

DESIGN AND TEST OF WIRE-SCANNERS FOR SwissFEL

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Abstract

The SwissFEL light-facility will provide coherent X-rays in the wavelength region 7-0.7 nm and 0.7-0.1 nm. In SwissFEL, view-screens and wire-scanners will be used to monitor the transverse profile of a 200/10 pC electron beam with a normalized emittance of 0.4/0.2 mm mrad and a final energy of 5.8 GeV. Compared to view screens, wire-scanners offer a quasi-non-destructive monitoring of the beam transverse profile without suffering from possible micro-bunching of the electron beam. The main aspects of the design, laboratory characterization and beam-test of the SwissFEL wire-scanner prototype will be presented.

INTRODUCTION

SwissFEL will provide coherent X-rays light in the wavelength region 7-0.7 nm and 0.7-0.1 nm [1]. Electron bunches with charge 200/10 pC and normalized emittance of 0.4/0.2 mm.mrad will be emitted by a photocathode at a repetition rate of 100 Hz according to a two-bunches train structure with a temporal separation of 28 ns. Thanks to a RF kicker switching the second electron bunch of the beam train into a magnetic switch-yard, the SwissFEL linac will simultaneously supply two distinct undulator chains at a repetition rate of 100 Hz: the hard-Xrays line Aramis and the soft-Xrays line Athos. The electron beam will be accelerated up to 330 MeV by a S-band RF Booster and to the final energy of 5.8 GeV by a C-band RF linac. Thanks to an off-crest acceleration in the RF Booster, the electron beam will experience a two-stages longitudinal compression from an initial bunch length of 3/1 ps (RMS) down to 20/3 fs (RMS) in two magnetic chicanes. Two X-band RF cavities will compensate the quadratic distortion of the longitudinal phase space due to the off-crest accelerating scheme of the beam and the non-linear contribution of the magnetic dispersion. A laser-heater in the Booster section will smooth down possible micro-structures affecting the beam longitudinal profile of the beam. Macro-bunching can be detrimental to the monitoring of the beam profile based on scintillator or OTR screens (Optical Transition Radiator) because of the emission of coherent OTR. As an alternative to view screens, wire-scanners (WSC) can be used to monitor the beam transverse profile. Moreover, the quasi-non-invasive feature of the WSCs - compared to view screens - can be beneficial to monitoring the beam transverse profile during FEL operations of the machine. In the following, results on design, characterization and beam test of wire-scanners for SwissFEL will be presented.

WSC DESIGN

Wire-Scanners can be used to measure the transverse profile of the electron beam in a particle accelerator [2–4]. Carbon or metallic wires with different diameter (D) - stretched on a wire-fork - can be vertically inserted at a constant velocity into the vacuum chamber by means of a motorized UHV linear stage to scan the beam transverse profile with an intrinsic resolution $D/4$ (rms). An encoder mounted on the linear stage allows the relative distance of the wire from the axis of the vacuum chamber to be measured at each machine trigger event. The interaction of the electron beam with the wire produces a "wire-signal" - scattered primary electrons and secondary particles (mainly electrons, positron and bremsstrahlung photons) - which is proportional to the number of the electrons sampled by the wire in the bunch. The Beam Synchronized Acquisition (BS-ACQ) - over a sufficient number of machine trigger shots - of the wire position and the wire-signal - detected by a loss monitors downstream the wire - allows the beam transverse profile along the horizontal or the vertical direction to be reconstructed. In SwissFEL, view screens and WSCs will be used to monitor the transverse profile of the electron beam which varies between $500\ \mu\text{m}$ and $5\ \mu\text{m}$ (rms) along the entire machine. View-screens will be mainly equipped with YAG crystals. In SwissFEL, only WSCs are in principle able to discriminate the 28 ns time structure of the two-bunches emitted at 100 Hz by the photocathode. The SwissFEL WSCs are designed according to the following criteria, see Fig. 1: use a single UHV linear stage to scan the beam profile in the X, Y and X-Y directions; use Tungsten wire with different diameters from 5 to $13\ \mu\text{m}$ to ensure a resolution in the range 1.5-3.5 μm ; equip each wire-scanner station with spare/different-resolution wires; detect the wire losses in the bunch charge range 10-200 pC and resolve the 28 ns time two-bunches structure of the electron beam; BS-ACQ of the read-out of both the encoder wire position and the loss-monitor; wire-fork suitably designed for routine scanning of the beam profile during FEL operations (no beam interception with the wire-fork); wire-fork equipped with different pin-slots where the wires can be stretched at different relative distances so that the scanning time can be minimized and optimized according to the WSC position in the machine, see Fig. 1. In SwissFEL the wire losses will be measured by means of scintillator fibers (PMMA, Poly-Methyl-Methyl-Acrylate, Saint Gobain BCF-20, emission in the green) winding up the beam pipe. The fiber are directly connected to a photomultiplier (PMT, Hamamatsu H10720-110). For more information on the detection sys-

