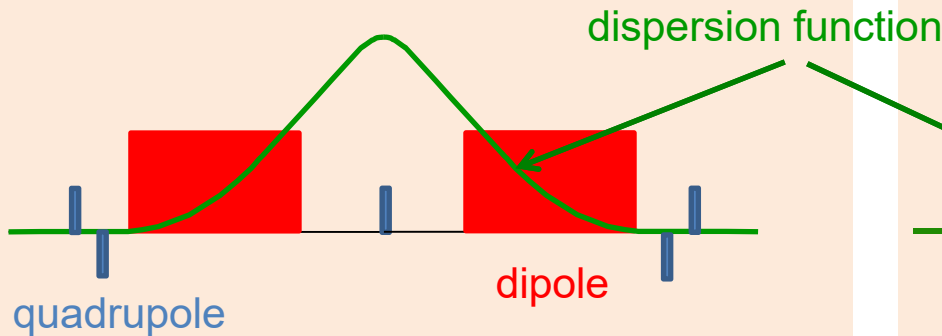
dispersion's betatron amplitude

$$\mathcal{H} = \frac{\eta^2 + (\alpha\eta + \beta\eta')^2}{\beta}$$

η : dispersion function
 β, α : Twiss functions

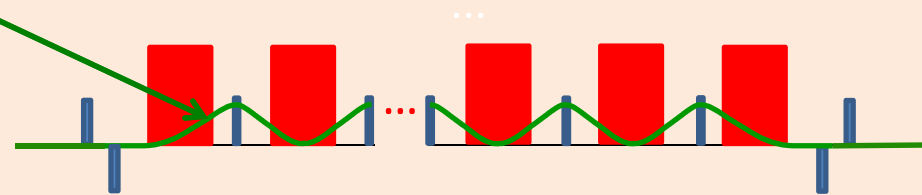
Emittance depends on optics at places where radiation is emitted (dipoles).

Double bend achromat - DBA



Multiple bend achromat - MBA

many small dipoles to keep horizontal focus in each dipole



Courtesy: Ricardo Rodrigues

MBA

Courtesy: Ricardo Rodrigues

MBA

Many short
cells

Courtesy: Ricardo Rodrigues

MBA

Many short
cells

Small beam
emittance

Courtesy: Ricardo Rodrigues

MBA

Many short
cells

Small beam
emittance

High orbit
stability

Courtesy: Ricardo Rodrigues

MBA

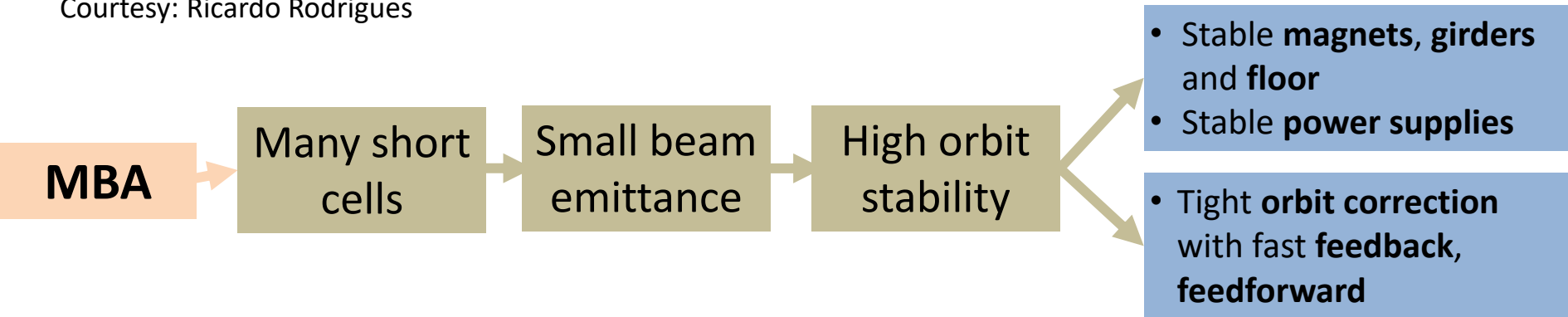
Many short
cells

Small beam
emittance

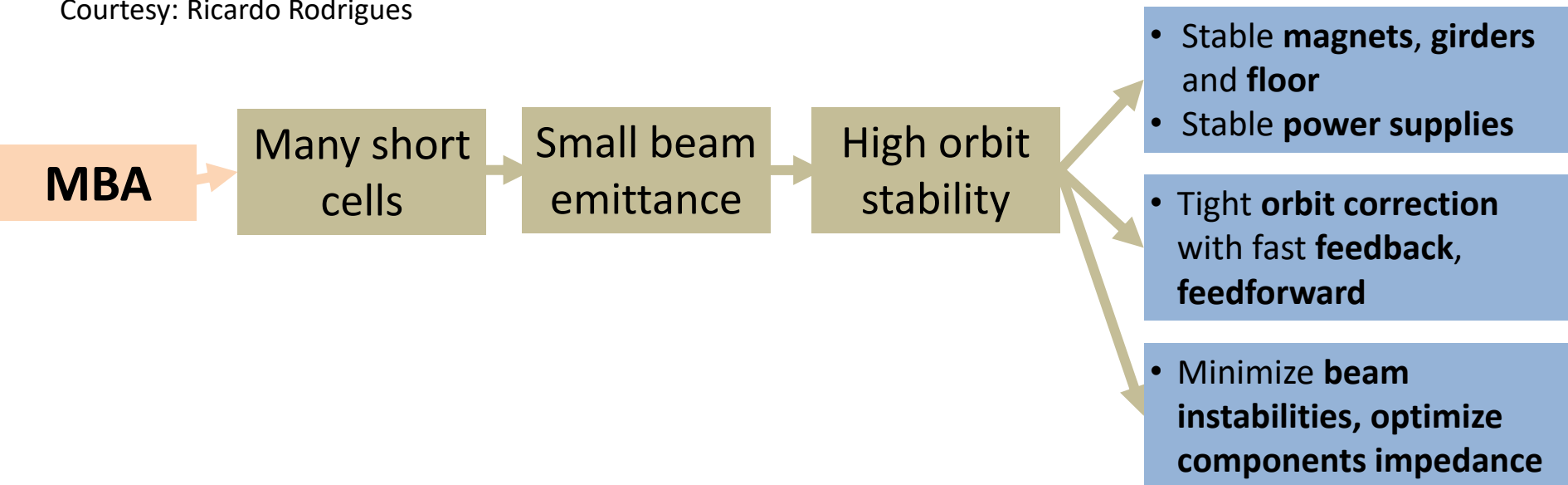
High orbit
stability

- Stable **magnets, girders and floor**
- Stable **power supplies**

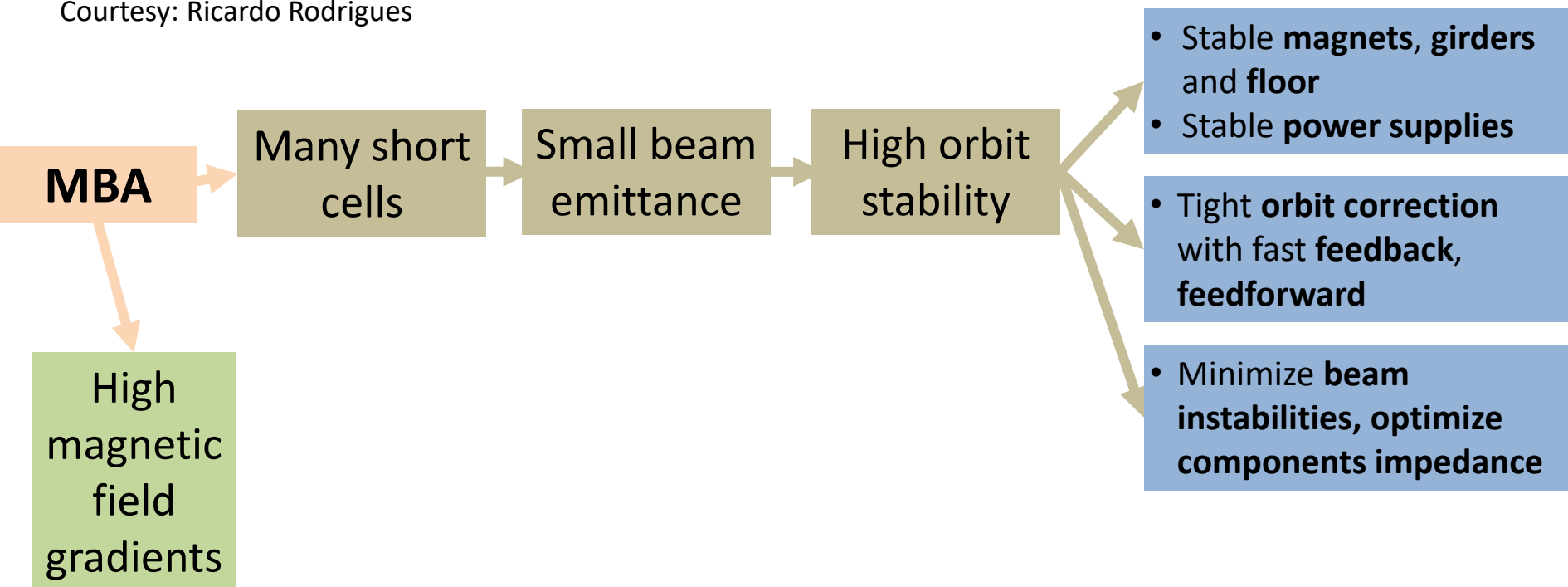
Courtesy: Ricardo Rodrigues



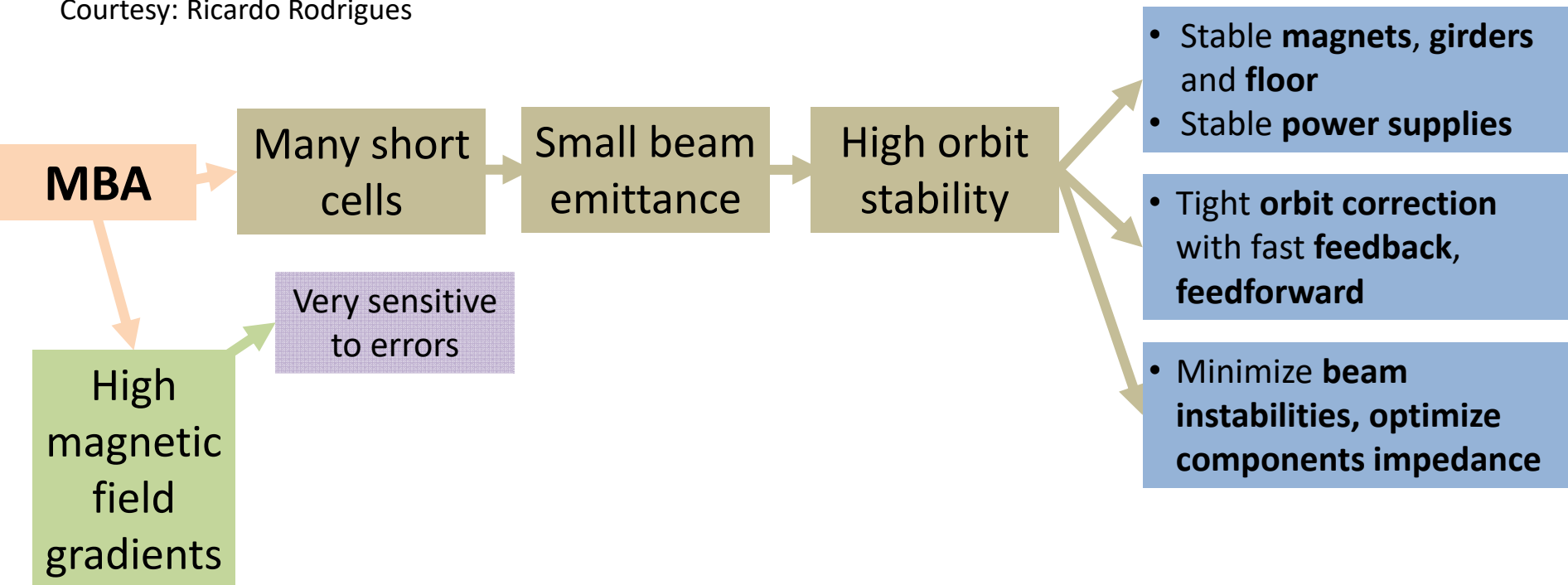
Courtesy: Ricardo Rodrigues



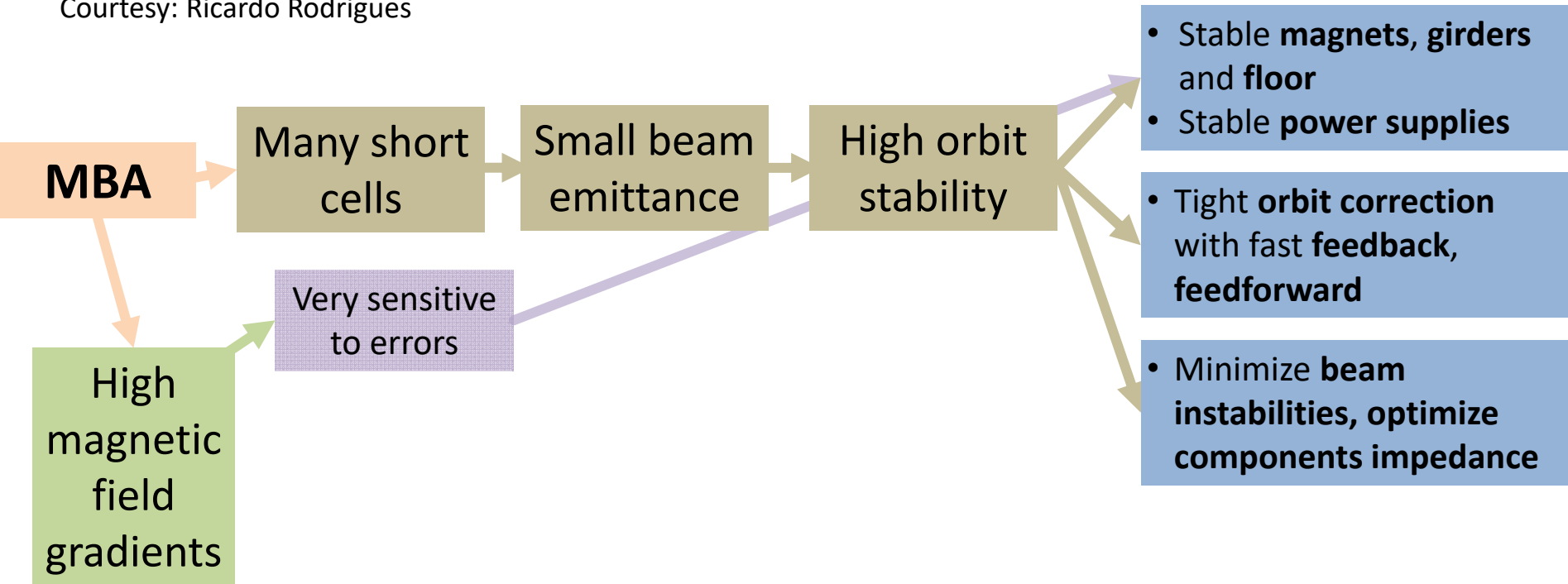
Courtesy: Ricardo Rodrigues



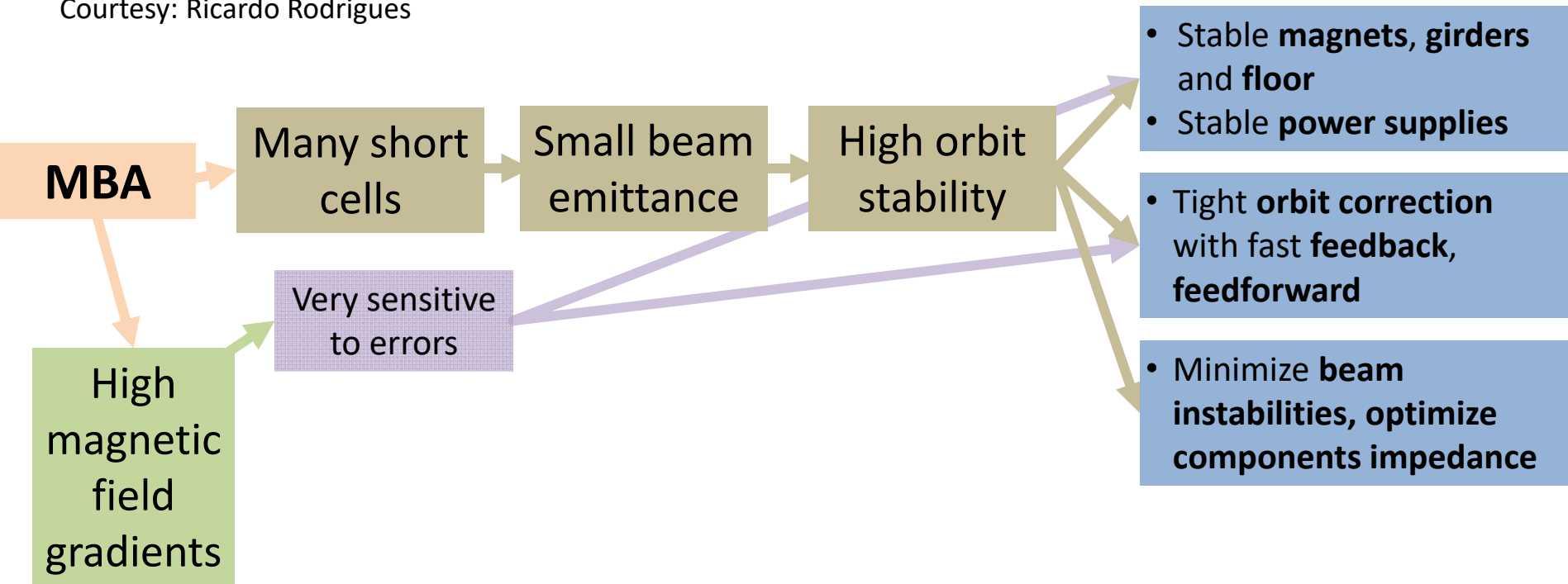
Courtesy: Ricardo Rodrigues



Courtesy: Ricardo Rodrigues



Courtesy: Ricardo Rodrigues



MBA

Many short cells

Small beam emittance

High orbit stability

- Stable **magnets, girders and floor**
- Stable **power supplies**

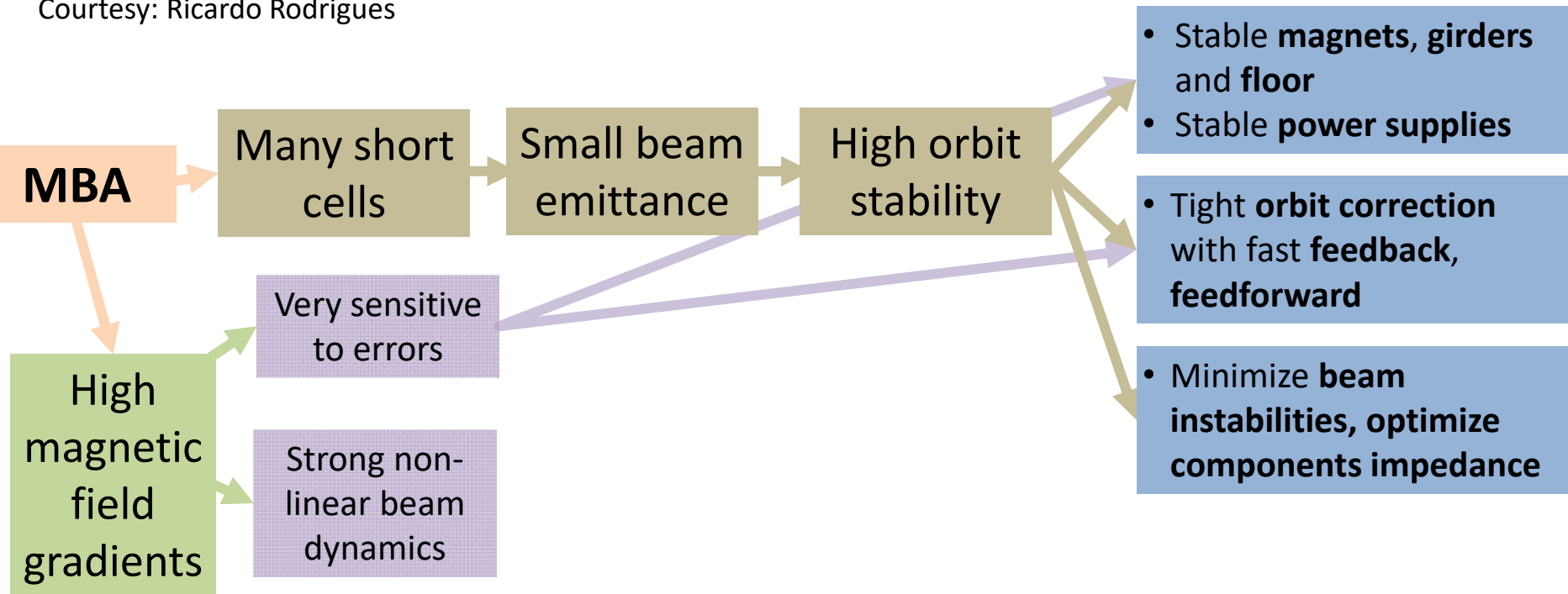
- Tight **orbit correction with fast feedback, feedforward**

- Minimize **beam instabilities, optimize components impedance**

High magnetic field gradients

Very sensitive to errors

Courtesy: Ricardo Rodrigues

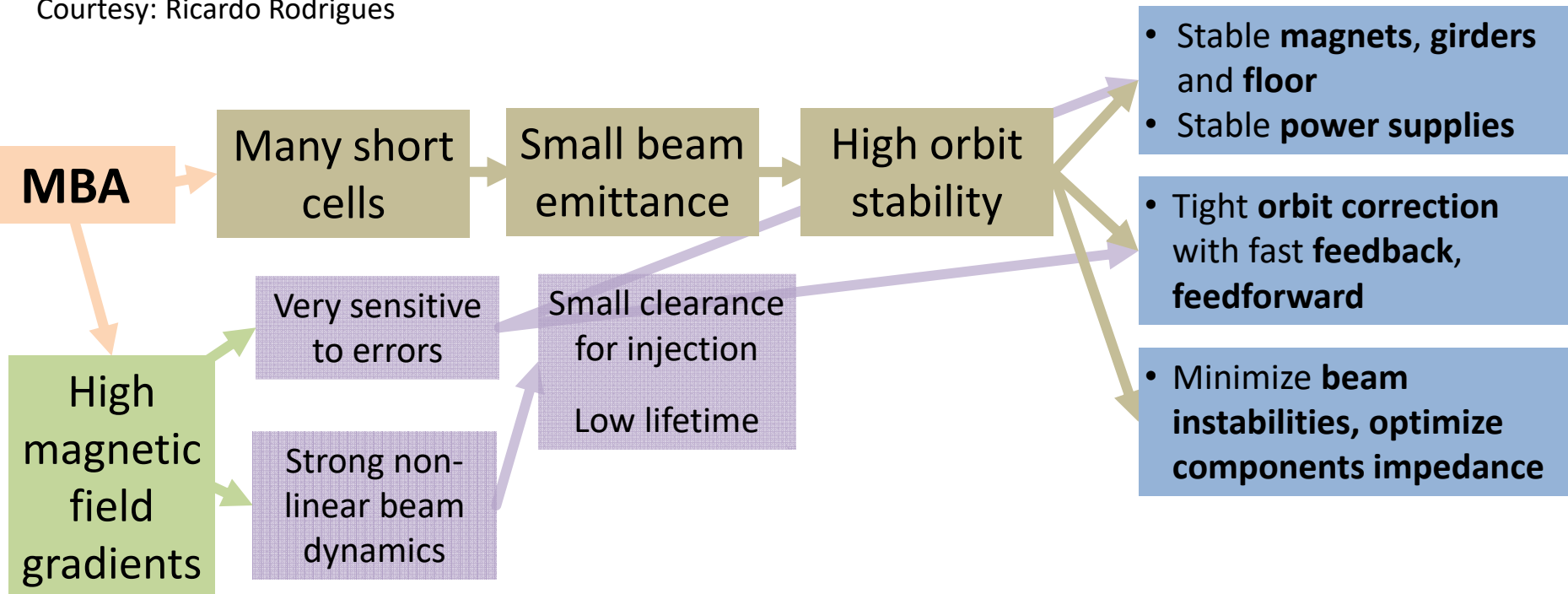


- Stable **magnets, girders and floor**
- Stable **power supplies**

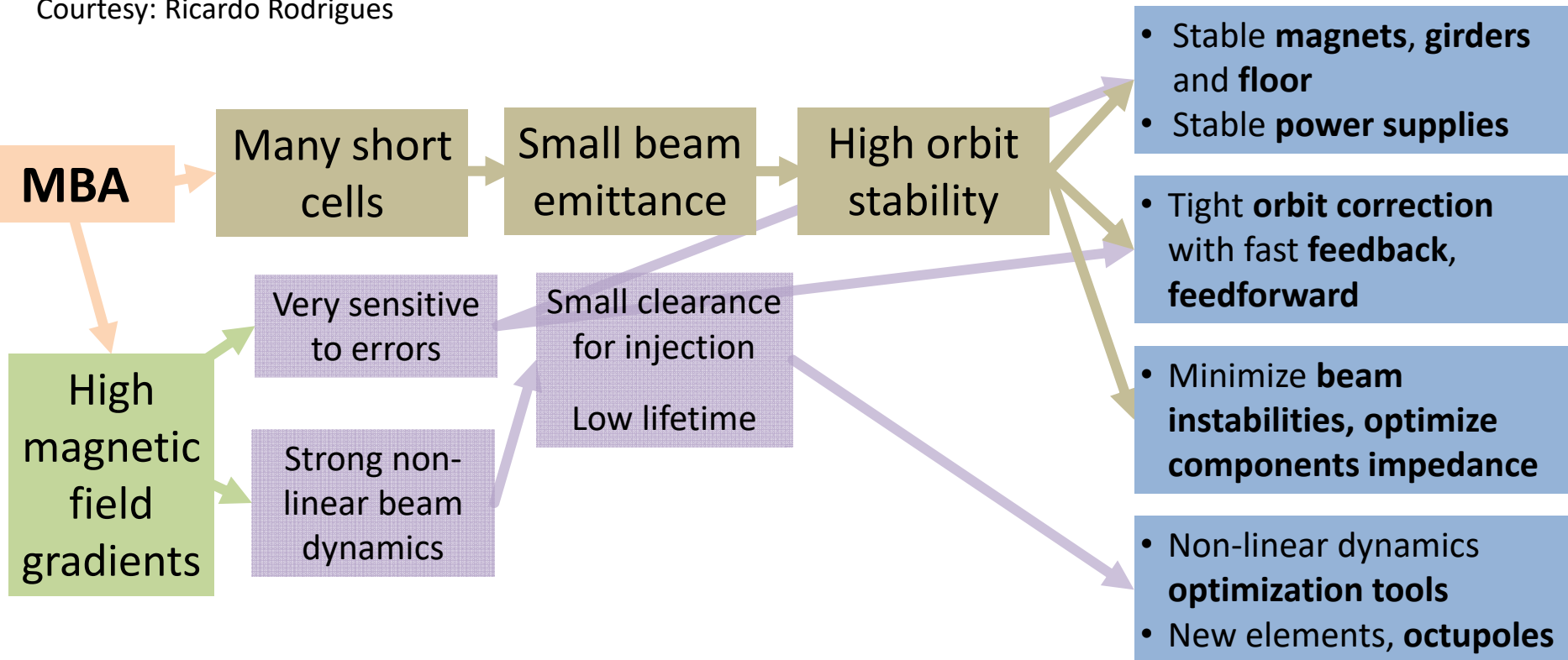
- Tight **orbit correction** with fast **feedback, feedforward**

