BEAM DYNAMICS IN g-2 STORAGE RING

W. Wu∗, B. Quinn, on behalf of the Muon g-2 Collaboration
University of Mississippi, University, MS, USA

Abstract

The muon anomalous magnetic moment has played an important role in constraining physics beyond the Standard Model. The Fermilab Muon g-2 Experiment has a goal to measure it to unprecedented precision: 0.14 ppm. To achieve this goal, we must understand the beam dynamics systematic effects in the muon storage ring. We will present the muon beam dynamics and discuss two specific topics here: the beam resonance which is related to the muon loss and the fast rotation analysis to determine the muon momentum distribution.

INTRODUCTION

Charged elementary particles with half-integer intrinsic spin have a magnetic dipole moment aligned with their spin:

\[ \vec{\mu} = \frac{g}{2m} \vec{s} \] (1)

where \( q \) is the electric charge and \( m \) is the mass. Dirac theory predicts the gyromagnetic ratio \( g = 2 \) [1, 2], but hyperfine structure experiments conducted in the 1940’s showed that \( g \neq 2 \) [3, 4]. Schwinger introduced radiative corrections in 1948 to resolve this discrepancy [5,6]. His calculation agreed with the experimental results and motivated the exploration of more corrections to the magnetic dipole moment anomaly defined by \( a \equiv (g - 2)/2 \).

Standard Model (SM) contributions of the muon anomaly \( a\mu \) come from the quantum electrodynamics, electroweak and quantum chromodynamics sectors. \( a\mu^{SM} \) has been calculated to a precision of around 0.42 parts per million (ppm) [7]. The most recent measurement of \( a\mu \) reported by the Brookhaven National Laboratory (BNL) E821 experiment achieved a precision of 0.54 ppm [8], which differs with SM prediction by about 3.6 standard deviation. This has motivated further theoretical and experimental investigation of \( a\mu \).

The upcoming E989 experiment at Fermi National Accelerator Laboratory (Fermilab) has restored and will use the E821 muon storage ring to measure the \( a\mu \) with a goal of 0.14 ppm [9]. To achieve this goal, we must have a detailed understanding of the muon beam dynamics in the storage ring to lower the associated systematic uncertainties. In the following sections, we will present the principles of the experiment and the muon beam dynamics in the g-2 storage ring. We will discuss beam resonance and how a fast rotation analysis is used to determine the muon momentum distribution.

∗ wwu1@go.olemiss.edu

PRINCIPLES OF THE EXPERIMENT

E989 will measure \( a\mu \) by using the spin precession resulting from the torque experienced by the muon magnetic moment when place in an external magnetic field. The rate of change of the component of spin \( \vec{s} \) parallel to the velocity \((\vec{\beta} = \vec{v}/c)\) is given by

\[ \frac{d}{dt}(\vec{\beta} \cdot \vec{s}) = -\frac{e}{mc} \vec{s}_\perp \cdot [\vec{g}(\vec{B} \times \vec{B}) + (\vec{g} \beta - \frac{1}{\beta}) \vec{E}] \] (2)

where \( \vec{\beta} \) is the unit vector in the direction of \( \vec{\beta} \) and \( \vec{s}_\perp \) is the component of \( \vec{s} \) perpendicular to the velocity [10].

The experiment uses four Electrostatic Focusing Quadrupoles (ESQ) to provide the vertical focusing of the muon beam. For a constant and purely vertical \( \vec{B} \), the anomalous precession frequency defined by the difference of the spin precession frequency and muon cyclotron rotation frequency is given by

\[ \omega_a = -\frac{q}{m}[a\mu \vec{B} - a\mu(\frac{\gamma}{\gamma + 1})(\vec{\beta} \cdot \vec{B}) \vec{B} - (a\mu - \frac{1}{\gamma^2 - 1}) \frac{\vec{\beta} \times \vec{E}}{c}] \] (3)

The second term in Eq. (3) is related to pitch correction (if \( \vec{\beta} \cdot \vec{B} \neq 0 \)) and the third term is related to electric field correction. The third term in Eq. (3) vanishes by choosing the “magic” momentum \( p_{magic} = m/\sqrt{\mu_a} \approx 3.09 \text{GeV/c} \). However, the muon beam has a momentum spread, and so an electric field correction is still needed. If the muon has a magic momentum and its velocity is also perpendicular to the magnetic field, Eq. (3) then becomes \( \omega_a = -a\mu q\vec{B}/m \).

The magnetic field is measured with nuclear magnetic resonance (NMR), where the calibrated Larmor precession frequency of a free proton is \( \omega_p = (g_p q \vec{B})/(2m_p) \). Combining \( \omega_a \) and \( \omega_p \) yields:

\[ a\mu = \frac{\omega_a/\omega_p}{\lambda - \omega_a/\omega_p} \] (4)

where \( \lambda = \mu_a/\mu_p \) is the muon-to-proton magnetic moment ratio externally determined from hyperfine splitting in muonium [11].

