HIGH-BRIGHTNESS BEAM DIAGNOSTICS FOR THE APS LINAC*

Advanced Photon Source, Argonne National Laboratory, Argonne, IL 60439 USA

Abstract

The Advanced Photon Source (APS) injector includes an S-band linac with the capability to accelerate beams to 650 MeV. The linac has recently been upgraded with the installation of an rf thermionic gun in addition to the standard DC thermionic gun. The rf gun is predicted to have lower emittance (5π mm mrad) and may be used to support the APS self-amplified spontaneous emission (SASE) experiments. The critical characterization of this gun’s beam has begun with a beam diagnostics station at the end of the linac that can address beam transverse size, emittance, and bunch length (peak current). This station uses both an optical transition radiation (OTR) screen at 45° to the beam direction and a Ce-doped YAG single crystal normal to the beam with a 45° mirror behind it. The visible light images are detected by a Vicon CCD camera and a Hamamatsu C5680 synchroscan streak camera. Spatial resolution of about 30 µm (σ) and temporal resolution of 1 ps (σ) have been demonstrated. Examples of rf gun beam characterization at 220 MeV are reported.

1 INTRODUCTION

The ongoing interest in high-brightness particle beams has been “stimulated” by the potential applications to self-amplified spontaneous emission (SASE) free-electron laser experiments [1,2]. At the Advanced Photon Source (APS) the injector linac has recently been upgraded with the installation of an rf thermionic gun [3,4] and a laser-driven photocathode (PC) rf gun [5]. The rf gun is predicted to have lower emittance (5π mm mrad) and sufficient peak currents (~100A) to be used in the APS SASE experiments [6]. The PC rf gun is expected to have even better brightness [5]. The capability to characterize such beams includes a beam diagnostic station with improved spatial and temporal resolution that can address beam transverse size, emittance, and bunch length (peak current). The station uses both optical transition radiation (OTR) and a Ce-doped YAG single crystal normal to the beam with a mirror behind it. The visible light images were detected by a charge-coupled device (CCD) camera and a dual-sweep streak camera. Spatial resolution of about 30 µm (σ) and temporal resolution of about 1 ps (σ) have been demonstrated in earlier studies [7,8]. Examples of beam measurements for the rf gun are reported.

2 EXPERIMENTAL BACKGROUND

The APS facility’s injector system uses a 250-MeV S-band electron linac and an in-line 450-MeV S-band positron linac. The original electron gun was a conventional, gated DC thermionic gun. The diagnostics station was initially commissioned using this gun’s beam [7]. More recently an rf thermionic gun was installed that is designed to generate low-emittance beams (< 5π mm mrad). It is configured with an alpha magnet that allows beam injection just after the first linac accelerating section [3,4]. Both in-line linacs can be phased to produce 50- to 650-MeV electron beams when the positron converter target is retracted.

The predicted lower normalized emittance of beams from the rf thermionic gun and PC rf gun versus the DC gun results in correspondingly smaller beam spot sizes. We have addressed this at selected diagnostics stations by supplementing or replacing the standard Chromox screen, with its approximate 200-µm spatial resolution and 300-ms decay time, with the OTR and Ce-doped YAG screens. For the station at the end of the linac the OTR screen was a molybdenum mirror from Melles Griot and oriented so that its surface was at 45° to the beam direction. The doped YAG crystal of 0.5 mm thickness (obtained from Startec) was mounted with its surface normal to the beam direction with a Zerodur mirror at 45° to the beam just behind it.

An alignment laser was injected on-axis into the bore of the linac just after the first accelerating structure and used to assess the relative orientation of the two mirrors involved at the station. The laser was reflected into the optical transport line and then used to simulate the path of the visible light images. In the case of OTR, the direction of the path is particularly critical since OTR is emitted around the angle of specular reflection. The visible light was transported out of the linac tunnel to an optics table via two 150-mm-diameter achromat lenses and two mirrors. For these experiments the OTR or YAG:Ce light could be viewed by a Vicon CCD camera and/or a Hamamatsu C5680 dual-sweep streak camera. The synchroscan unit was phase-locked to 119.0 MHz, the 24th subharmonic of the 2856-MHz linac frequency. A low-jitter countdown circuit has been built using Motorola ECLIN PS Logic to generate the 24th subharmonics with sub-ps jitter [9]. The synchroscan unit was critical to the bunch length measurements due to the lower charge/micropulse in early experiments. The model M5677 slow-sweep unit was used in the measurement of the YAG:Ce response time, which was previously...


