STUDY OF A DETUNED ACCELERATING SECTION WITH THE COMPUTER PROGRAM PROGON

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Abstract

The longitudinal coupling impedance for a number of lower passbands, bunch to bunch energy variation due to longitudinal wake fields, the beam loading compensation, some effects of production errors, and the rf pulse transmission through a detuned disk loaded accelerating section with finite wall conductivity have been studied using the computer program PROGON.

MOTIVATION AND METHOD

The detuned accelerating section [1] designed to diminish the undesirable effects of the transverse wake field has been studied previously using approximate methods [2][3]. Here these results are confirmed by employing the new computer code PROGON [4] which is based on the field matching technique for the Fourier harmonics of the longitudinal electromagnetic (EM) travelling waves. The geometry of the considered 204 cell section is presented in Fig. 1a.

ACCELERATING FIELD

First, the EM fields in the structure, the corresponding stored energy \( W(N) \) per cell and the power flow \( S_z(N) \) through the cavity cross section are found. These quantities define the relative group velocity \( v_g/c = S_z/d/W \), where \( d \) is the cell length. The stored energy \( W \) and the group velocity \( v_g \) are plotted in Figs. 1b,c, respectively.

ERRORS AND DAMPING

Next, the effect of fabrication errors for infinite and finite wall conductivity on the excitation of and propagation through the section of a cylindrical accelerating wave with frequency \( f \) has been studied. The absolute values of the reflection \( R(f) \) and the transmission \( T(f) \) coefficients are plotted in Fig. 2 for "small" relative random radi-

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Figure 1. (a) The radii of the \( N \)th cavity \( b(N) \) and iris \( a(N) \) for the considered 204 cell section. The gaps \( g(N) = 0.729 \) cm, the iris thicknesses \( l(N) = 0.149 \) cm. (b) The stored energy in the \( N \)th cell \( W(N) \). (c) The relative group velocity in the \( N \)th cell \( v_g(N)/c \).

Figure 2. The absolute values of the reflection \( R(f) \) and the transmission \( T(f) \) coefficients: (a),(b) "small" relative random radii errors in the range \( \Delta a/a = \Delta b/b = \pm 2.5 \cdot 10^{-4} \); (c),(d) "large" errors \( \Delta a/a = \Delta b/b = \pm 2.5 \cdot 10^{-5} \); (e),(f) the finite conductivity of the copper walls.
The real part of the longitudinal coupling impedance $Z_c(f)$ is presented in Fig. 3 for a few lower passbands.

**COUPLING IMPEDANCE**

The real part of the longitudinal coupling impedance $Z_c(f)$ is presented in Fig. 3 for a few lower passbands.

**PULSE TRANSMISSION**

The distortion of an rf pulse transmitted through the detuned section with finite wall conductivity is illustrated in Fig. 4a where the envelopes of the rf pulse at the entrance (dashed curve) and the exit (solid curve) are plotted [5].

**BEAM LOADING**

The beam loading [6] for a long train of bunches is illustrated in Fig. 4b [7]. The energy gain by a bunch in an accelerating wave with power 70.56 MW (curve 1), the energy loss due to wake fields of the previous bunches (curve 3) and their difference (curve 2) are plotted versus the bunch number for the number of electrons $N_p = 0.7 \cdot 10^{10}$ per bunch. The impedance of the first passband only is taken into account for the wake field calculations. Fig. 4c is a close up of part of curve 2 in Fig. 4b. Fig. 4d illustrates the effect of the random bunch population in the range $\Delta N_p/N_p = \pm 3\%$.

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**REFERENCES**