

MODELING OF BEAM LOSS IN TEVATRON AND BACKGROUNDS IN THE BTeV DETECTOR *

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Abstract

Detailed STRUCT simulations are performed on beam loss rates in the vicinity of the BTeV detector in the Tevatron $C\emptyset$ interaction region due to beam-gas nuclear elastic interactions and out-scattering from the collimation system. Corresponding showers induced in the machine components and background rates in BTeV are modeled with the MARS14 code. It is shown that the combination of a steel collimator and concrete shielding wall located in front of the detector can reduce the accelerator-related background rates in the detector by an order of magnitude.

MODELING WITH STRUCT AND MARS14

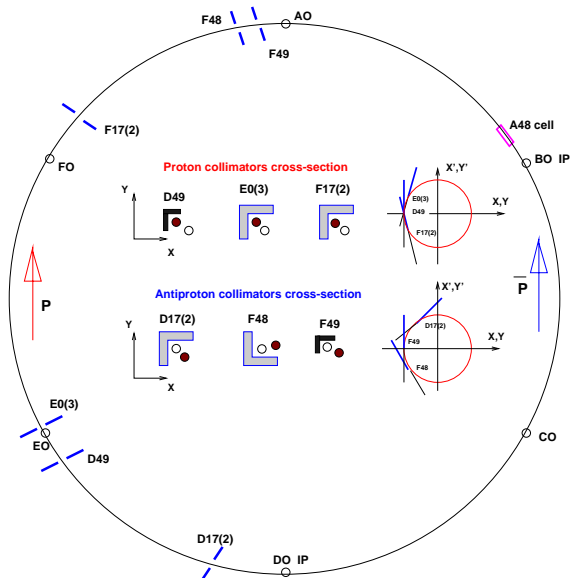


Figure 1: Tevatron Run II beam collimation system.

As a result of halo interactions with limiting apertures, hadronic and electromagnetic showers are induced in accelerator and detector components causing excessive backgrounds in the CDF and D \emptyset collider detectors these days and in the future BTeV detector. A two-stage collimation system has been developed for the Tevatron Run II [1] to reduce uncontrolled beam losses in the machine to an allowable level. About 0.1% of halo particles hitting the collimators are scattered back into the beam pipe. These particles are lost mostly in the high- β regions upstream of the experimental halls producing the background rates in the detector on the level of a few percent of those due to $\bar{p}p$ collisions.

To evaluate these rates for the BTeV detector, a multi-turn proton beam tracking through the Tevatron lattice with

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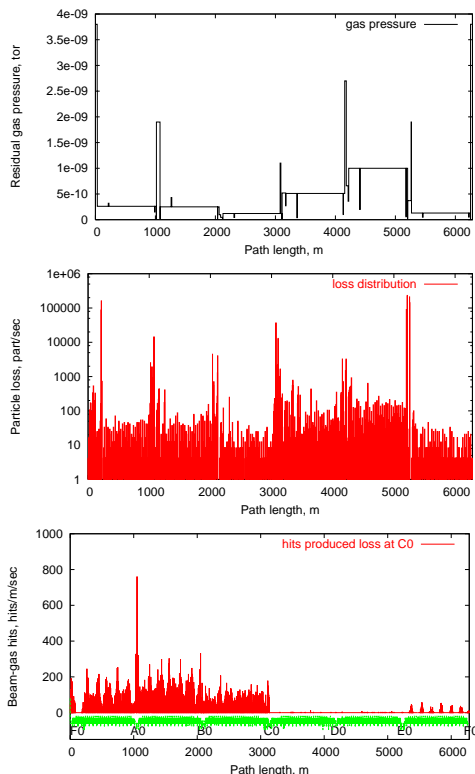


Figure 2: Residual gas pressure in the Tevatron Run-II (top), beam-gas induced loss distributions (middle) and beam-gas hit distribution for protons lost at $C\emptyset$ (bottom).

elastic beam scattering on the residual gas [2] and halo interactions with the collimators was conducted with the STRUCT code [3]. All the accelerator components with their real strengths and aperture restrictions were taken into account. Using the beam loss distributions calculated this way in the vicinity of the $C\emptyset$, detailed hadronic and electromagnetic shower simulations with the MARS14 code [4] were performed in the machine, detector and tunnel components with a cutoff energy for hadrons, leptons and photons of 0.1 MeV. Two protective measures – a short steel collimator at the B48 location and a concrete shielding wall at the tunnel/hall interface on the proton side – were considered as a way to reduce the machine-related backgrounds in the BTeV detector.