- Minimize **beam instabilities, optimize components impedance**

Courtesy: Ricardo Rodrigues



Courtesy: Ricardo Rodrigues



MBA

Many short cells

Small beam emittance

High orbit stability

High magnetic field gradients

Very sensitive to errors

Strong non-linear beam dynamics

Small clearance for injection
Low lifetime

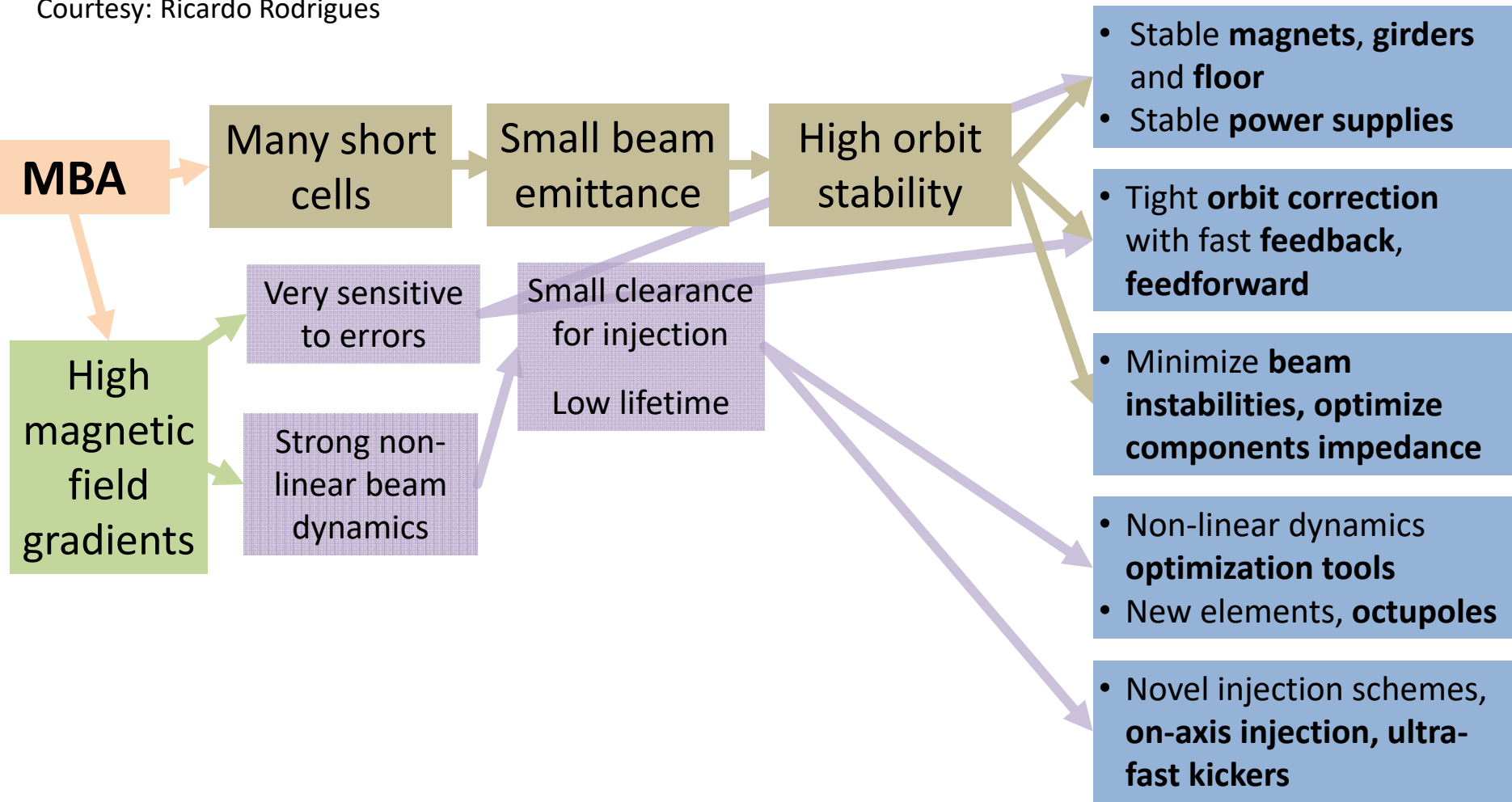
- Stable **magnets, girders and floor**
- Stable **power supplies**

- Tight **orbit correction with fast feedback, feedforward**

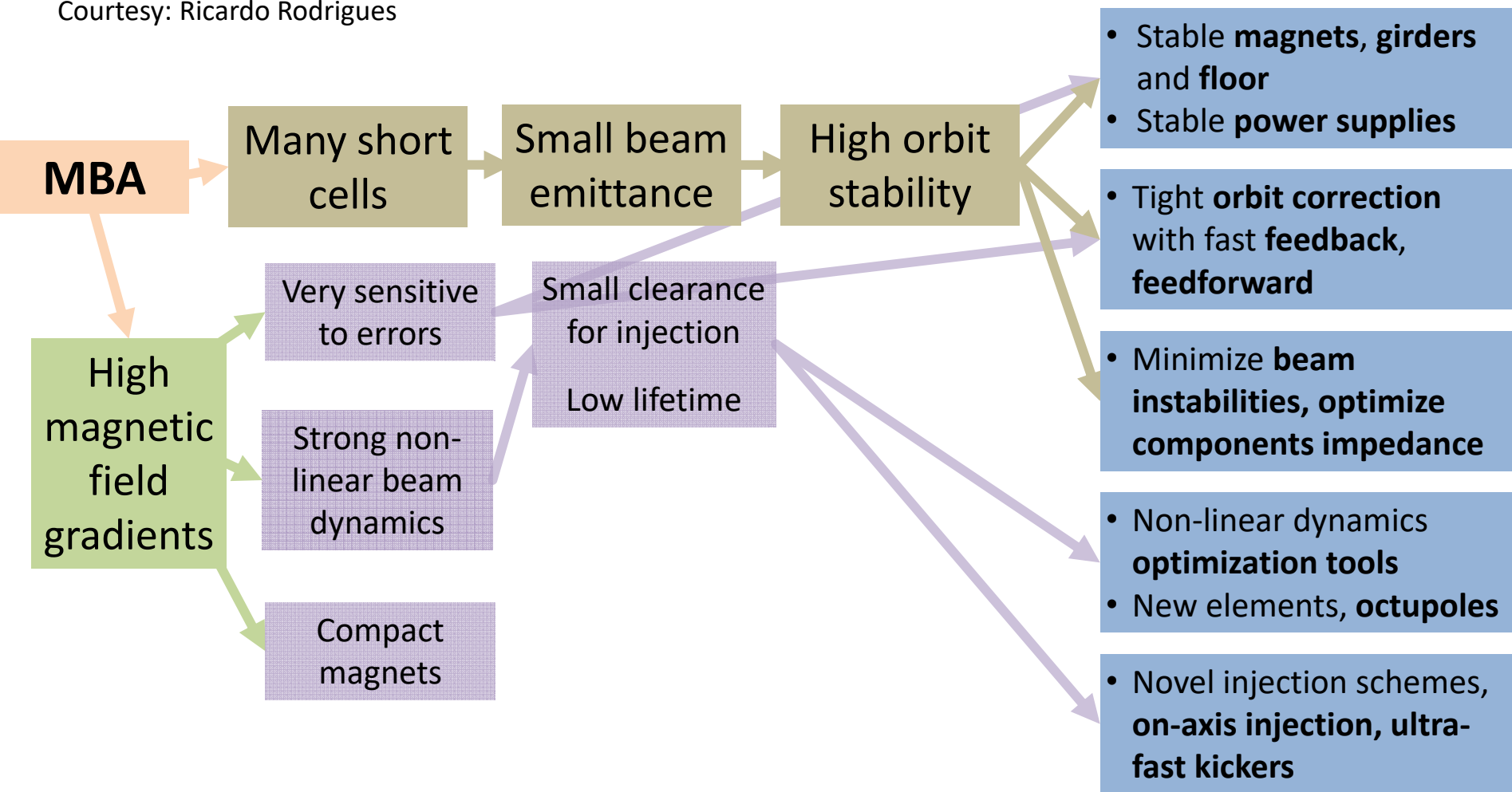
- Minimize **beam instabilities, optimize components impedance**

