EMITTANCE GROWTH STUDY USING 3DE CODE FOR THE ERL INJECTOR CAVITIES WITH VARIOUS COUPLER CONFIGURATIONS

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Abstract

We compute emittance growth in electron bunches in high gradient accelerating cavities due to field non-linearities and phase dependence associated with the finite bunch length. The determination of the emittance is based on tracking of individual particles through tabulated, 3-dimensional, electromagnetic fields. No symmetry is assumed. The fields for the structure, including input coupler, are calculated using MAFIA or MWS[1]. We examine particle trajectories RF focusing and emittance growth of injector cavities with various coupler configurations under study for Cornell Energy-Recovery-Linac proposal [2].

INTRODUCTION

A study of beam emittance growth in ERL injector cavities [2] was done using the new code 3DE[3]. The study was done for four different coupler configurations over the particle energy range of 0.5 MeV – 10 MeV. All four structures have two identical RF cells. The input and output beam pipe dimension of the first structure is slightly different than the pipe dimensions of the other three structures. The arrangements of the couplers in each structure is listed below:

- **Structure I**: Two cells cavity - NO couplers
- **Structure II**: Two cells + ONE vertical coupler at the bottom of the output pipe located at z = 0.281 m from the input, see Figure 1.
- **Structure III**: Two cells + TWO vertical symmetric couplers at the bottom and top of the output pipe located at z = 0.281 m from the input, See Figure 2.
- **Structure IV**: Two cells + TWO vertical asymmetric couplers, the same as structure III, except there is a 1 mm vertical shift of the bottom coupler relative to the top one.

The top coupler in structures III & IV was added in order to compensate the effect of the transverse kick of the bottom coupler on the emittance. In structure IV a 1 mm shift was added to simulate practical imperfection in manufacturing.

THREE DIMENSIONAL ELECTROMAGNETIC FIELD

The three dimensional Electromagnetic field for all the structures were calculated using the code MWS[1]. The fields were calculated for both electric and magnetic wall boundary conditions and were combined, in order to simulate propagating waves, using the following expressions [4]:

\begin{align}
E &= (-E_e + j \cdot k_n \cdot E_m) e^{-j \omega t + \varphi} \\
H &= (H_e - j \cdot k_n \cdot H_m) e^{-j \omega t + \varphi}
\end{align}

$E_e$, $H_e$, are fields calculated with the electric wall boundary condition

$E_m$, $H_m$, are fields with calculated the magnetic wall boundary condition

$k_n$ is a normalization coefficient. The value of $k_n$ is determined from the characteristic impedance of the coaxial line of the RF couplers. The sign of $k_n$ is chosen to satisfy the a condition for a propagating waves into the structure.

TRACKING

The 3DE code calculates the new kick on the particle due to the RF field at each coordinate and calls BMAD’s tracking routine [5], which uses a Runge-Kutta algorithm, to calculate the new particle location. Using the new location 3DE code calculates the new angles and particle energy. A comparison of the trajectories of three particles with vertical displacement of $y = -2$, 0, 2 mm...
and horizontal displacement of \( x = 0, 0, 0 \) mm entering structure II (two cells with ONE vertical coupler) and structure IV (two cells with TWO, 1 mm shifted, vertical couplers) is seen in Figure 3. The particles traveling in the ONE coupler structure are being kicked in the \(-y\) direction after passing the location of the coupler at \( z = 0.281 \) m, due to the transverse kick.

Figure 3: Particles trajectories in structures II (dashed line) & IV (solid line) (These are the one-coupler & two-asymmetric-couplers structures)

**RF FOCUSING**

We observe the expected transverse focusing[3], and compute its strength as \( \frac{\Delta y'_{out}}{\Delta y_{in}} \). A comparison of the RF focusing versus Input-Particle-Energy for all four structures can be seen in Figure 4. The RF phase was optimized to get maximum acceleration at each input energy. The RF focusing of structures II, III, IV (two cells WITH couplers) is essentially identical and about 90% of the focusing with no couplers. When the gradient in the structure was raised from 1 MV/l to 3 MV/l (\( l = 0.54 \) m, total length of the structure), the RF focusing strength was increased by a factor of 3.4 for Input-Particle-Energy of 1 MeV in the structure of two cells and ONE coupler.