The Tevatron lattice, that corresponds to the BTeV operation, with $\beta_{x,y}^* = 0.35$ m at $C\emptyset$ was used without collisions at B \emptyset and D \emptyset ($\beta_{x,y}^* = 1.7$ m). The BTeV pixel detectors aperture radius is 2.75 mm ($31\sigma_{x,y}$), the $C\emptyset$ triplet quadrupoles aperture radius is 31.5 mm ($14\sigma_{x,y}$), and all other machine components with their strength and apertures were implemented in the model. The luminosity at $C\emptyset$ is assumed to be $\mathcal{L} = 2 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$. The collimation system (Fig. 1) and residual gas pressure distribution (Fig. 2) of the Run-II [1, 2] were assumed in the model-

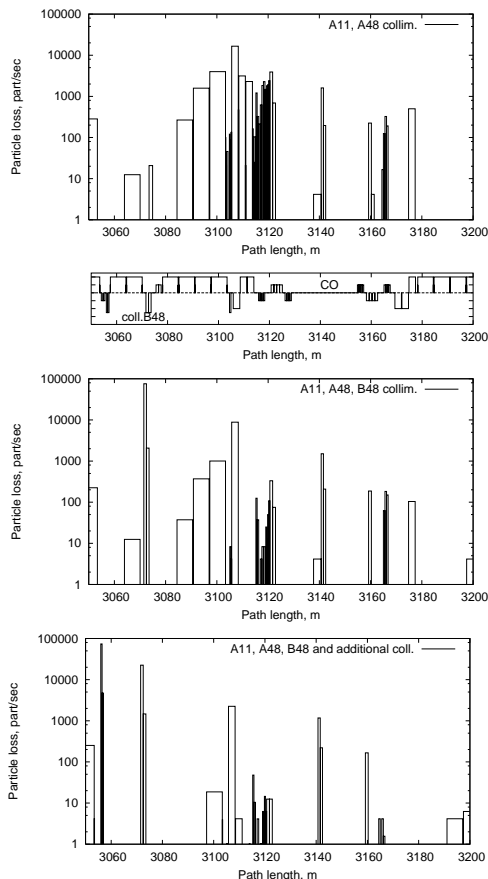


Figure 3: Beam-gas induced beam loss distributions in the CØ region: baseline (top), with B48 collimator (3-d line) and with additional collimator (bottom).

ing. Detailed 3D geometry, magnetic field and materials description in a 70-m region upstream of the CØ IP were implemented in the MARS14 model for lattice and tunnel components along with a few meters of the dirt surrounding the tunnel.

RESULTS

For the current vacuum conditions, the nuclear elastic beam-gas interactions is a dominant source of beam loss on the electrostatic separators and low- β quadrupoles as shown in Table 1.

Calculated beam loss distributions in the CØ region due to elastic beam-gas interactions and collimation system inefficiency are shown in Figs. 3 and 4 for the baseline layout and the case with a 1.2-m long stainless steel collimator at the B48 warm region upstream the last four SC dipoles preceding the IP. The collimator jaws are at $12\sigma_{x,y}$ from the beam axis, with their rectangular full aperture of 15.6 mm (hor.) and 4.6 mm (ver.). A further substantial reduction is possible with an additional short collimator located at $12\sigma_{x,y}$ in ~ 25 m upstream of the B48 one. Unfortunately this collimator should be placed in a short region between superconducting magnets, that makes it difficult for realization.