- Non-linear dynamics **optimization tools**
- New elements, **octupoles**

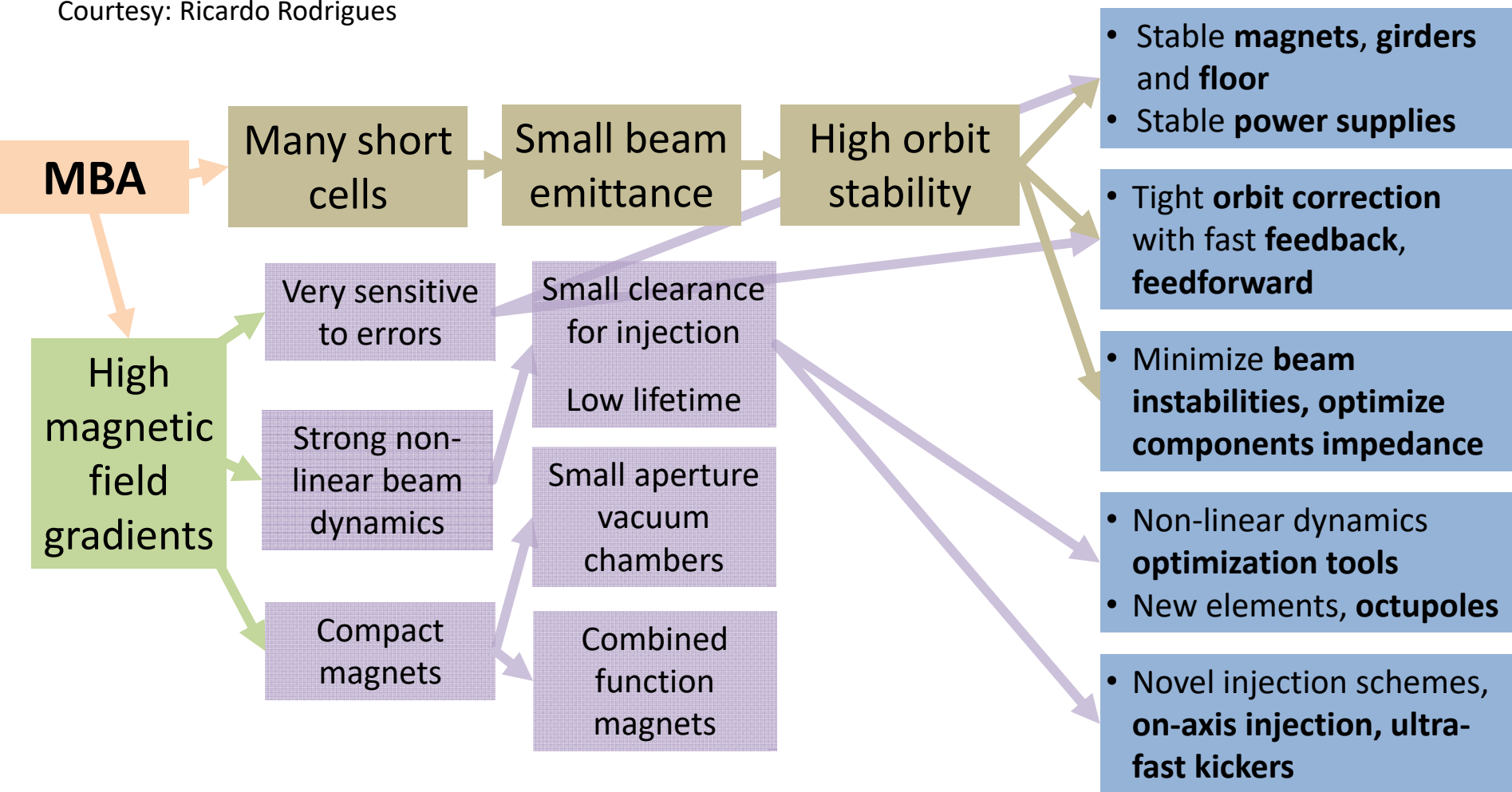
Courtesy: Ricardo Rodrigues



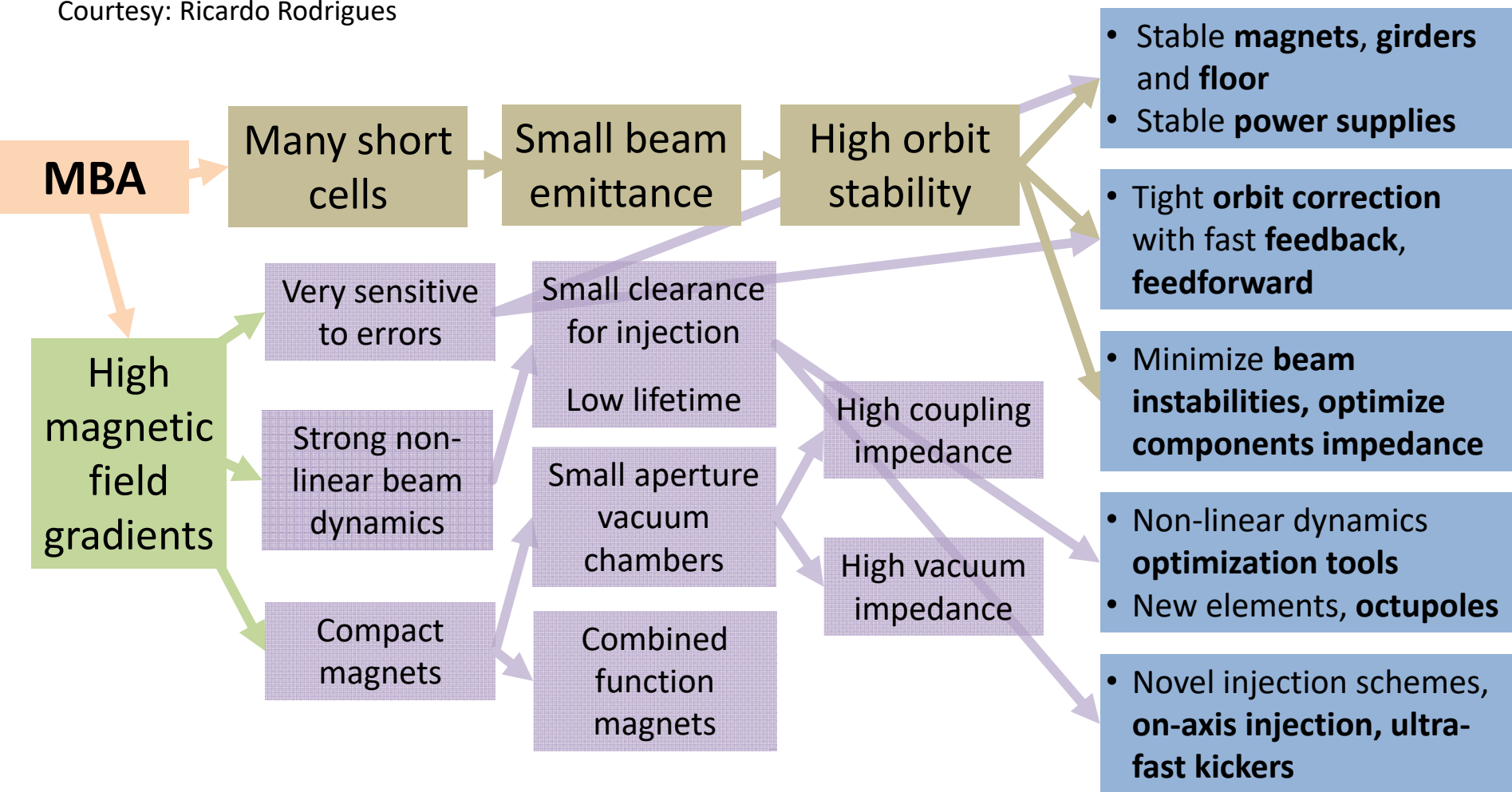
Courtesy: Ricardo Rodrigues



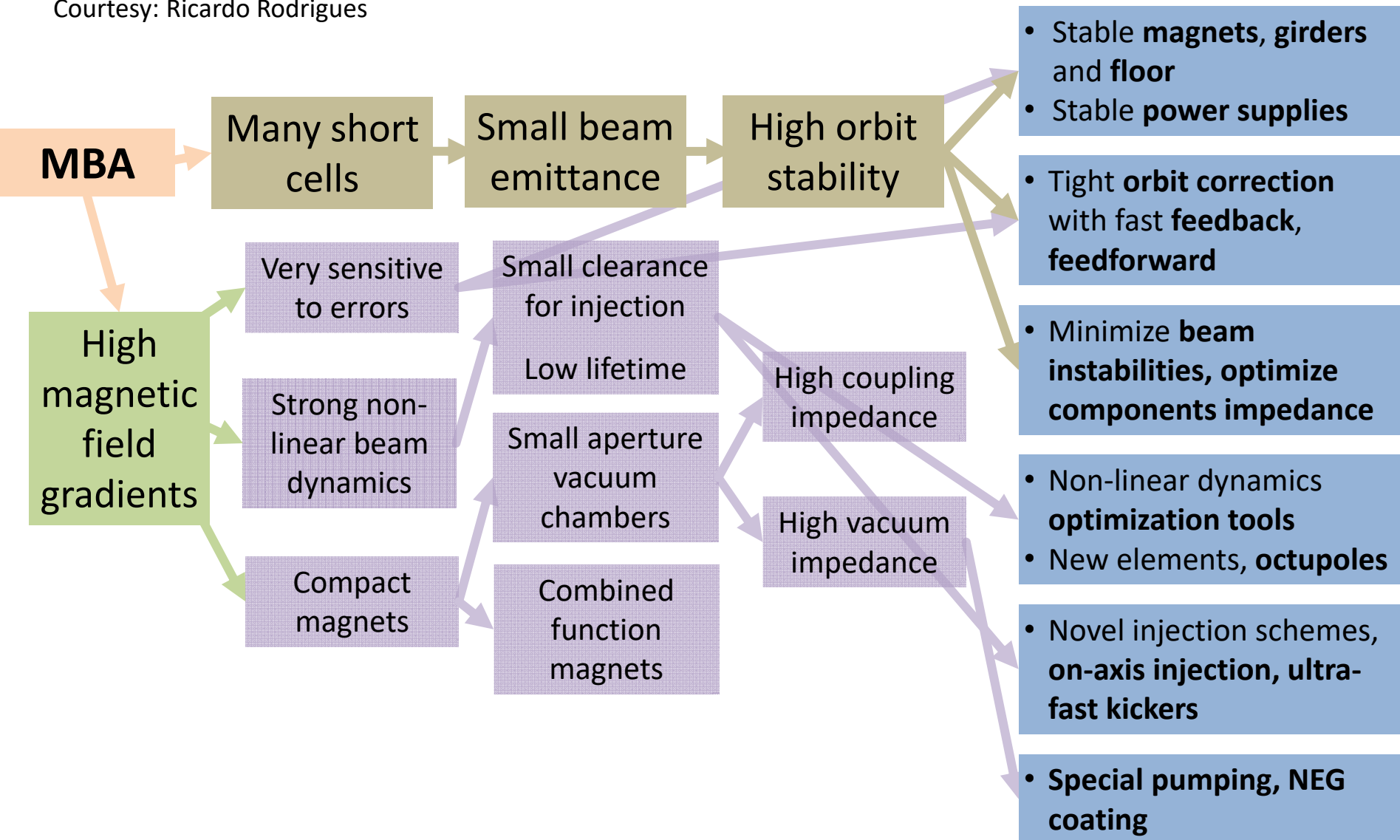
Courtesy: Ricardo Rodrigues



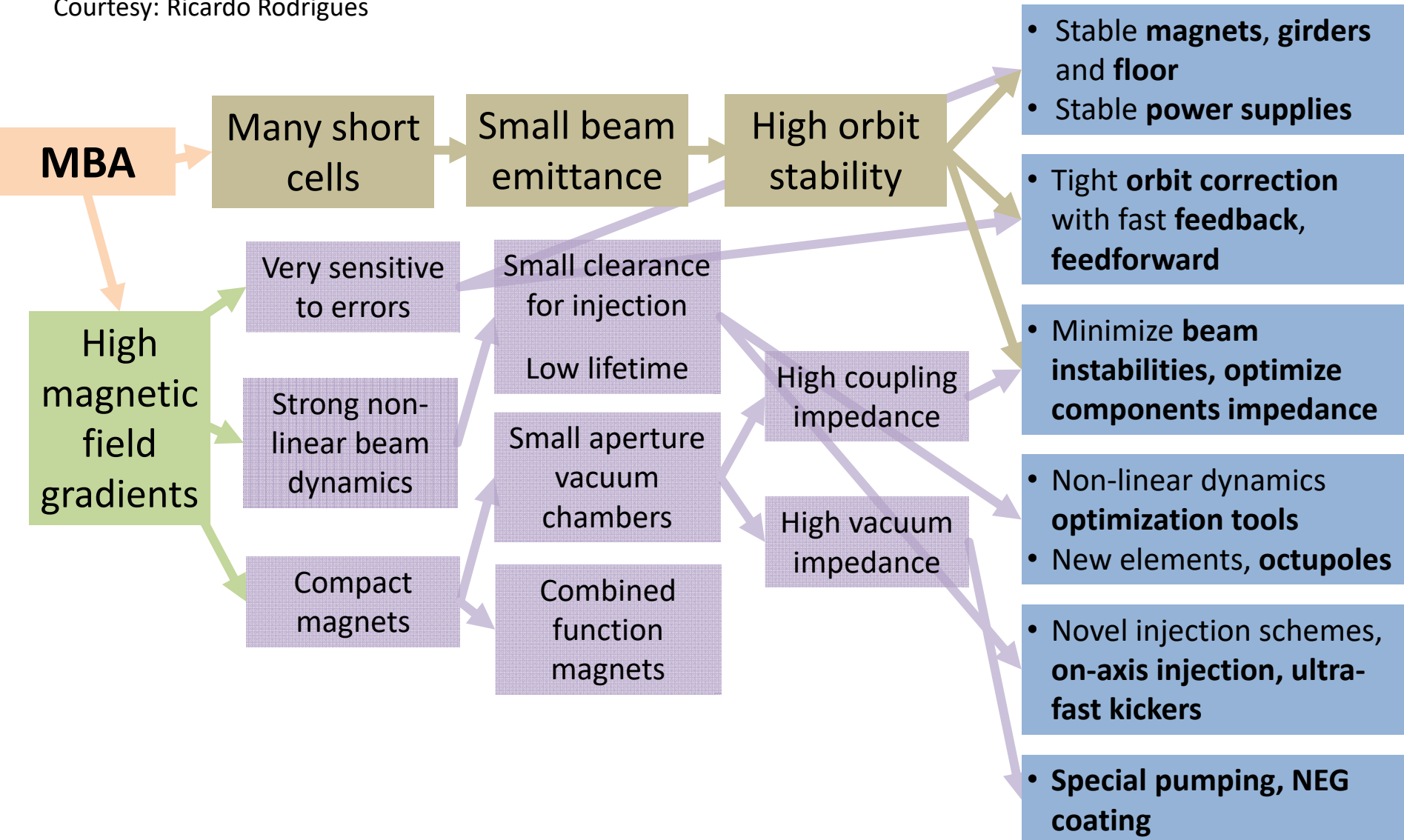
Courtesy: Ricardo Rodrigues



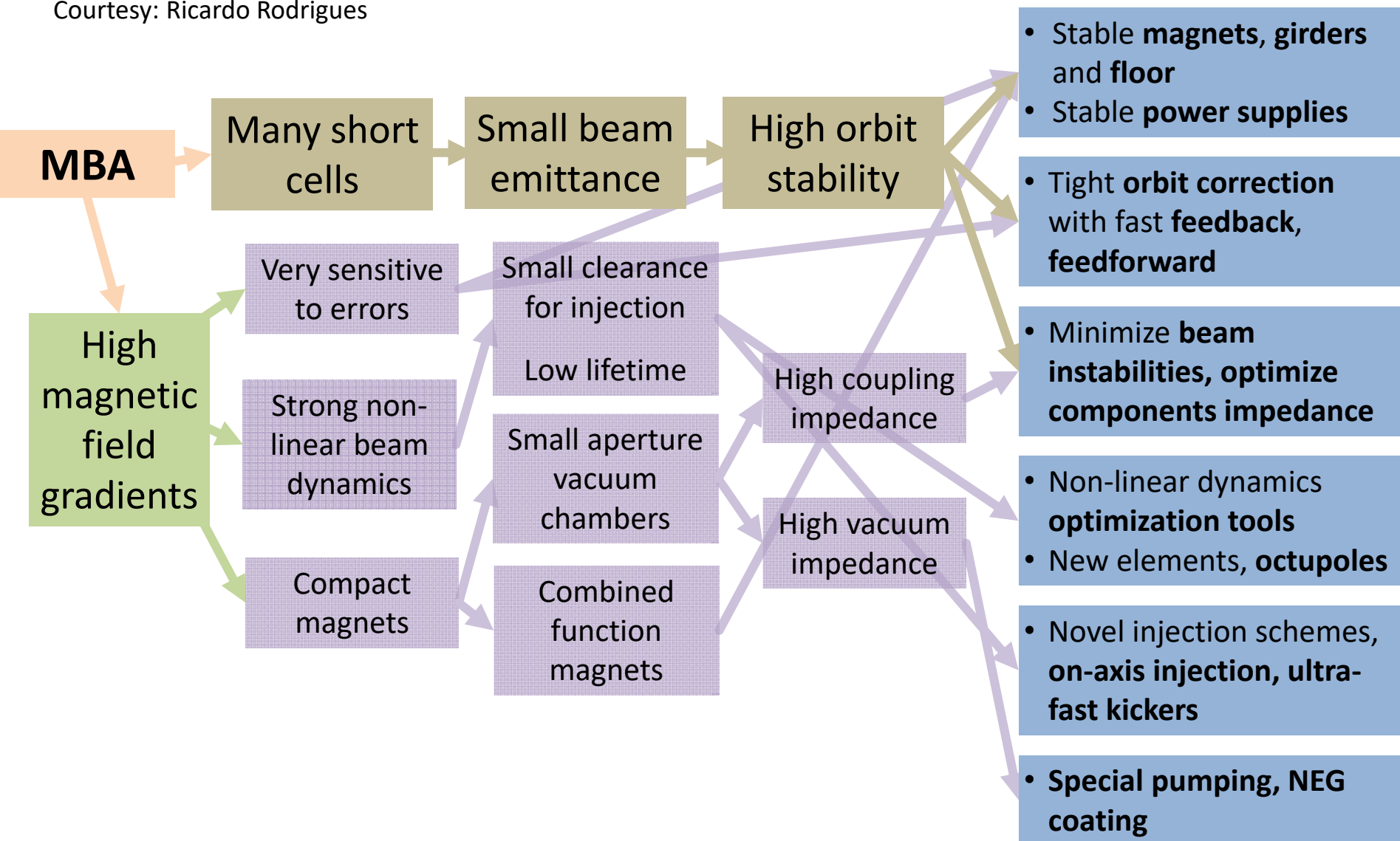
Courtesy: Ricardo Rodrigues



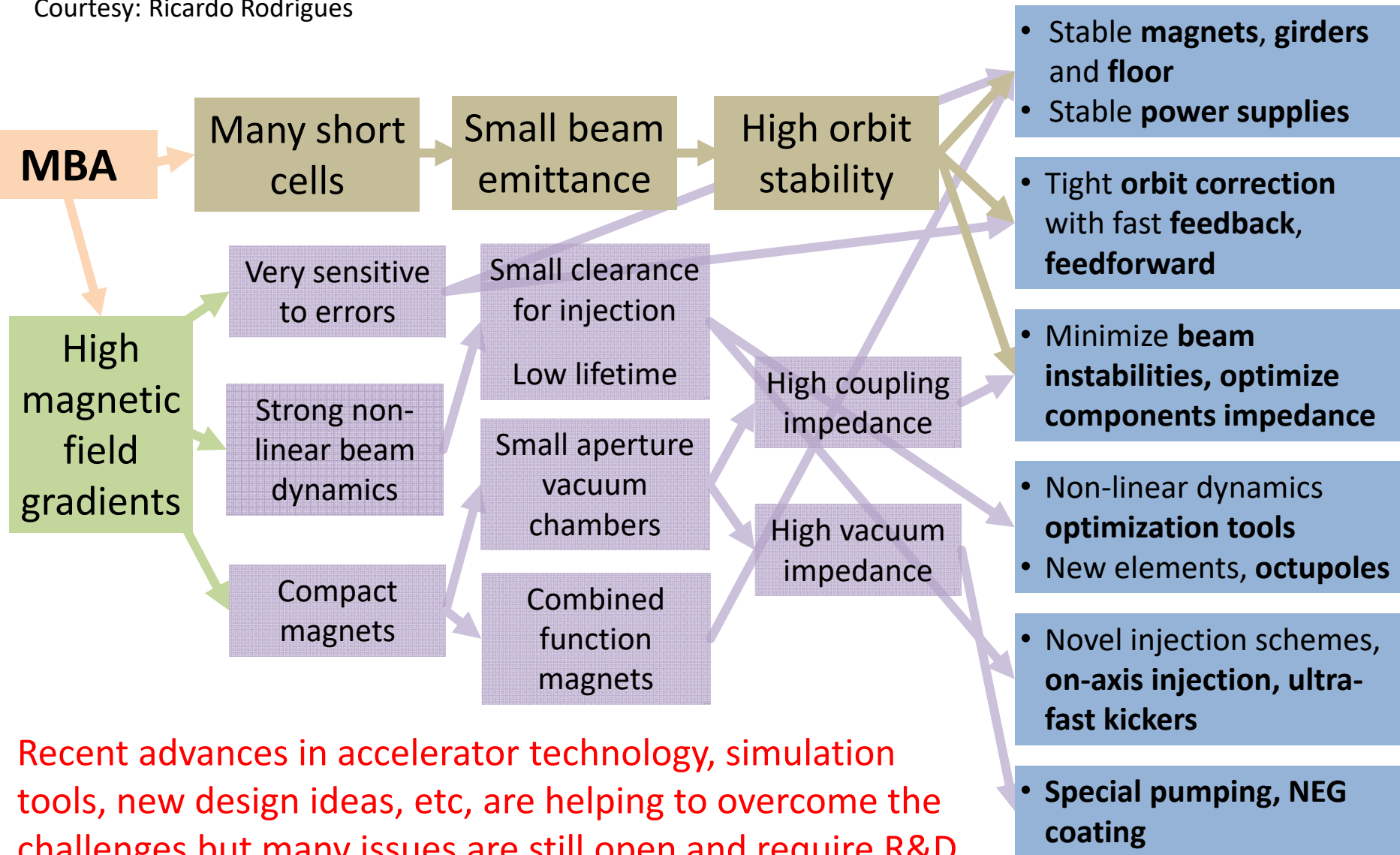
Courtesy: Ricardo Rodrigues



Courtesy: Ricardo Rodrigues



Courtesy: Ricardo Rodrigues



Recent advances in accelerator technology, simulation tools, new design ideas, etc, are helping to overcome the challenges but many issues are still open and require R&D.

- Stable **magnets, girders and floor**
- Stable **power supplies**
- Tight **orbit correction with fast feedback, feedforward**
- Minimize **beam instabilities, optimize components impedance**
- Non-linear dynamics **optimization tools**
- New elements, **octupoles**
- Novel injection schemes, **on-axis injection, ultra-fast kickers**
- **Special pumping, NEG coating**

- Increase damping partition number J_x by adding transverse field gradient in dipoles.

$$J_x = 1 - \frac{\oint (1 - 2n)\eta |h(s)|^3 ds}{\oint h(s)^2 ds}$$

- Increase damping partition number J_x by adding transverse field gradient in dipoles.

$$J_x = 1 - \frac{\oint (1 - 2n)\eta |h(s)|^3 ds}{\oint h(s)^2 ds}$$

- Longitudinal dipole gradient

$$\epsilon_x \propto \oint \mathcal{H}(s) h(s)^3 ds$$

- Curvature function $h(s) = 1/\rho(s)$
- To keep product small: compensate variation in $\mathcal{H}(s)$ with variation in $h(s)$
- Radiate more (high curvature) where $\mathcal{H}(s)$ is small.

- Increase damping partition number J_x by adding transverse field gradient in dipoles.

$$J_x = 1 - \frac{\oint (1 - 2n)\eta |h(s)|^3 ds}{\oint h(s)^2 ds}$$

- Longitudinal dipole gradient

$$\epsilon_x \propto \oint \mathcal{H}(s) h(s)^3 ds$$

- Curvature function $h(s) = 1/\rho(s)$
- To keep product small: compensate variation in $\mathcal{H}(s)$ with variation in $h(s)$
- Radiate more (high curvature) where $\mathcal{H}(s)$ is small.
- Achromatic cells and low field dipoles to enhance emittance reduction with Insertion Devices.

- Increase damping partition number J_x by adding transverse field gradient in dipoles.

$$J_x = 1 - \frac{\oint (1 - 2n)\eta |h(s)|^3 ds}{\oint h(s)^2 ds}$$

- Longitudinal dipole gradient

$$\epsilon_x \propto \oint \mathcal{H}(s) h(s)^3 ds$$

- Curvature function $h(s) = 1/\rho(s)$
- To keep product small: compensate variation in $\mathcal{H}(s)$ with variation in $h(s)$
- Radiate more (high curvature) where $\mathcal{H}(s)$ is small.
- Achromatic cells and low field dipoles to enhance emittance reduction with Insertion Devices.
- Different dipole lengths, shorter dipoles at cell ends, where $\eta = \eta' = 0$.

- Increase damping partition number J_x by adding transverse field gradient in dipoles.

$$J_x = 1 - \frac{\oint (1 - 2n)\eta |h(s)|^3 ds}{\oint h(s)^2 ds}$$

- Longitudinal dipole gradient

$$\epsilon_x \propto \oint \mathcal{H}(s) h(s)^3 ds$$

- Curvature function $h(s) = 1/\rho(s)$
- To keep product small: compensate variation in $\mathcal{H}(s)$ with variation in $h(s)$
- Radiate more (high curvature) where $\mathcal{H}(s)$ is small.
- Achromatic cells and low field dipoles to enhance emittance reduction with Insertion Devices.
- Different dipole lengths, shorter dipoles at cell ends, where $\eta = \eta' = 0$.
- Increase damping with Damping Wigglers \rightarrow energy spread increases.

- Increase damping partition number J_x by adding transverse field gradient in dipoles.

$$J_x = 1 - \frac{\oint (1 - 2n)\eta |h(s)|^3 ds}{\oint h(s)^2 ds}$$

- Longitudinal dipole gradient

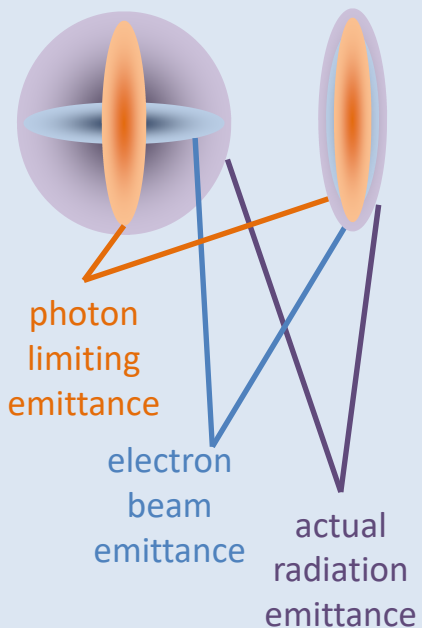
$$\epsilon_x \propto \oint \mathcal{H}(s) |h(s)|^3 ds$$

- Curvature function $h(s) = 1/\rho(s)$
- To keep product small: compensate variation in $\mathcal{H}(s)$ with variation in $h(s)$
- Radiate more (high curvature) where $\mathcal{H}(s)$ is small.
- Achromatic cells and low field dipoles to enhance emittance reduction with Insertion Devices.
- Different dipole lengths, shorter dipoles at cell ends, where $\eta = \eta' = 0$.
- Increase damping with Damping Wigglers \rightarrow energy spread increases.
- Anti-bends (SLS). Disentangle dispersion η and beta function β_x .

Electron beam and radiation phase-space

mismatched

matched



Matching condition

$$\beta_e = \frac{\sigma_e}{\sigma'_e} = \frac{\sigma_r}{\sigma'_r}$$

Highest brilliance from undulator of length L is achieved when

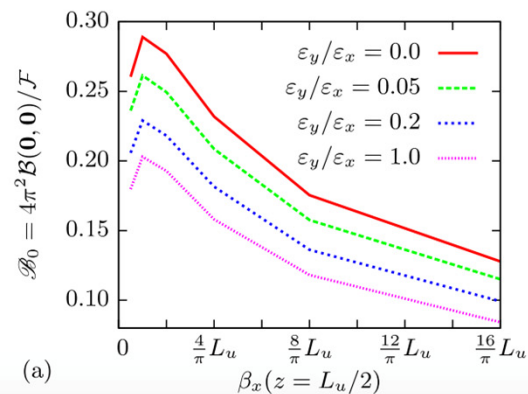
$$\beta_{x,y}^{opt} \approx \frac{L}{\pi}$$

$$\beta_{x,y}^{opt} \sim 1 - 2m$$

Ryan R. Lindberg and Kwang-Je Kim (2015)

Phys. Rev. ST Accel. Beams **18**, 090702 (2015)

COMPACT REPRESENTATIONS OF PARTIALLY ...

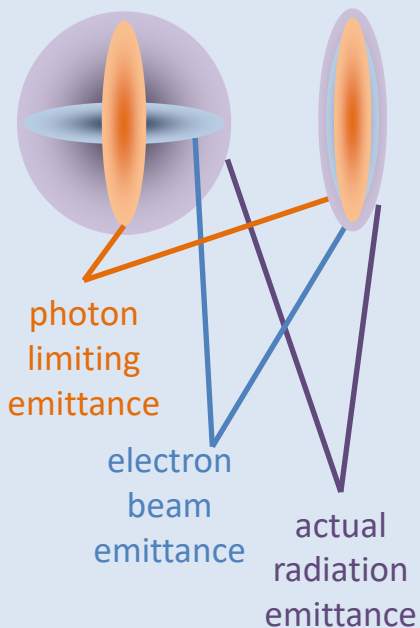


IV. CONCLUSIONS

In this paper we have described three different coherent mode representations of partially coherent undulator radiation. We began with the well-known Gauss-Schell decomposition in terms of Gauss-Hermite modes, which is valid provided the electron beam emittance is much larger than the natural radiation emittance $\lambda/4\pi$. In this largely incoherent case the specifics of the single-electron undulator field are unimportant. We then refined our analysis to include the situation when the electron beam emittance ϵ_y in one direction is arbitrary, and found that the modes along y are determined by solving a matrix

Electron beam and radiation phase-space

mismatched matched



Matching condition

$$\beta_e = \frac{\sigma_e}{\sigma'_e} = \frac{\sigma_r}{\sigma'_r}$$

Highest brilliance from undulator of length L is achieved when

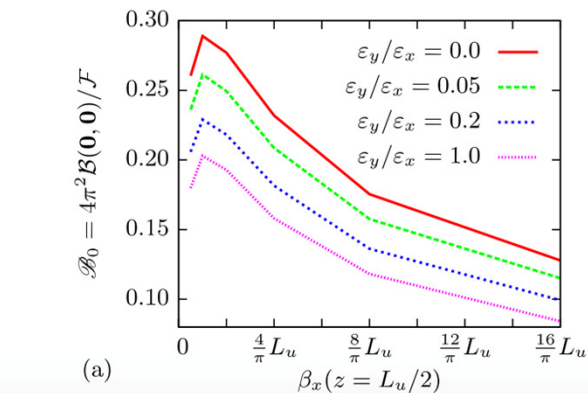
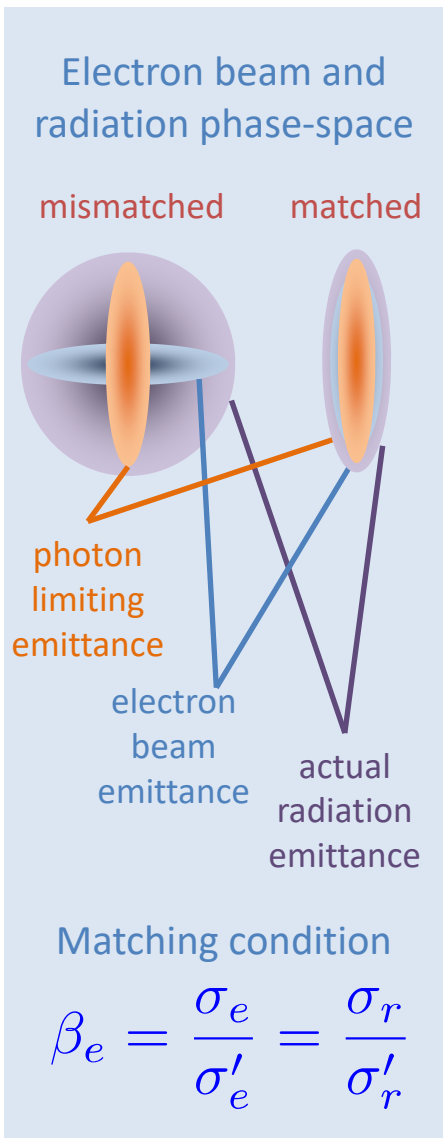
$$\beta_{x,y}^{opt} \approx \frac{L}{\pi}$$

$$\beta_{x,y}^{opt} \sim 1 - 2m$$

Ryan R. Lindberg and Kwang-Je Kim (2015)

COMPACT REPRESENTATIONS OF PARTIALLY ...

Phys. Rev. ST Accel. Beams **18**, 090702 (2015)



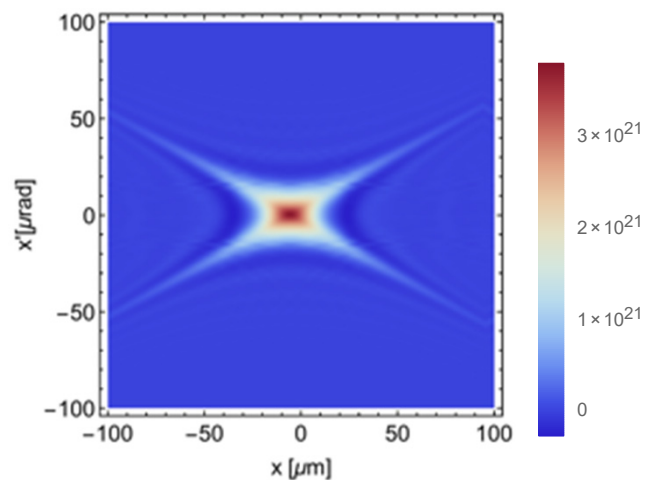
IV. CONCLUSIONS

In this paper we have described three different coherent mode representations of partially coherent undulator radiation. We began with the well-known Gaussian-Schell decomposition in terms of Gauss-Hermite modes, which is valid provided the electron beam emittance is much larger than the natural radiation emittance $\lambda/4\pi$. In this largely incoherent case the specifics of the single-electron undulator field are unimportant. We then refined our analysis to include the situation when the electron beam emittance ϵ_y in one direction is arbitrary, and found that the modes along y are determined by solving a matrix

Figure 8(a) shows that the coherence is maximized when $\beta_x \approx L_u/\pi$ (or $\hat{\beta}_x \approx 1$), which indicates that the “natural” Rayleigh range of undulator radiation $Z_R \approx \beta_x \approx L_u/\pi$. Unfortunately, it is nearly impossible for lattice designers to make the beta functions in both x and y be simultaneously that small, and typically $\beta_x > 3L_u/\pi$.

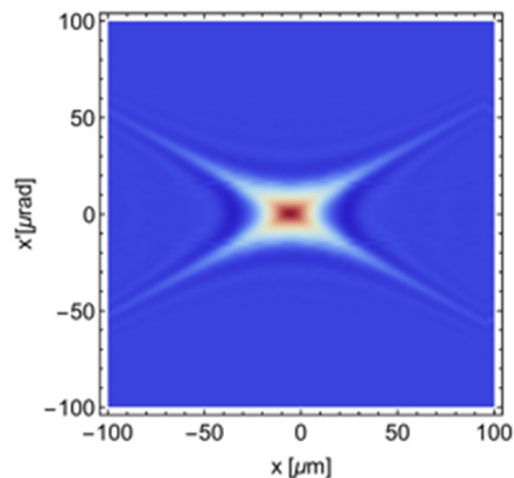
very similar: the profiles when β_x is increased by a factor of 16 can be approximately obtained by multiplying those in Fig. 8(b) by 1.3. Figure 8(c) plots the profiles along x as we vary β_x . Each plot is approximately Gaussian, and we see that the angular spread decreases as β_x increases. More careful inspection shows that the width of the angular

SINGLE PHOTON'S "PHASE SPACE" (900 eV)
(Wigner Distribution Function)
from single electron

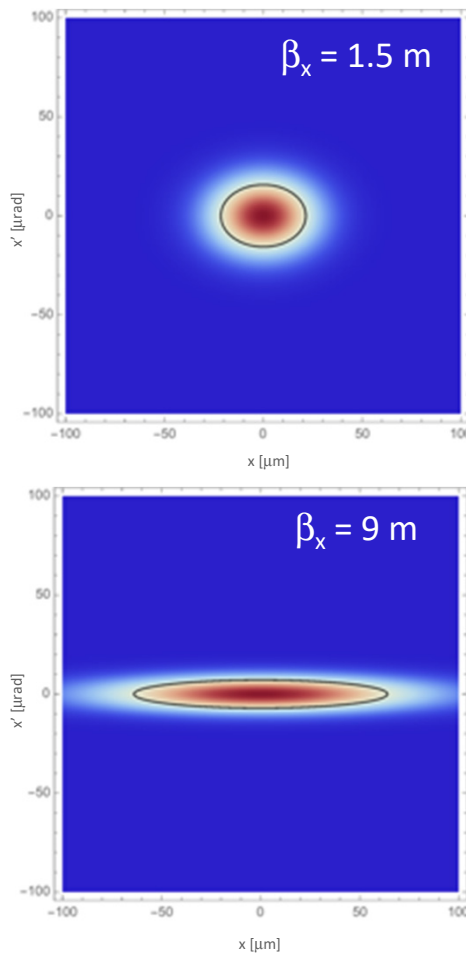


Courtesy: Harry Westfahl Jr.

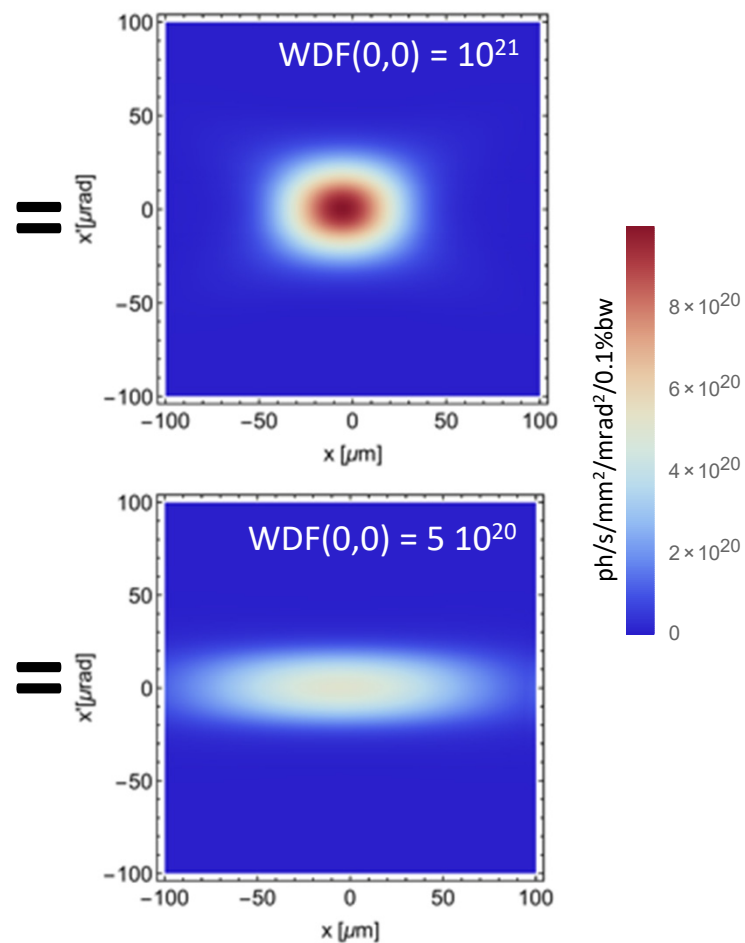
SINGLE PHOTON'S "PHASE SPACE" (900 eV)
(Wigner Distribution Function)
from single electron



ELECTRONS PHASE SPACE
(HORIZONTAL)



AVERAGE PHOTON'S "PHASE SPACE"
(Wigner Distribution Function=Brilliance)



Courtesy: Harry Westfahl Jr.

First beam 2018 – Open in 2019



Budget

- Accelerators 100 M US\$
- 13 beamlines 140 M US\$
- Building 213 M US\$
- Human Res 57 M US\$
- Total 510 M US\$


Schedule

- Jan.2015 start of building construction
- Oct.2017 start of machine installation
- Jul.2018 start of SR commissioning
- Sep.2018 phase 1 operation (20mA, NCC)
- Feb.2019 phase 2 operation (100mA, SCC)

City of Campinas (population: 1.100.000)



- UVSX**
- 1.37 GeV
 - 100 nm.rad
 - 18 beamlines
 - Over 1200 users/yr



UNICAMP
40.000
students



CTBE



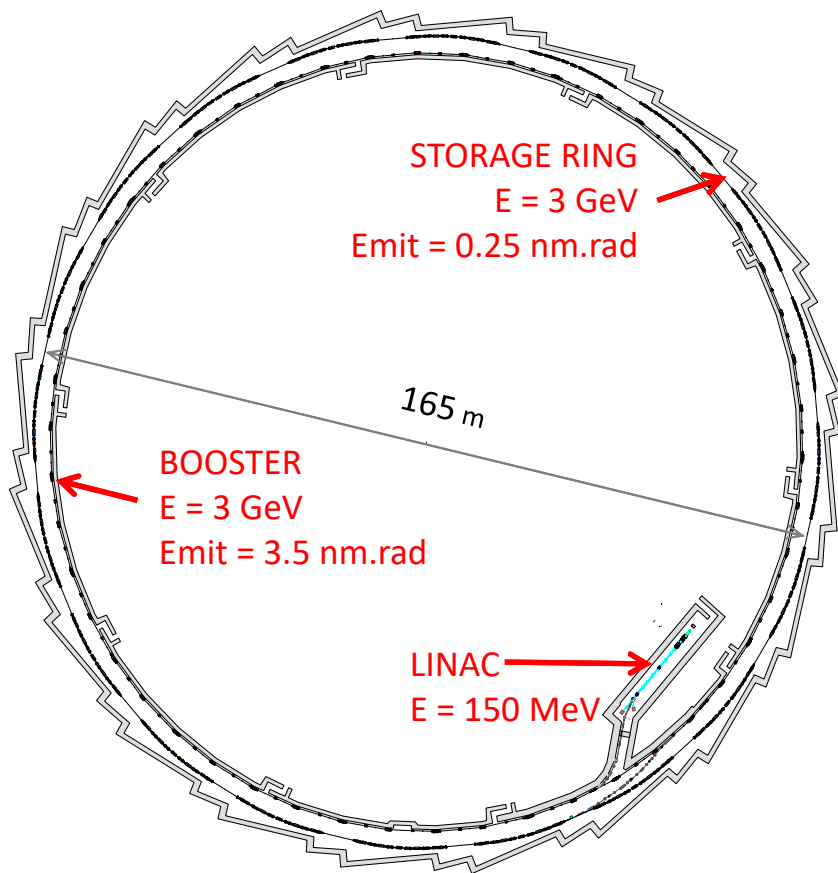
LNBIQ



LNLS
200 employees
80 students &
trainees



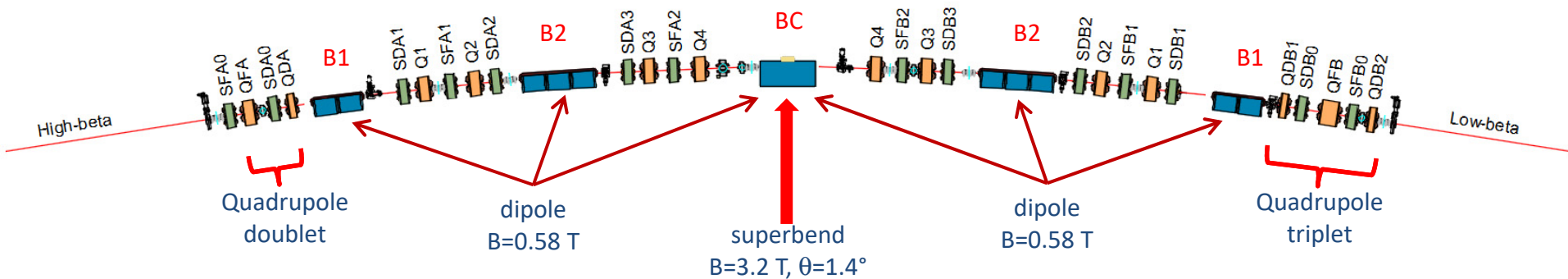
LNNANO

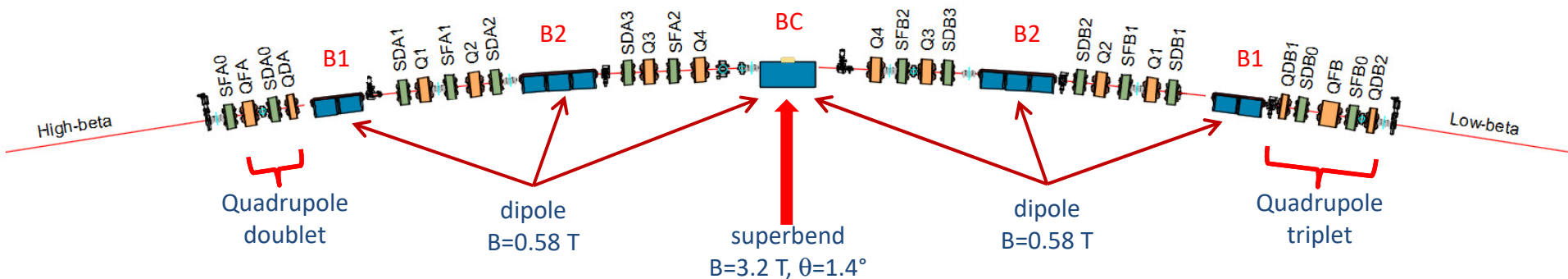


Storage Ring	
Beam energy	3.0 GeV
Circumference	518.4 m
Lattice	20 x 5BA
Hor. emittance (bare lat.)	0.25 nm.rad
Hor. emittance (with IDs)	→ 0.15 nm.rad
Betatron tunes (H/V)	49.11 / 14.17
Natural chrom. (H/V)	-119.0 / -81.2
rms energy spread	0.85×10^{-3}
Energy loss/turn (dipoles)	473 keV
Dam. times (H/V/L) [ms]	16.9 / 22.0 / 12.9
Nominal current, top up	350 mA