**EMITTANCE**

The emittance obtained for structure I with two cells, NO coupler for a bunch with a uniform square distribution, at input particle energy of 0.5 MeV gradient of 1 MV/l with \( \sigma_{yin} = 1 \) mm, \( \sigma_{y'in} = 0 \) is 2.8334 \( 10^{-9} \) m - rad. Note that the input emittance is nominally zero where the output emittance \( \epsilon \) is:

\[
\epsilon_{RMS} = \sqrt{\sigma_y^2 \cdot \sigma_{y'}^2 - \sigma_{y,y'}^2} \tag{3}
\]

\( y \) and \( y' \) are the vertical displacement and angle, and \( \sigma_y \) and \( \sigma_{y'} \) are the rms beam sizes. The corresponding phase space distribution is seen in Figure 5. The emittance growth calculated for a bunch with Gaussian distribution \( \sigma_{yin} = 2 \) mm, \( \sigma_{y'in} = 0 \), at input energy of 0.5 MeV and Gradient of 1 MV/l, are summarized in Table I. The output phase space for bunch with Gaussian distribution for structures II (ONE coupler) & IV (TWO asymmetric couplers) are shown in Figure 6 and a histogram of the transverse angle, \( y'_{out} \), of the particles at the output are shown in Figure 7.

Calculation of emittance versus energy for II (ONE coupler) & IV (TWO asymmetric couplers) is shown in Figure 8. At low energy where the emittance is dominated by the RF focusing there is not much difference between one coupler and two couplers. At high energy, where the non-uniformity of transverse kick due to the coupler is the dominant contribution, we can see the effect of the cancellation in structure IV due to the second opposite coupler.
CONCLUSIONS

Calculation of the emittance (not including space charge effect) shows that the highest emittance obtained in these structures will be \(4.74 \times 10^{-8} \, \text{m} - \text{rad}\) for the first cavity after the gun (input particle energy of 0.5 MeV). At this energy the different arrangements of couplers makes very little effect. On the other hand, when the particle input energy increases the non uniformity of the transverse kick is dominant. At input energy of 2 MeV the output emittance of structure IV (TWO asymmetrical couplers) is \(\frac{1}{16}\) of the emittance of structure II (ONE coupler). See Figure 8. Having two couplers also keeps the bunch along the center of the beam pipe, where in the case of one coupler the bunch is displaced by 0.5 mm in the \(-y\) direction of the center, as seen in Figure 6. Also note, transversely displacing the beam by 1 mm in structure IV compensated for the asymmetry between the couplers, causing to a decrease in the emittance growth as seen in Table 1.

Table 1: Emittance Growth calculated by the Code 3DE for a Gaussian bunch. Input Particle Energy 0.5 MeV. The Gradient in the Cavity is 1 MV/m. (l=0.5418m)

<table>
<thead>
<tr>
<th>Structure</th>
<th>Emittance</th>
<th>Notes</th>
</tr>
</thead>
<tbody>
<tr>
<td>II</td>
<td>(4.397 \times 10^{-8})</td>
<td>Two cells ONE coupler</td>
</tr>
<tr>
<td>II off center</td>
<td>(5.354 \times 10^{-8})</td>
<td>bunch is 1 mm displaced in y</td>
</tr>
<tr>
<td>III</td>
<td>(4.6916 \times 10^{-8})</td>
<td>Two cells – TWO symmetric couplers</td>
</tr>
<tr>
<td>IV</td>
<td>(4.74 \times 10^{-8})</td>
<td>Two cells – TWO asymmetric couplers</td>
</tr>
<tr>
<td>IV off center</td>
<td>(4.34 \times 10^{-8})</td>
<td>bunch is 1 mm displaced in y</td>
</tr>
</tbody>
</table>

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REFERENCES

[2] V. Shemelin, SRF notes 021028-08