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tem of the wire-signal in SwissFEL, see these conference proceedings [5].

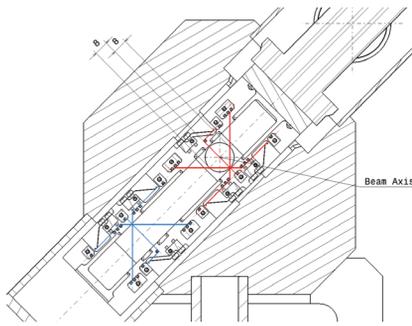


Figure 1: View of the Transverse section of the WSC vacuum chamber and, in particular, of the wire-fork and of the CF16 vacuum chamber. The wire-fork is equipped with 3 different pin-slots where the wire can be stretched. The distance of the "wire-vertex" from the vacuum-chamber axis can be set at 8, 5.5 and 3 mm in correspondence of the 3 different pin-slots.

WSC SIMULATIONS

Most part of the primary and secondary particles resulting from the interaction of a relativistic electron beam with a metallic wire are emitted within a narrow cone with respect to the direction of incidence. Since, only a small fraction of the solid angle of emission of the wire-signal can be covered by a loss-monitor, the efficiency of the wire-losses detection strongly depends on the relative distance between the wire and loss-monitors. In order to determine the most suitable distance of a loss-monitor from a wire in SwissFEL, the particle composition and the related energy and angular distribution of the particle shower has been evaluated with the Monte Carlo code FLUKA [6]. In FLUKA simulations, Tungsten wires with a diameter of 13 μm and 25 μm and beam energies of 340 MeV, 1.33 GeV, 2.99 GeV and 5.2 GeV are considered. The electron beam is modelled as a pencil beam (no energy spread) travelling along the axis of a CF16 vacuum-chamber (inner diameter 16 mm, outer diameter 18mm). Both the angular and the energy distributions - expressed as number of particles per unit primary particle - of the primary and secondary particles are calculated in FLUKA, see Figs. 2,3. Electrons and photons mainly contribute to the wire-signal, see Fig. 2. Moreover, 95% of the particle shower is emitted within a polar angle less than 0.1 rad, i.e., intercepts the CF16 vacuum chamber within 80 mm and 4000 mm from the wire, see Fig. 3.

Beam tests of WSCs and loss-monitor response were carried out at the 250 MeV SwissFEL Injector Test Facility (SITF) [7]. For such a purpose, a scintillator fiber winding around the vacuum chamber was installed just before the high energy bending dipole of the machine [5]. The loss-monitor response was tested intercepting alternatively with several OTR screens - 300 μm thick Si Wafer with a 200 nm Al coating - a 180 pC electron beam with energy 245 MeV and recording the corresponding time-integration