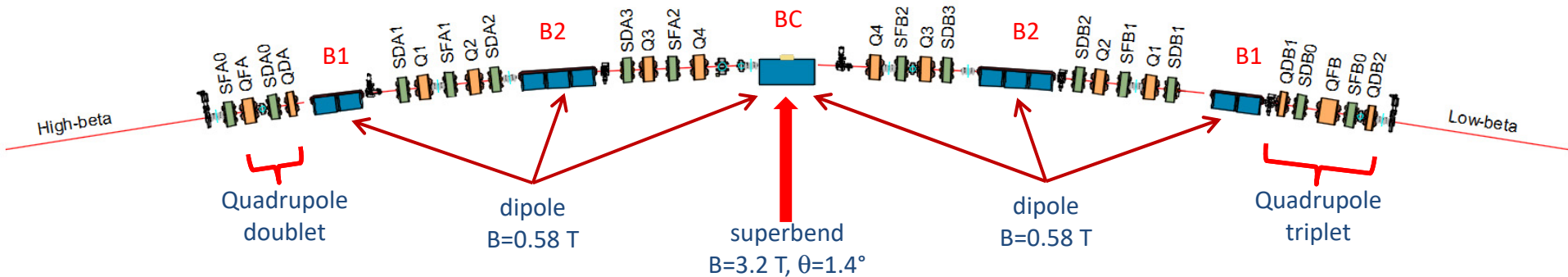
Booster	
Circumference	496.8 m
Emittance @ 3 GeV	3.5 nm.rad
Lattice	50 Bend
Cycling frequency	2 Hz

The Sirius 5BA magnet lattice

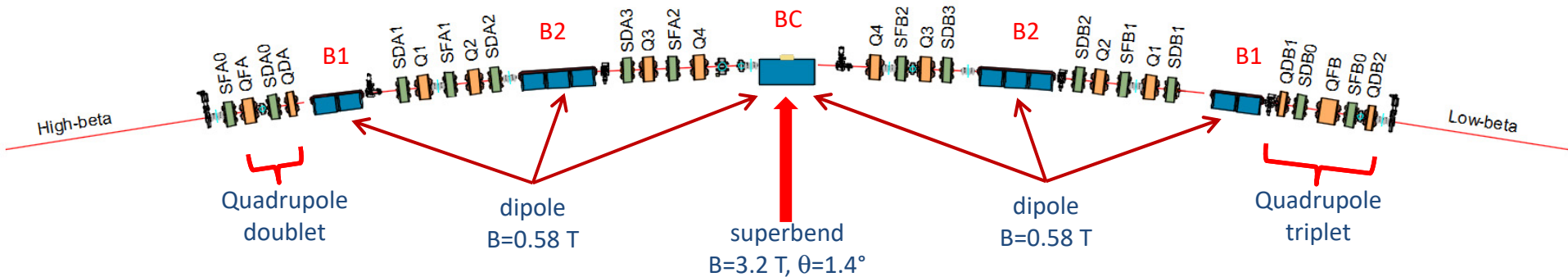




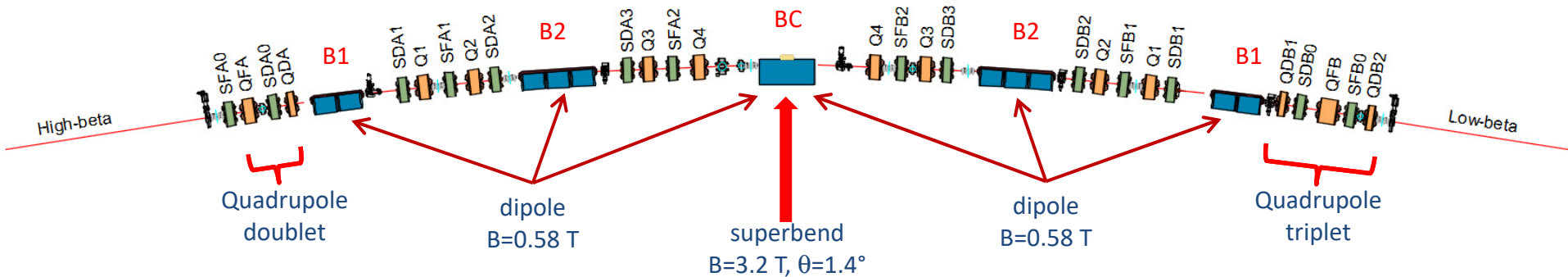
- 20 - 5BA arcs and 2 types of straight sections for insertion devices:
 - 5 **high β_x** straight sections of **7.0 m** – matching with quad **doublets**.
 - 15 **low β_x** straight sections of **6.0 m** – matching with quad **triplets**.



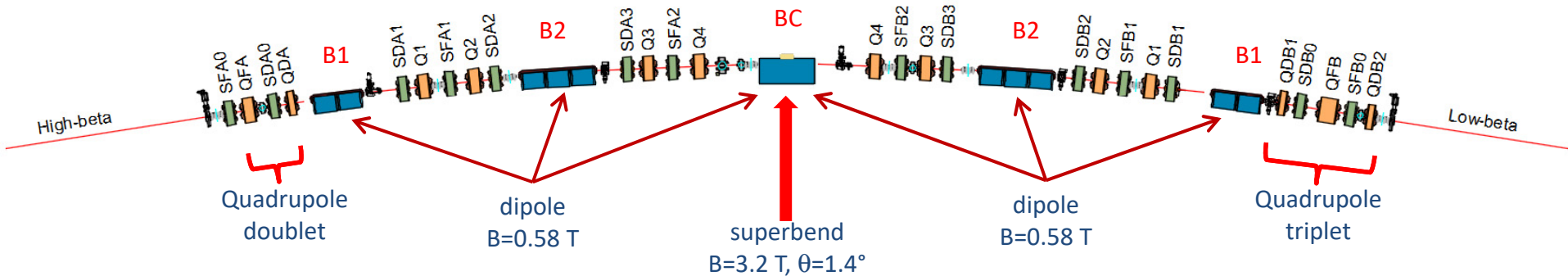
- 20 - 5BA arcs and 2 types of straight sections for insertion devices:
 - 5 **high β_x** straight sections of **7.0 m** – matching with quad **doublets**.
 - 15 **low β_x** straight sections of **6.0 m** – matching with quad **triplets**.
- 20 PM longitudinal gradient **superbends**
 - sharp peak field of **$B_p = 3.2$ T** in the center \rightarrow critical photon energy of **$e_c = 19.2$ keV**



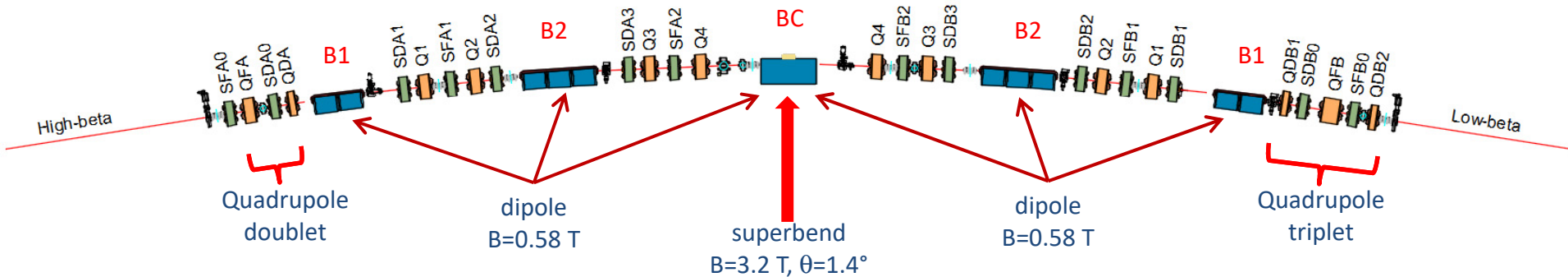
- 20 - 5BA arcs and 2 types of straight sections for insertion devices:
 - 5 **high β_x** straight sections of **7.0 m** – matching with quad **doublets**.
 - 15 **low β_x** straight sections of **6.0 m** – matching with quad **triplets**.
- 20 PM longitudinal gradient **superbends**
 - sharp peak field of **$B_p = 3.2$ T** in the center \rightarrow critical photon energy of **$e_c = 19.2$ keV**
- Low field (0.58 T) EM and PM dipoles with transverse field gradient (7.8 T/m)



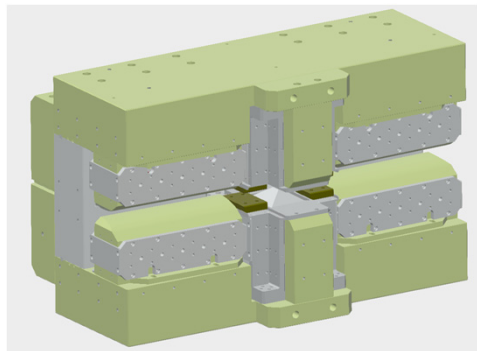
- 20 - 5BA arcs and 2 types of straight sections for insertion devices:
 - 5 **high β_x** straight sections of **7.0 m** – matching with quad **doublers**.
 - 15 **low β_x** straight sections of **6.0 m** – matching with quad **triplets**.
- 20 PM longitudinal gradient **superbends**
 - sharp peak field of **$B_p = 3.2$ T** in the center \rightarrow critical photon energy of **$e_c = 19.2$ keV**
- Low field (0.58 T) EM and PM dipoles with transverse field gradient (7.8 T/m)
- 13 quadrupoles/cell, 14 sextupoles/cell, no octupoles



- 20 - 5BA arcs and 2 types of straight sections for insertion devices:
 - 5 **high β_x** straight sections of **7.0 m** – matching with quad **doublers**.
 - 15 **low β_x** straight sections of **6.0 m** – matching with quad **triplets**.
- 20 PM longitudinal gradient **superbends**
 - sharp peak field of **$B_p = 3.2$ T** in the center \rightarrow critical photon energy of **$e_c = 19.2$ keV**
- Low field (0.58 T) EM and PM dipoles with transverse field gradient (7.8 T/m)
- 13 quadrupoles/cell, 14 sextupoles/cell, no octupoles
- Different lengths for B1 and B2 (but same unit block)



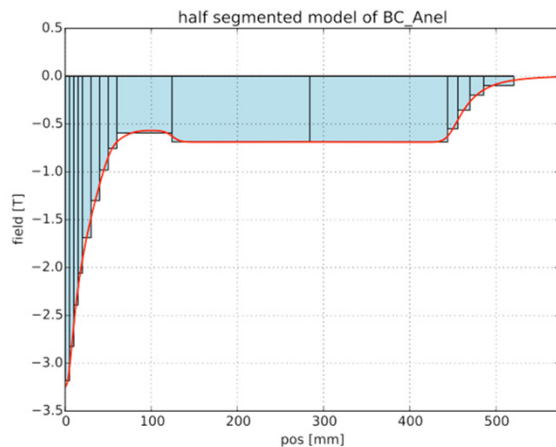
- 20 - 5BA arcs and 2 types of straight sections for insertion devices:
 - 5 **high β_x** straight sections of **7.0 m** – matching with quad **doublets**.
 - 15 **low β_x** straight sections of **6.0 m** – matching with quad **triplets**.
- 20 PM longitudinal gradient **superbends**
 - sharp peak field of $B_p = 3.2\text{ T}$ in the center \rightarrow critical photon energy of $e_c = 19.2\text{ keV}$
- Low field (0.58 T) EM and PM dipoles with transverse field gradient (7.8 T/m)
- 13 quadrupoles/cell, 14 sextupoles/cell, no octupoles
- Different lengths for B1 and B2 (but same unit block)
- NEG coated copper beam pipe $\varnothing = 24\text{ mm}$ (internal)

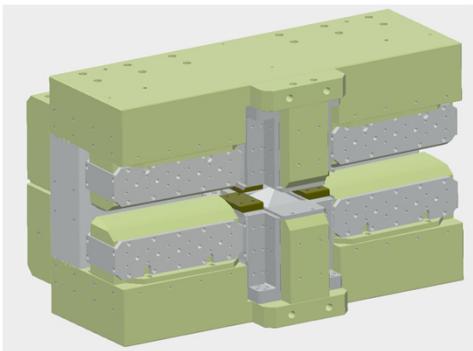


Permanent magnet (NdFeB)

High field insert (3.2 T) superbend

- 19 keV critical energy at peak
- Hard X-rays produced only at beamline exit
- Total energy loss/turn from dipoles = 473 keV

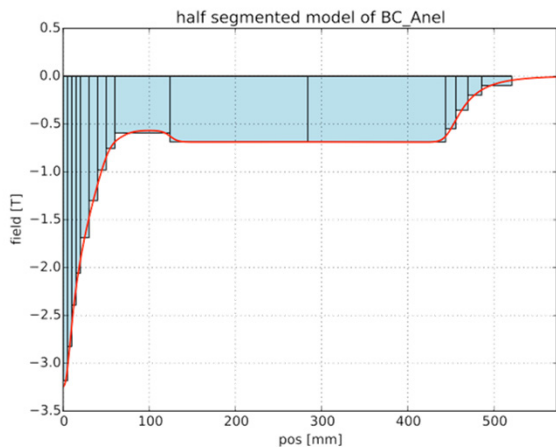




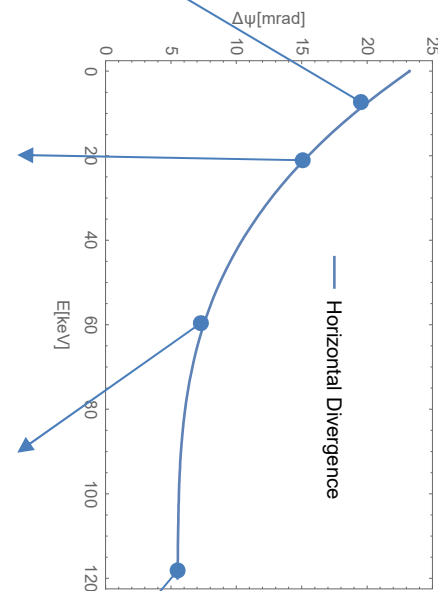
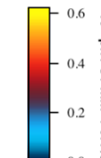
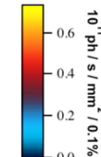
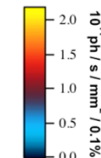
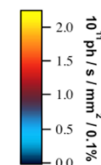
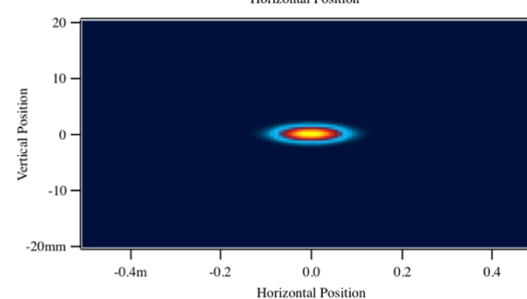
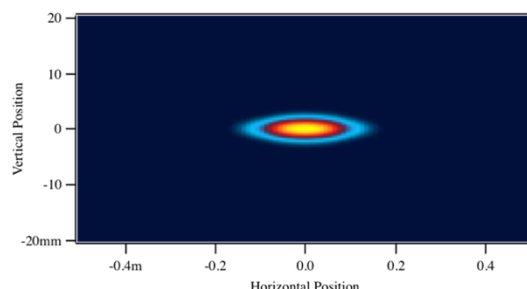
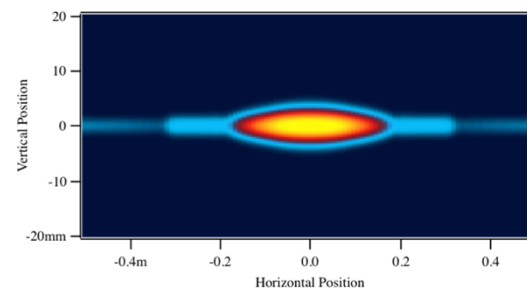
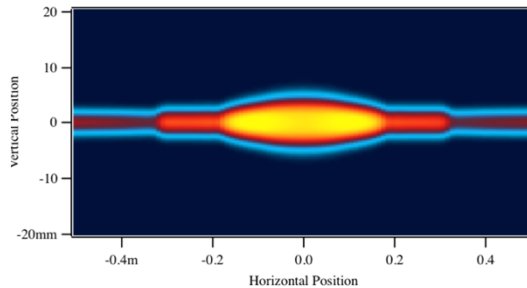
Permanent magnet (NdFeB)

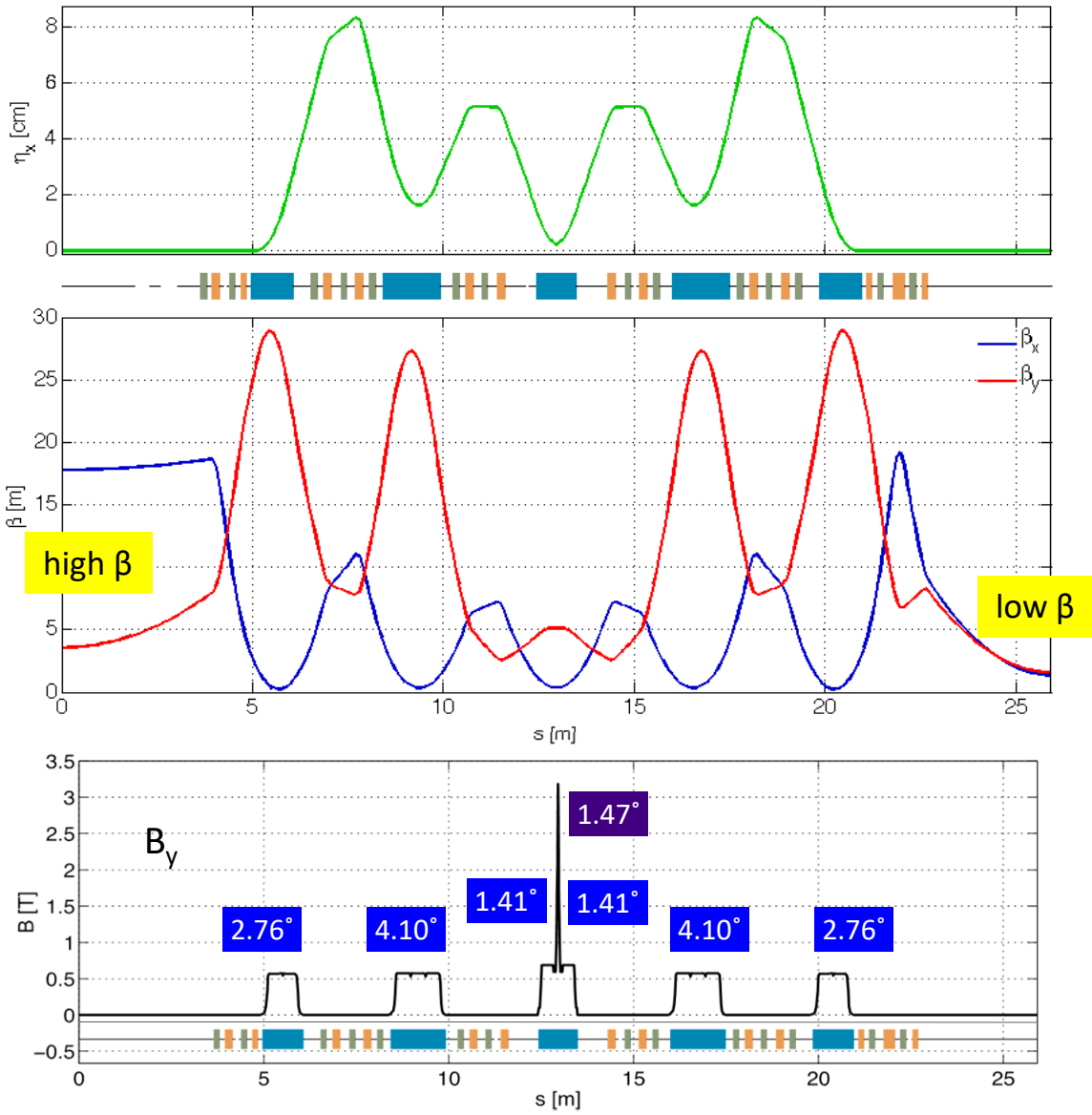
High field insert (3.2 T) superbend

- 19 keV critical energy at peak
- Hard X-rays produced only at beamline exit
- Total energy loss/turn from dipoles = 473 keV

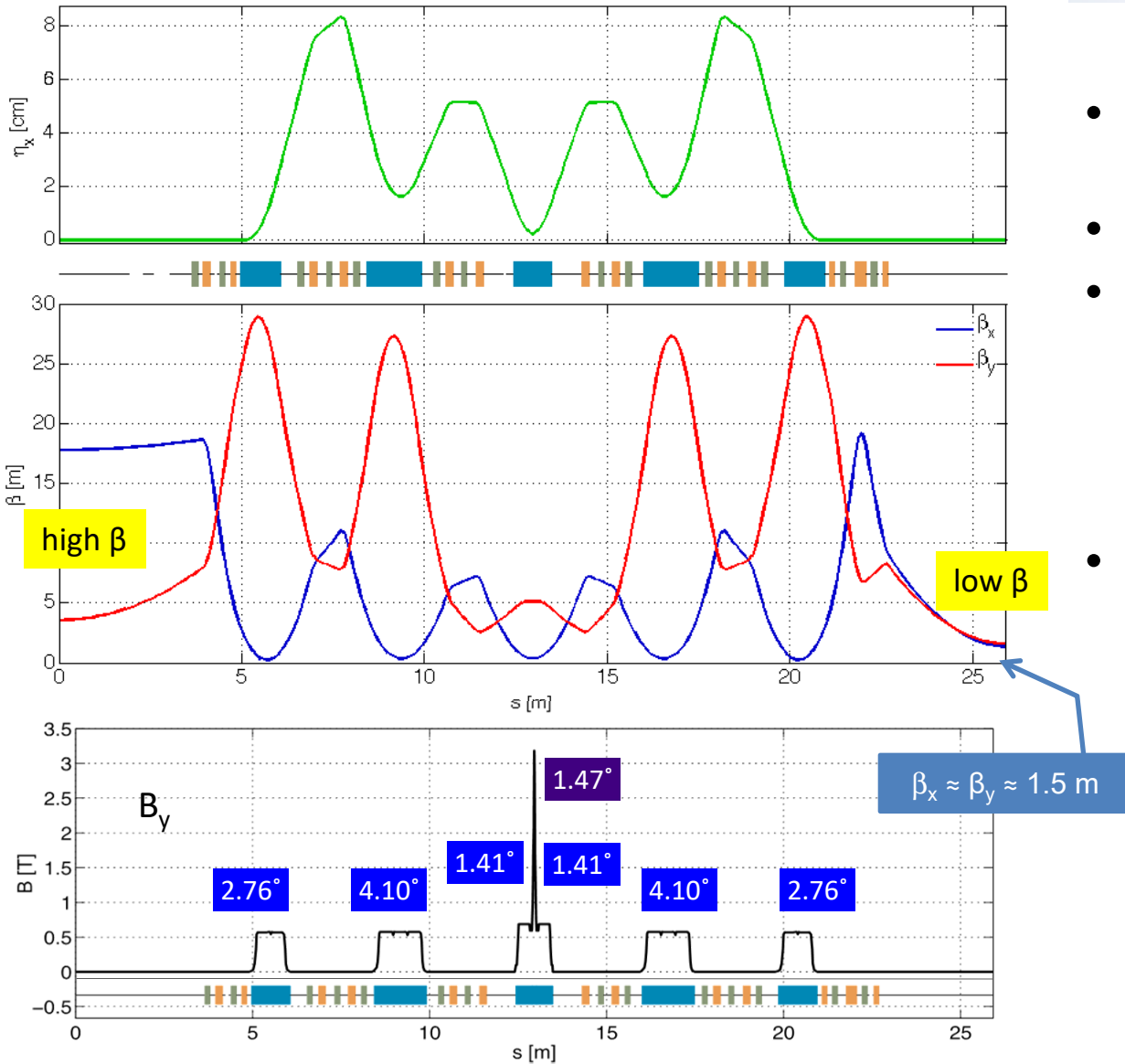


From SRW:

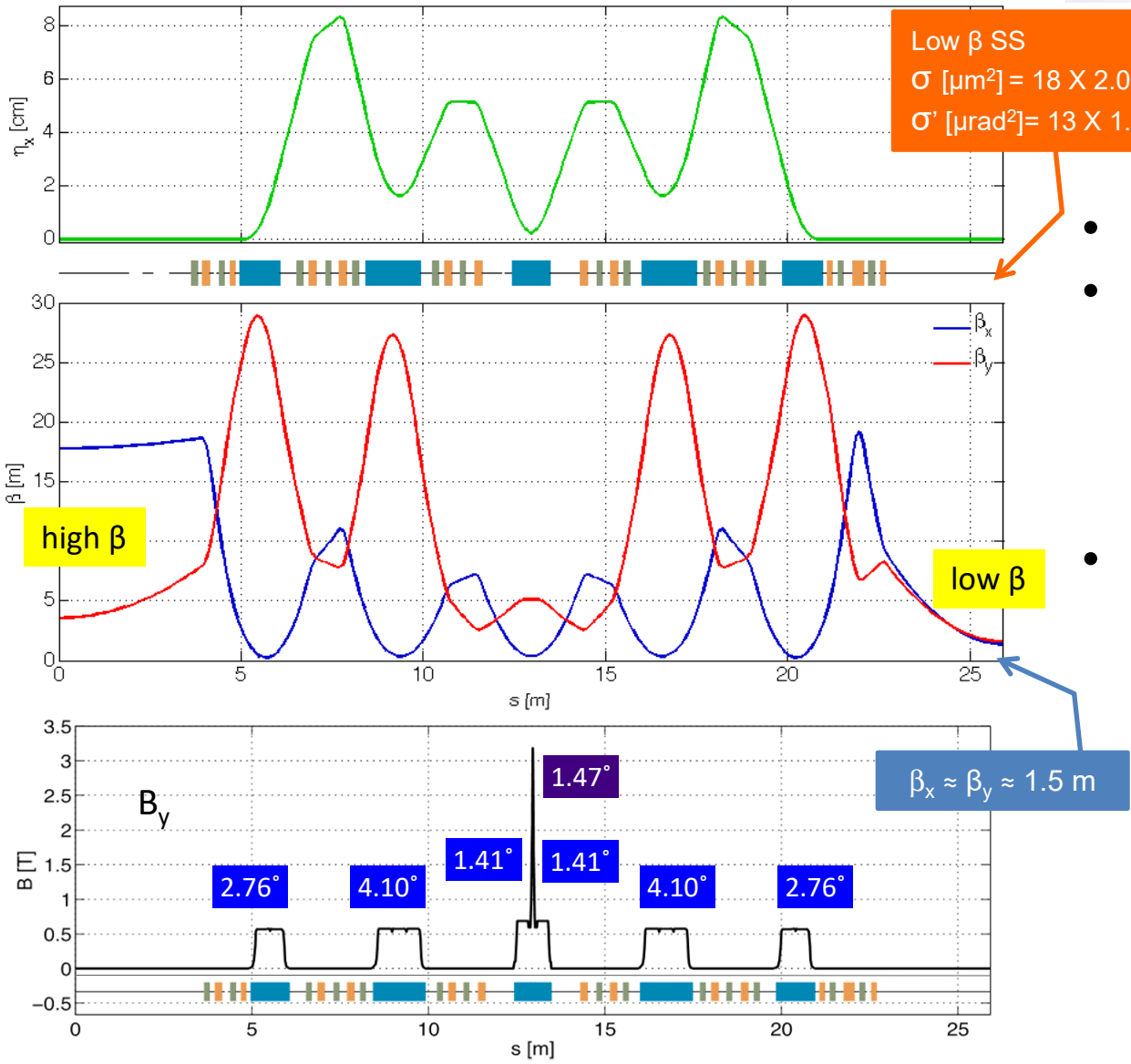




- 5-fold symmetric optics with 5 high and 15 low β sections.
- Achromatic cells.
- At low β sections
 - $\beta_x \approx \beta_y \approx 1.5$ m
 - Optimized electron and photon beam phase-space matching for undulators.
- At superbend
 - Strong focusing of dispersion and β_x functions
 - Beam size: $9.6 \times 3.6 \mu\text{m}^2$

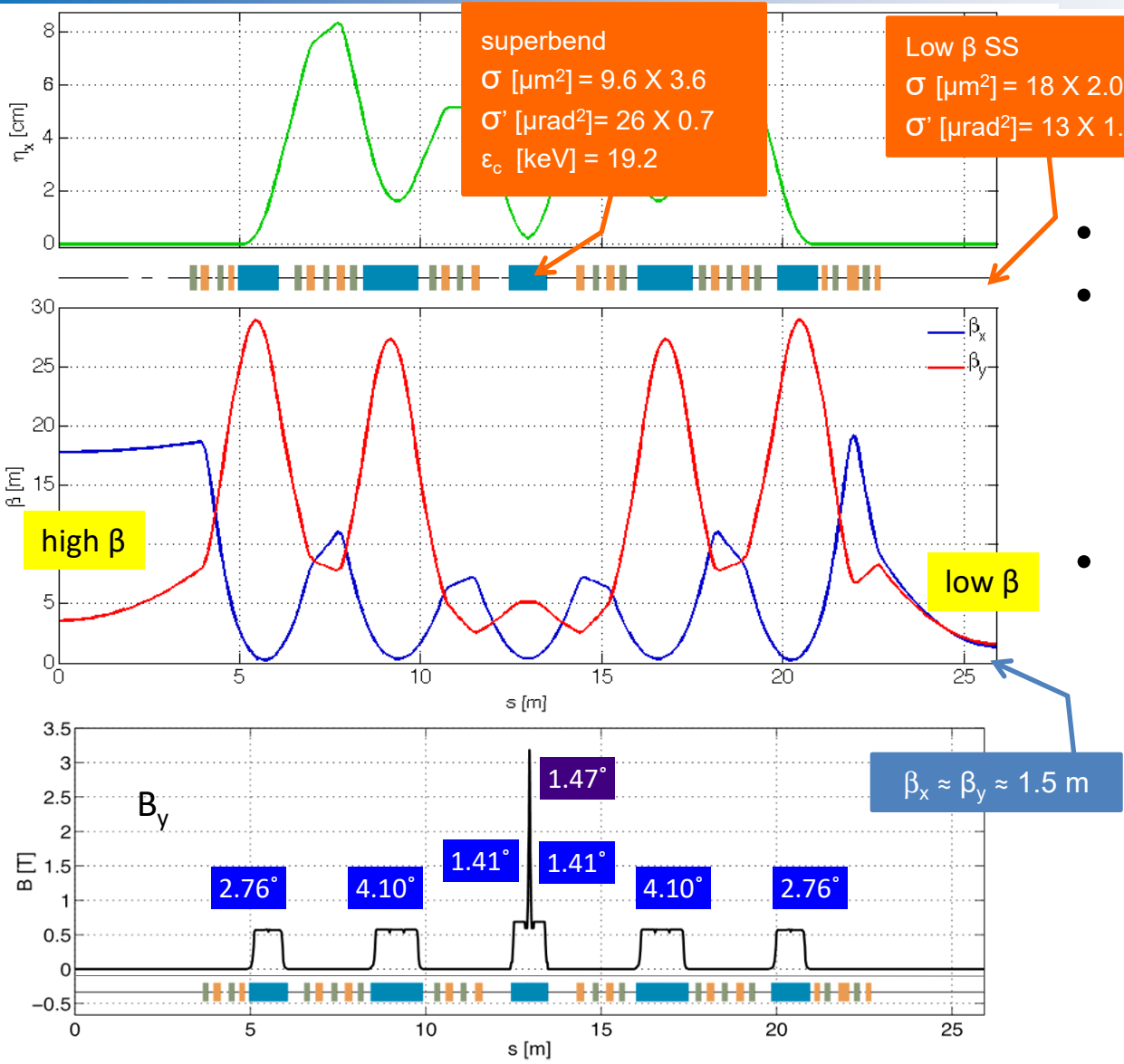


- 5-fold symmetric optics with 5 high and 15 low β sections.
- Achromatic cells.
- At low β sections
 - $\beta_x \approx \beta_y \approx 1.5$ m
 - Optimized electron and photon beam phase-space matching for undulators.
- At superbend
 - Strong focusing of dispersion and β_x functions
 - Beam size: $9.6 \times 3.6 \mu\text{m}^2$



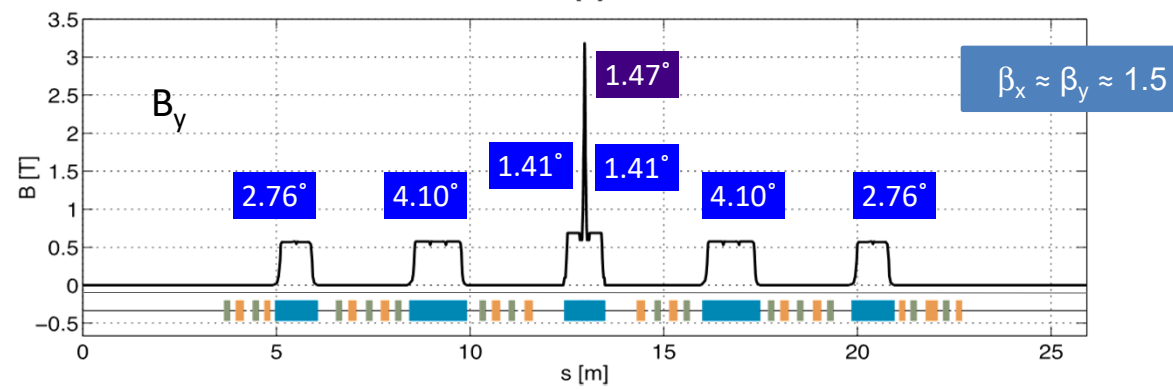
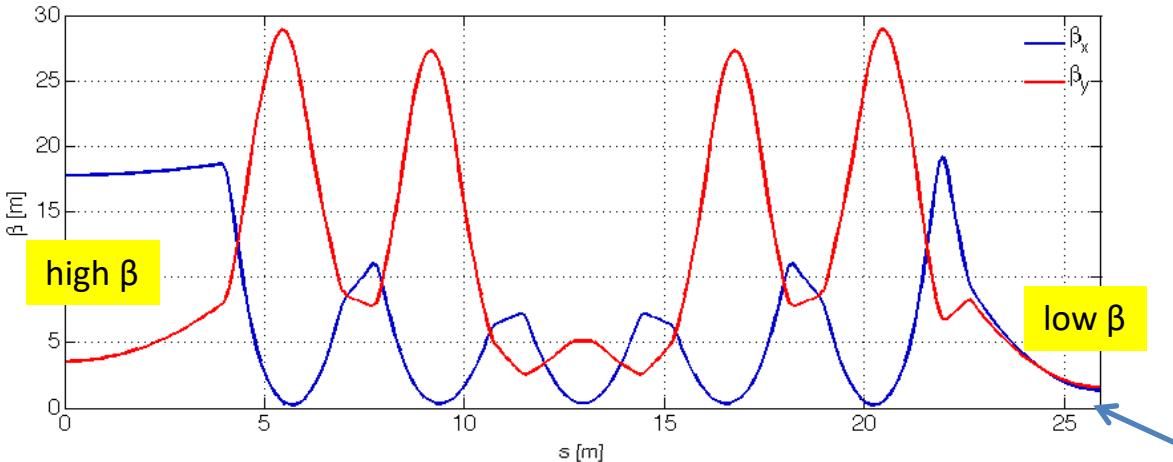
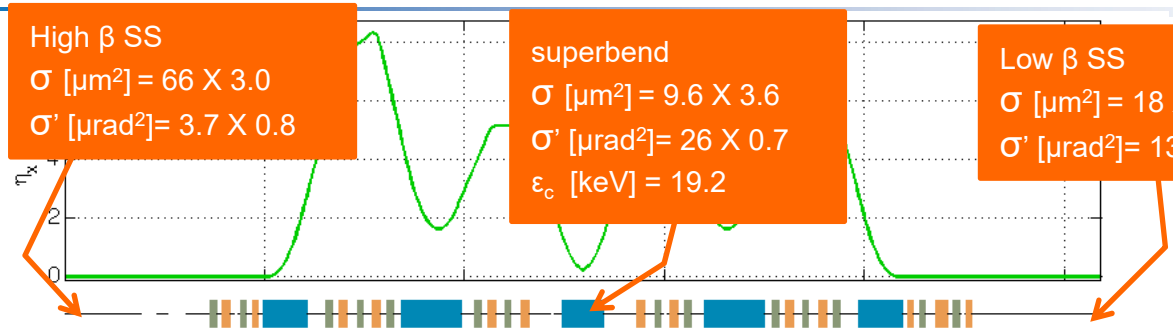
and symmetric optics with 5 high and 15 low β sections.

- Achromatic cells.
- At low β sections
 - $\beta_x \approx \beta_y \approx 1.5$ m
 - Optimized electron and photon beam phase-space matching for undulators.
- At superbend
 - Strong focusing of dispersion and β_x functions
 - Beam size: $9.6 \times 3.6 \mu\text{m}^2$



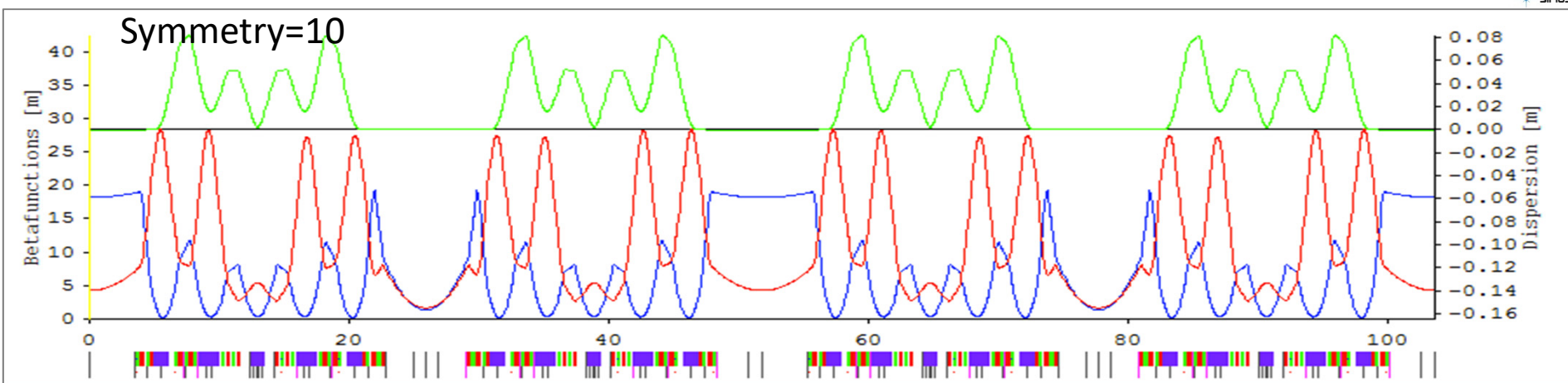
and symmetric optics with 5 high and 15 low β sections.

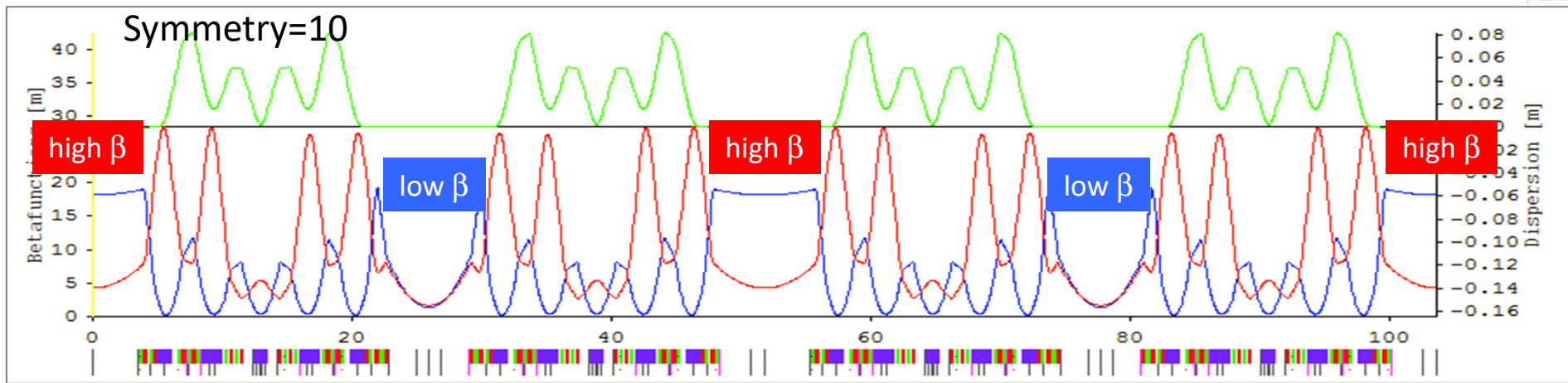
- Achromatic cells.
- At low β sections
 - $\beta_x \approx \beta_y \approx 1.5$ m
 - Optimized electron and photon beam phase-space matching for undulators.
- At superbend
 - Strong focusing of dispersion and β_x functions
 - Beam size: $9.6 \times 3.6 \mu\text{m}^2$

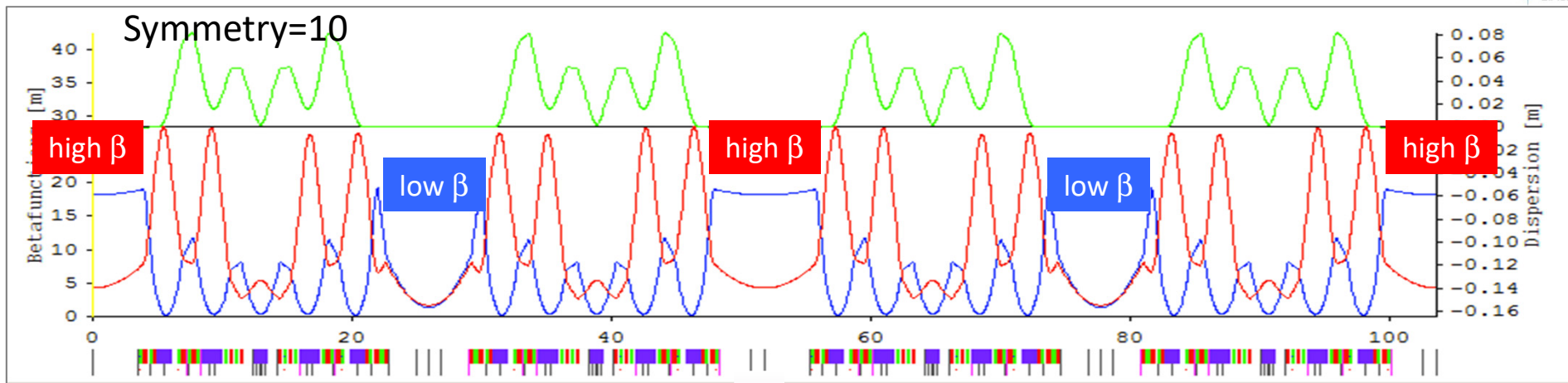


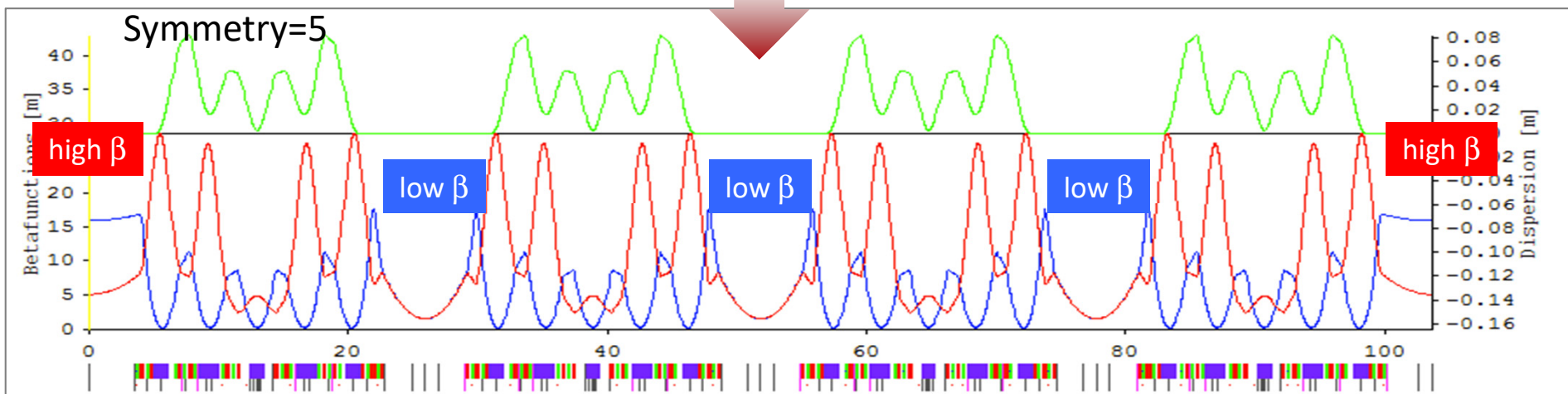
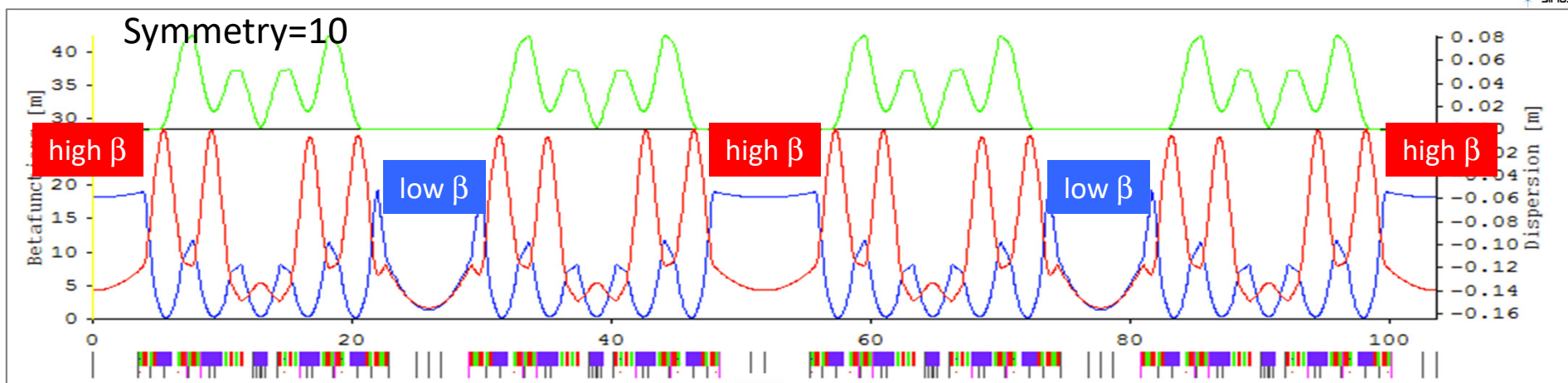
and symmetric optics with 5 high and 15 low β sections.

- Achromatic cells.
- At low β sections
 - $\beta_x \approx \beta_y \approx 1.5$ m
 - Optimized electron and photon beam phase-space matching for undulators.
- At superbend
 - Strong focusing of dispersion and β_x functions
 - Beam size: 9.6 x 3.6 μm^2



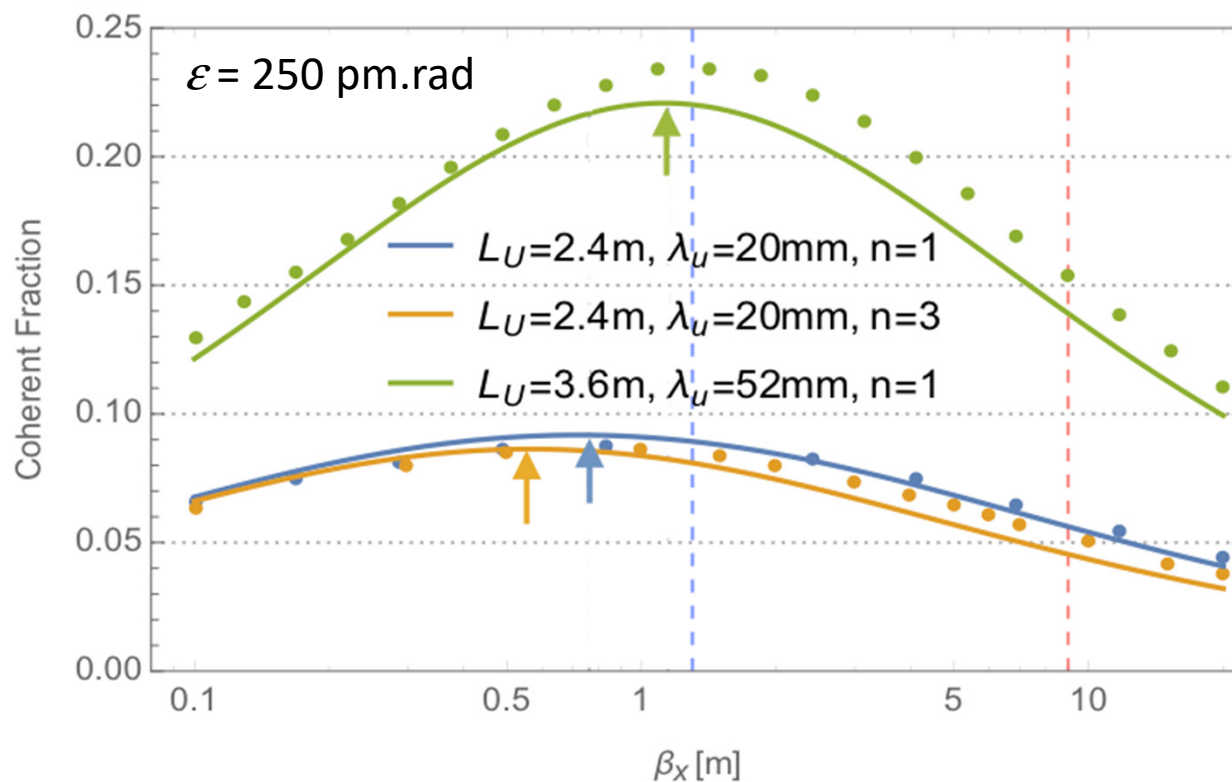






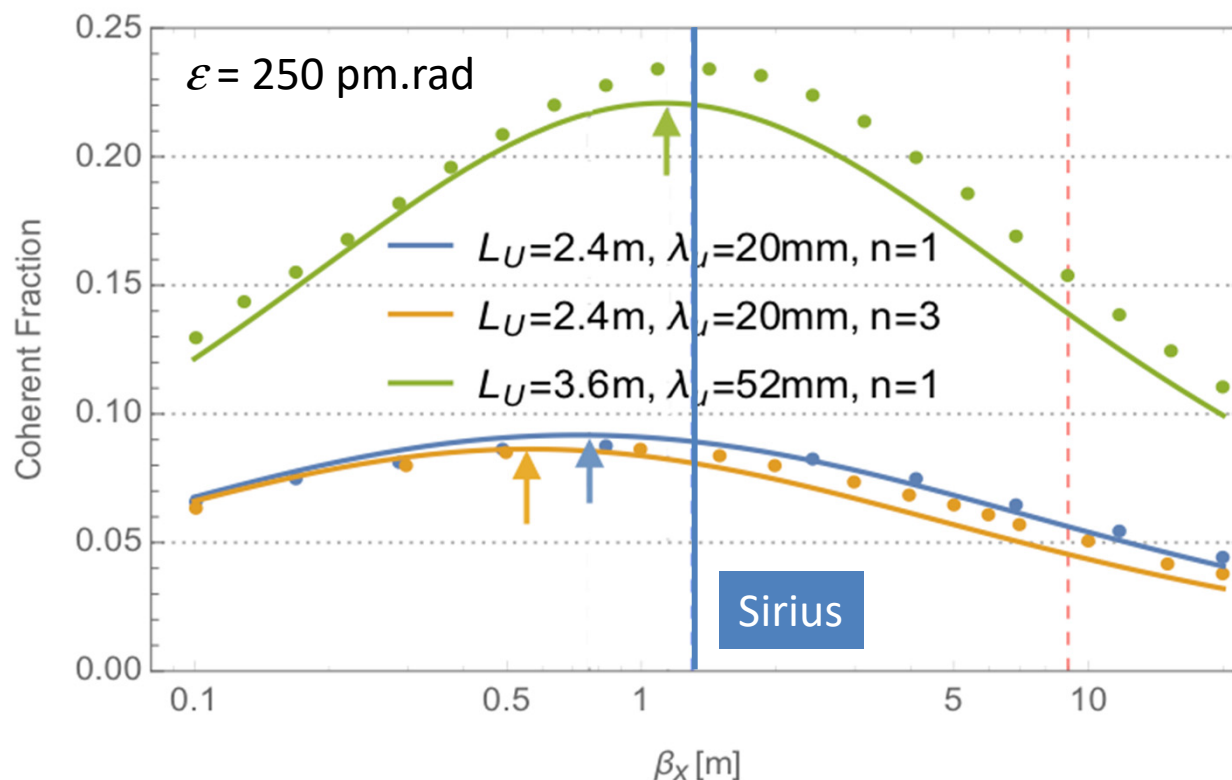
Symmetry 5: 5 high β_x + 15 low β_x

Courtesy: Harry Westfahl Jr



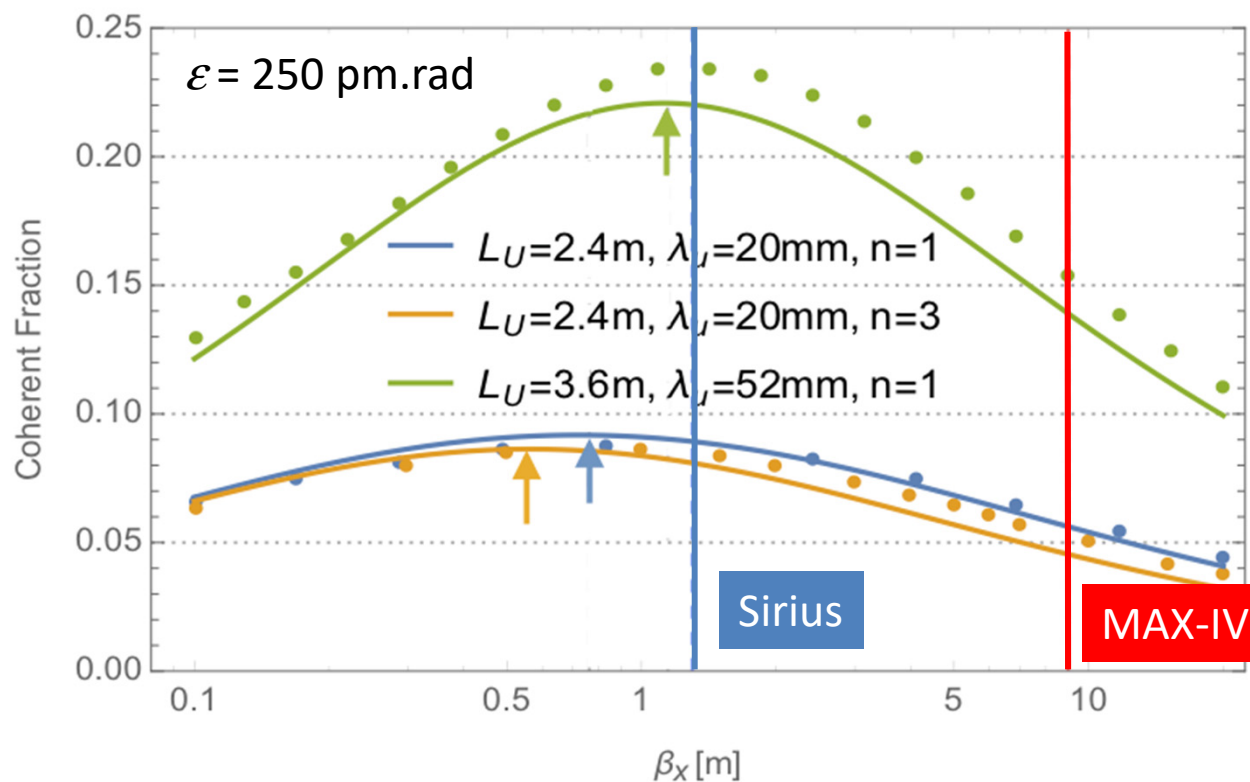
- Numerical integration of Wigner Distribution Function
- Gaussian approximation of reference [H. Westfahl Jr *et al*, JSR, 24, 2017]

