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## Introduction

It is the purpose of this paper to discuss a new type of construction consisting of slotlines transverse to the beam (not to be confounded with transverse slots = Faltn type pick-up).<sup>1</sup> The slot pick-up of Faltn may be described as a line propagating a TEM-like wave in the beam direction coupled to the beam by slots in the outer conductor. In contrast to this the slotline pick-up to be discussed here propagates a beam-excited TEM-like wave along a slot between conducting planes perpendicular to the beam direction. The main advantage should be the rather simple fabrication process (double layer planar structure) and thus offering well defined mechanical tolerances. Also due to the photolithographic process very small structures and thus high frequencies can be reached.

Consider a charged particle beam (Fig. 1) propagating in +z-direction between two infinite

Fig. 1 - Slotlines and beam coordinates.

$$V_s = J_z \cdot dx_0 \cdot z_L/2 \quad (2)$$

assuming  $J_z$  not to be perturbed by the slot.

As a first approximation we assume that for a single current element at a source point  $y = d, x_0, 0$ ,  $J_z$  may be written as:

$$\vec{J}(d, x_0, 0, t) = \delta(x_0) \cdot f(t) \vec{u}_z = g(y=d, x, z, t) \cdot \vec{u}_z. \quad (3)$$

Let us remark that:

$$J(d, x_0, 0, t) = J(d, x_0, z, t + z/v_{beam})$$

whatever J is, an observer at any position  $y=d, x, z=0$ , t sees:

$$v_s(x, t) = \frac{Z_L}{2} \int_{-\infty}^{+\infty} J \left( x_0, t - \frac{|x - x_0|}{v_{s\text{lot}}} \right) dx_0 \quad (4)$$

rewritten as:

$$V_S(x,t) = \frac{Z_L}{2} \int_{-\infty}^x dx_0 J \left( x_0, t - \frac{x - x_0}{v_{\text{slot}}} \right) + \frac{Z_L}{2} \int_x^{+\infty} dx_0 J \left( x_0, t + \frac{x - x_0}{v_{\text{slot}}} \right) \quad (5)$$

With the present assumption the image current distribution  $g(y=\pm d, x, z, t)$  and the transversely ( $\pm x$ ) travelling waves in the slotlines are propagating both with  $v = c$  ( $v_{\text{beam}} = v_{\text{slot}} = c$ ) perpendicular to each other. The observer at  $y$  will measure a voltage across the slotline which has been raised from an image current distribution along the line:

$$x - x_0 = z \cdot v_{\text{slot}}/v_{\text{beam}}$$

(dashed line in Fig. 1) for slotline waves in  $+x$  direction and

$$x - x_0 = -z \cdot v_{\text{slot}}/v_{\text{beam}}$$

(dotted line in Fig. 1) for slotline waves in  $-x$  direction.

Thus by separating the  $+x$  and  $-x$  travelling slotline waves one obtains at a given time  $t$  the superposition of an image current distribution along a  $\pm 45^\circ$  line (Fig. 1) if  $v_{\text{beam}} = v_{\text{slot}}$  or more generally:

$$\tan \alpha = \frac{v_{\text{slot}}}{v_{\text{beam}}} \quad (6)$$

Under the further assumption that each surface current element can be attributed only to charged particles travelling perpendicularly above or below it, one might interpret the signal  $V_S(x,t)$  as being raised from a particle population in two sample disks of the beam tilted by  $\pm \alpha$  to the  $z$ -axis. Separating the  $+x$  and  $-x$  propagating slotline wave (e.g. by measuring  $V_S(x,t)$  at a distance  $x$  where  $|j_z| \approx 0$ ) a pick-up signal corresponding to a single tilted sample disk of the beam can be found.

#### Experimental Slotline Pick-Up Structures

The first mask, which was used to investigate the properties of a slotline pick-up, is reproduced in Fig. 2. It contains one line with  $Z_L = 150\Omega$ , 3 lines  $Z_L = 100\Omega$  and 3 lines  $Z_L = 50\Omega$  on a 3 mm alumina substrate. The impedances were calculated for  $f = 1.5$  GHz. Slotline-coaxial cable transition are provided at each end of the slotline, soldering a semirigid cable on the metallised substrate surface with the inner conductor across the slot next to the circle (= open end). For the realisation of all masks shown here, the "MICROS 3" CAD programme<sup>2</sup> has been used. It turned out that already for this very simple approach reasonable electrical properties could be obtained and that this structure is well suited to investigate the mutual coupling between slotlines, effects of mismatches, radiation and attenuation<sup>3,4</sup> with two transitions to 0.141" semirigid cable. The S-parameter  $S_{11}$  (reflexion coefficient) and  $S_{21}$  or  $S_{12}$  (transmission coefficient) are defined as the complex ratio of the reflected wave/incident wave and transmitted wave/incident wave, respectively. In the frequency range 1-4 GHz the transmission characteristic is reasonably flat and about half the transmission loss (ca 1 dB) can be explained by mismatch (-10 dB).<sup>4</sup> Since the transmission loss

amounts to roughly 3 dB/m here<sup>5</sup>, the radiation loss would be below 0.2 dB (2 dB/m). It should be pointed out that the intended interaction mechanism with the beam is not due to radiation (since slotlines or slots are often used as radiating elements in antenna arrays). It rather works like a transformer, where the image current produces some voltage in a given load.