of the fiber-signal read-out (instead of WSCs, OTR screens were used in this measurement in order to avoid any dependence of the loss-monitor response on the variation of the beam profile along the machine). The loss-monitor response measured as a function of the different OTR screen - i.e., as a function of the relative distance OTR screen loss-monitor - is shown in Fig. 4. The curve in Fig. 4 shows a maximum of the signal for a distance between OTR-screen and loss-monitor of about 2.5 m in agreement with previous measurements [8]. Taking into account the radius of the CF40 vacuum-chamber of SITF, from this measured value of the optimum distance screen-loss-monitor a mean emission polar angle of about $\theta_{Si,300 \mu\text{m}} = 7.6$ mrad can be estimated for the particle shower produced by the interaction beam-screen. In order to rescale this result to the case of a Tungsten foil, the Rossi-Greisen formula [9] can be used. The Rossi-Greisen formula allows to calculate the rms scattering angle of a charged particle experiencing multiple Coulomb interaction while crossing a block of material. The resultant average scattering angle scales down with the charge energy and is proportional to the square-root of the normalized material thickness per radiation length of the material. According to the Rossi-Greisen formula, the mean emission polar angle for Tungsten foils 13 μm and 25 μm thick can be extrapolated from the previous result $\theta_{Si,300 \mu\text{m}} = 7.6$ mrad: $\theta_{Si,300 \mu\text{m}}/\theta_{W,25 \mu\text{m}} = 0.67$ and $\theta_{Si,300 \mu\text{m}}/\theta_{W,13 \mu\text{m}} = 0.93$. Taking into account the scaling formulae above, the following optimum distances of a loss-monitor from a 25 μm and a 13 μm thick Tungsten wire can be estimated: 0.7 m and 1 m, respectively. Under the limit of applicability of this extrapolation method, an optimum distance between wire and loss-monitor of about 1m can be estimated (in agreement with FLUKA results).

WSC BENCH AND BEAM TESTS

Reliable WSC measurements require a precise knowledge of the relative position of the wire with respect to the centroid of the beam at each machine trigger event. For such a purpose, the read-out of both the encoder wire-position and the loss-monitor must be acquired in a BS-ACQ mode. The encoder read-out can provide a precise position of the wire provided that appreciable vibrations are not affecting the wire during the scan. The mechanical stability of a scanning wire can be measured in a test-bench by imaging the moving wire with a high speed camera. The wire vibration during a scan can be evaluated from the analysis of the centroid and the sigma of the projected images of the moving wire. Wire vibration measurements have been performed for different WSC set-ups, i.e., for different wire diameters, stepping motors and motor controllers. The mechanical stability of a Tungsten wire (25 μm diameter) stretched on a wire-scanner fork was measured for different velocities (0.1-10 mm/s) of a 3-phase stepping motor. Measurements were performed by imaging the wire by means of a high speed camera (camera frame rate 1kHz) and a microscope 10X (resolution 1px=1 μm). Appreciable vibrations of the wire

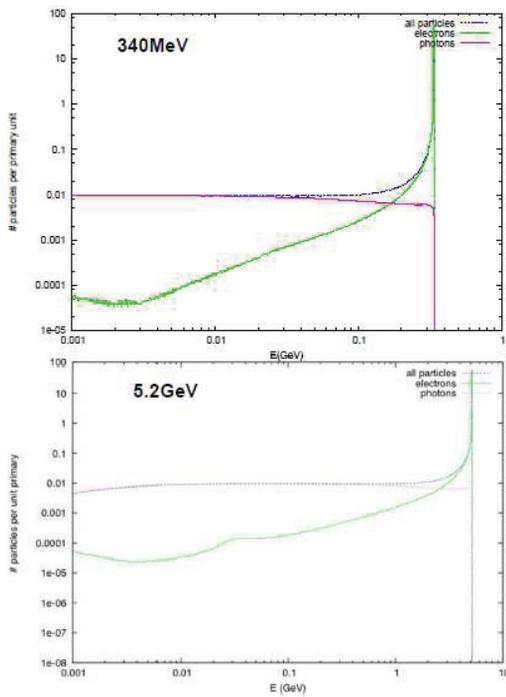


Figure 2: FLUKA results of the energy distribution of the wire losses for a beam energies of 340MeV and 5.2GeV.