Courtesy: Harry Westfahl Jr

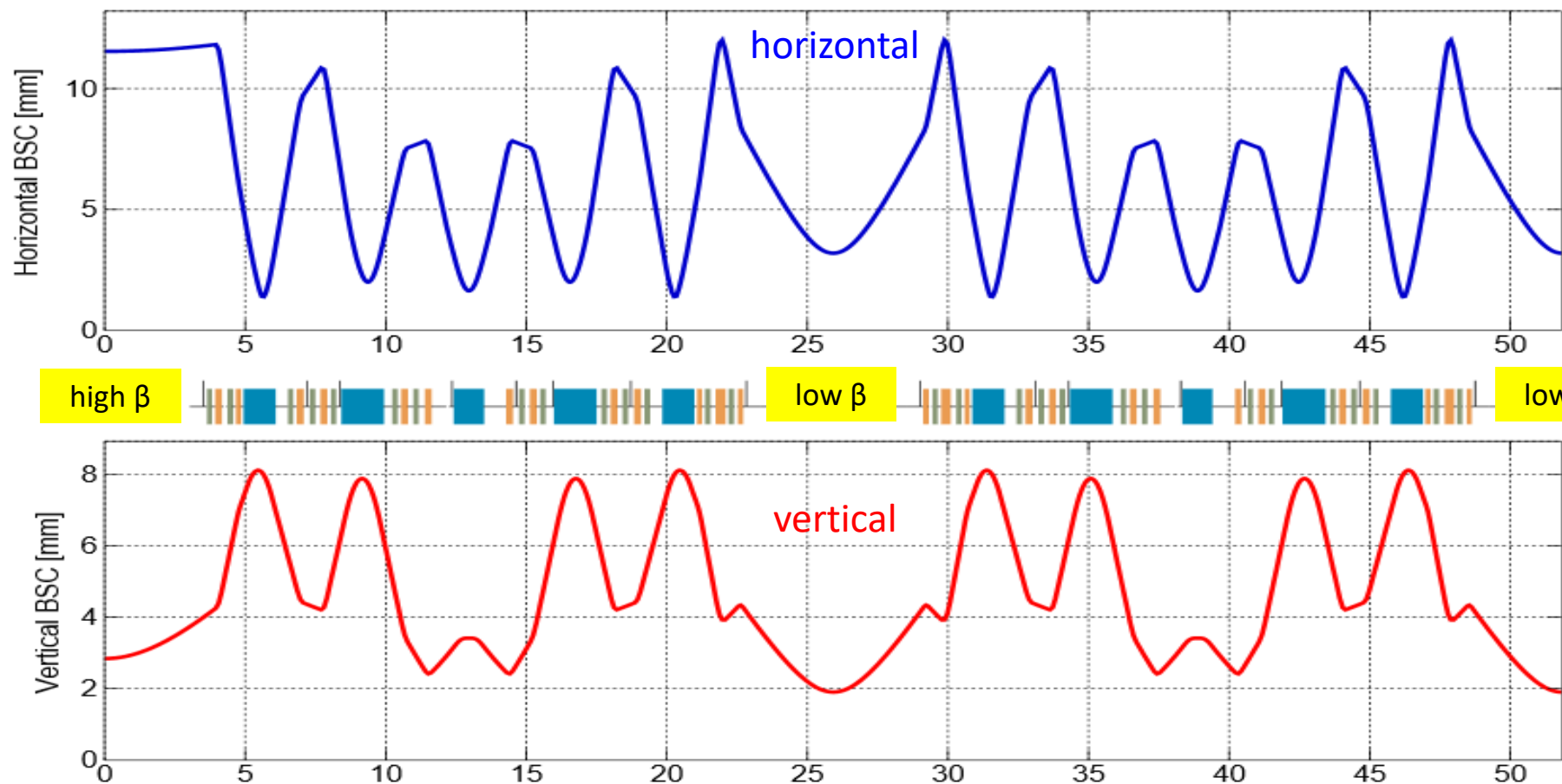


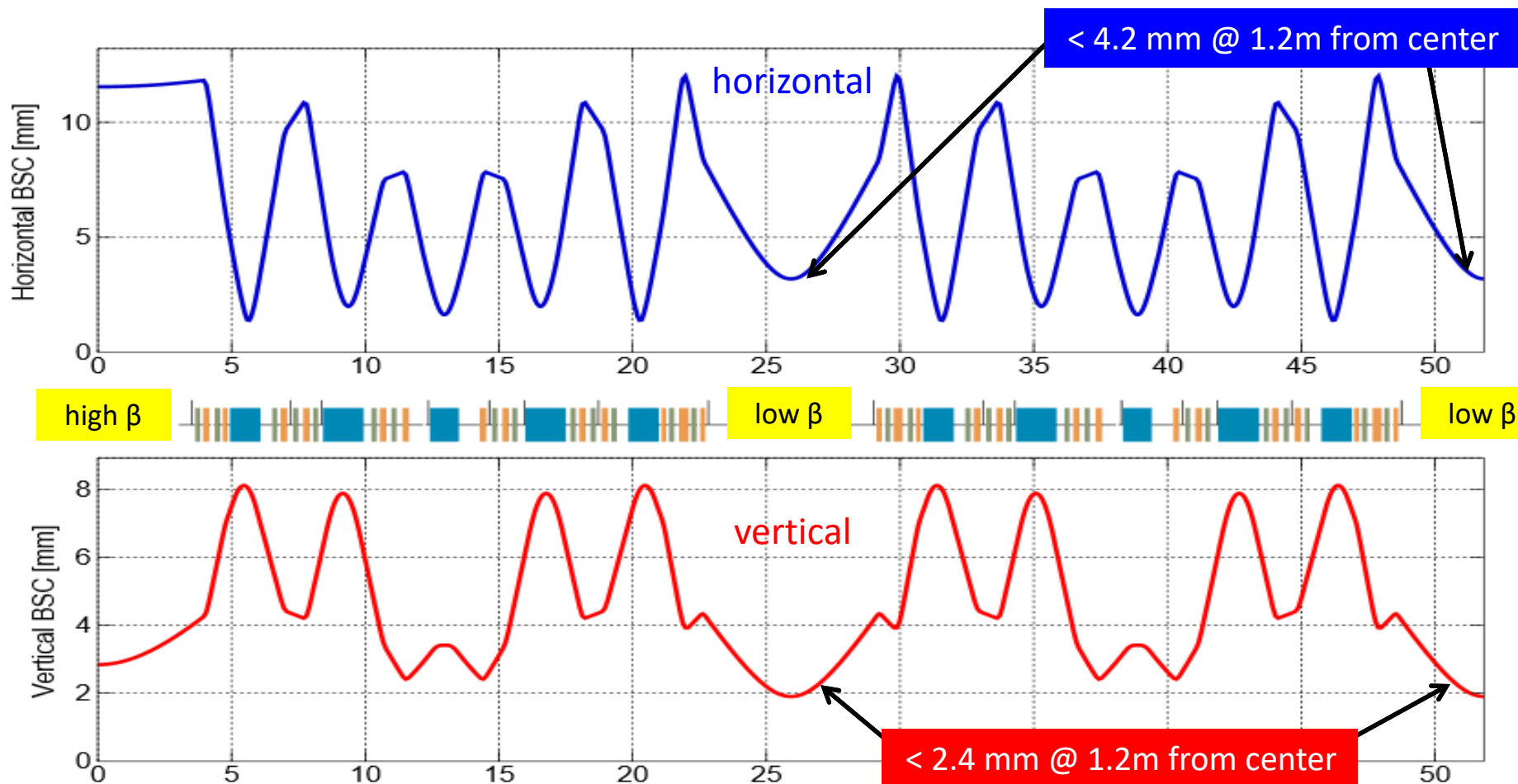
- Numerical integration of Wigner Distribution Function
- Gaussian approximation of reference [H. Westfahl Jr *et al*, JSR, 24, 2017]

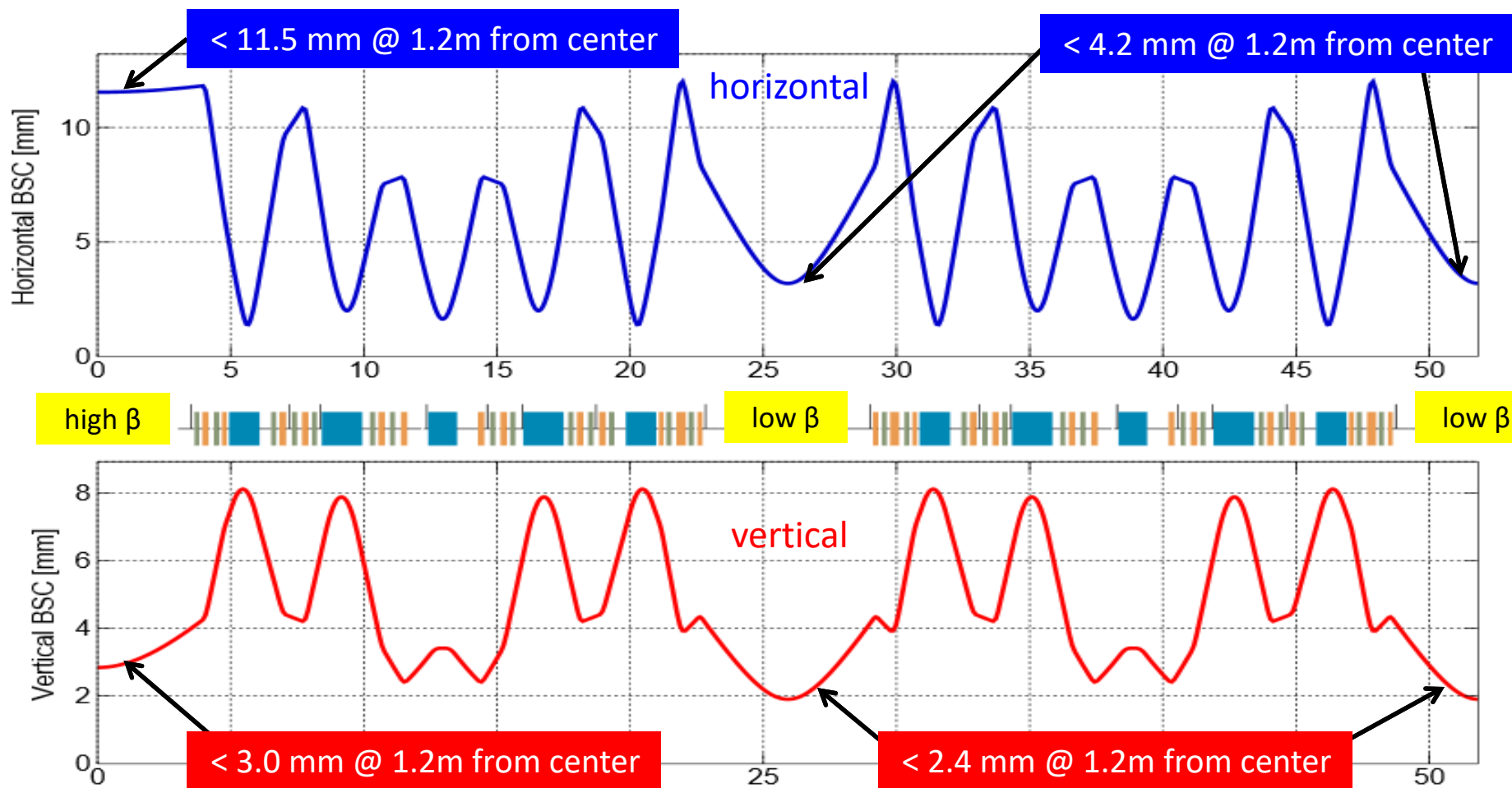
Courtesy: Harry Westfahl Jr



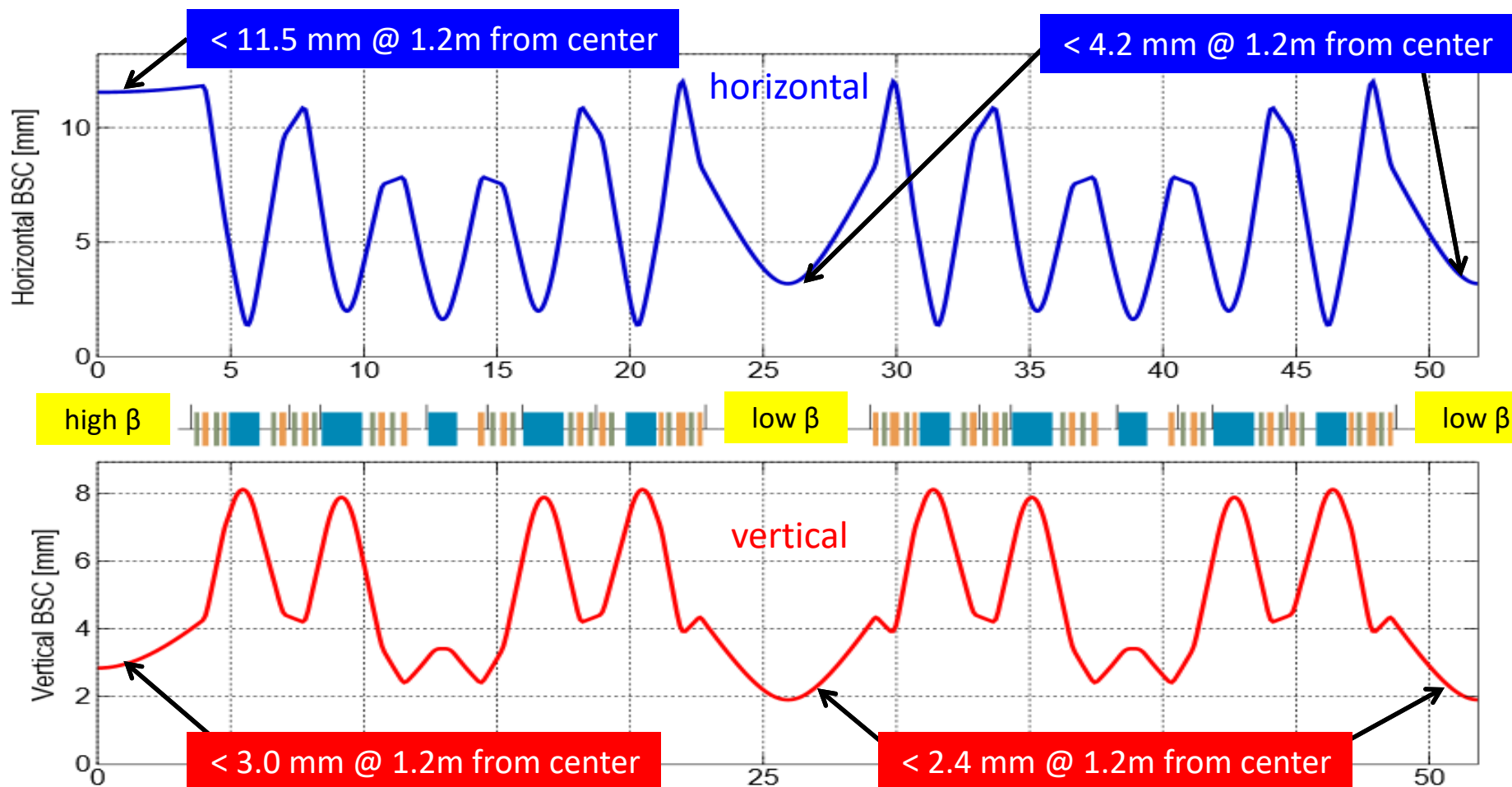
- Numerical integration of Wigner Distribution Function
- Gaussian approximation of reference [H. Westfahl Jr *et al*, JSR, 24, 2017]





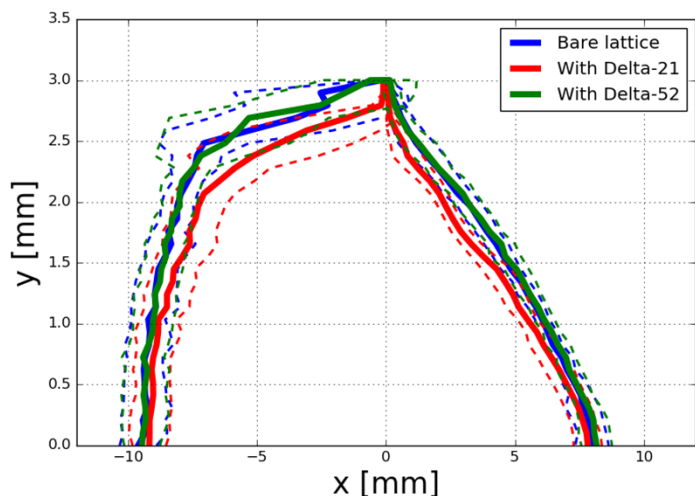


Sirius IDs will be based on Delta and APU undulators.

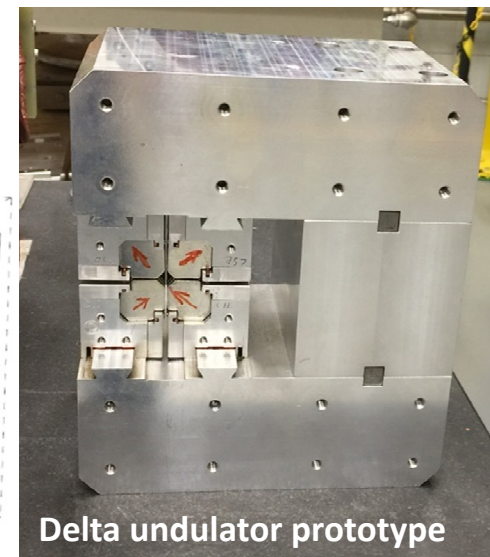
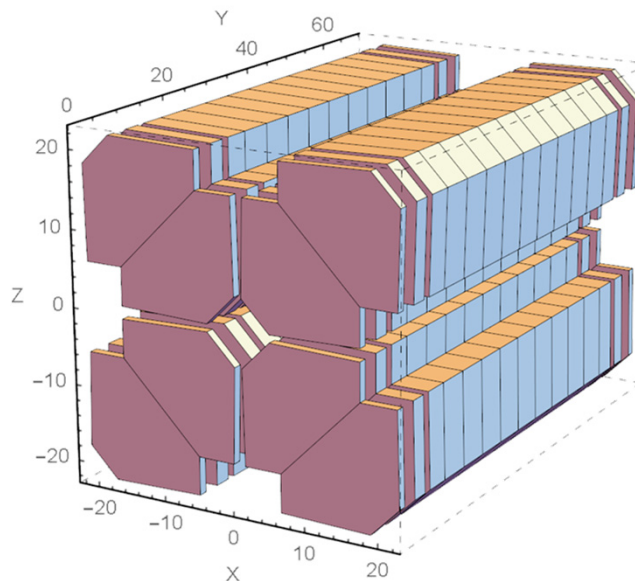
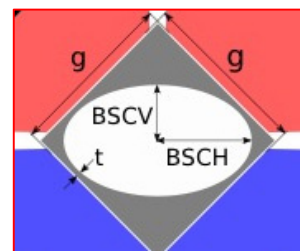


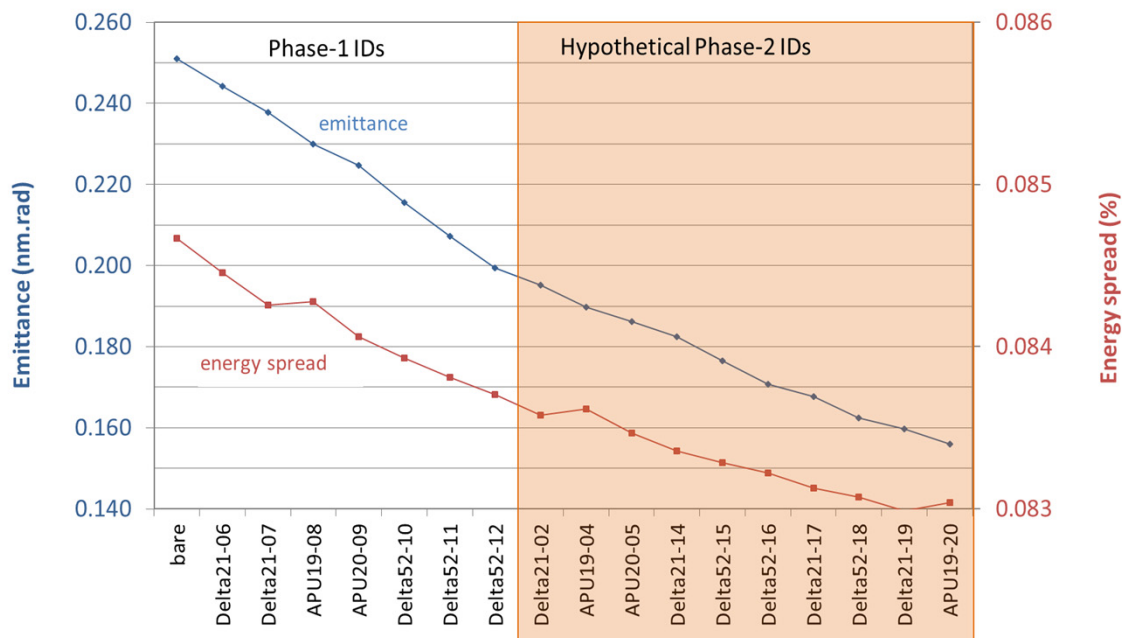
- Delta Undulators
 - Possible for Sirius due to small Hor. BSC
 - Smoother K changes
 - Horizontal, Vertical and Circular x-ray polarizations on the same energy range
 - Do not introduce strong harmful multipoles

	B_0 [T]	λ [mm]	L [m]	K_{max}	Diag. gap [mm]
Delta21	1.12	21	2.4	2.2	7.0
Delta52	1.19	52	3.6	5.85	13.85



Effect of 14 Delta21 and 14 Delta52 on DA
 7 in each polarization (H/V)
 distributed in low beta sections



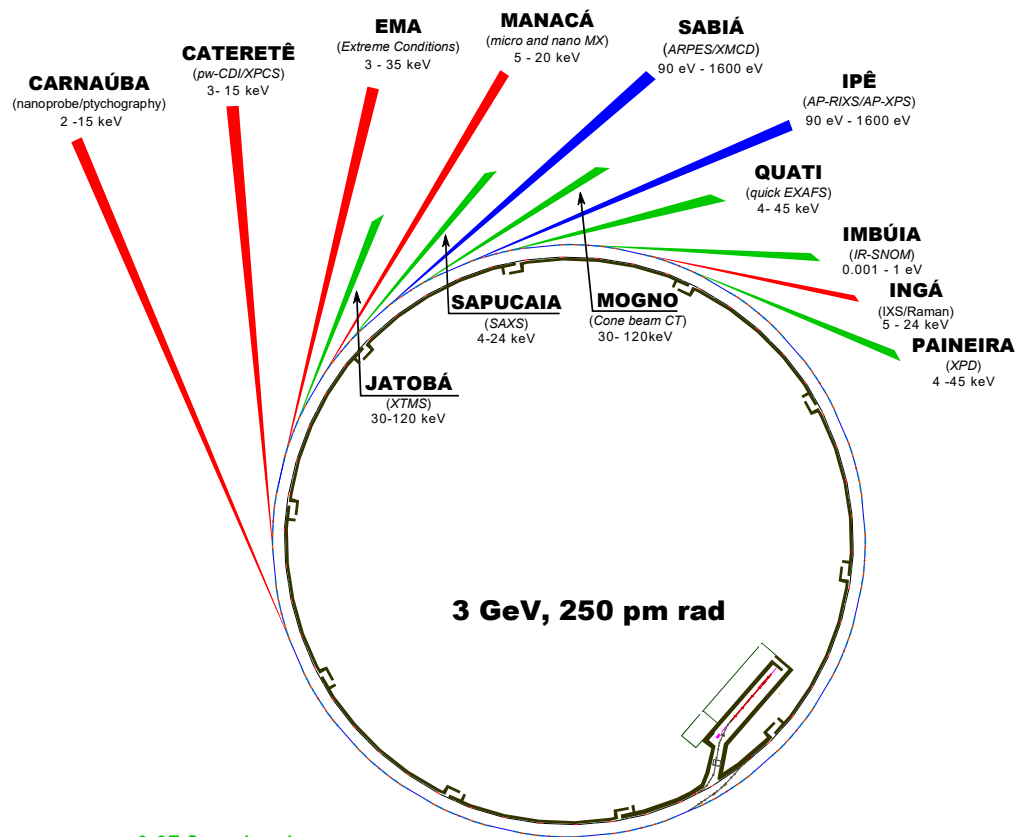


Sirius Phase-1 Beamlines

Beamline	ID Type	SS	β_x
CARNAÚBA	Delta21	06	low
EMA	APU19	08	low
CATERETÊ	Delta21	07	low
IPÊ	Delta52	11	low
SABIÁ	Delta52	10	low
MANACÁ	APU20	09	high
PGM++	Delta52	12	low