The calculations show that the B48 collimator in a combination with a recently installed A48 collimator protects

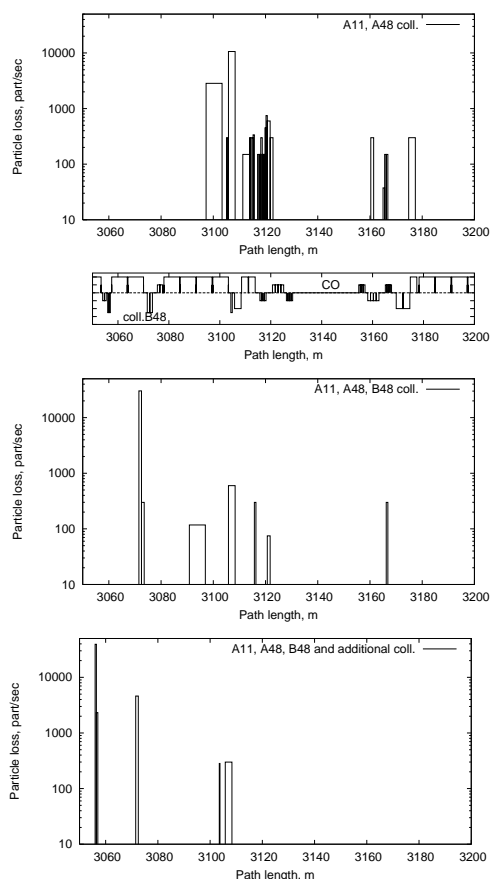


Figure 4: Collimation system induced beam loss distributions in the CØ region: baseline (top), with B48 collimator (3-d line) and with additional collimator (bottom).

the BTeV pixel detectors and the low- β quads at an abort kicker prefire, reducing a number of particle lost by three orders of magnitude.

Particle flux isocontours ($E_{th}=0.1$ MeV) in the orbit plane in the 60-m long region preceding the BTeV experimental hall are presented in Fig. 5. Shown are neutrons in the baseline configuration and charged hadrons for the case with the B48 collimator and a 2-m concrete wall at 12.7 m upstream of the CØ IP. Fig. 6 shows hadron flux XY-isocontours at the entrance to the hall (12.2 m from IP) for the case with the B48 collimator and shielding wall.

Total background rates are summarized in Table 2. The dominant component is photons: about 10^8 soft photons per second (baseline) entering the hall around the beamline. There is no wall effect at $R < 0.25$ m. The B48 collimator

Table 1: Beam loss rates (10^4 s $^{-1}$) in the 70-m regions upstream of DØ and BØ (now) and CØ (2009) with Run-II vacuum parameters.

Source	DØ	BØ	CØ
Nuclear elastic beam-gas	8.8	8.0	9.4
Large angle Coulomb beam-gas	0.12	0.06	0.1
Tails from collimators	2.4	3.5	0.99
Elastic $p\bar{p}$ at two IPs	0.144	0.105	-

