Correlation Method of Measurements of Ion Beam Parameters

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Abstract

A correlation method of nonperturbative control on the ion energy spectrum, beam profile and transverse emittance in a bending transport line area, is suggested. The method is based on measurements of the cross-correlation function between a flux of photons or particles from a probing target pseudorandomly modulated in time, and that of fast information-carrying convoy particles produced in the target and recorded at a drift distance. Characteristics of the apparatus used to realize the proposed diagnostic method by means of time-integrated correlometers based on charge-coupled devices, are considered for a source of $H^-$ ions.

1. INTRODUCTION

For high-brightness ion accelerators it is important to obtain information on beam parameters not affecting them appreciably during measurements (nonperturbative diagnostics). For this purpose in a bending transport line area, fast neutral or charged convoy particles can be used. These particles are produced as a result of ion destruction or the charge-exchange process in a specially shaped target which is practically transparent for a beam (for $H^-$ beams see [1-5]). The target is formed so that these information-carrying convoy particles (IN-particles) follow the ion velocity in magnitude (in relative units) and in direction (in rad) with accuracies required for measurements. These accuracies can be estimated by $\leq (\mu_0 I_0/M_i E_i)^{0.5}$ where $\mu_0$ is the reduced mass of the IN-particle and the remaining part of the ion in its destruction or the ion and electrons in their recombination, $I_0$ is the affinity energy, $M_i$ and $E_i$ are the ion mass and energy, respectively. In sources, for example of negative ions, the probability of IN-particle generation ($\eta$) in residual gas can be quite considerable (for the neutral IN-particles $\eta_r \approx 0.2 - 0.4$). In this case, using the well-known methods [1-5] for any density of a probing target, it is impossible to separate directly information on beam parameters from a flux of background IN-particles on a detector. The correlation method of nonperturbative measurements of ion beam parameters considered in this paper allows one to overcome these difficulties.

2. METHOD

The correlation method of nonperturbative measurements of the ion energy spectrum has been previously proposed [6,7]. It is based on the use of test IN-particles produced in a target, pseudorandomly modulated in time, and detected at drift distance $L$. To measure the transverse beam emittance, for example in the $(Y, Y')$-plane (see Figure), one or a few thread-type targets parallel to the $(X, Z)$-plane can be formed in front of a bending transport line area. If the ion beam current is invariable during measurements, the spatial $X$-dimension of the target must be required for reproducing target time modulation by the flux $\Psi_p^*$ of test IN-particles

$$\Psi_p^*(t) = const_{(n)} \cdot I_p^*(t),$$

where $I_p^*$ is the flux of photons or particles in the $n$-target. The targets are fixed in space.
and separated from each other along the Y-axis. When one target is used, it moves in parallel along the Y-axis. Taking into account (1), the autocorrelation function of the flux of test IN-particles on a m-detector is equal to

$$R_{nm}^{pp}(\tau) = \int_{-\infty}^{+\infty} \Psi_{pn}(t) \cdot \Psi_{pn}(t - \tau) \, dt =$$

$$= \sum_{k=-\infty}^{+\infty} \delta(\tau - kT). \quad (2)$$

The pulsed characteristic \( h_{nm}(t) \) of the drift distance from the n-target to the m-band-type detector \((Y' = fix)\) is related to the velocity \((V)\) distribution of IN-particles \((t = L/V)\) in the \(n-m\) direction and, hence, to the energy spectrum of ions. The fluxes of IN-particles in the n-target area \((\Psi_n^o)\) and on the m-detector \((f_{nm}^o)\) are related by the convolution

$$f_{nm}^o(t) = C_m \cdot \int_{0}^{+\infty} h_{nm}(\tau) \cdot \Psi_n^o(t - \tau) \, d\tau,$$

$$\sum_m C_m = 1, \quad (3)$$

where \(\Psi_n^o = \Psi_{bn}^o + \Psi_{pn}^o\), \(\Psi_{pn}^o\) is the flux of background IN-particles produced in the residual gas. Taking into account the independence of \(\Psi_{bn}^o\) and \(\Psi_{pn}^o\) and measuring the cross-correlation function between the fluxes of target particles or photons and IN-particles on the detector