Sirius IDs

ID Type	B_0 [T]	λ [mm]	L [m]	K_{max}	gap [mm]
Delta21	1.12	21	2.4	2.2	6.92
Delta52	1.19	52	3.6	5.85	13.85
APU19	1.28	19	2.4	2.3	5.0
APU20	1.07	20	2.4	2.0	6.2

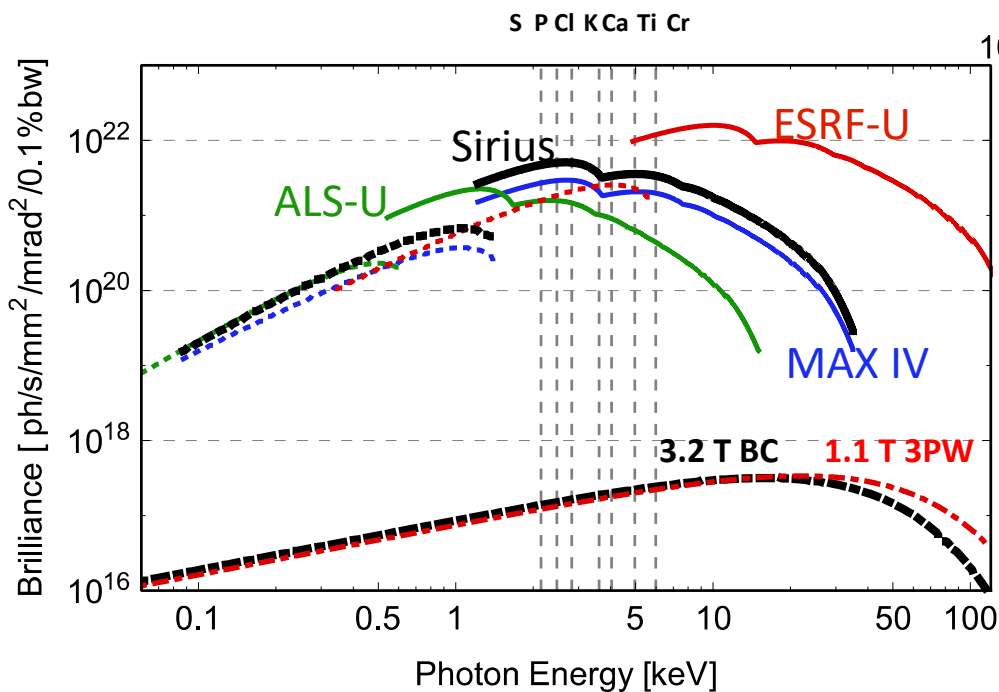
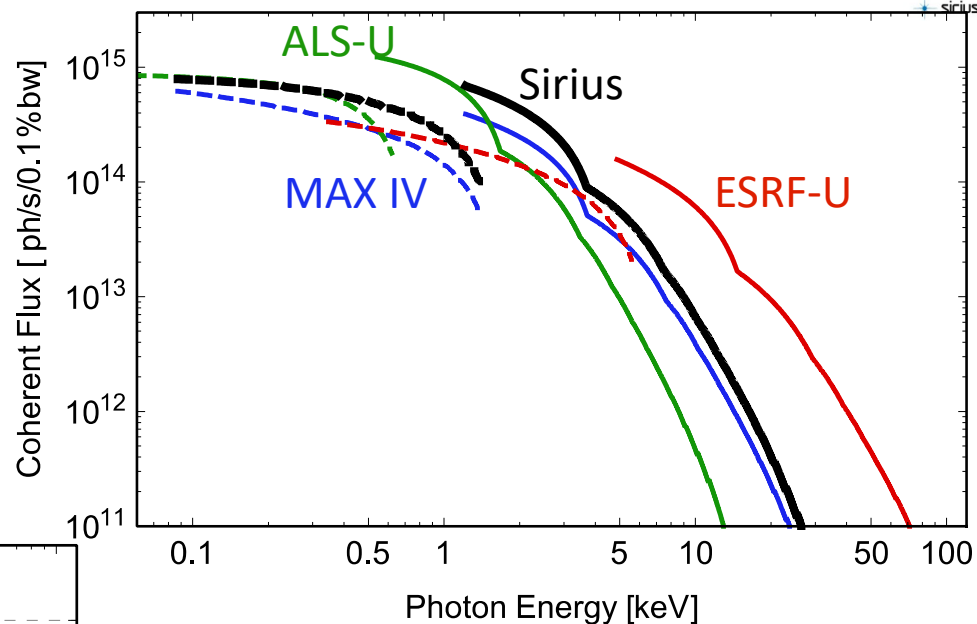


— 3.2T Superbend
— Short period undulator
— Long period undulator

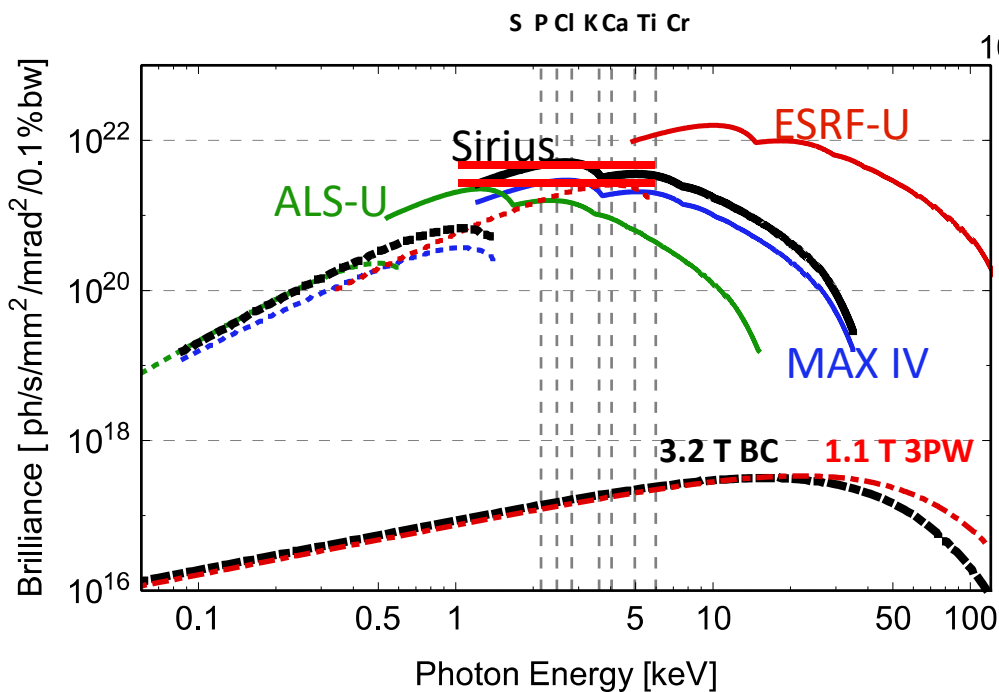
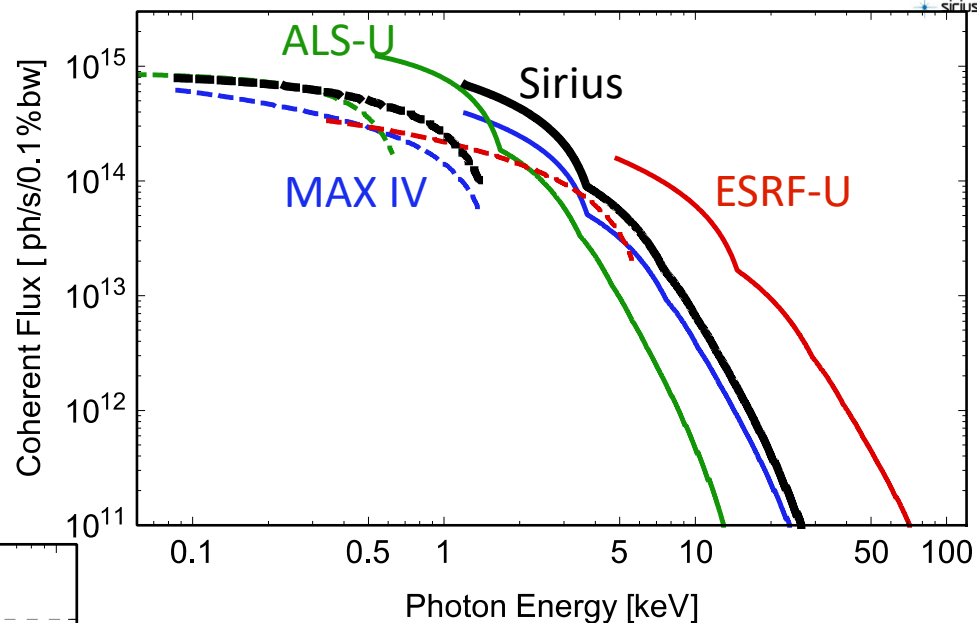
— Experimental Programs

- Tender nano-probe for spectro-ptychography
- Large FOV (30 μm) Coherent Diffraction Imaging
- Bragg CDI/XRD/XAFS under extreme conditions
- Serial micro and nano MX
- Tender x-ray RIXS
- AP-RIXS/XPS
- ARPES/PEEM
- Cone beam High Energy Tomography
- Quick-EXAFS
- 3D X-Ray Diffraction Microscopy
- High-Throughput SAXS
- Time Resolved Powder Diffraction
- nano-FTIR

- Even with future upgrades, Sirius will be competitive in the energy range of tender X-rays.

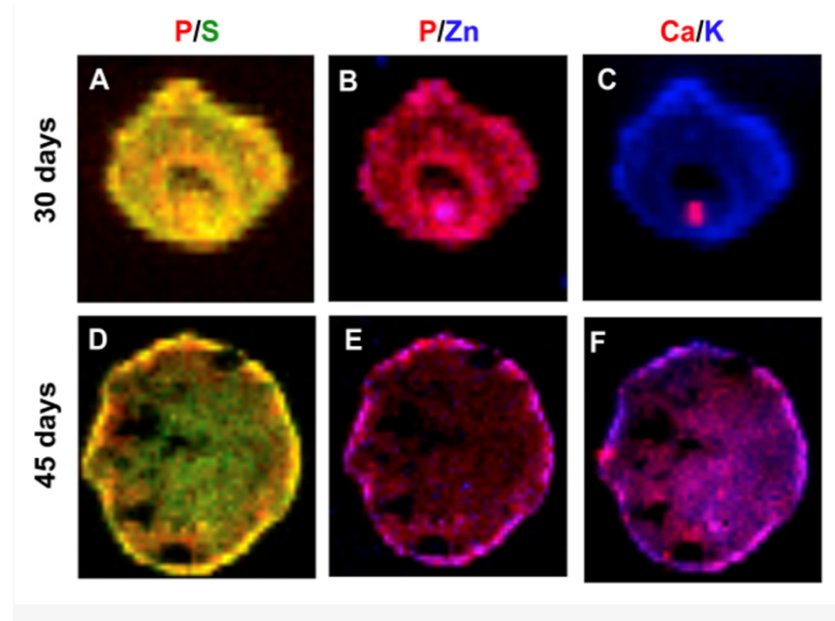
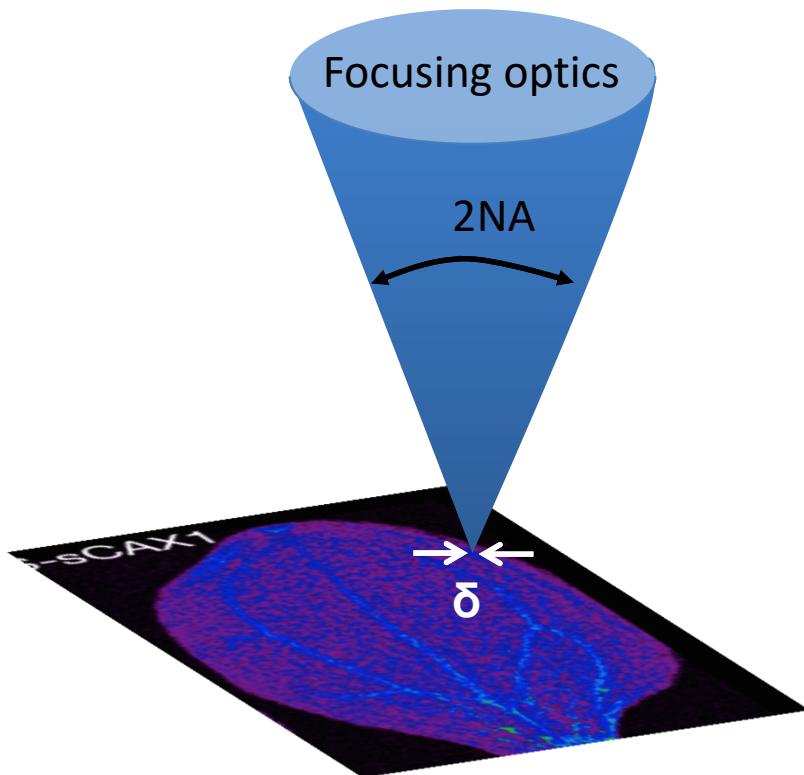


- Even with future upgrades, Sirius will be competitive in the energy range of tender X-rays.



factor ~ 2 comes from betatron function matching

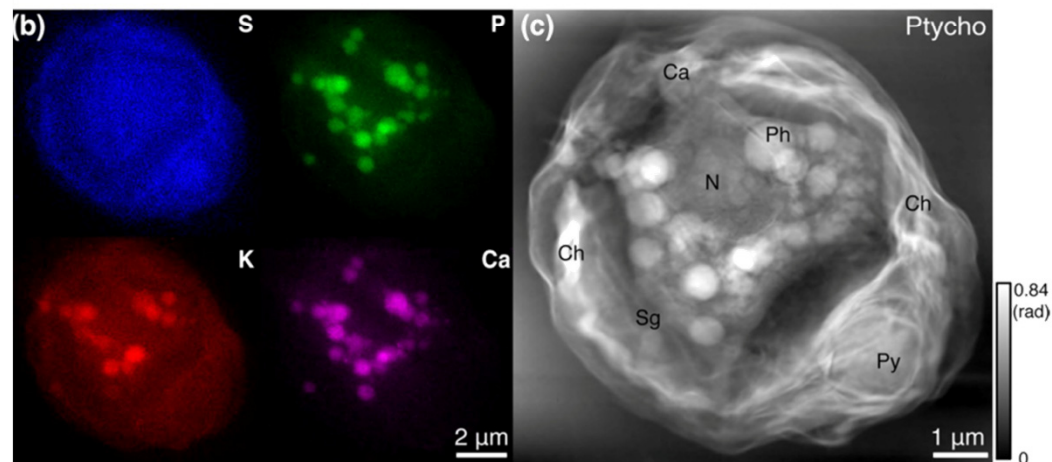
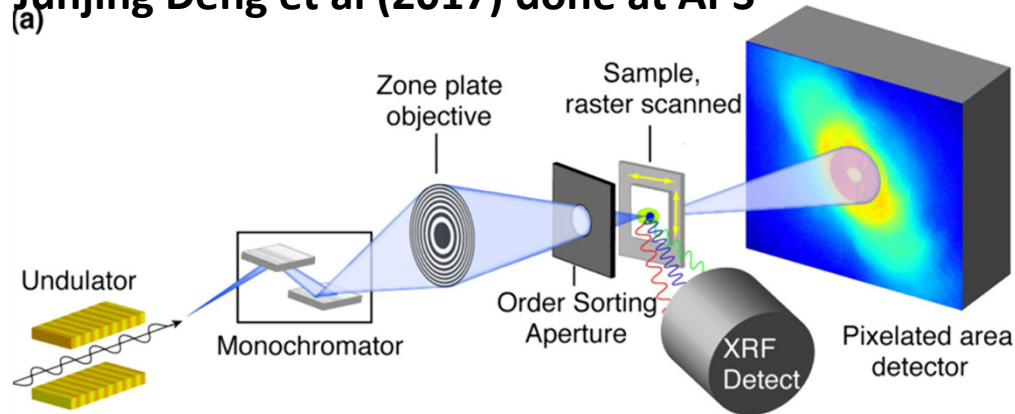
$$\text{Photon Flux} \sim \text{Brilliance} (NA \times \delta)^2$$



@ LNLS today: micronutrients during brain Formation 1 mm cerebral organoids
 Rafaela C. Sartore *et al.* (2017)

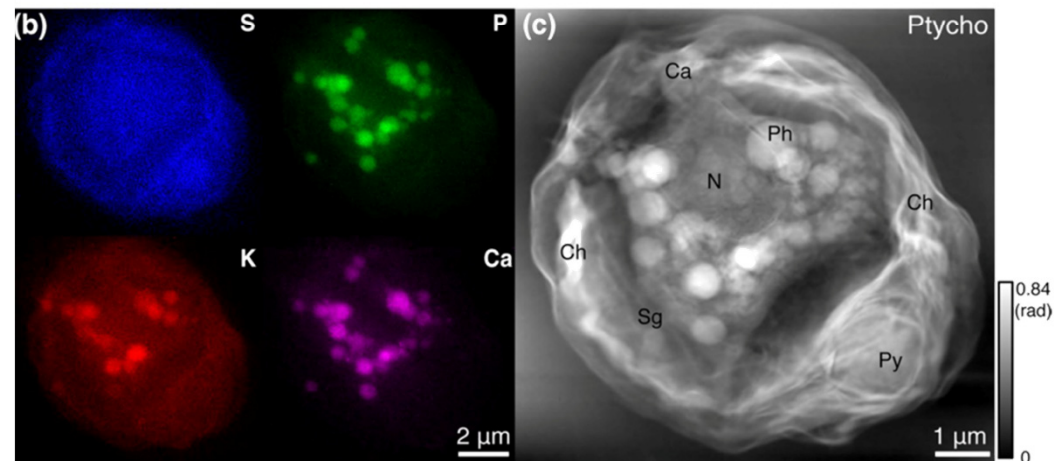
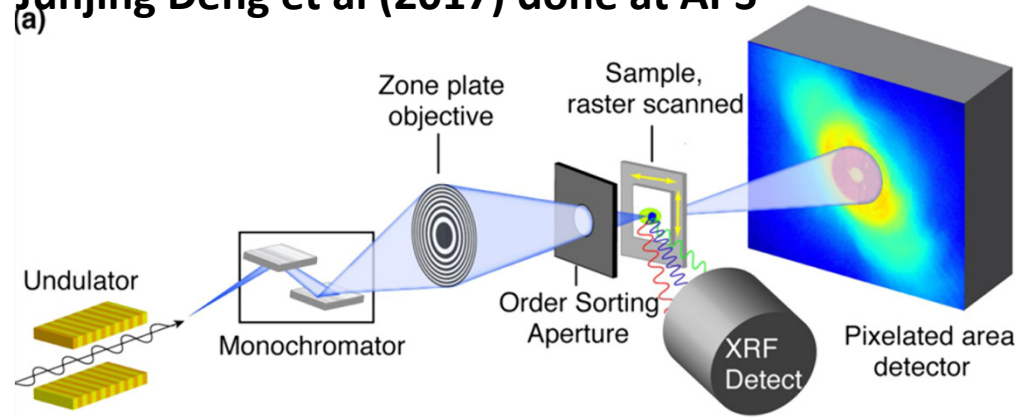
20 μm (@ LNLS today) \longrightarrow **20 nm** (@ Sirius in 2019)

Junjing Deng et al (2017) done at APS



green algae Chlamydomonas reinhardtii
model cell for studying photosynthesis

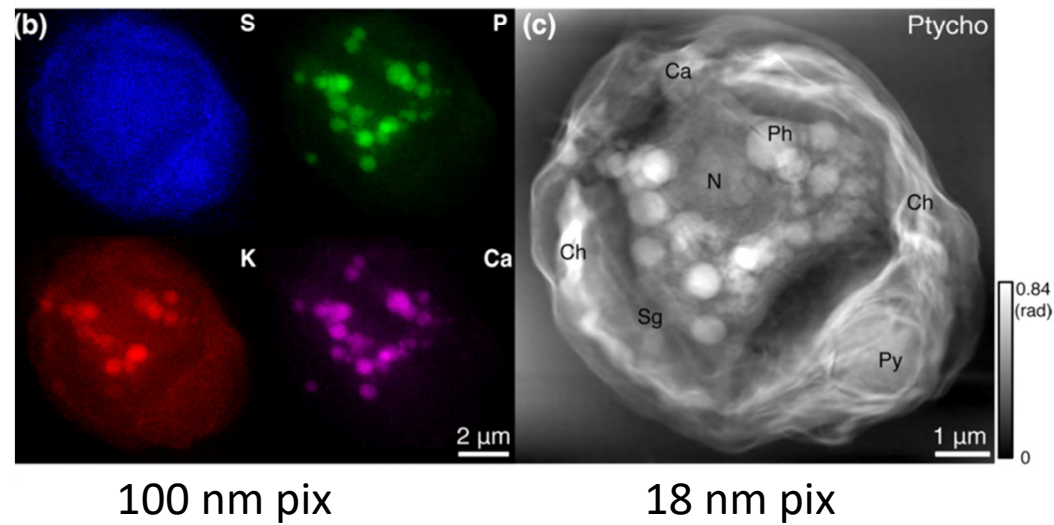
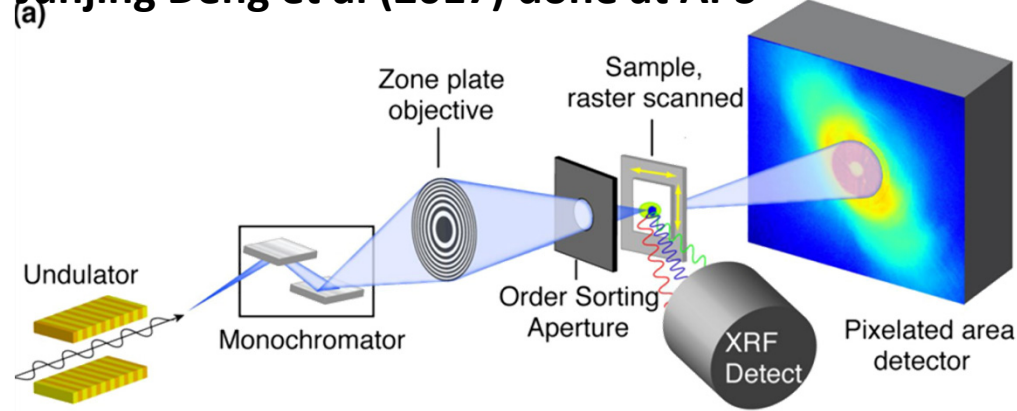
(a) Junjing Deng et al (2017) done at APS



100 nm pix

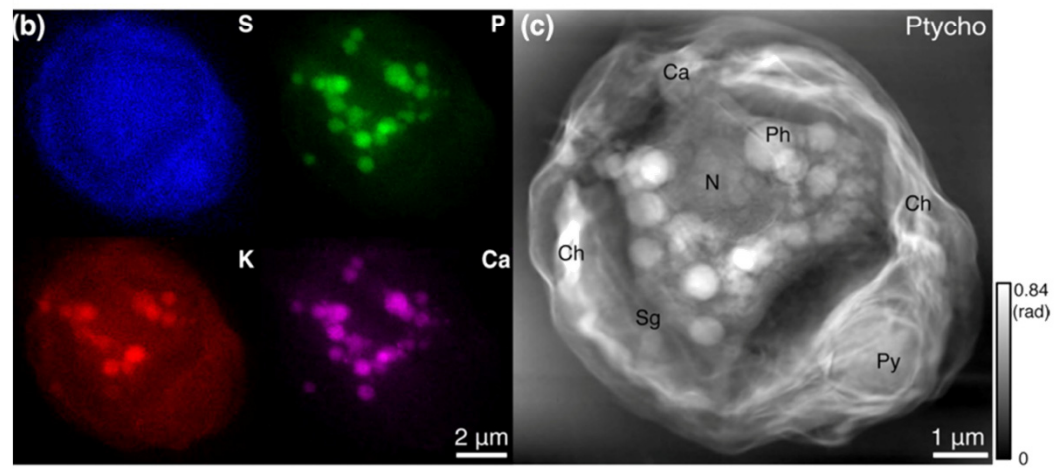
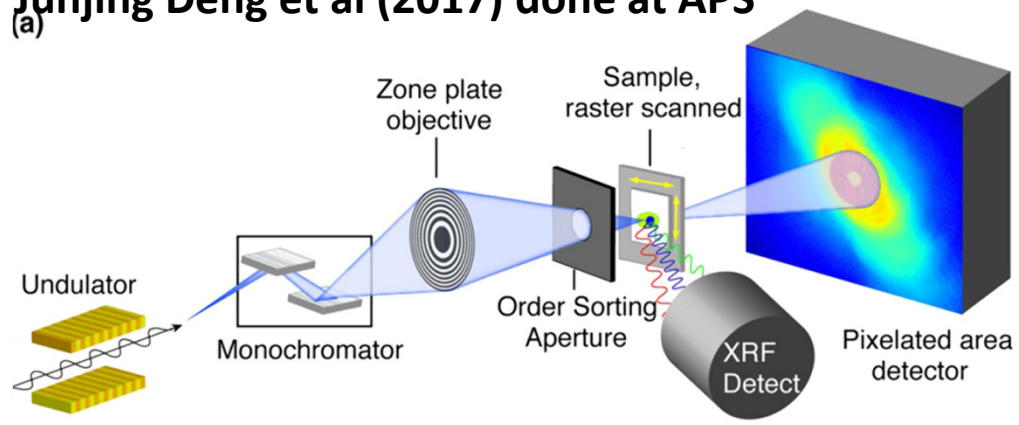
green algae Chlamydomonas reinhardtii
model cell for studying photosynthesis