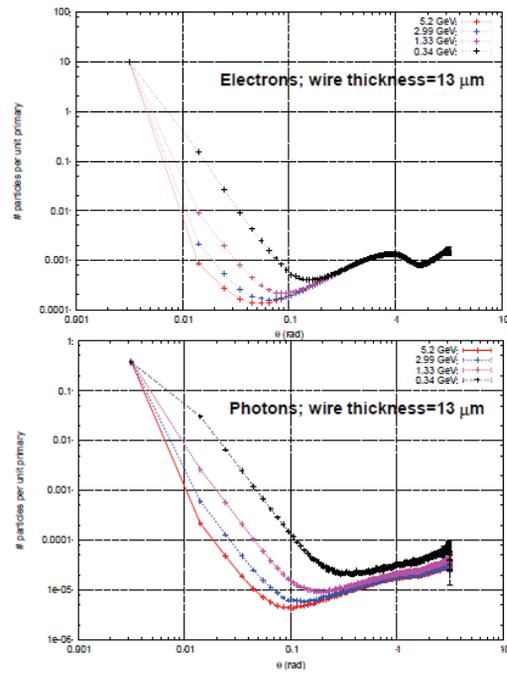


Figure 3: FLUKA results of the angular distribution of photons and electrons composing the wire losses.

were observed only for a wire velocity larger than 1mm/s. For a wire speed of 2mm/s, a vibration amplitude of about $0.6 \mu\text{m}$ (rms) was measured. A second series of vibration measurements have been performed in vacuum (10^{-4}mbar) on a prototype of the SwissFEL WSCs (2-phase stepping motor) for a velocity range of the wire 0.2-2 mm/s, see Fig. 1. A back-illuminated Tungsten wire ($13 \mu\text{m}$) was imaged by a camera with a frame rate of 500 fps equipped with a 200mm lens (projected pixel size $9.7 \mu\text{m}$). In Figs. 5, the relative variation of the centroid and the sigma of the projected image of a $240 \mu\text{m}$ long portion of the wire in motion at a velocity of 2 mm/s is shown. After subtracting in quadrature from the rms values of the centroid and sigma distributions - Figs. 5 - the corresponding rms values measured with the wire at rest, it results that, for a wire velocity of 2 mm/s, the centroid vibration is about $1.3 \mu\text{m}$ and the apparent enlargement of the sigma due to possible oscillation of the wire at a frequency higher than 500 Hz is about $0.1 \mu\text{m}$. In conclusion, the results of the measurements of the mechanical stability of the SwissFEL WSCs indicate that, for the wire velocity range of interest for SwissFEL 0.2-2 mm/s, the measured vibration of the wire is less than the resolution limit which can be achieved in a measurement of the beam profile with a Tungsten wire with a diameter of $5 \mu\text{m}$.

Several WSC tests on the electron beam have been also carried out in SITF. Two different techniques were adopted to detect the wire-signal produced by a Tungsten wire with a diameter of $25 \mu\text{m}$. In the former case, the wire-signal was retrieved as the difference of the charge read-out of two Beam-Position-Monitors (BPM), the one upstream the wire

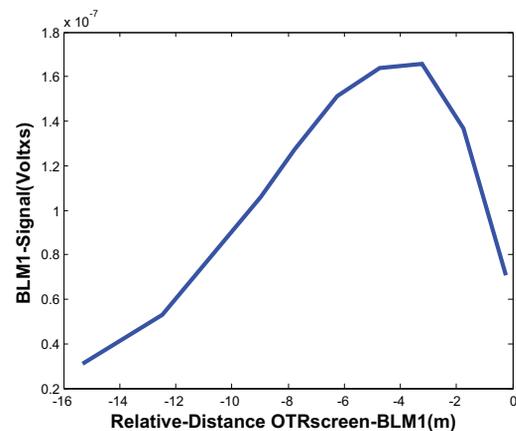


Figure 4: Measured beam losses as a function of the distance of the loss monitor from the single OTR screen intercepting the beam (charge 180 pC, energy 245 MeV).

and the other downstream the wire just behind the bending dipole of the high energy spectrometer. In this case a 10 Hz BS-ACQ of both the BPMs and encoder position of the wire was possible. In the latter case, the wire-signal was directly measured by a loss-monitor (scintillator fiber); no BS-ACQ available in this case. In both cases, the vertical profile of the beam measured by the WSC was compared with the beam profile measured by an OTR screen placed at the same longitudinal position of the wire. Results of both campaigns of measurements are shown in Fig. 6 (BPM read-out of the wire-signal) and in Fig. 7 (loss-monitor read-out of the wire-signal). Taking into account that the WSC data are not corrected by the transverse jitter of the beam, the OTR