$$R_{nm}^{co}(\tau) = \int_{-\infty}^{+\infty} I_{nm}(t) \cdot f_{nm}(t - \tau) \, dt =$$

$$= B_{nm} \cdot \int_{0}^{+\infty} h_{nm}(t) \cdot R_{nm}^{pp}(\tau + t) \, dt =$$

$$- B_{nm} \cdot \sum_{k=0}^{+\infty} h_{nm}(\tau - kT), \quad (4)$$

we obtain the pulsed characteristic of the drift distance in the \(n-m\) direction \((I_{nm}^o \propto I_n^o)\). Using normalization \(\int h_{nm}(\tau) \, d\tau = 1\), we get from the \(B_{nm}\)-matrix information on the ion distribution in the \((Y, Y')\)-plane and thus on the beam Y-profile and transverse emittance.

In reality, we must form such targets when convolution (4) of \(h_{nm}\) and \(R_{nm}^{pp}\) does not change the supposed \(h_{nm}(t)\)-function. In accordance with [8], this condition means that a periodically replicating element of the autocorrelation function of the \(I_{n}^o\)-flux must have a sufficiently narrow shape in time with width \(\Delta < \tau_{max}\), where \(h_{nm}(\tau) = 0\) for \(|\tau| \geq \tau_{max}\), and its period \(T\) must meet the condition \(T > 2 \cdot \tau_{max}\). Correlation methods measure a useful signal with a background which is several orders of magnitude more than the signal. Thus, measuring \(R_{nm}^{co}(\tau)\) by \(n \times m\) correlometers, the energy spectrum and distribution of ions in the \((Y, Y')\)-plane can be controlled without perturbation of beam parameters.

### 3. APPARATUS

Nonperturbative measurements of ion beam parameters, for example in a source of \(H^-\) ions, can be realized according to the scheme shown in Figure. It is analogous to the previously proposed one [7], but it contains "n" identical photon targets \((I_n^o)\) and photon detectors \((D_n^o)\), \(n \times m\) correlometers \((C_m)\) and band-type detectors \((D_{nm}^o)\) of fast \(H^o\) atoms. When probing targets are formed by diaphragming radiation with an optimum polarization and a wavelength of \(\lambda = 10600\AA\) from the Nd:YAG laser with synchronized modes (see Fig.1), the test \(H^o\) atoms follow the \(H^-\) velocity in magnitude (in relative units) and in direction (in rad) with accuracies of \(-4 \cdot 10^{-3} \cdot (E_i[keV])^{-0.5}\). The series duration of pseudorandom radiation pulses is \(T_s \approx 100\) ns and the width of the autocorrelation function is \(\Delta \approx 50\) ps [9]. Thus, such photon targets due to \(H^-\) photodetachment can efficiently generate test IN-particles \((H^o)\) and allow one to measure pulsed characteristics of the drift distance \(h_{nm}(\tau)\) which are fairly short in time. At present, possibilities of the above diagnostics are mainly limited by the fast action of correlometers. The cross-correlation function \(R_{nm}^{co}(\tau)\) between the
The $R^o_{nm}$-functions can be read out during intervals between target switchings or between ion beam pulses and taking into account the guiding frequency of modulating structures.

As estimates for the source of $H^-$ ions with an energy of $E_i \approx 20$ keV, $L \approx 100$ cm and the average power density of laser radiation within the duration of a series of pulses $I^r_T \approx 4 \cdot 10^5 W/cm^2$ ($\Psi_{bn}/\Psi_{pn} \approx 10^2$), the proposed apparatus allows one to realize nonperturbative measurements of beam parameters during $T_m \geq T_s \approx 100$ ns with accuracies, e.g., of $\approx 0.4\%$ in energy and of $\approx 2 \cdot 10^{-4}$ rad in $Y'$.

References