Junjing Deng et al (2017) done at APS



green algae Chlamydomonas reinhardtii
model cell for studying photosynthesis

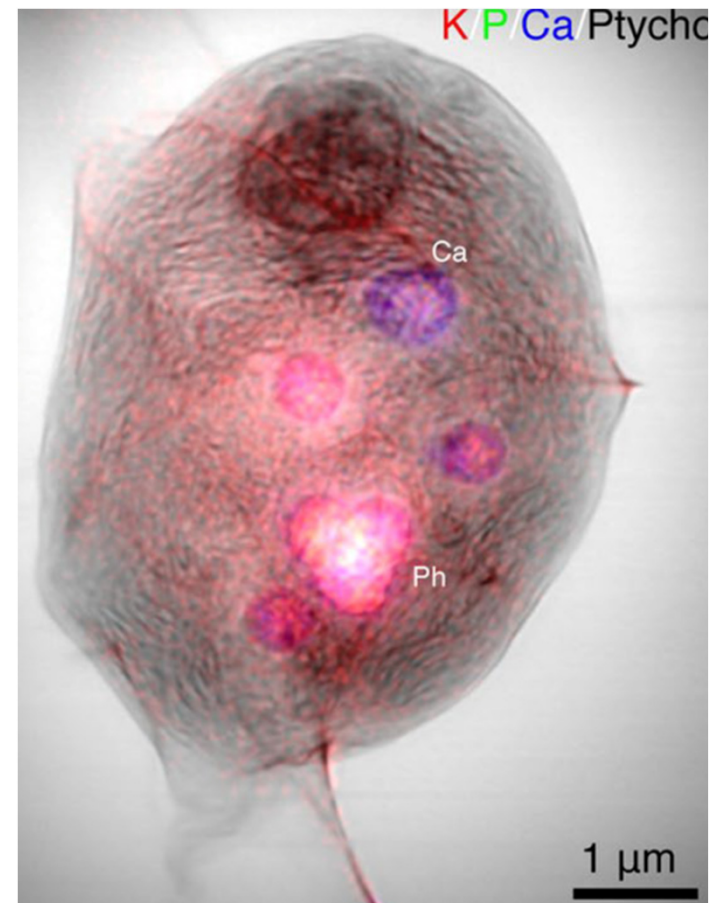
(a) Junjing Deng et al (2017) done at APS



100 nm pix

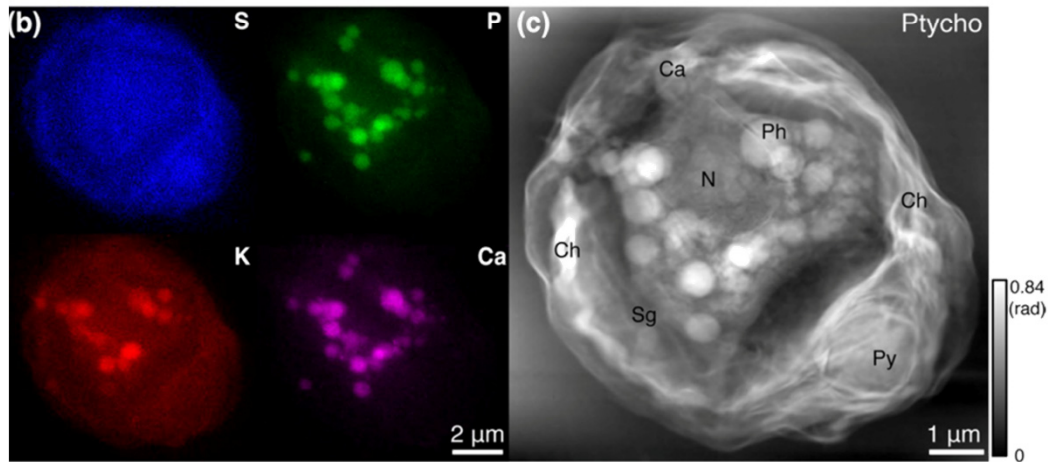
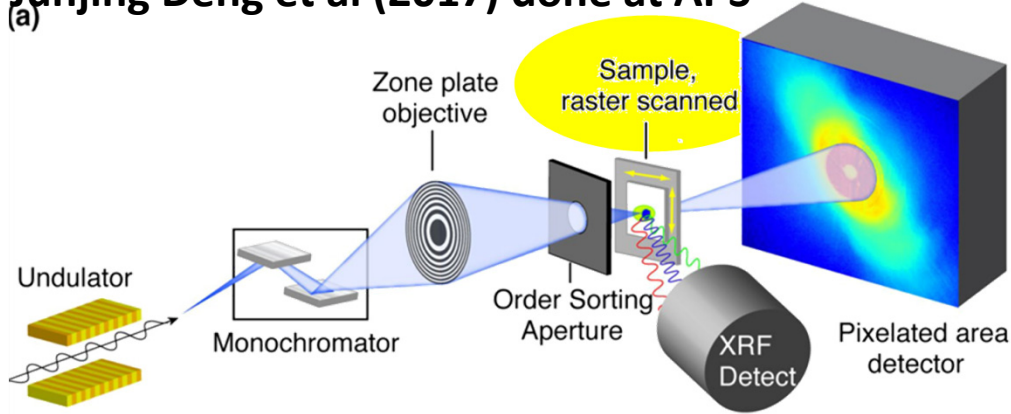
18 nm pix

green algae Chlamydomonas reinhardtii
model cell for studying photosynthesis



Overlay of the two measurements

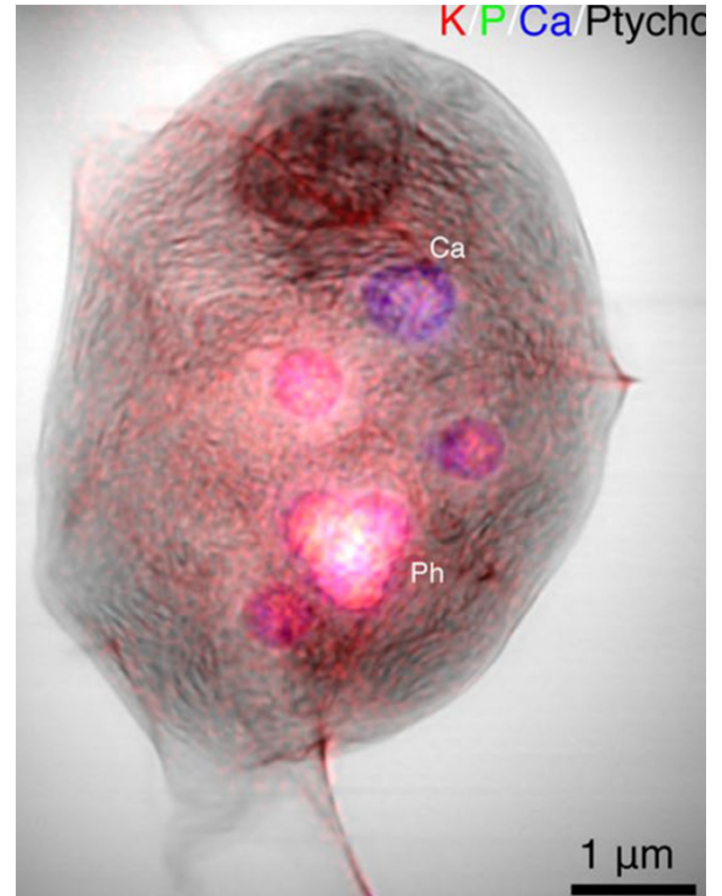
(a) Junjing Deng et al (2017) done at APS



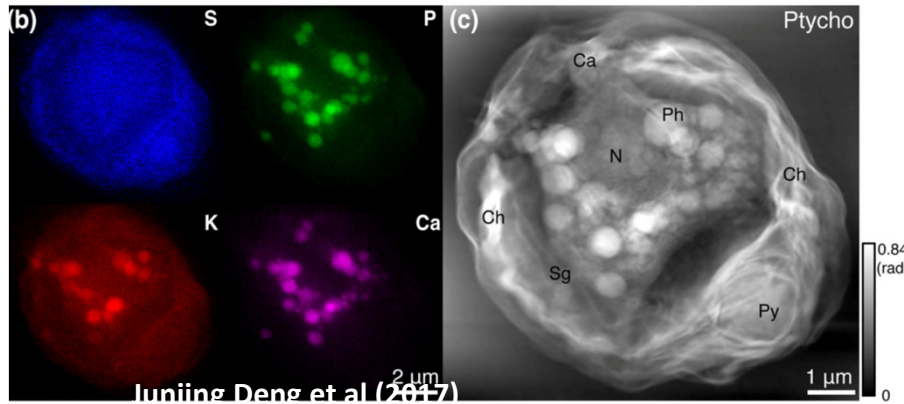
100 nm pix

18 nm pix

green algae Chlamydomonas reinhardtii
model cell for studying photosynthesis



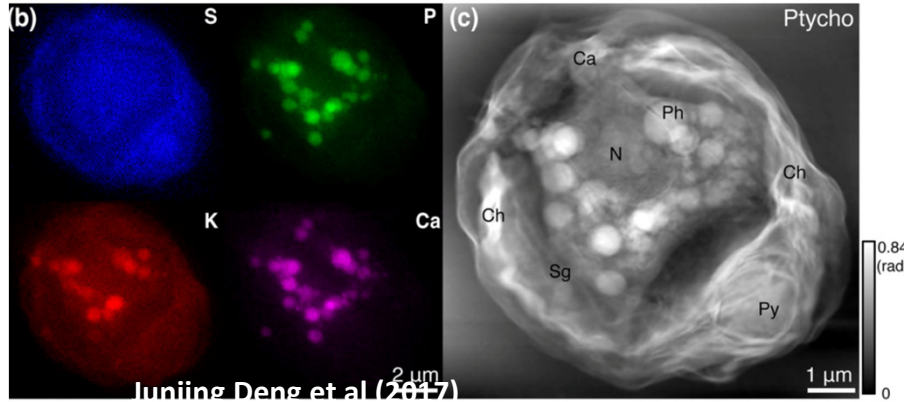
Overlay of the two measurements



Fluorescence: ~100 nm pix Lenseless imaging: ~18 nm pix

@ APS today:
Acquisition time (2D):
 ~0.1 s/pixel & ~75 min/image
 ~10⁴ photons/nm² at 5.2 keV

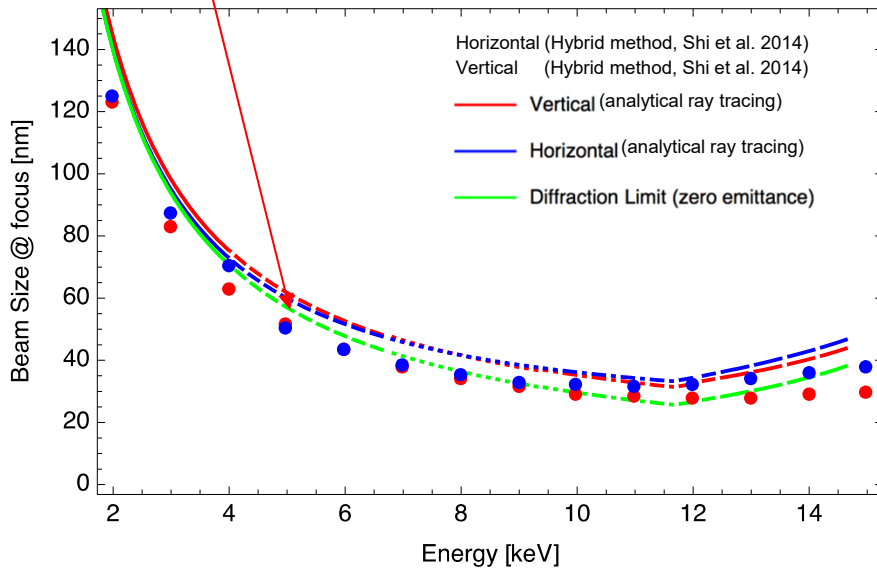


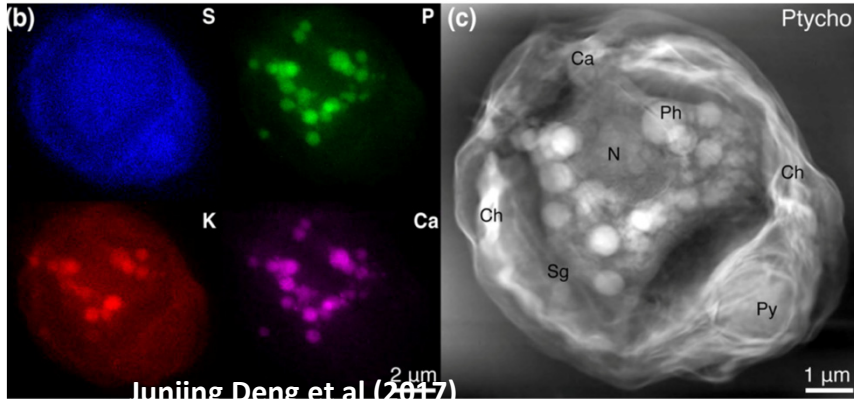


Sirius
50 nm pix

@ APS today:
Acquisition time (2D):
 ~0.1 s/pixel & ~75 min/image
 ~10⁴ photons/nm² at 5.2 keV

Fluorescence: ~100 nm pix Lenseless imaging: ~18 nm pix





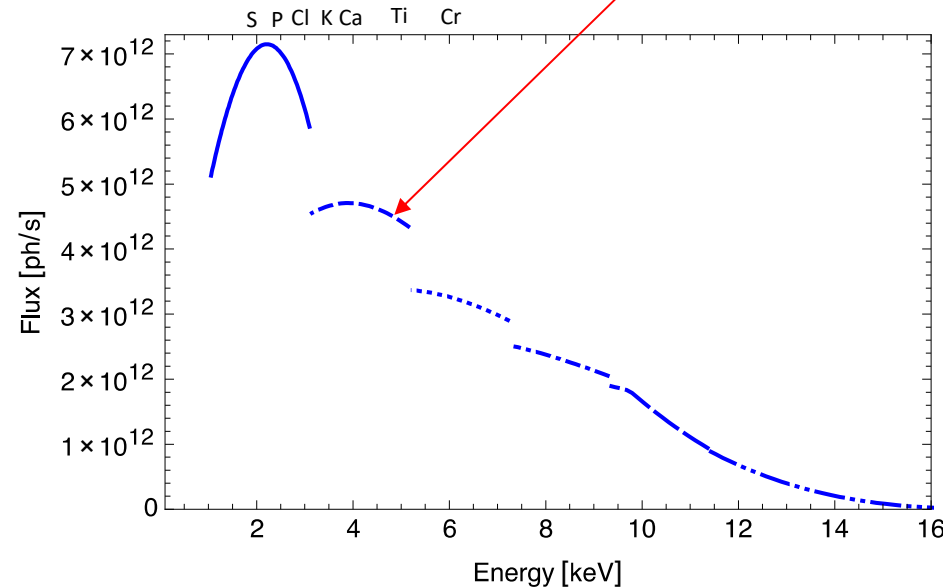
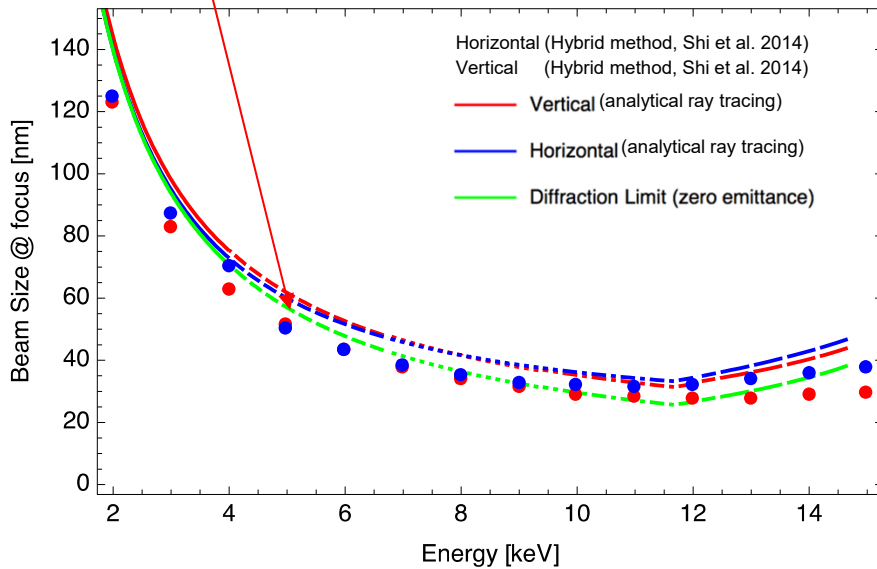
Sirius
50 nm pix

Fluorescence: ~ 100 nm pix Lenseless imaging: ~ 18 nm pix

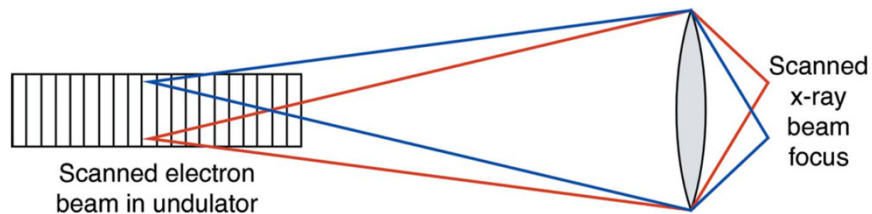
@ APS today:

Acquisition time (2D):
 ~ 0.1 s/pixel & ~ 75 min/image
 $\sim 10^4$ photons/nm² at 5.2 keV

Sirius
 $\sim 10^8$ photons/nm² at 5.2 keV
 ~ 10 μ s/pixel & ~ 1 s/image

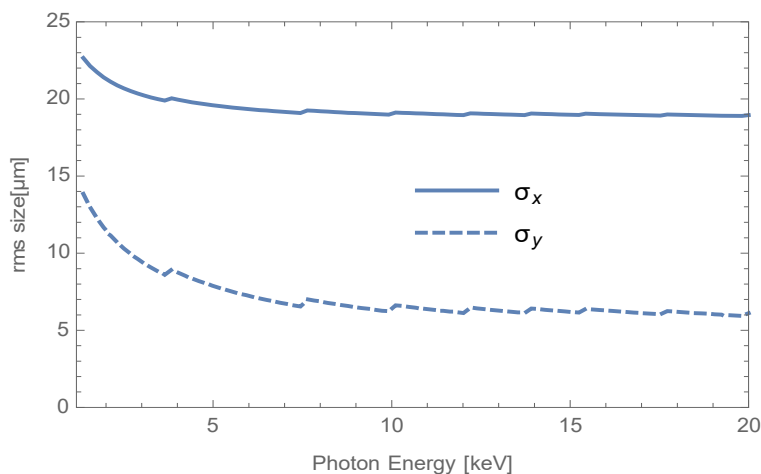
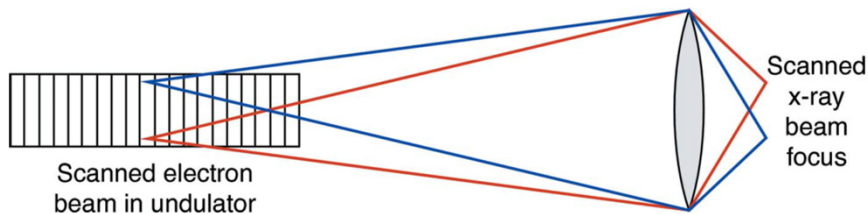


M. D. de Jonge et al. J. Sync. Rad. (2014)



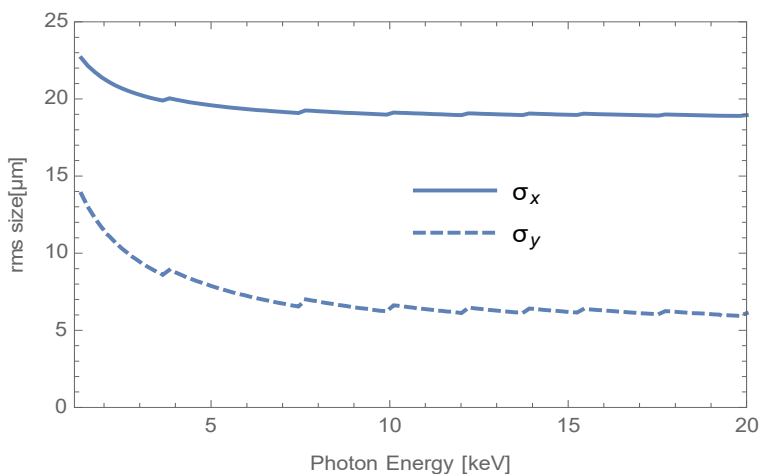
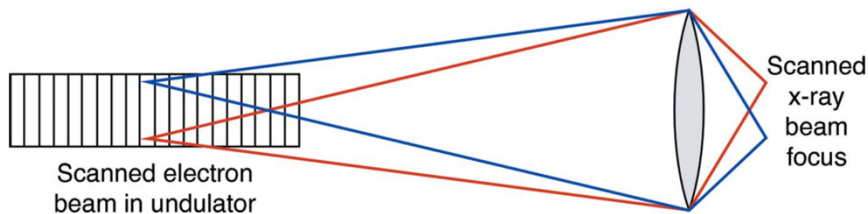
- Scanning the source position with corrector strengths of $\pm 400\mu\text{rad}$ can result in scanning ranges of $\pm 400\mu\text{m}$ in each direction.

M. D. de Jonge et al. J. Sync. Rad. (2014)

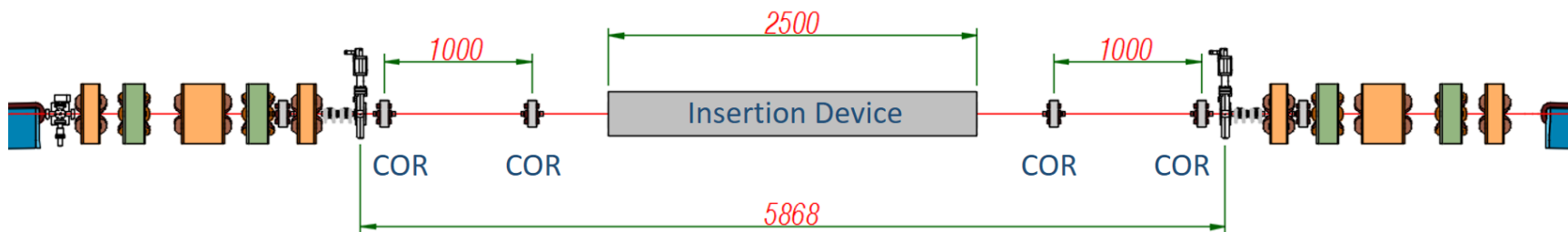


- Scanning the source position with corrector strengths of $\pm 400 \mu\text{rad}$ can result in scanning ranges of $\pm 400 \mu\text{m}$ in each direction.
- Step sizes of σ_x and σ_y result in $\sim 40 \times 100$ overlapping scanning points for ptychography.

M. D. de Jonge et al. J. Sync. Rad. (2014)



- Scanning the source position with corrector strengths of $\pm 400\mu\text{rad}$ can result in scanning ranges of $\pm 400\mu\text{m}$ in each direction.
- Step sizes of σ_x and σ_y result in $\sim 40 \times 100$ overlapping scanning points for ptychography.
- Local beam bumps can be created with 4 correctors in the low beta straight section.



- The Synchrotron Radiation Light Source community is going through a very exciting time, with many new developments under way both in the machine and scientific application sides. Many new machines and machine upgrades are expected for next years.

- The Synchrotron Radiation Light Source community is going through a very exciting time, with many new developments under way both in the machine and scientific application sides. Many new machines and machine upgrades are expected for next years.

If you are a student or a young person starting in this field:
you'll have lots of fun! 😊

- The Synchrotron Radiation Light Source community is going through a very exciting time, with many new developments under way both in the machine and scientific application sides. Many new machines and machine upgrades are expected for next years.

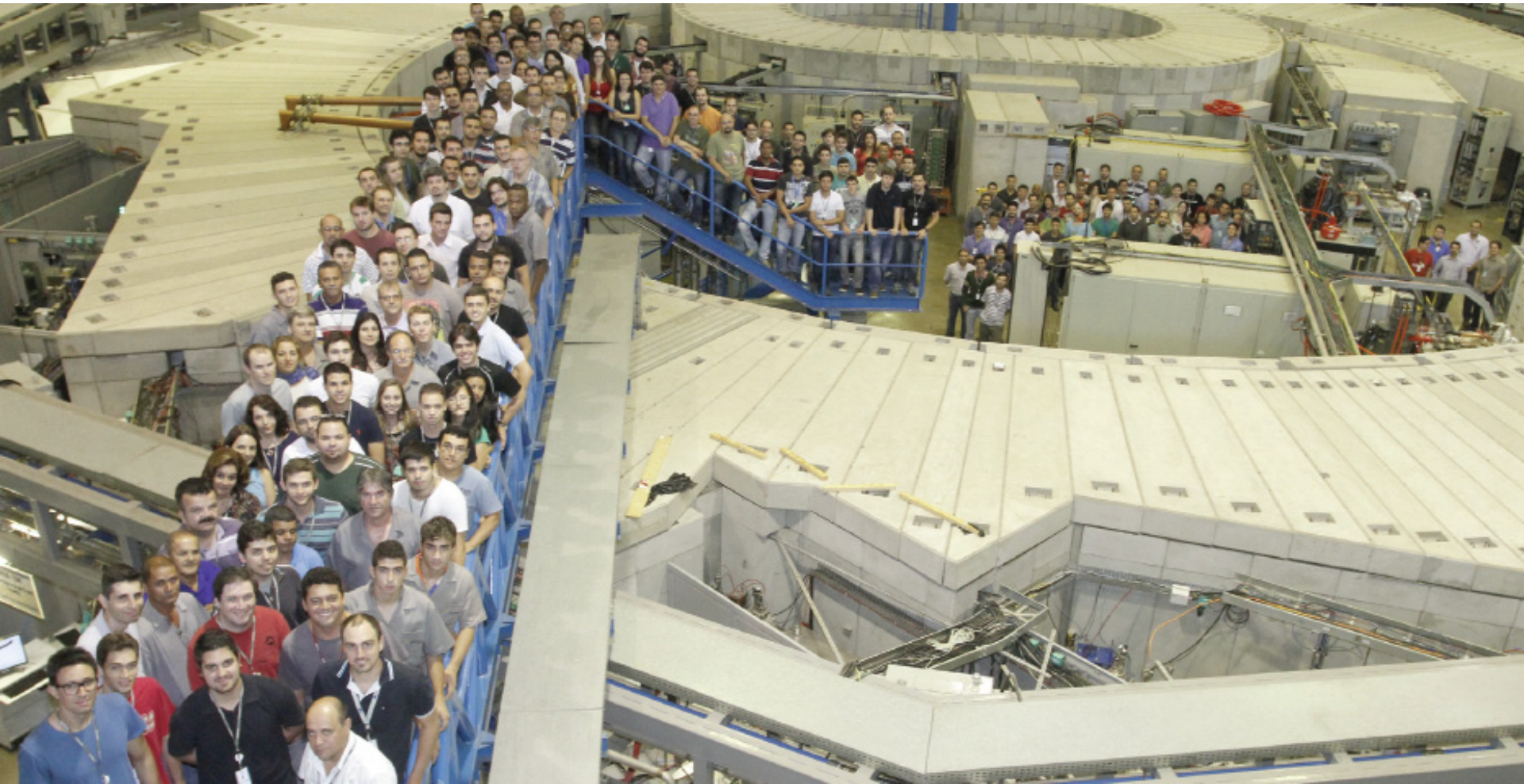
If you are a student or a young person starting in this field:
you'll have lots of fun! 😊

- It is important to integrate machine and beamline teams in the optimization of experiments.

- The Synchrotron Radiation Light Source community is going through a very exciting time, with many new developments under way both in the machine and scientific application sides. Many new machines and machine upgrades are expected for next years.

If you are a student or a young person starting in this field:
you'll have lots of fun! 😊

- It is important to integrate machine and beamline teams in the optimization of experiments.
- This is an open community and international cooperation is one of the most important sources for learning and advancing in this area.



Sirius Team – a small but highly motivated and integrated team!