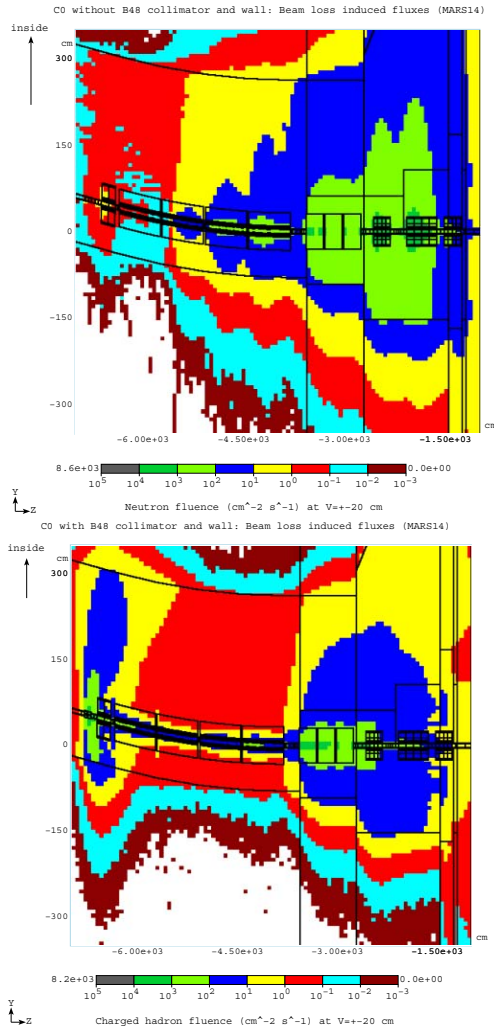


Figure 5: Particle isofluxes in the $C\emptyset$ region: neutrons, baseline (top) and charged hadrons with B48 collimator and 2-m concrete wall (bottom).

alone reduces the backgrounds by a factor of two compared to the baseline configuration. Installation of the shielding wall results in a combined reduction effect of a factor of ten. The numbers in the table are to be increased by 10% to account for tails from the Tevatron main collimators.

CONCLUSIONS

A combined reduction of particle flow by a factor of 10 is obtained from the B48 collimator and a 2-m concrete shielding wall at 12.7 m from the IP. Machine-related backgrounds at the last muon plane of the BTeV detector are quite low: 3.5 charged hadrons and 160 photons per cm^2 per second. With a 5-GeV cut this makes the machine-related backgrounds in the BTeV pixel detectors at a per-

Table 2: Number of particles (10^5 s^{-1}) above 0.1 MeV entering the BTeV hall at $z=-12.2 \text{ m}$ and $R < 3.5 \text{ m}$.

Scenario	n	h^\pm	e^\pm	γ	μ^\pm
No B48, no wall	24.2	14.5	58.9	1147	2.80
B48, no wall	11.0	9.29	42.4	730	1.81
B48, 2-m wall	6.29	2.48	7.55	132	1.00

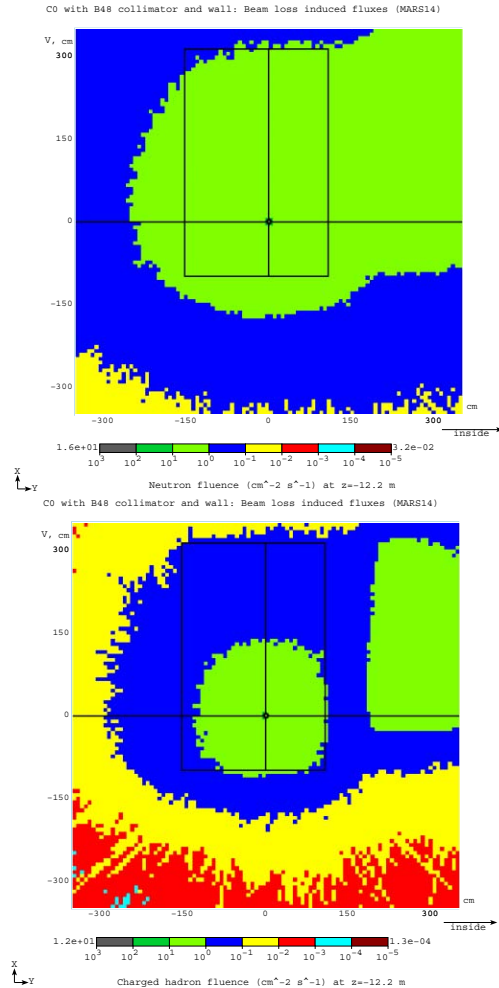


Figure 6: Neutron (top) and charged hadron (bottom) isofluxes at the entrance to the $C\emptyset$ hall, with B48 collimator and 2-m concrete wall.

cent level of those from $p\bar{p}$ collisions. Calculations show that a further substantial reduction is possible with an additional short collimator (unfortunately in the cold region) 25 m upstream of the B48 one. Simulations show that B48 collimator in a combination with a recently installed A48 collimator protects the BTeV pixel detectors and the low- β quads at an abort kicker prefire.

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