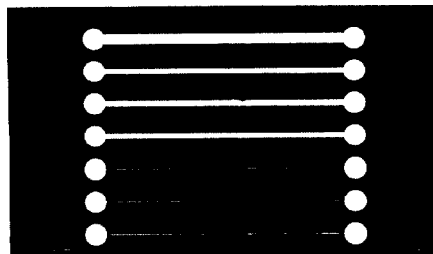


Fig. 2 - Mask for experimental (transverse travelling wave) slotline pick-up.<sup>4</sup>

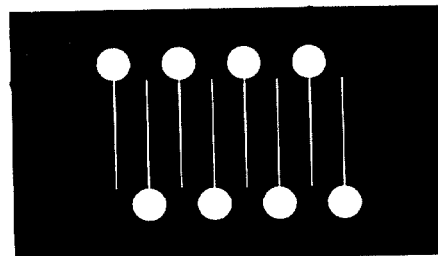


Fig. 3 - Mask for experimental (transverse standing wave) slotline pick-up ( $f_0 = 1.5$  GHz).<sup>4</sup>

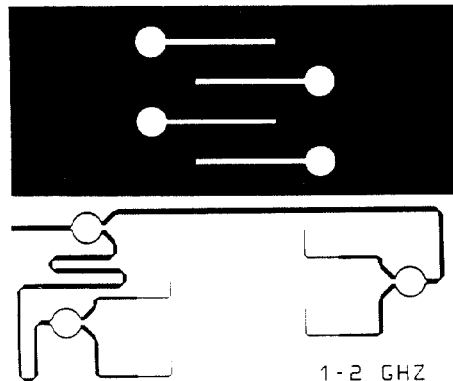


Fig. 4 - Mask for staggered standing wave slotline pick-up ( $4 \times 100\Omega$ ;  $f_0 = 1.5$  GHz); microstrip combiner board (right).

If more than one slotline is provided, mutual coupling will take place. The directivity turns out to be about 10-15 dB (1-4 GHz) and the forward coupling gives a relatively flat response in the same frequency range. A good theoretical treatment of slotline couplers can be found in references [5], [6] and [7].

For measurements of the pick-up response the field of a beam is simulated by a wire closely mounted ( $\approx 2$  mm) above a conducting ground plate which forms a 50 $\Omega$  transmission line, where most of the energy is concentrated between the wire and the ground plate. For quantitative calibration reference measurements were carried out with 50 $\Omega$  and 100 $\Omega$   $\lambda/4$  loop (or stripline) pick-ups positioned at the same distance  $h$  ( $h > 10$  mm) from the ground plate as the slotline metallisation. Here the structure from Fig. 3 was mainly looked at, because resonant slotlines give a response more similar to loop pick-ups than travelling wave slotlines. The length of these slots in Fig. 3 approaches  $\lambda/2$  at 1.5 GHz,  $Z_0 = 50\Omega$ . For an image current element passing in the middle of that array, each short circuit trans-

## Conclusion

The theoretical and measured results on planar slotline pick-ups indicate that these structures might become an interesting alternative to both the well known stripline or  $\lambda/4$  loop pick-ups as well as the slot pick-ups on Falin type structures.

Present fabrication methods allow very small dimensions in the surface pattern on the alumina substrate. The process involved is relatively easy and returns a good reproducibility with tight tolerances.

Resonant slotlines (TSWS) could have a higher beam impedance than stripline  $\lambda/4$  loops for a centred beam, but for an offset beam a faster lateral decay in sensitivity should be taken into account. Resistive losses are about comparable to loop pick-ups. Traveling wave slotlines (TTWS) offer the possibility of various sample choices on the beam. For a transverse slotline with  $v(\text{slot}) = c/2$  and a beam ( $v(\text{beam}) = c$ ) the normal vector of the sample disk would be at  $26.6^\circ$  to the beam axis ( $45^\circ$  for  $v(\text{slot}) = c$ ). By changing the angle of the slotline a wide range of angles of the sample disk can be set. Since the slotline must not necessarily be straight a variety of sample disks may be chosen. For a TTWS the beam diameter is not limited to about  $\lambda/2$  at the center frequency since the slotlines interact mainly with the TEM-like wakefield of the beam. Possible waveguide modes in the beam pipe may be suppressed by means of absorbing material.