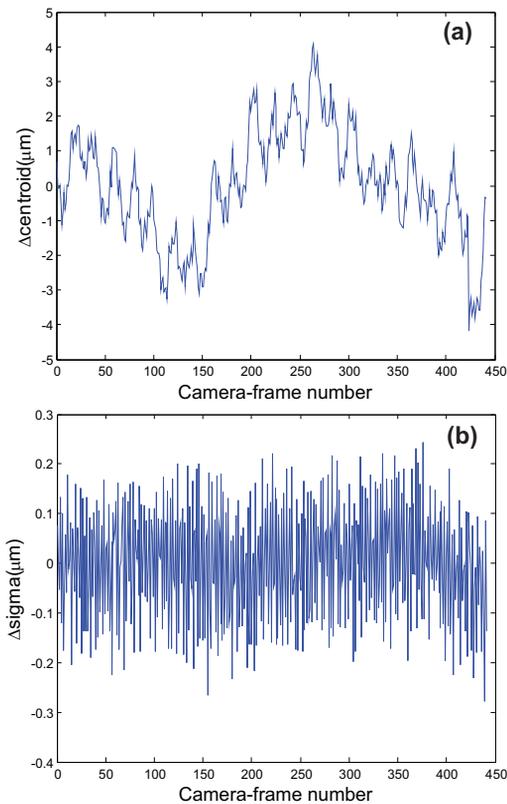


Figure 5: Wire vibration measurements results: relative variation of the centroid (a) and the sigma (b) of the projection of the camera image of the wire vs. camera frames number.

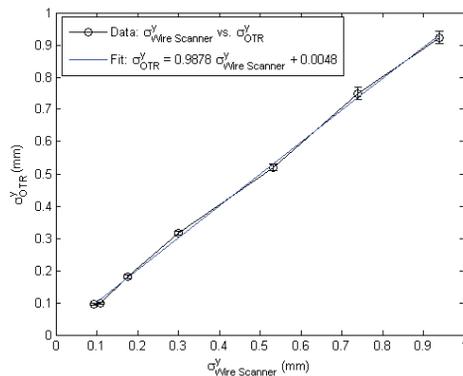


Figure 6: Measurements of the vertical profile of the electron beam: OTR vs. WSC results.

measurements vs. WSC measurements - see Fig. 6 - show an excellent linearity (less than 2% deviation) as well as the comparison of the WSC-loss-monitor and OTR measurements shows a good agreement within 8%, $\sigma_Y = 0.244 \pm 0.002$ mm (OTR), $\sigma_Y = 0.265 \pm 0.008$ mm (WSC).

CONCLUSIONS

Design and tests of the SwissFEL WSCs are presented. Results of the test-bench measurements of the mechanical stability of the prototype system indicate that, for the wire ve-

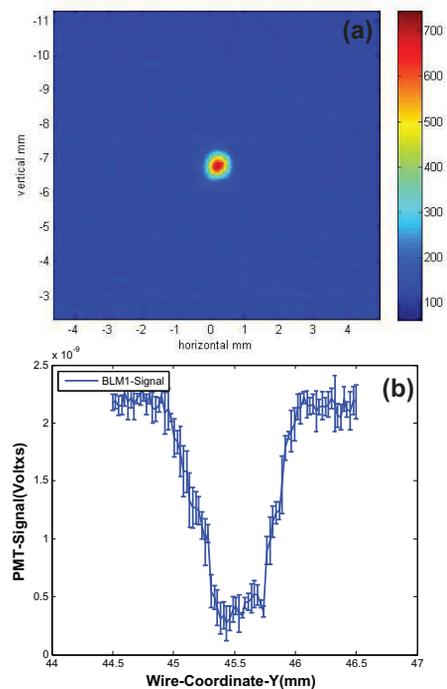


Figure 7: Measurements of the vertical profile of the electron beam with OTR screen and loss-monitor read-out of the WSC signal.

locities of interest for SwissFEL (0.2-2 mm/s), the measured wire vibration is less than the intrinsic resolution expected for a $5 \mu\text{m}$ thick Tungsten scanning the transverse profile of the electron beam. WSC tests on the beam indicate that WSC measurements of the beam transverse profile are consistent with analogous measurements performed with OTR screens.

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