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reported [10] as about 80 ns (FWHM). Although beam energies as high as 540 MeV have been used, one set of data is at 220 MeV, the energy that would be used for SASE experiments at 517 nm.

3 EXPERIMENTAL RESULTS

Preliminary results are reported in this section on rf gun beam transverse size, emittance, and bunch length. Additionally, the measurement of the YAG:Ce crystal response time is reported.

3.1 Beam Transverse Size

Recently measurements at 220 MeV were performed on the beam at the end of the linac. Beam images were obtained with both the YAG:Ce crystal and the OTR converter. Figure 1 shows an example of the beam vertical profile using OTR. The MAC-TA algorithm calculated a spot size of ~170 µm (FWHM) that corresponds to about 72 µm (σ), assuming a Gaussian-shaped profile. The aspect ratio (H:V) was about 4:1 for this beam focus.

3.2 Preliminary Emittance

Having obtained reasonable transport and focus of the beam, an evaluation of the beam emittance was performed by the standard tracking of beam size with the quadrupole field strength of an upstream magnet. In this screen position there were two quadrupoles that were available to be adjusted. A script developed by M. Borland was adapted that automatically stepped the fields and logged the measured beam sizes via the EPICS process variables [11]. As shown in Fig. 2, the horizontal beam size can be minimized by selecting a combination of fields from the linac-to-PAR magnets LTPQ9 and LTPQ10. The Max Video-20 digitized data were fit to Gaussian profiles in this case. An example of emittance data is shown in Fig. 3. The data are consistent with normalized emittances, $\varepsilon_x = 8.8 \pi$ mm mrad and $\varepsilon_y = 9.5 \pi$ mm mrad for a 25-mA macropulse current. Errors are estimated as 20-30%. Simulations of this particular rf gun setup with elegant [12] have reproduced the emittance asymmetry between the two planes. It is noted that the measured and computed emittance ratios are similar although the measured values are 60% higher. These measurements were made without an energy filter in the alpha magnet. As a result, the beam contains a trailing, low-energy component that is expected to significantly increase the projected (but not slice) emittance. Although this low-energy “tail” contains only a small fraction of the total charge (approximately 10% of the total bunch charge by the time the bunch reaches the end of the linac), it is the greatest contribution to the projected (longitudinal whole-beam) emittance. Considerably lower measured emittances are expected once the alpha magnet energy filter has been installed.

Figure 1: An example of the rf thermionic gun’s vertical beam profile: 170 µm (FWHM).

Figure 2: Variations of the observed horizontal beam size with the two upstream quadrupoles Q9 and Q10.

Emittances of $\varepsilon_{x,y} < 10 \pi$ mm mrad were determined from the fitted curves.
3.3 YAG:Ce Crystal Response Time

The rf gun’s macropulse is typically 30 to 40 ns, shorter than the 80-ns (FWHM) response time of YAG:Ce reported by Graves et al. [10]. This would preclude any submacropulse imaging with this converter. We have used the slow sweep streak module operating at 1-µs full range to assess the actual response times of our specific crystals. The electron-beam macropulse duration is not a delta function impulse to the system, but it is still sufficiently short that a reasonable measurement is possible. In these data we actually used the DC gun beam with a shortened macropulse of ~20 ns. As shown in Fig. 4, the measured YAG:Ce response is 105 ns (FWHM). A preliminary assessment of the 1/e time is also about 112 ns. The streak unit’s calibration factor is based on the original factory calibrations.

![Figure 4: Streak camera data showing the response time of the YAG:Ce crystal to the incident 20-ns electron beam. The effective response is about 105 ns (FWHM). The horizontal axis spans 1 µs.](image)

3.4 Bunch Length Measurements

Assessments of micropulse-averaged bunch length using the synchroscan streak camera have begun. Preliminary results at 540 MeV and for macropulse currents of less than 150 mA are 3-4 ps (σ). Low signal levels required averaging over about 30 macropulses (few seconds) and so may involve some phase slew blurring the image and lengthening the apparent size. Measurements with the optimized gun and an optimized optical path at 220 MeV are planned. Use of a fifth-harmonic cavity signal to optimize bunch length was previously reported by Lewellen et al. [4].

4 SUMMARY

In summary, a diagnostics station that can support the characterization of bright electron beams at the end of the APS linac has been tested. It has also been used to attain preliminary data for the rf thermionic gun beam quality at low beam currents. Further experiments are planned this year at high macropulse currents of the rf gun and on the PC rf gun.

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6 REFERENCES