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## References

- [1] L. Falin, Slot-Type Pick-up and Kicker for Stochastic Beam Cooling, CERN/ISR-RF/77-47, 1977.
- [2] J.F. Zurcher, "MICROS 3" A CAD/CAM Program for Microstrip Circuits EPFL Lausanne-LEMA, 1985.
- [3] F. Caspers, A Broadband Slotline Pick-up, PS/AA/ACOL Note 31, CERN, Feb. 1985.
- [4] F. Caspers, Planar Slotlines, Pick-ups and Kickers for Stochastic Cooling, CERN/PS/85-48 (AA), Aug. 1985-June 1986.
- [5] R.K. Hoffmann, Integrierte Mikrowellenschaltungen, Springer Verlag 1983.
- [6] K.C. Gupta, R. Garg and I.J. Bahl, Microstrip Lines and Slotlines, Artech House 1979.
- [7] M. Aikawa and H. Ogawa, A New MIC Magic-T Using Coupled Slotlines, IEEE-MTT-28 No 6, June 1980, pp. 523-528.
- [8] D. Suddeth, Argonne National Laboratory, Private Communication.

forms into an open one at  $\lambda/4$  distance from the short circuit and the entire structure should have an impedance of  $8 \times 50\Omega$  for this image current element. Considering a beam close to the surface in the middle of the array, one may assume that nearly all the image current is passing over the slotline at  $x = 0$  (Fig. 1), thus resulting in a power extraction of  $I_{\text{beam}}^2/4 \times 400\Omega$  which is equivalent to a beam coupling impedance of  $100\Omega$  for a single pick-up plate. The results of these measurements (Fig. 5) prove that a single resonant slotline (TSWS) has a pick-up response shape comparable to a single  $\lambda/4$  loop (centred beam). If the outputs of either side are combined in a 4 to 1 power combiner (proper delay for each output provided), a smoothing effect can be found (Fig. 6) and slight losses due to the combining network have to be accepted. For an off-centred beam, however, one might obtain deep notches in the response characteristic due to interference effects on the slotline ( $v_{\text{slot}}$  is only  $\approx 0.5 c$ ). The shape of the curves in Fig. 6 were also confirmed by measurements carried out at Argonne National Laboratory with a 22 MeV electron beam.<sup>8</sup> For these measurements the distance of the beam from the pick-up plate was 15 mm and with the combined signals of the 8 slotlines the maximum coupling impedance was found to be about  $50\Omega$ . The missing factor 2 (or 3 dB) can be attributed partly to losses in the 8:1 power combiner ( $\approx 1$  dB) and 8 times 2 metre coaxial cable. The other part is due to the fact that not all the image current passes over the centre of the slotline since the beam is already 15 mm apart.

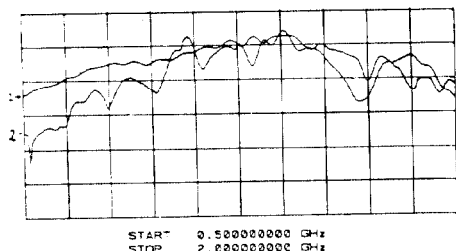


Fig. 5 - Pick-up response of a  $50\Omega$   $\lambda/4$  loop (2) and a single  $50\Omega$  resonant slotline (1) simulated beam centred (vert. scale: 5dB/div).

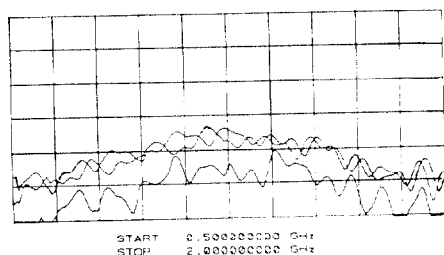


Fig. 6 - Pick-up response of  $\lambda/4$  loop (lower trace) and 4 combined resonant slotlines (Fig. 2) of either side (upper trace); simulated beam centred (vert. scale: 5dB/div).

The staggered slotlines ( $100\Omega$ ) in Fig. 4 have been realised in conjunction with a microstrip combiner board containing Wilkinson power combiners. Microstrip-slotline transitions with a through hole were applied in this case. A distinct ripple and rather poor sensitivity (compared with results from Fig. 6) is a difficulty linked to staggered slotlines. But with an average signal level 6 dB above the response of a single  $\lambda/4$ ,  $100\Omega$  loop approximately the same beam coupling impedance per unit length in the frequency range 1-2 GHz as such a loop has been obtained.