BEAM DYNAMICS

Polarized muons with positive charge are injected into the storage ring through the “Inflector” magnet. The “Kicker” magnet will kick muons onto the closed orbit. The muon storage ring uses ESQ for weak vertical focusing, of which the field index \( n \) with dipole magnetic field \( B \) and electric gradient \( \partial E_R/\partial R \) is defined as:

\[ n(s) = \frac{R}{\beta B} \frac{\partial E_R(s)}{\partial R}, \quad \beta = v/c \] (5)

5 Beam Dynamics and Electromagnetic Fields
D01 Beam Optics - Lattices, Correction Schemes, Transport

ISBN 978-3-95450-182-3

Copyright © 2017 CC-BY-3.0 and by the respective authors
where $s$ is the longitudinal coordinate with $\theta = s/R$ as the azimuthal angle. The muon equations of motion are then approximated by $\frac{d^2x}{dt^2} + (1 - n)x = 0$ and $\frac{d^2y}{dt^2} + ny = 0$.

There are four symmetrically located quadrupole regions, where the electric potential map of the cross section is as shown in Fig. 1. As an approximation, the horizontal and vertical tunes are given by $v_x = \sqrt{1 - n}$ and $v_y = \sqrt{n}$ for an ideal weak-focusing storage ring, where details about the lattice design of the $g$-2 storage ring and beam dynamics are discussed in Ref. [12].

![Figure 1: Electric equipotential map of quadrupoles from OPERA3D [13] with ±27.2 kV on the quadrupole plates and azimuthal average.](image)

### Beam Resonance

Deviation from the electric quadrupole and magnetic dipole fields in the storage region will result in periodic forces that perturb the muon orbits. If the periodicity of the forces falls on some resonance, the horizontal or vertical oscillations increase and may cause a loss of muons. The beam resonances can be studied from the muon equations of motion. The condition for the resonance takes the form $Lv_x + Mv_y = N$, where $L$, $M$, and $N$ are integers.

The field index $n$ increases for higher quadrupole voltage. The operating quadrupole voltage should be chosen to avoid the beam resonances, where a weak-focusing storage ring operates using the approximate condition $v_x^2 + v_y^2 = 1$. Therefore, the operating point should not intersect any of the resonance lines shown in Fig. 2.

### Fast Rotation Analysis

Muons are injected into the storage ring as a bunch and not all of them sit on the magic momentum. Eq. (3) shows an electric term to $\omega_d$ and the muon momentum distribution is required to analyze and estimate the electric field correction to systematic errors.

The so-called Fast Rotation Analysis (FRA) is a technique for measuring the muon radial momentum distribution that uses the time evolution of the bunch structure. Muons in a bunch with momentum $p > p_{\text{magic}}$ will naturally assume higher orbits ($r > r_{\text{magic}}$), which take longer to complete one cyclotron revolution. After some time (around 100 ~ 1000 revolutions), a muon with lower momentum will lap a muon with high momentum in the same bunch and the bunch will stretch. The stored muon momentum distribution as well as radial distribution can be determined by analyzing the debunching beam.

There are two techniques used to perform a fast rotation analysis: minimized $\chi^2$ method [14] and Fourier Transform method [15]. Both methods were used in the BNL experiment [8], although only the minimized $\chi^2$ method is described here.

The measurement of $\omega_d$ uses the decay position time spectra (e.g., Fig. 3) seen by the electromagnetic calorimeters sitting on the inside radius of the storage ring. Because of the parity violation in the weak decay $\mu^+ \to e^+ + \nu_e + \bar{\nu}_\mu$, there is a correlation between the muon spin and the direction of the high-energy daughter positron. After applying a high energy cut around 1.8 GeV, the number of daughter positrons can be modulated by the frequency $\omega_d$ (i.e., $N(t) = N_0\exp(-t/\tau_\mu)[1 - A\cos(\omega_d t + \phi)]$ [8].

The expected decay positron count for the $j$th time bin in the minimized $\chi^2$ method is given by

$$C_j = \sum_{i=1}^{I} f_i \beta_{ij}$$

where $i$ is the radial momentum bin index, $I$ is the total number of radial bins, $f_j$ is the fraction of muons in the $i$th bin, and $\beta_{ij}$ are the muon temporal-radial evolution coefficients. The $f_j$ are the quantities being solved for, and the $\beta_{ij}$ are muon beam bunch geometrical factors that can be calculated separately.
ACKNOWLEDGMENT

The author thanks the organizers of this conference and his many colleagues on the Fermilab E989 g-2 experiment, in particular W. Morse for useful discussions and J. D. Crnkovic for help with editing the manuscript. The author is supported by the U.S. Department of Energy Office of Science, Office of High Energy Physics, award DE-SC0012391, and a Universities Research Association Visiting Scholar award. His travel to IPAC’17 is supported by the Division of Physics of the U.S. National Science Foundation (Accelerator Science Program) and the Division of Beam Physics of the American Physical Society.

REFERENCES